

QED in 2 spatial dimensions, A study of confinement with quantum computers



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Work in collaboration with Arianna Crippa and Karl Jansen (DESY) → [arxiv:2411.05628](https://arxiv.org/abs/2411.05628)

Enrico Rinaldi – 2024/11/13 – Lead R&D Scientist at

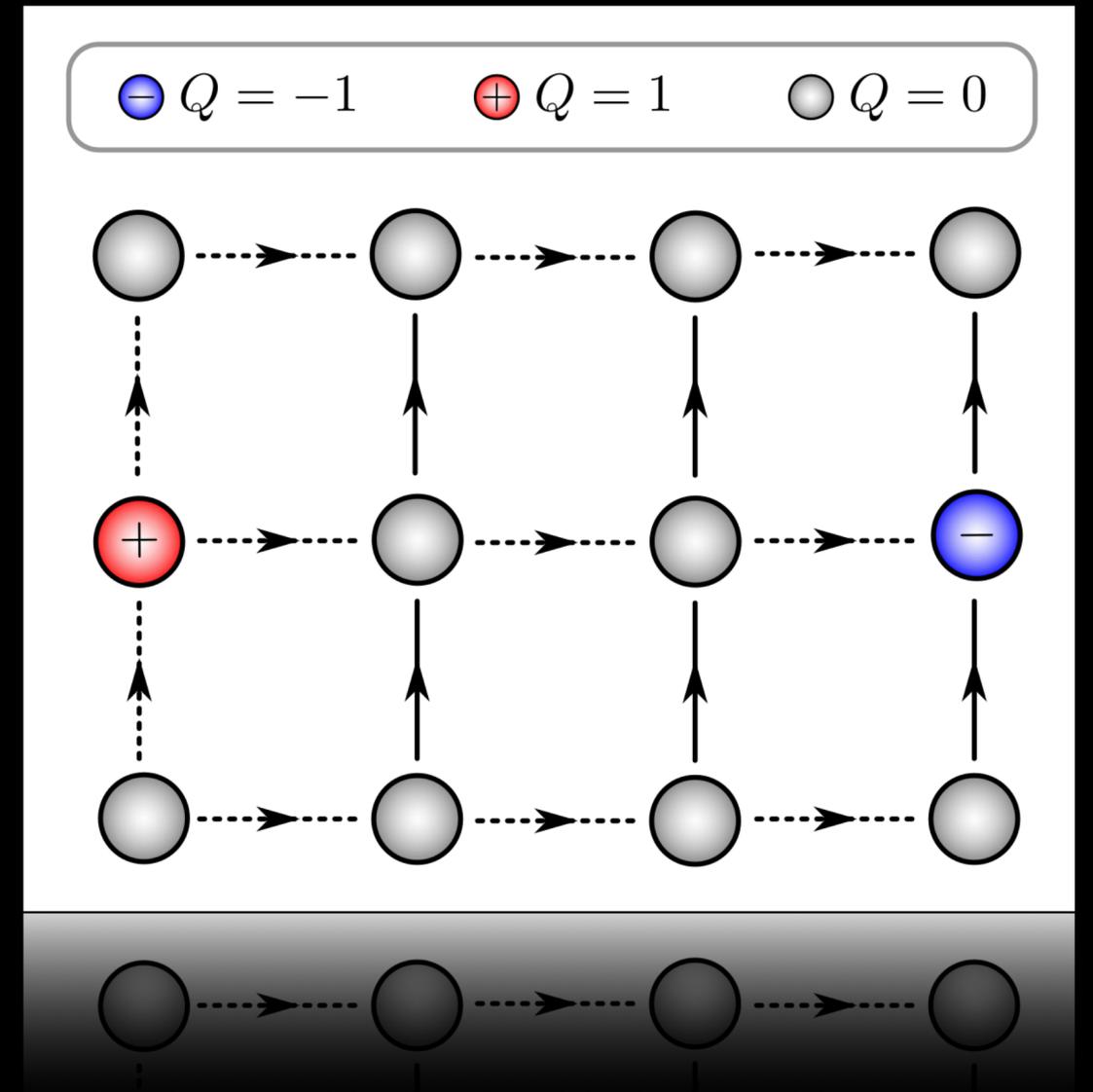


QED in 2 spatial dimensions ...

Hamiltonian Lattice QED

In 2 spatial dimensions

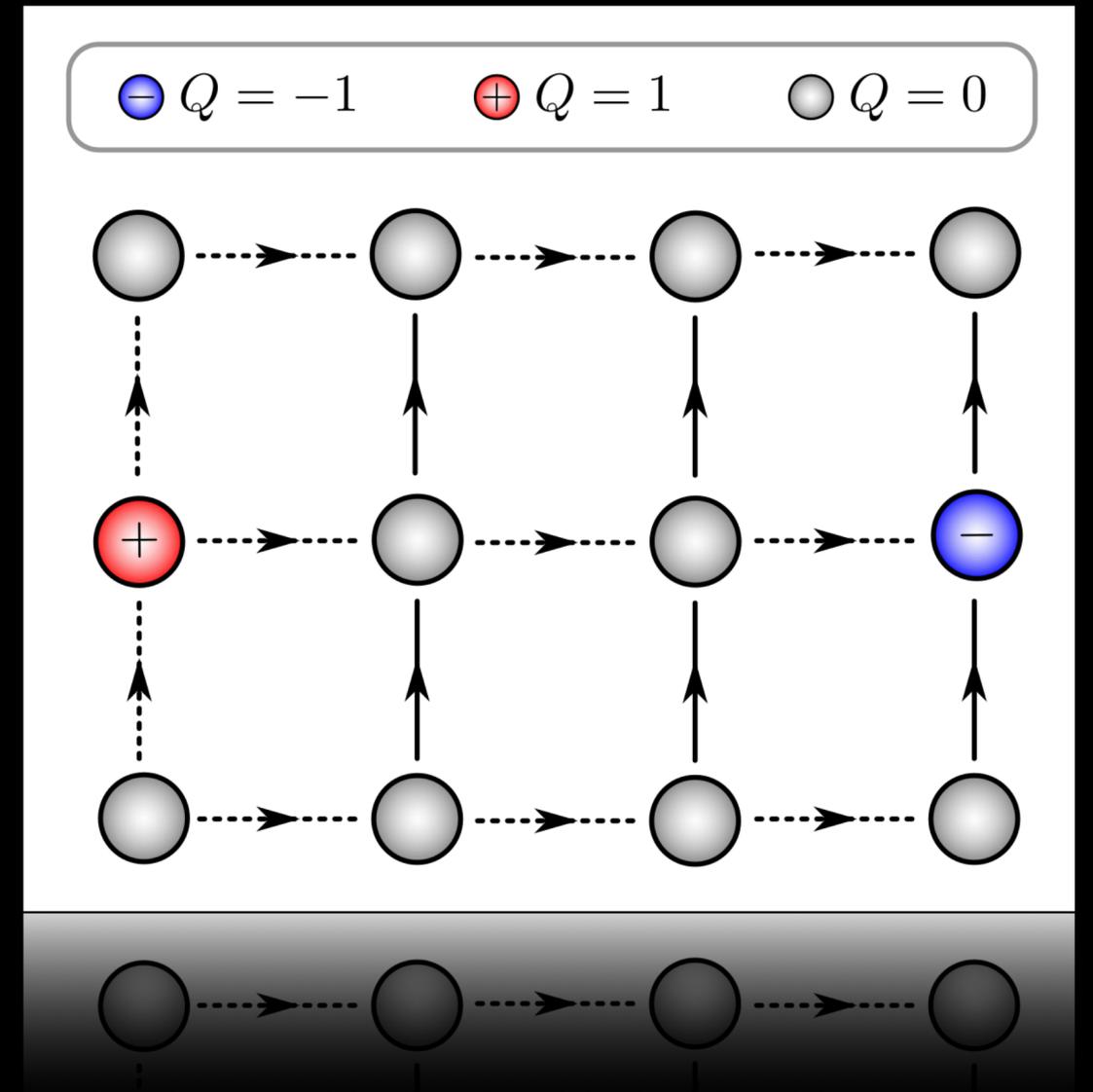
- We use the **Kogut–Susskind Hamiltonian** formalism of lattice gauge theory. **Time is continuous.**
- The Hilbert space is defined as the tensor product of the **local Hilbert spaces of each degree of freedom on the lattice**
- A state is a superposition of amplitudes for each possible **configuration of degrees of freedom on the lattice**



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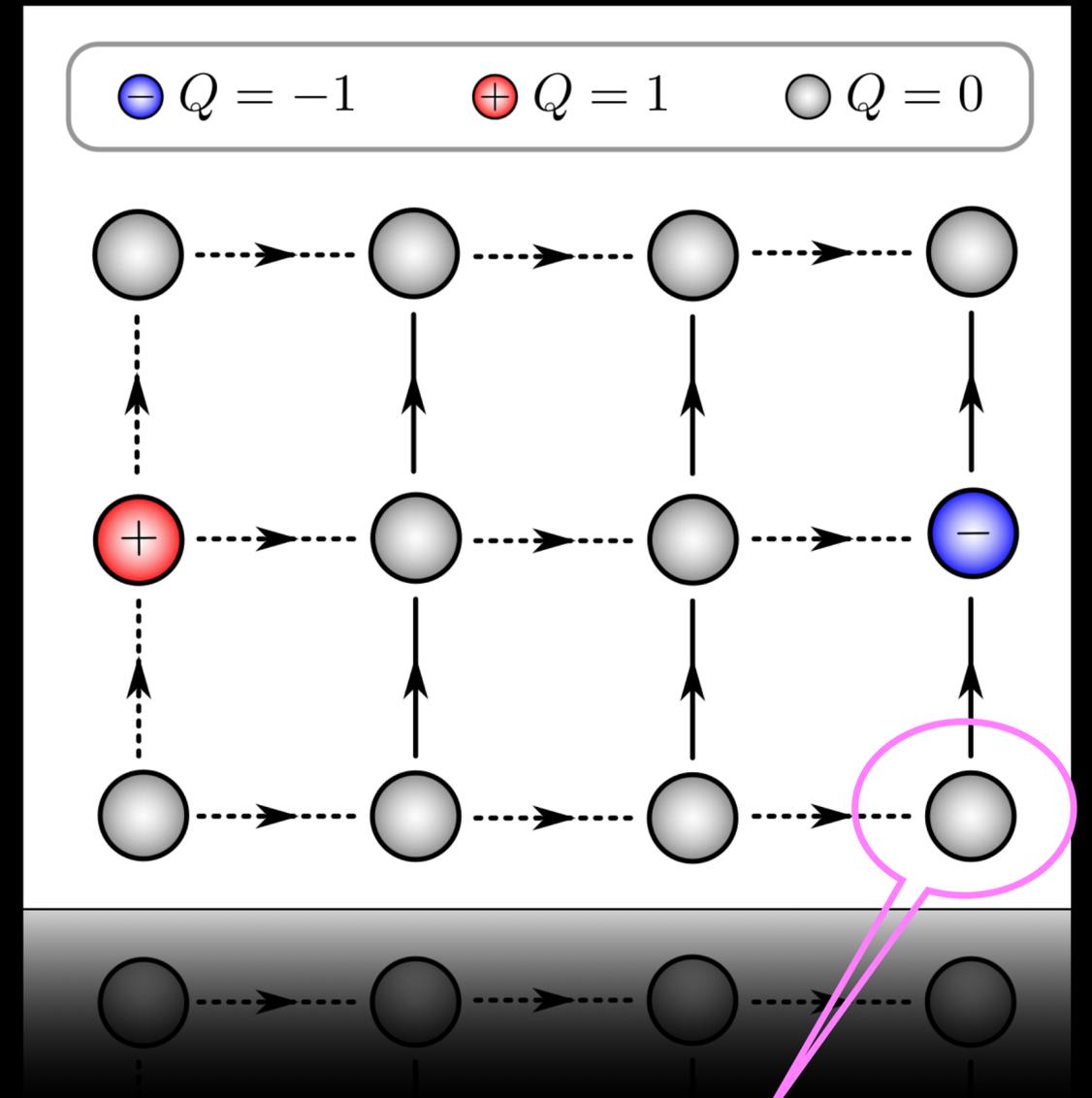


- Site: fermion – Electron
- Link: gauge – Electric field

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Fermion

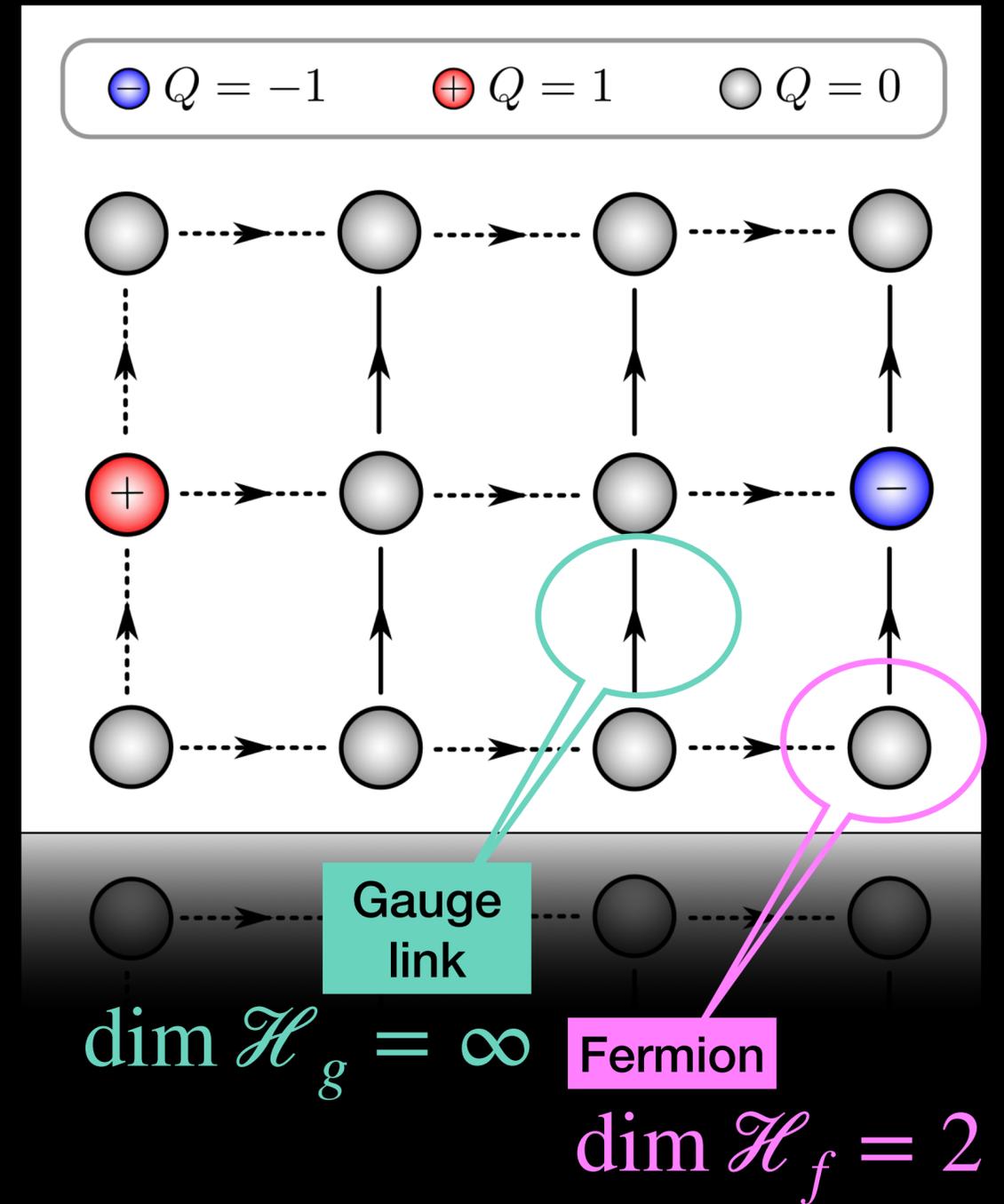
$$\dim \mathcal{H}_f = 2$$

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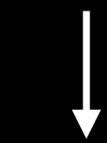


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QED on qubits

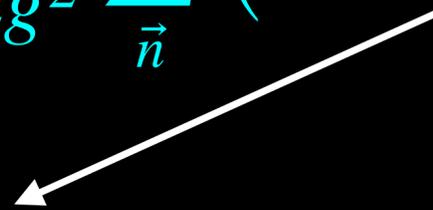
Electric and Magnetic terms

$$\hat{H}_E = \frac{g^2}{2} \sum_{\vec{n}} \left(\hat{E}_{\vec{n},x}^2 + \hat{E}_{\vec{n},y}^2 \right)$$

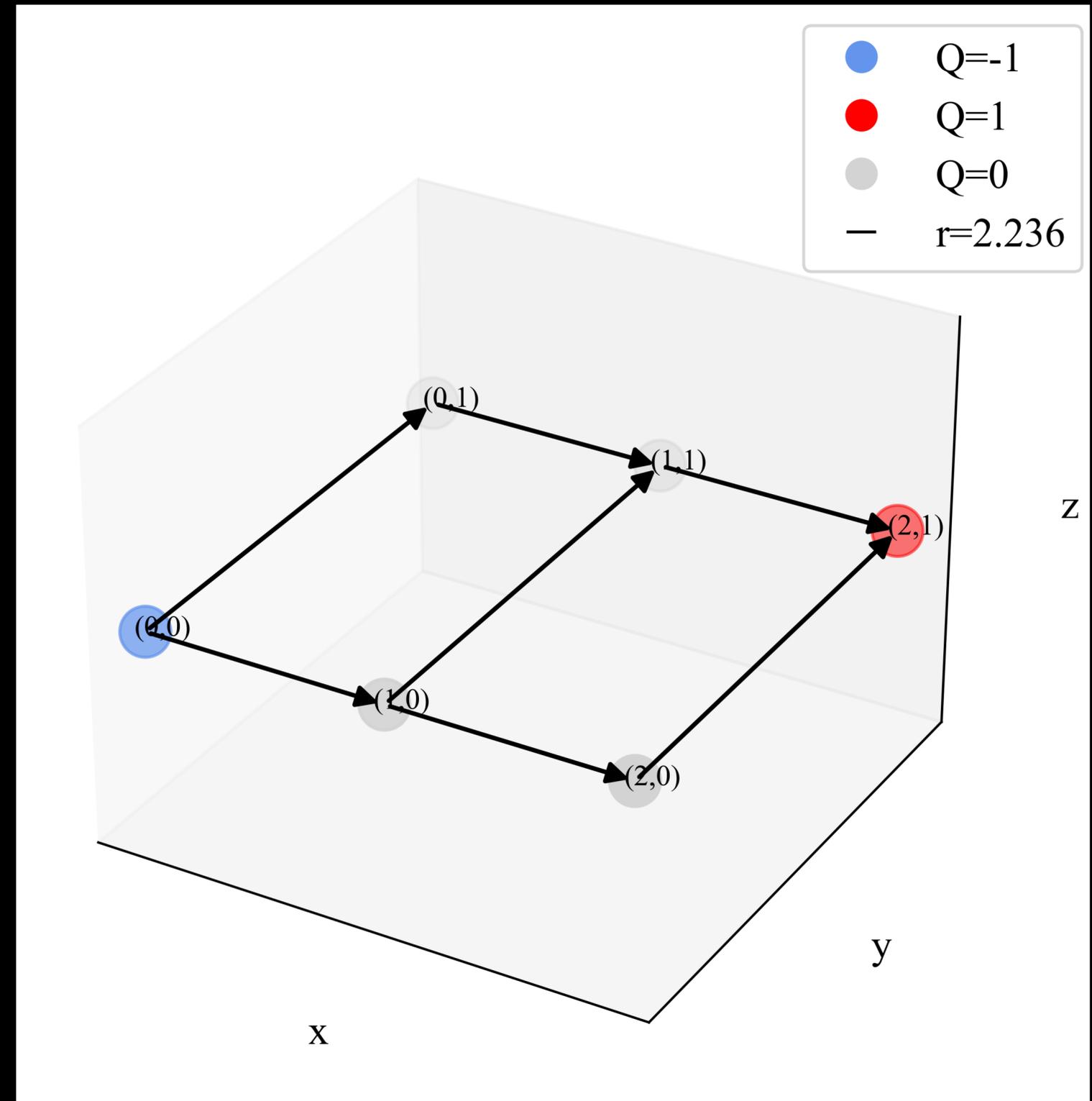


$$\hat{E}_{\vec{n},\mu} |e_{\vec{n}}\rangle = e_{\vec{n}} |e_{\vec{n}}\rangle$$

$$\hat{H}_B = -\frac{1}{2g^2} \sum_{\vec{n}} \left(\hat{U}_{\vec{n},x} \hat{U}_{\vec{n}+x,y} \hat{U}_{\vec{n}+y,x}^\dagger \hat{U}_{\vec{n},y}^\dagger + \dots \right)$$



$$\hat{U}_{\vec{n},\mu} |e_{\vec{n}}\rangle = |e_{\vec{n}} - 1\rangle$$



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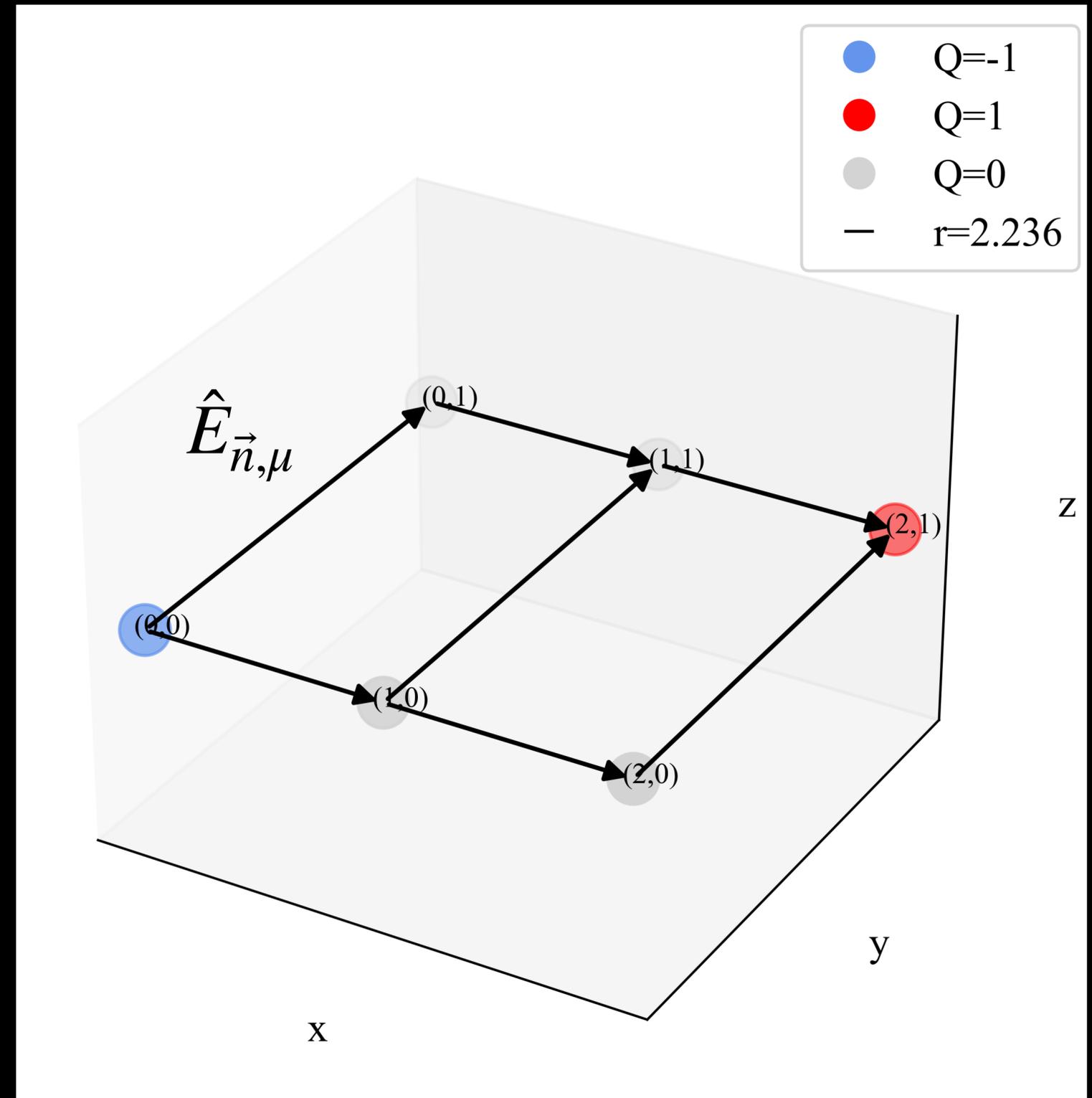
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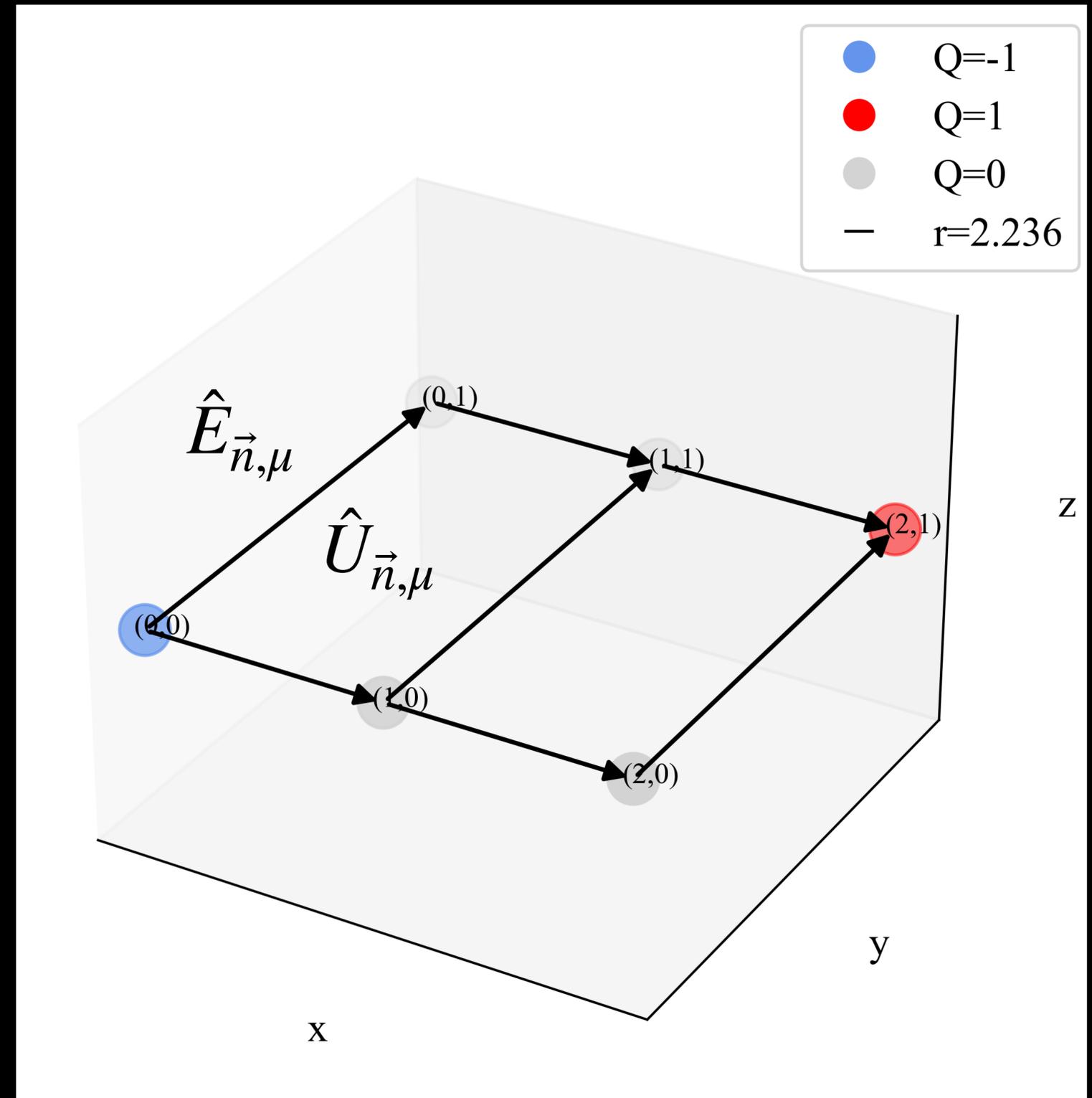
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Magnetic term

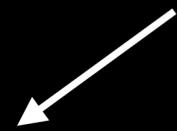
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QED on qubits

Discretize: $U(1) \rightarrow \mathbb{Z}_{2l+1}$

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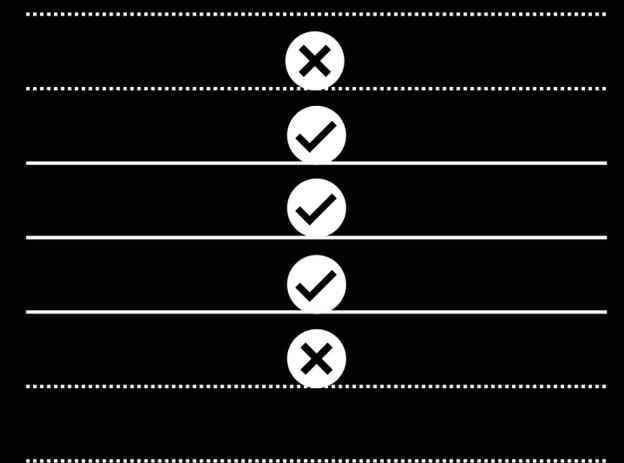
$$|e_{\vec{n}}\rangle = |-l_{\vec{n}}\rangle, |-l+1_{\vec{n}}\rangle, \dots, |-1_{\vec{n}}\rangle, |0_{\vec{n}}\rangle, |+1_{\vec{n}}\rangle, |l-1_{\vec{n}}\rangle, |l_{\vec{n}}\rangle$$

$$l = 1$$

$$i = +1$$

$$i = 0$$

$$i = -1$$



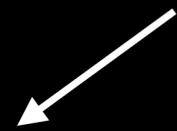
Encoding to qubits:

$l = 1$ We need 2 qubits to represent 4 states. 1 state is “unphysical”

QED on qubits

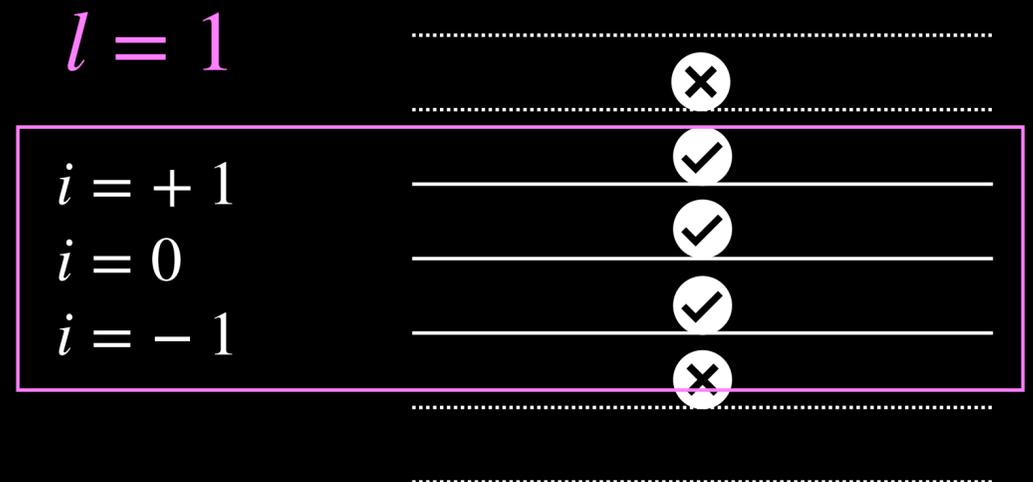
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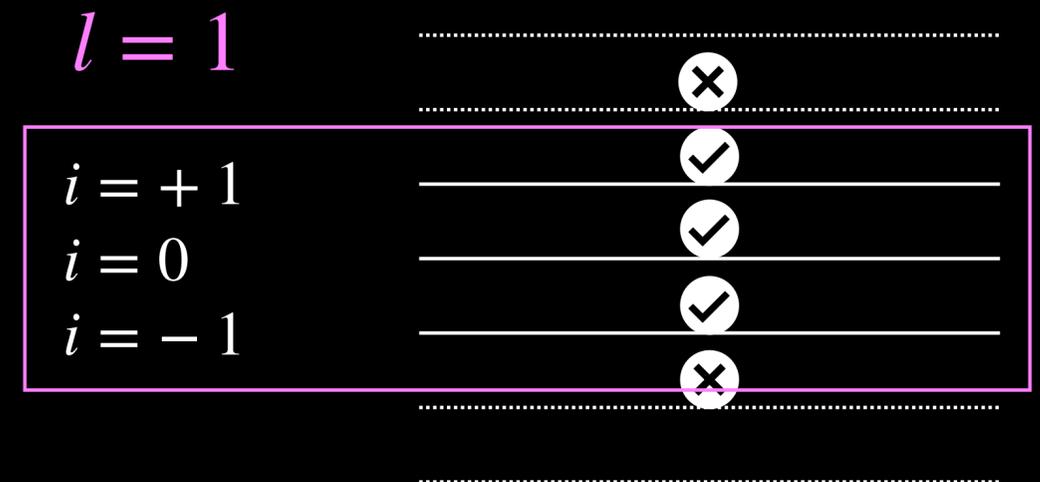
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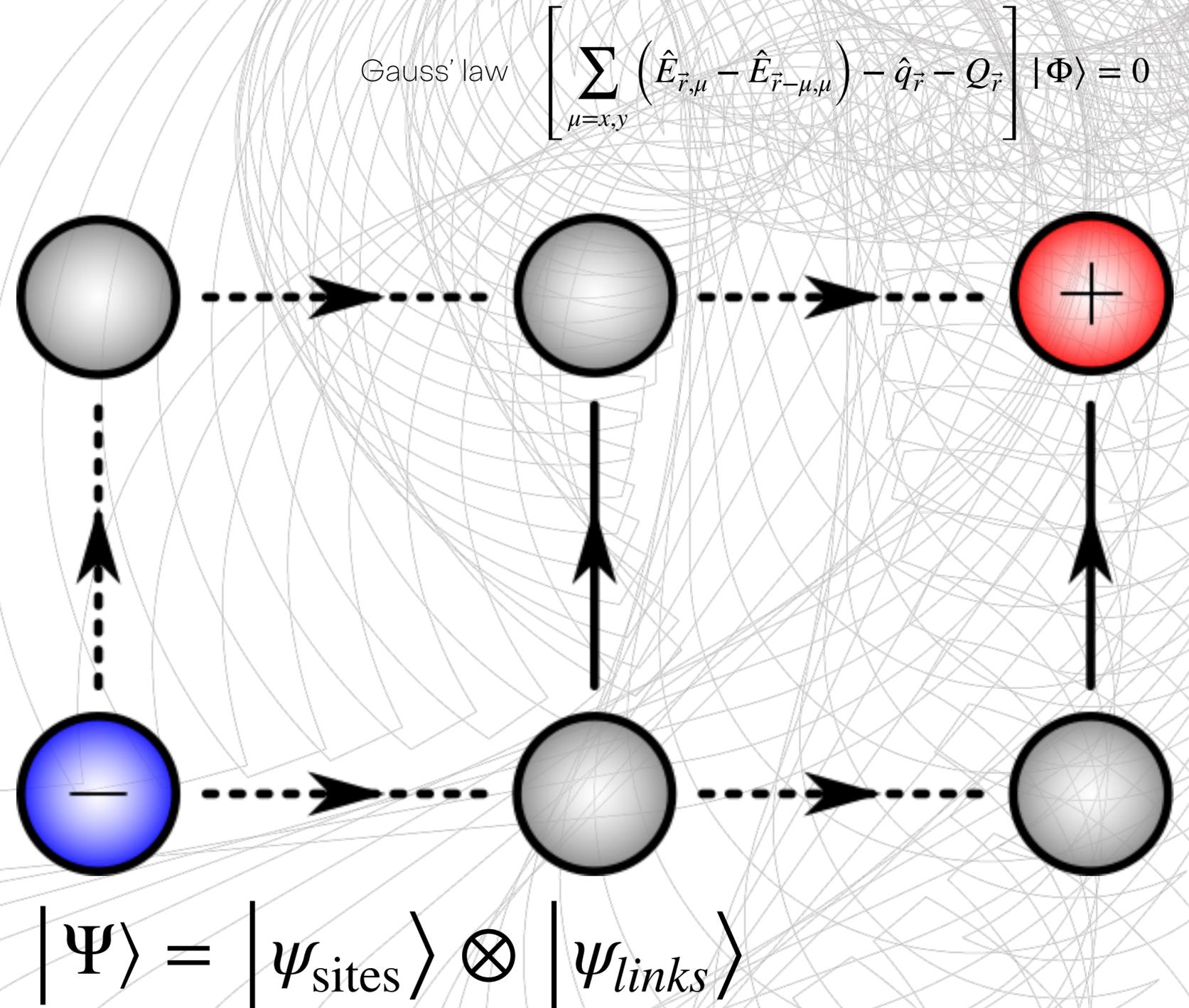


Gray Encoding

$ i\rangle_{\text{phys}}$	$ i\rangle$
$ -1\rangle_{\text{phys}}$	$ 00\rangle$
$ 0\rangle_{\text{phys}}$	$ 01\rangle$
$ +1\rangle_{\text{phys}}$	$ 11\rangle$
Unphysical	$ 10\rangle$

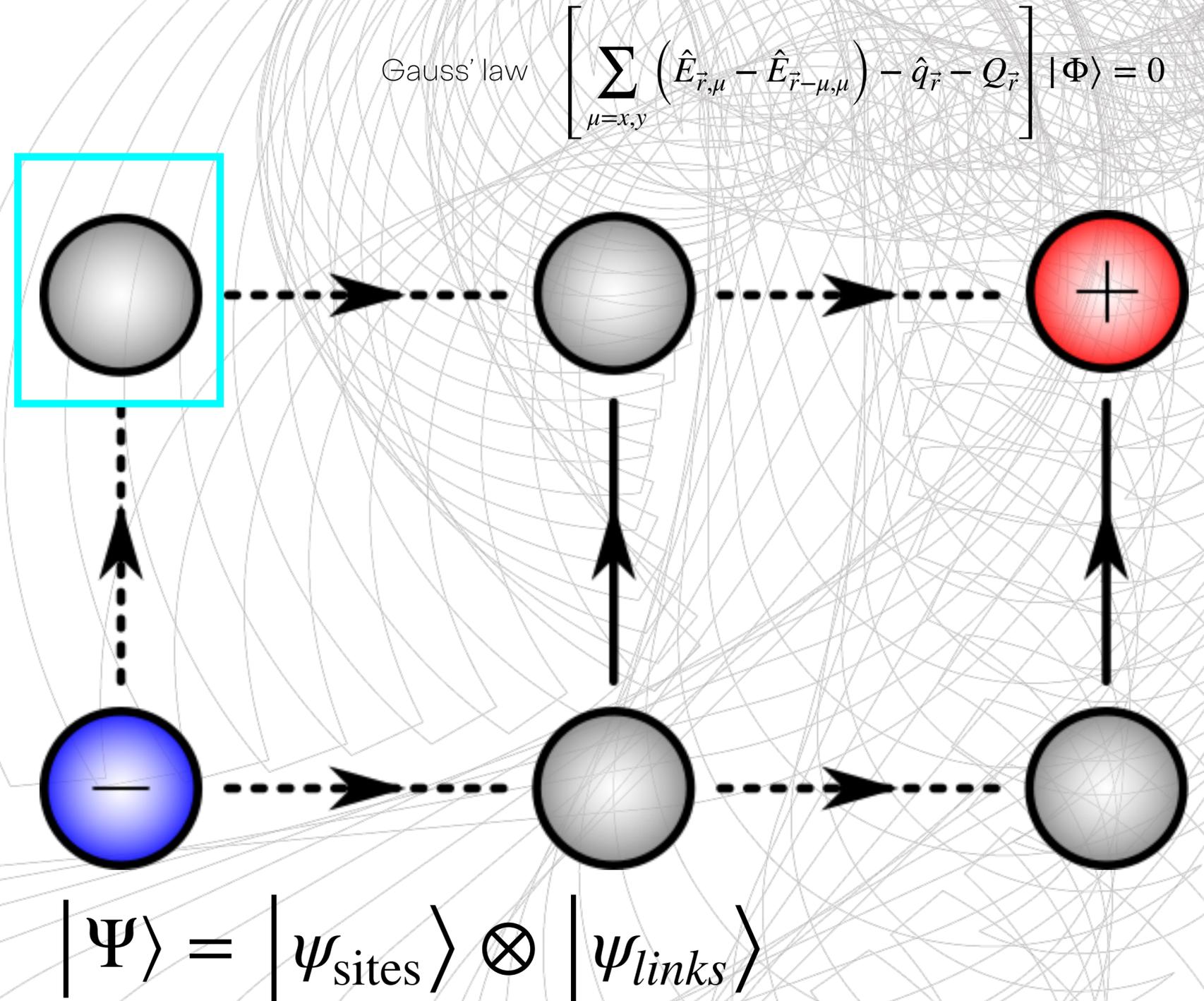
Example on a 3×2 lattice

- Using Gauss' law we reduce the number of dynamical degrees of freedom: **6 sites**, **2 links**
- We use **1 qubit for each site**
- We use **2 qubits for each link**
- Any state of this lattice QED theory is defined on **10 qubits**
- This classically requires manipulating a vector with $2^{10} = 1024$ complex components



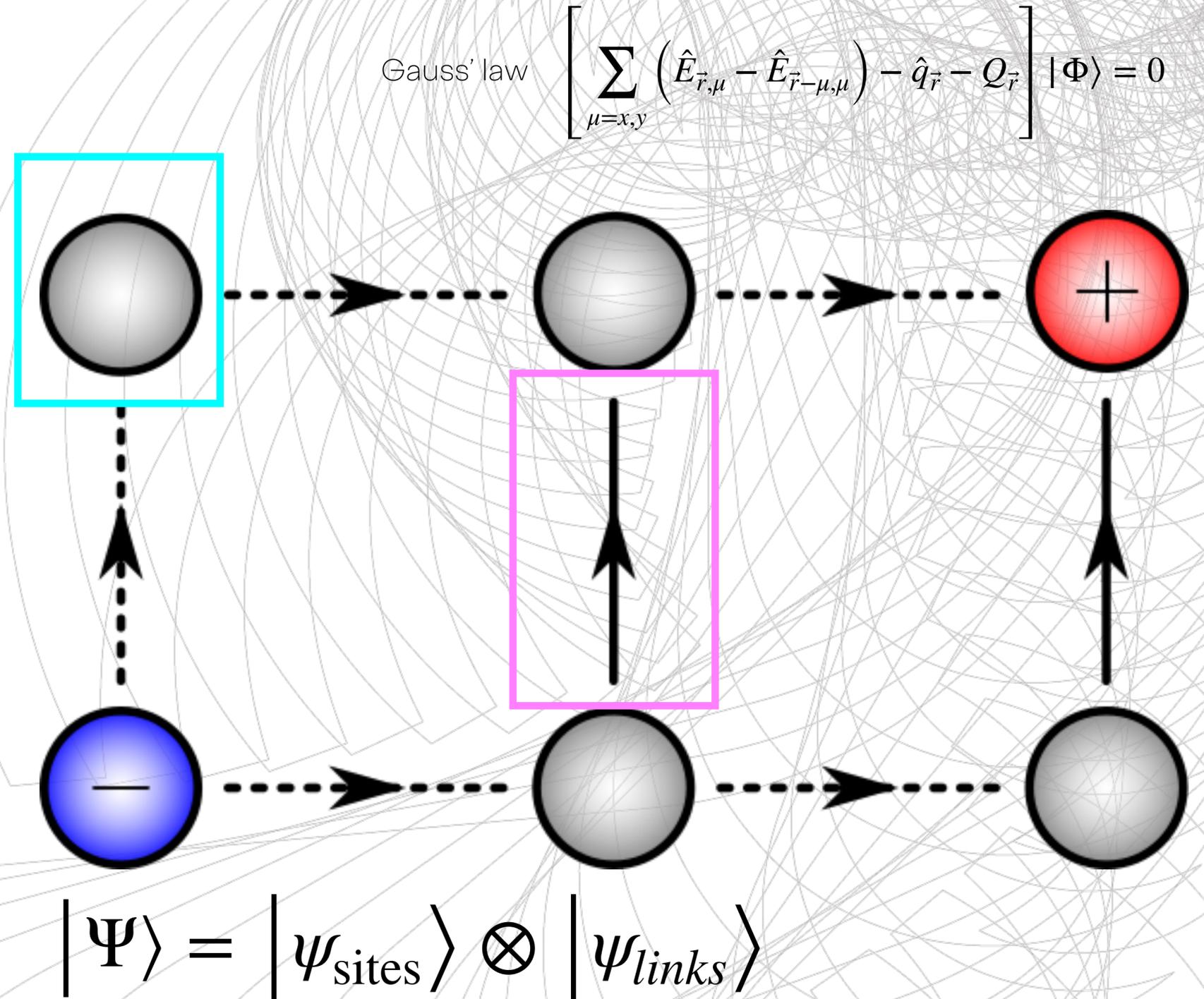
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A study of confinement ...



QUANTINUUM

The potential energy

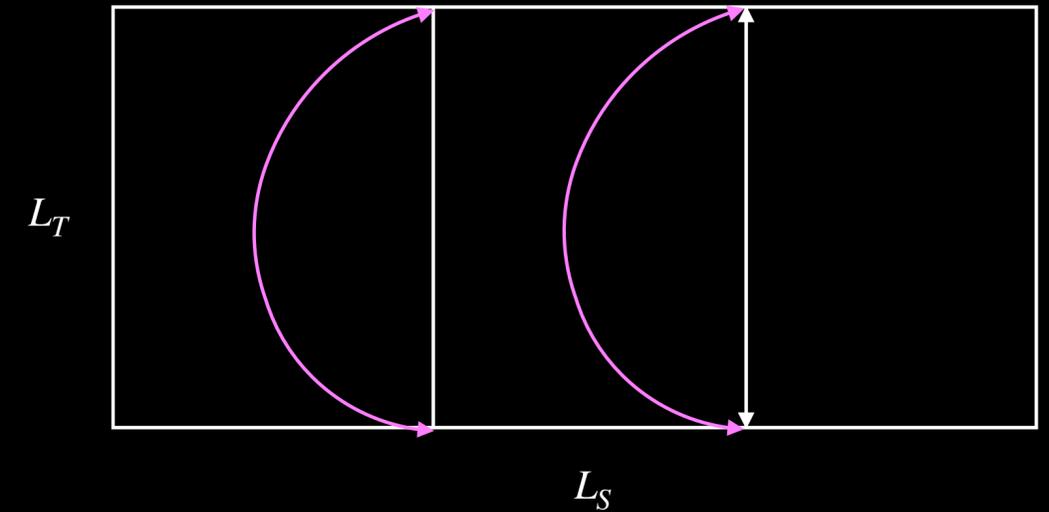
Between two static charges

- In Path Integral Monte Carlo we extract this by computing **Wilson Loops** of various lengths
- In quantum computing we have direct access to the **Hamiltonian and the states of the system!**
- The potential **$V(r)$** is the energy of the **ground state** with 2 opposite static charges
- By **changing the distance** between static charges we can study the force between them

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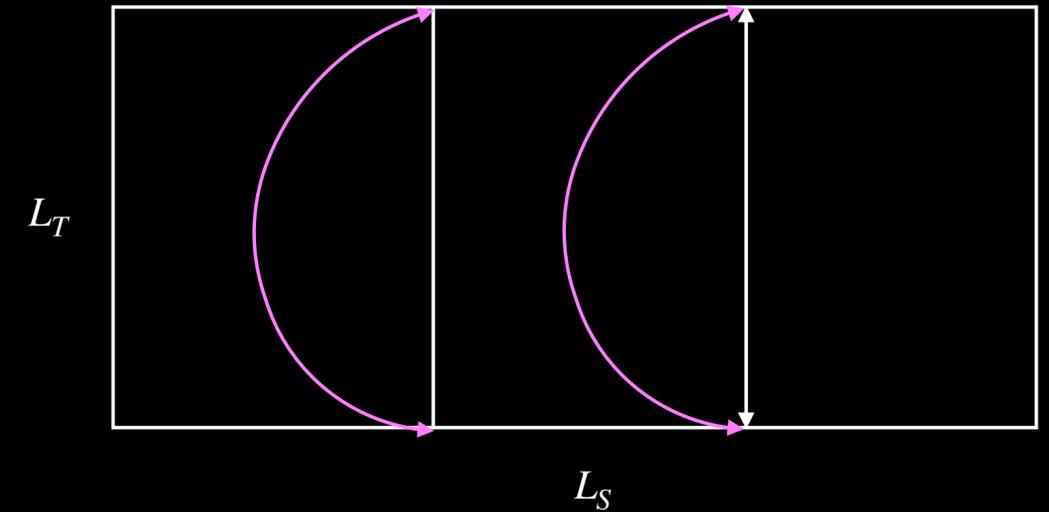
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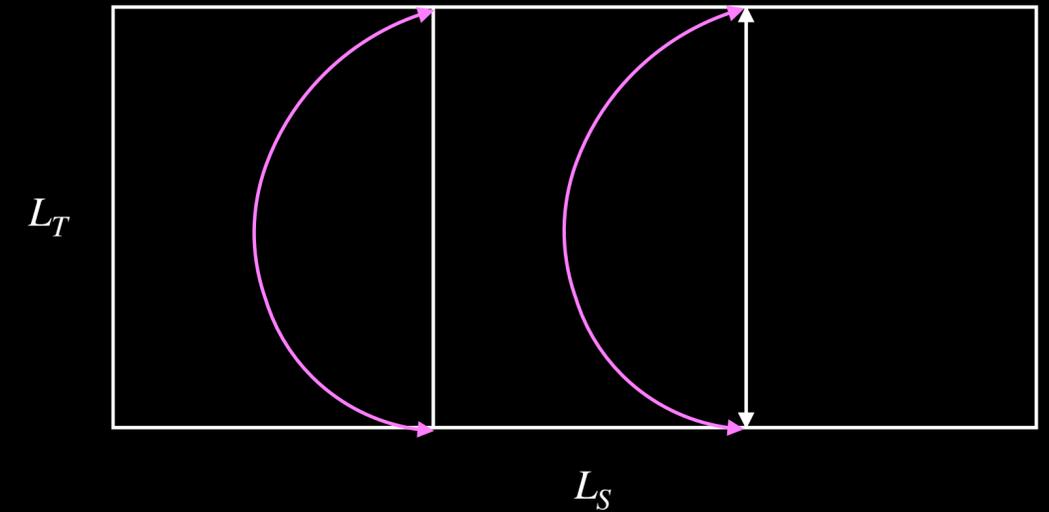


$$V(r) = \langle \Psi_0(r) | \hat{H} | \Psi_0(r) \rangle$$

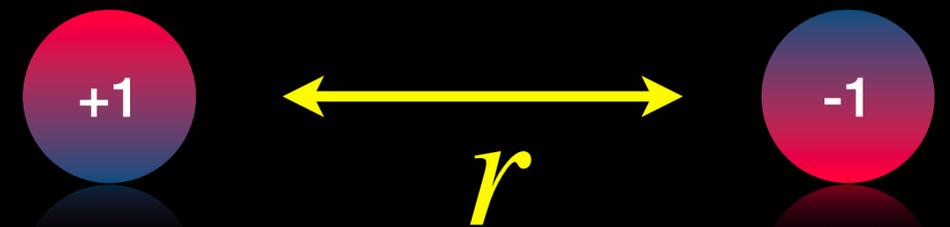
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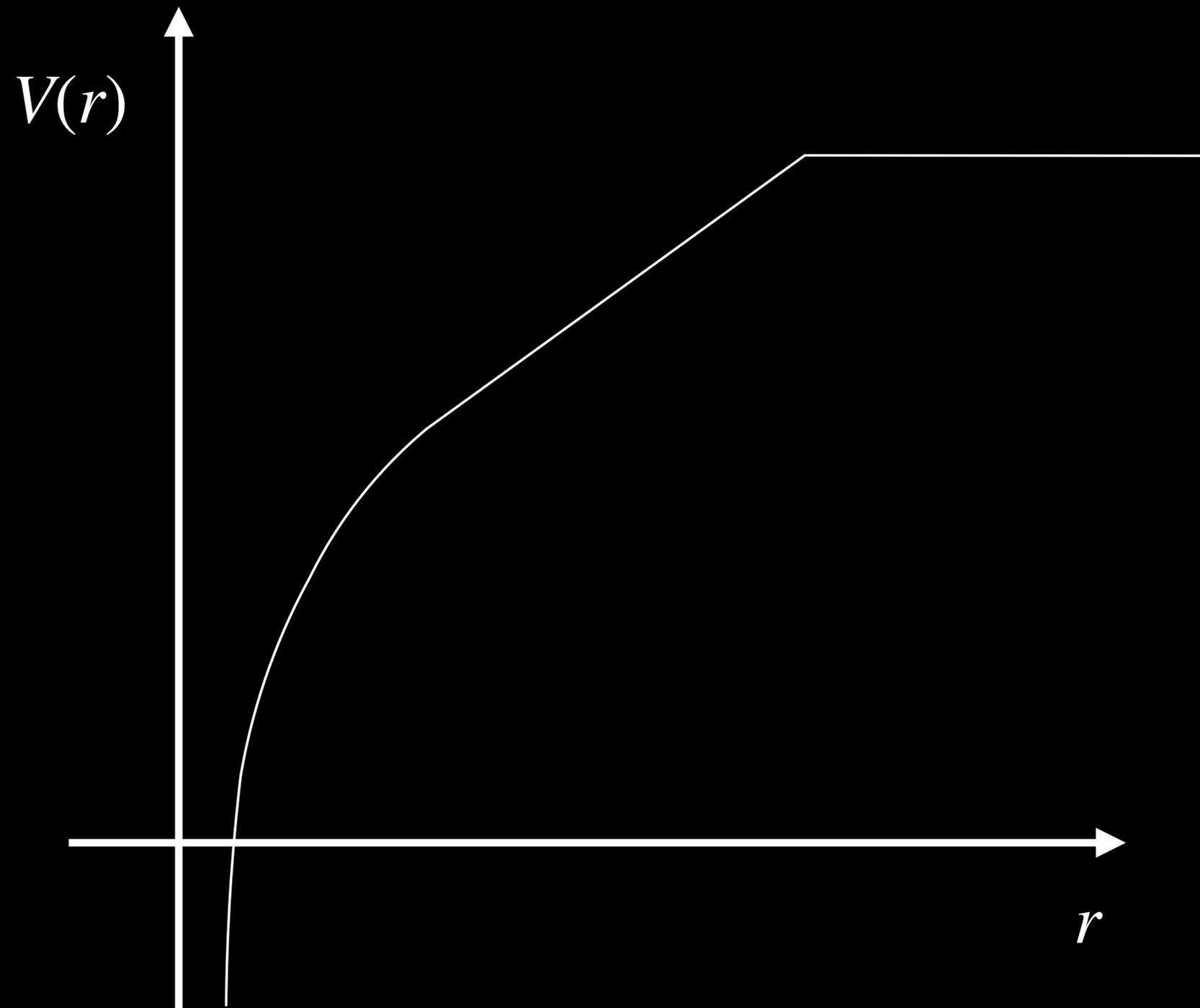
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Static Potential

In (2+1)D QED

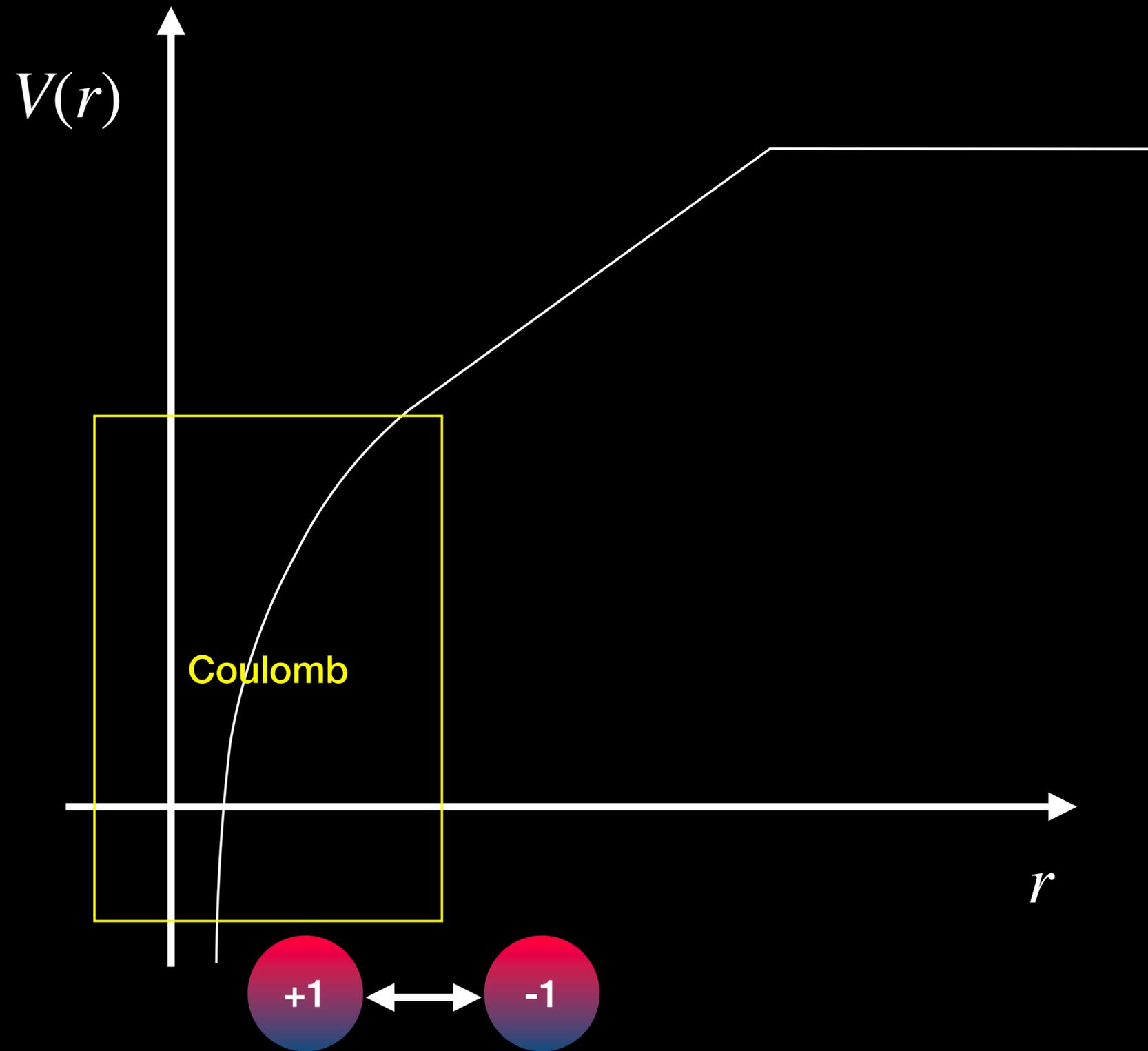
- $V(r) = V_0 + \alpha \log r + \sigma r$
- On the lattice we can change r by changing the lattice spacing a
- The lattice spacing a depends non-perturbatively on the coupling constant g
 - $V(r) \rightarrow V(g)$



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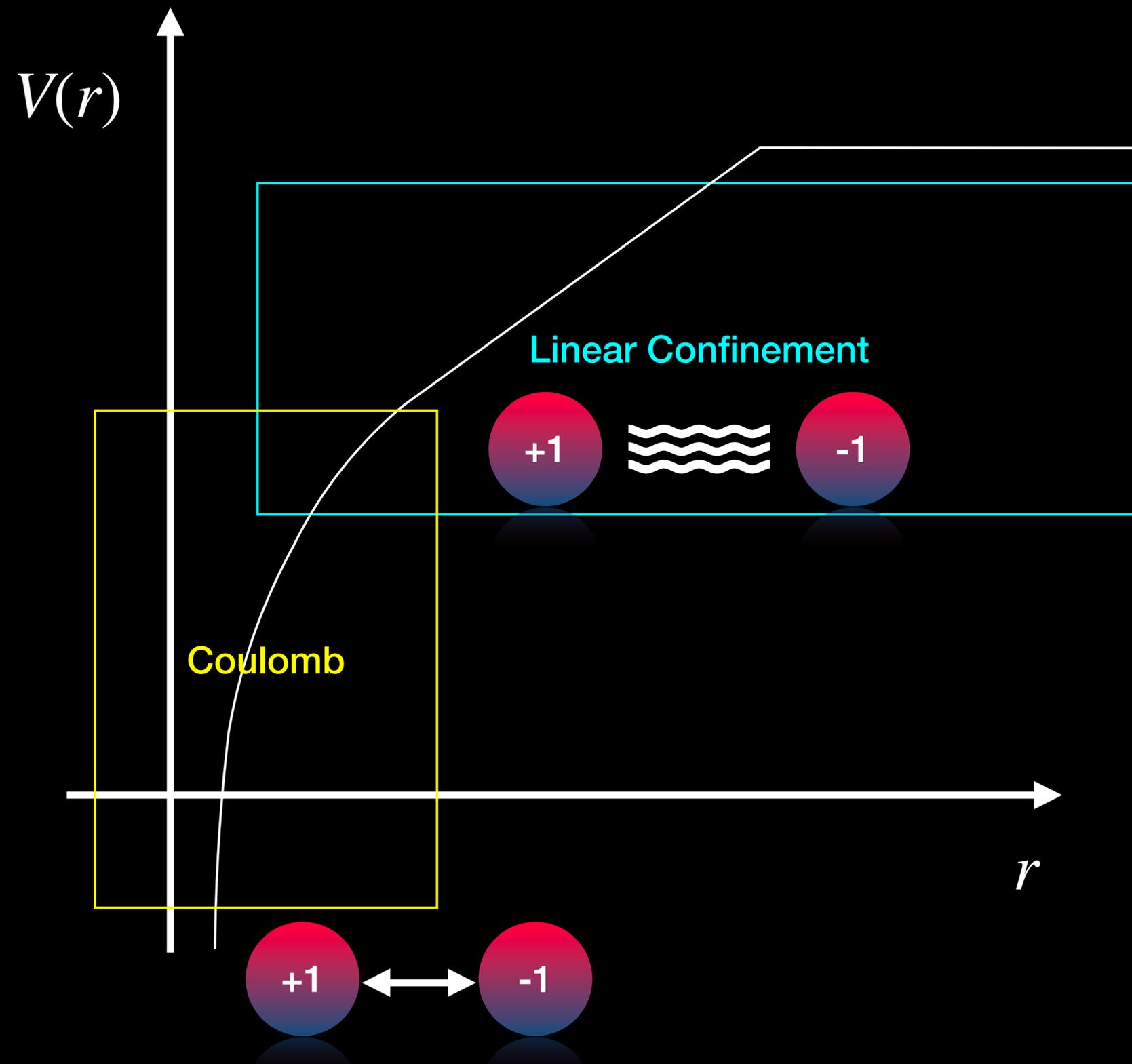
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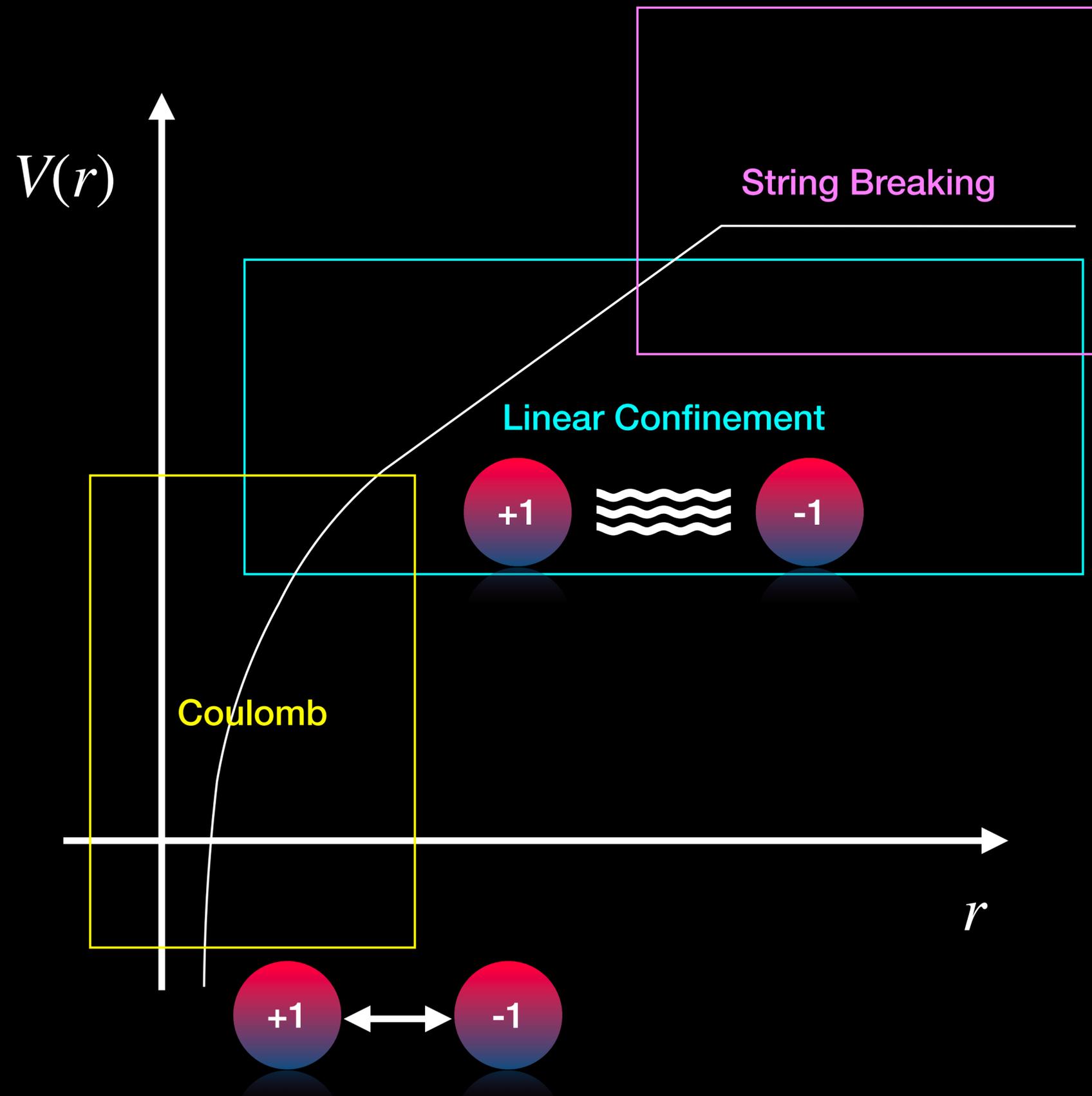
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... with quantum computers



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Quantum computing

For Lattice Gauge Theory

- A new paradigm for scientific computing
(many beautiful reviews online)
- Quantum algorithms work by manipulating quantum states in Hilbert space and measuring them
- Represent the full wavefunction of a quantum many-body system
- Can do unitary time evolution of such wavefunction
- Digital quantum computing implements unitary operators as sequence of gates

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$$U = e^{-iHt} \equiv \prod_i U_i$$

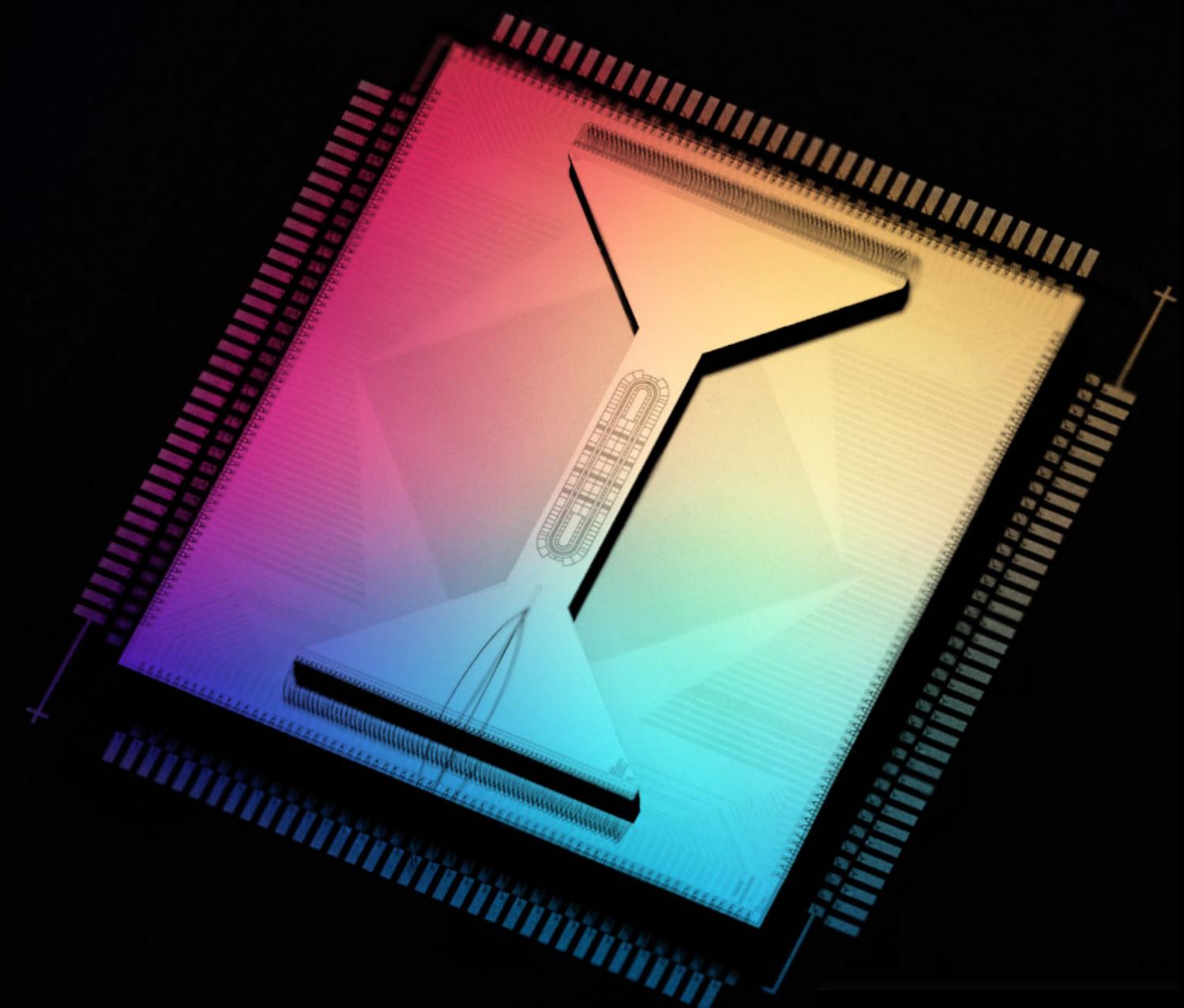
Quantinuum

H-series Quantum Hardware

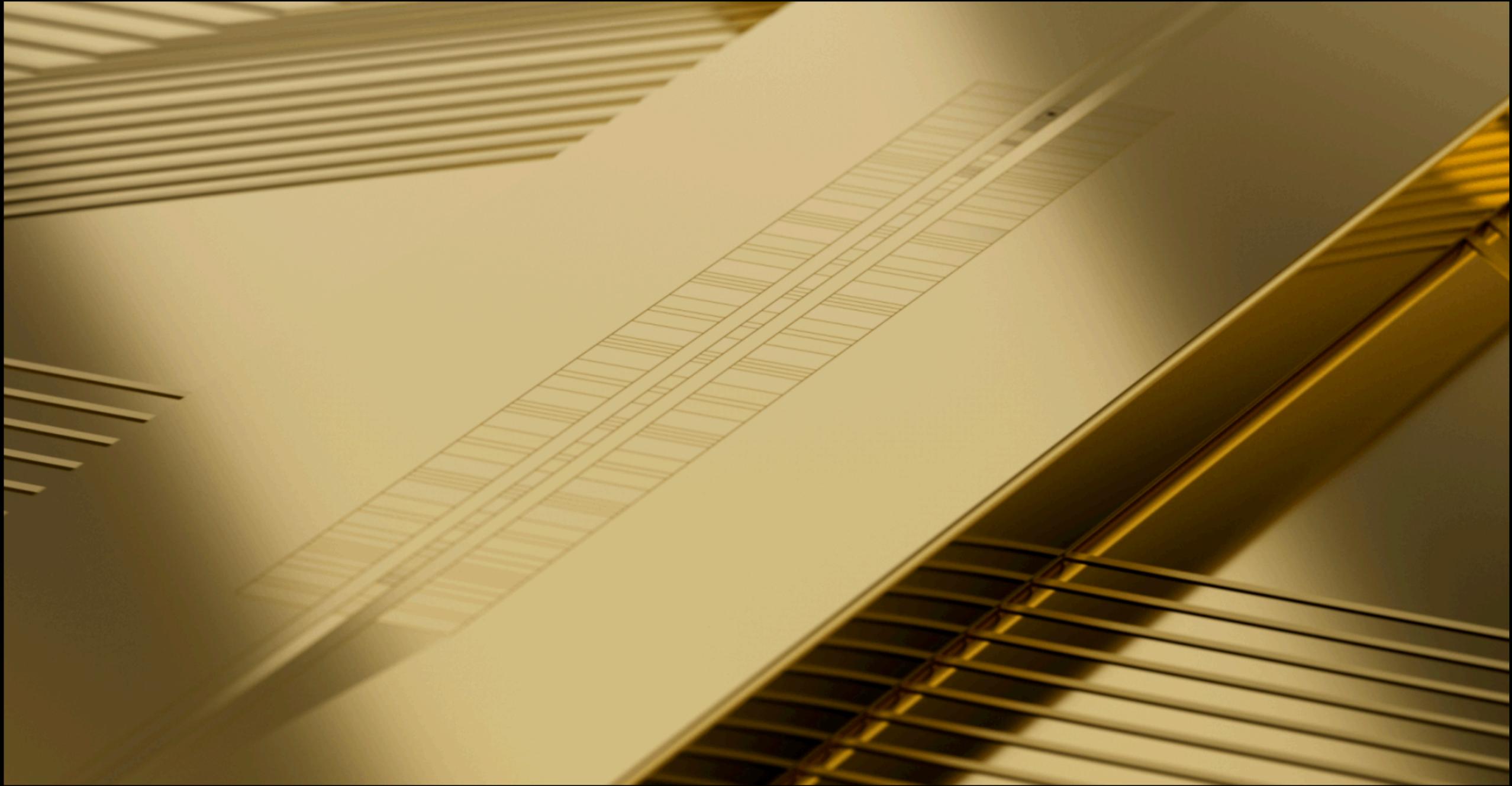
Most benchmarked quantum computer

Lowest-error commercial quantum device

20 and 56 qubits on trapped ions



H1-1



20

FULLY-CONNECTED QUBITS

1,048,576 (2^{20})

QUANTUM VOLUME

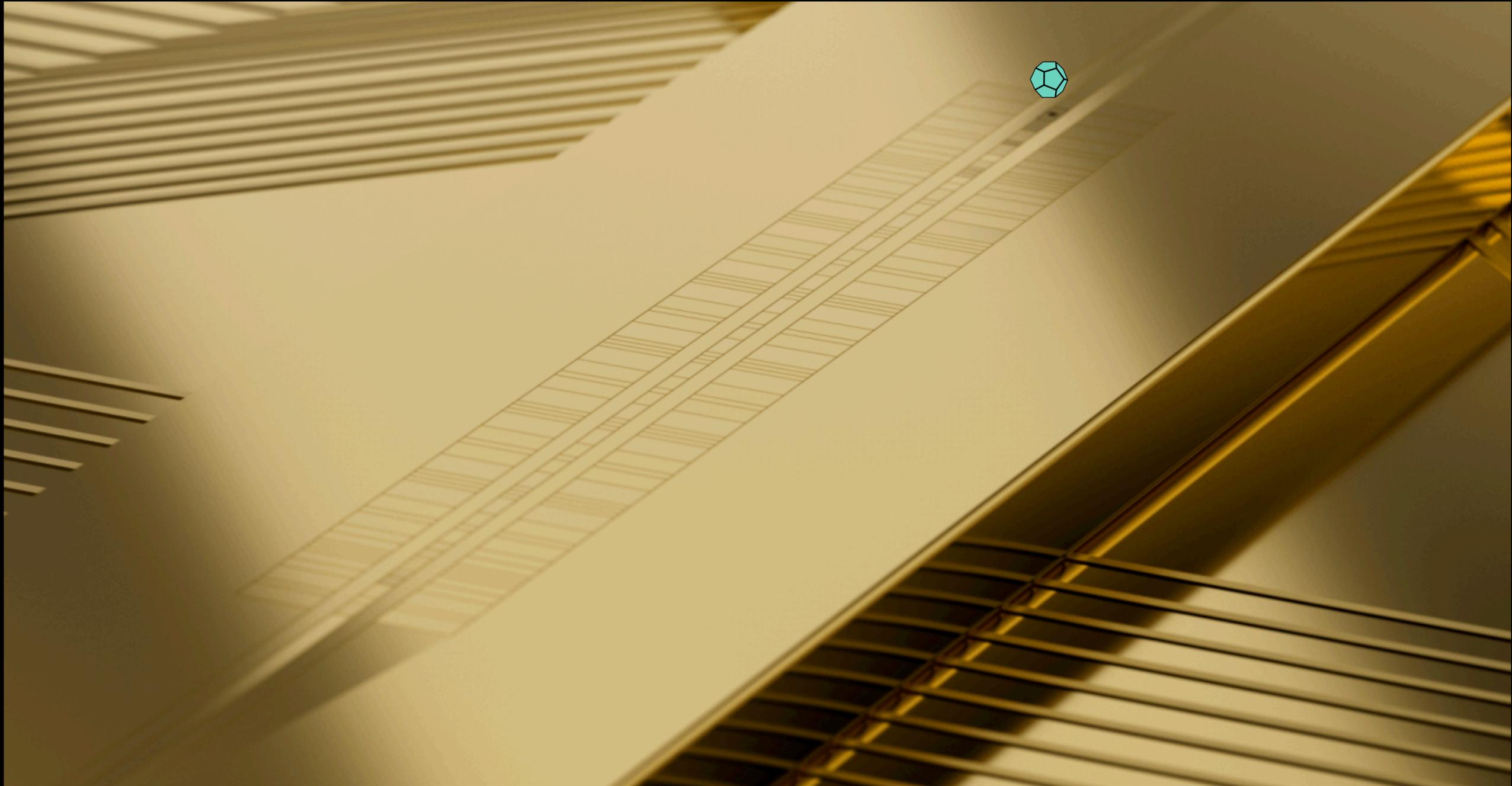
99.998%

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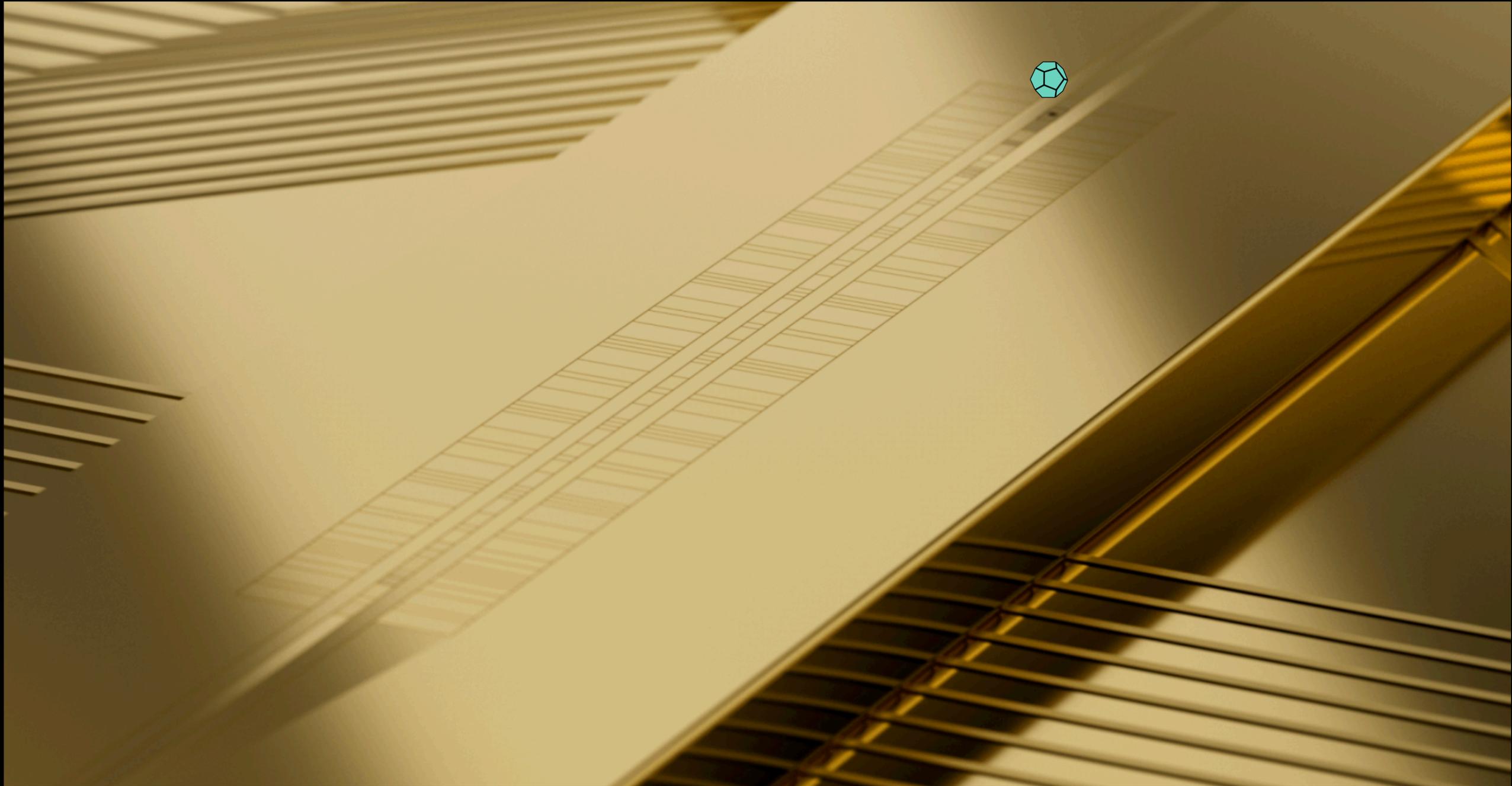
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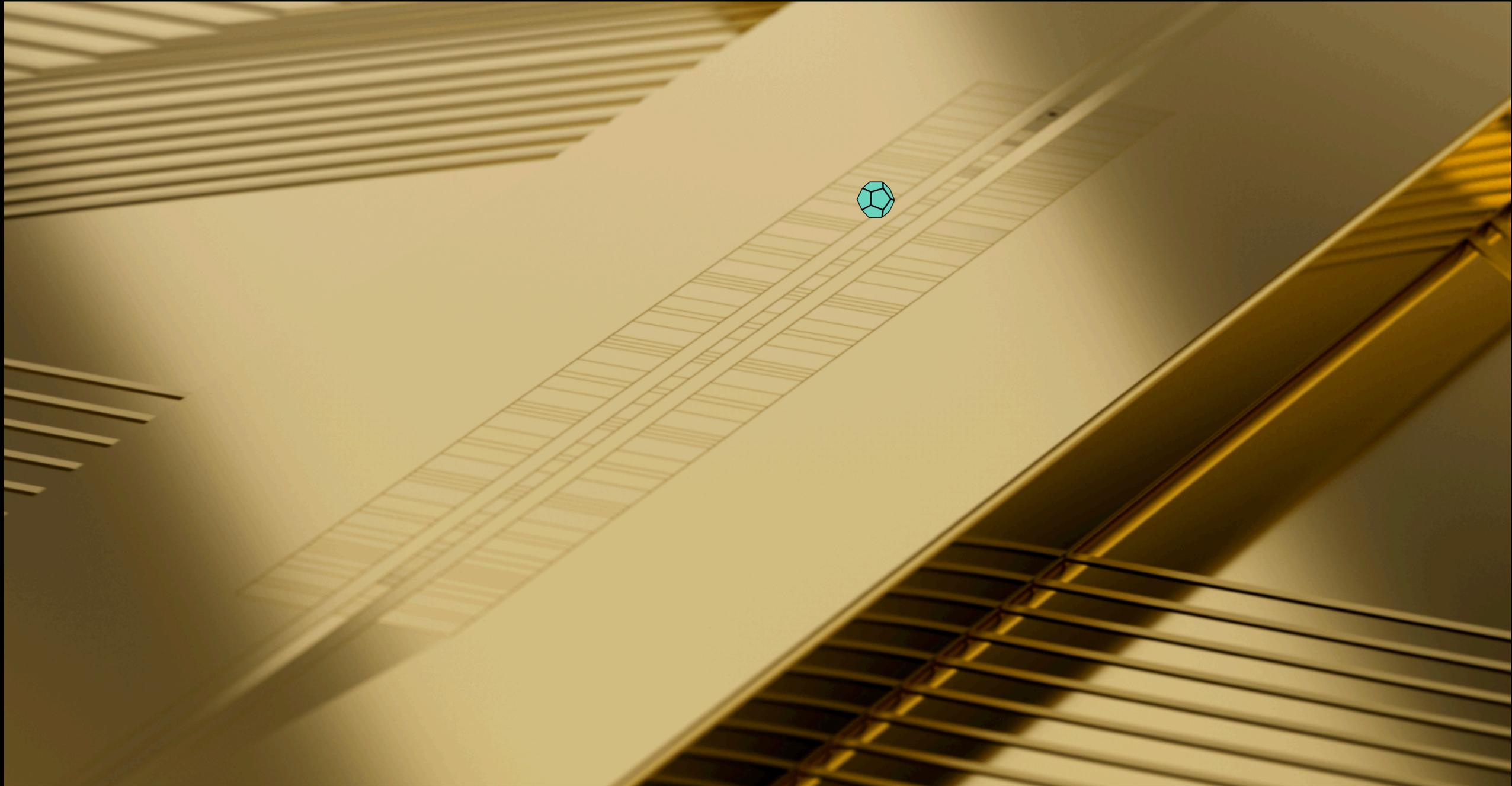
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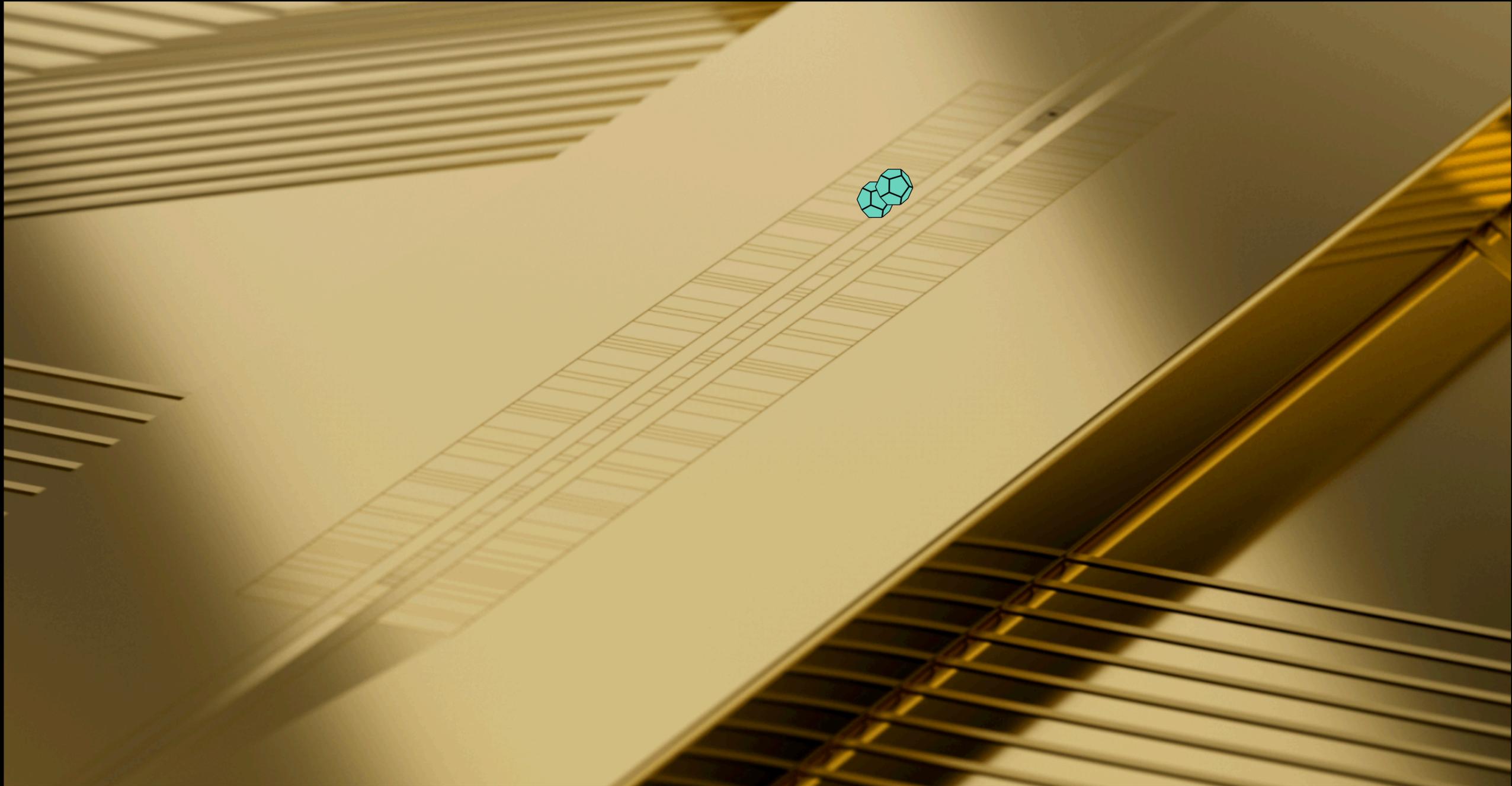
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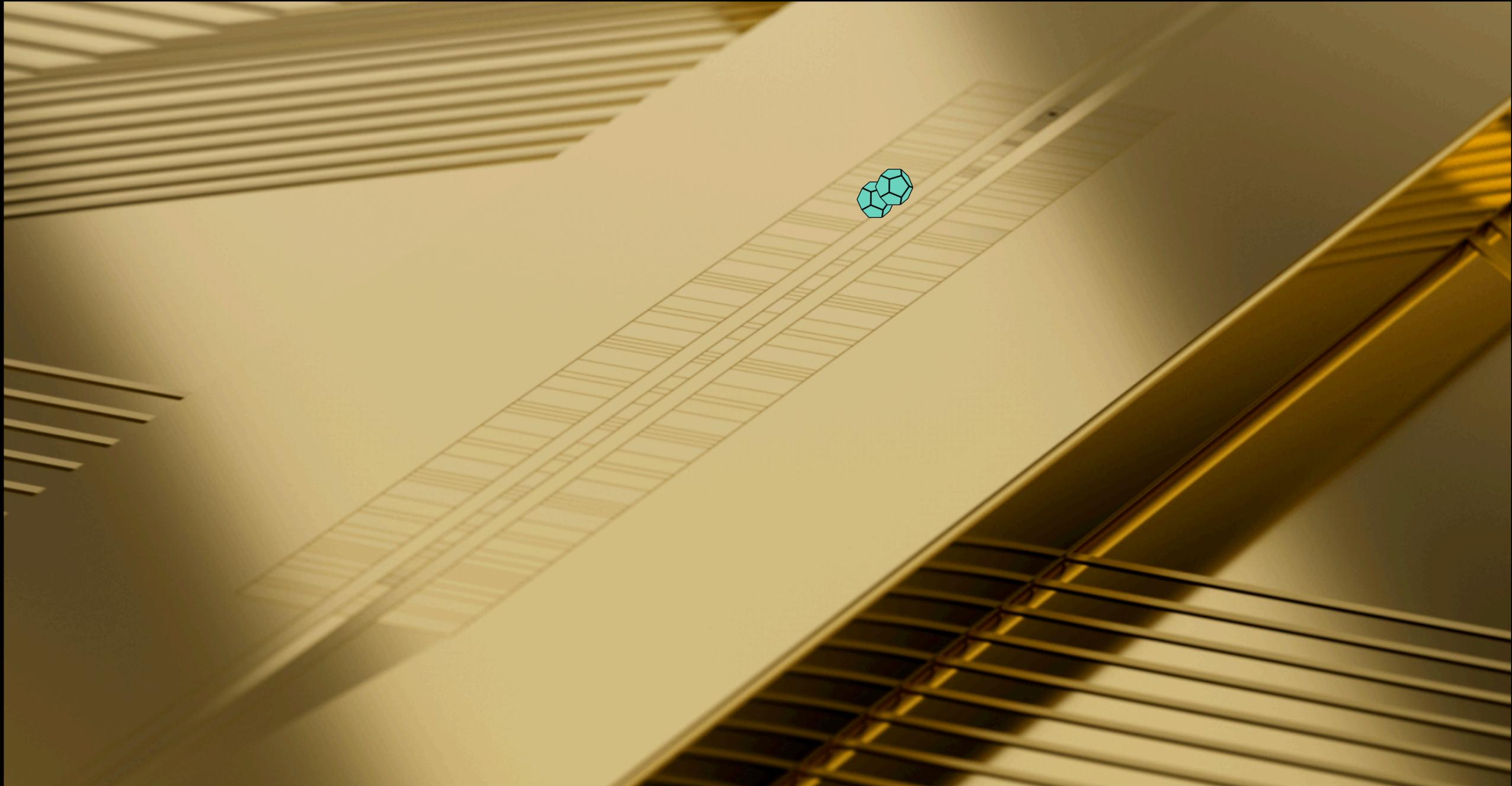
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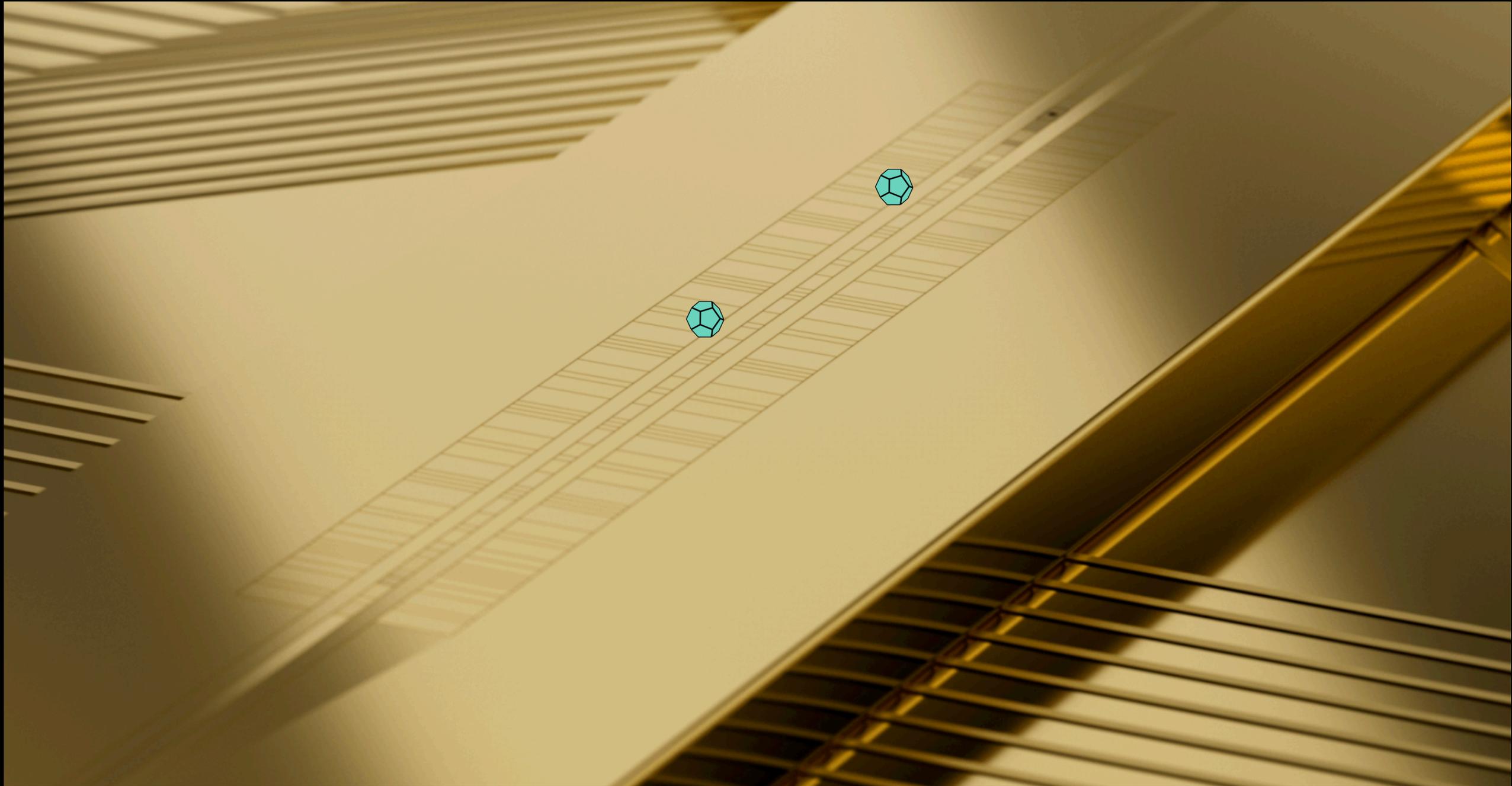
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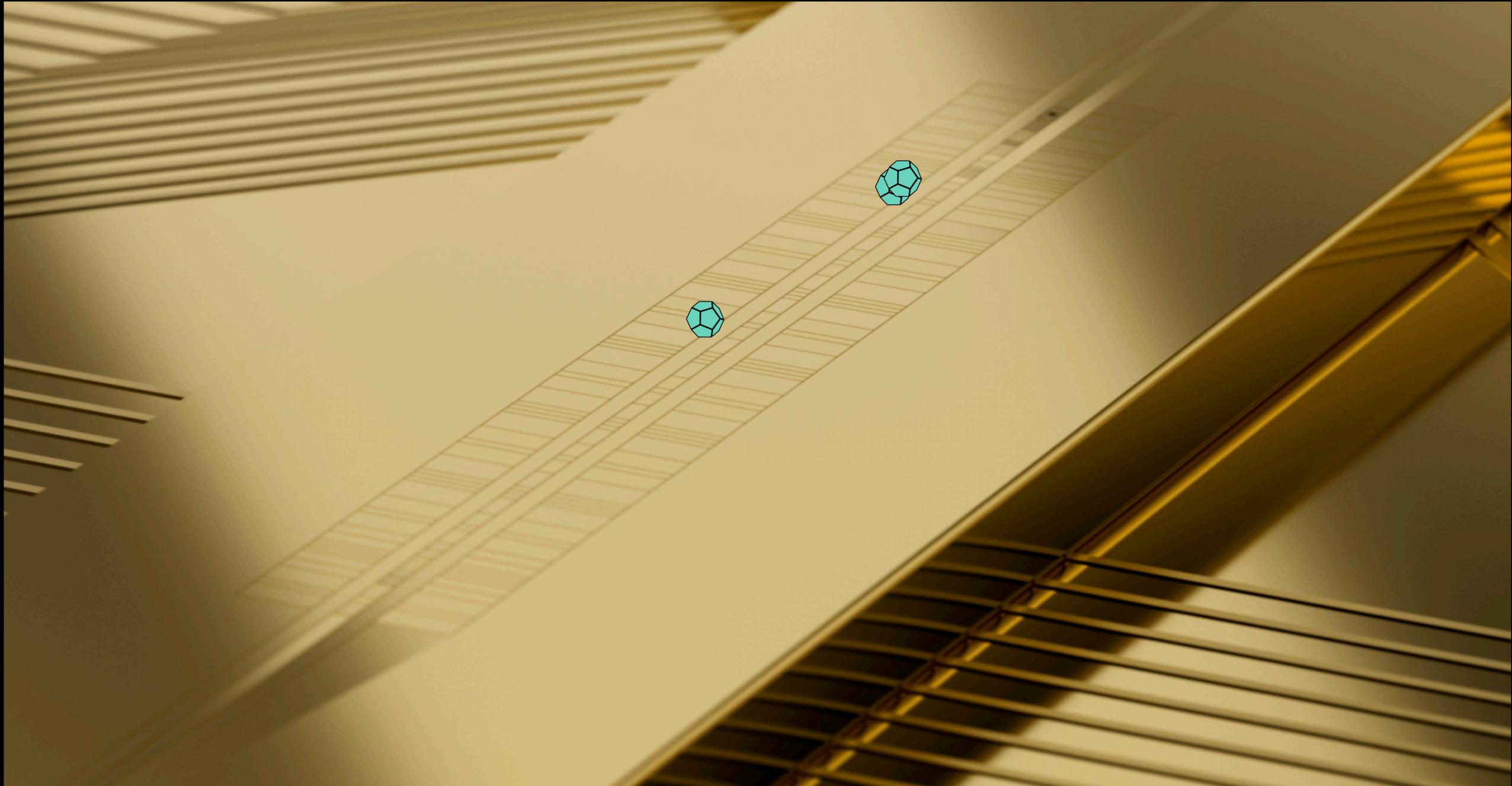
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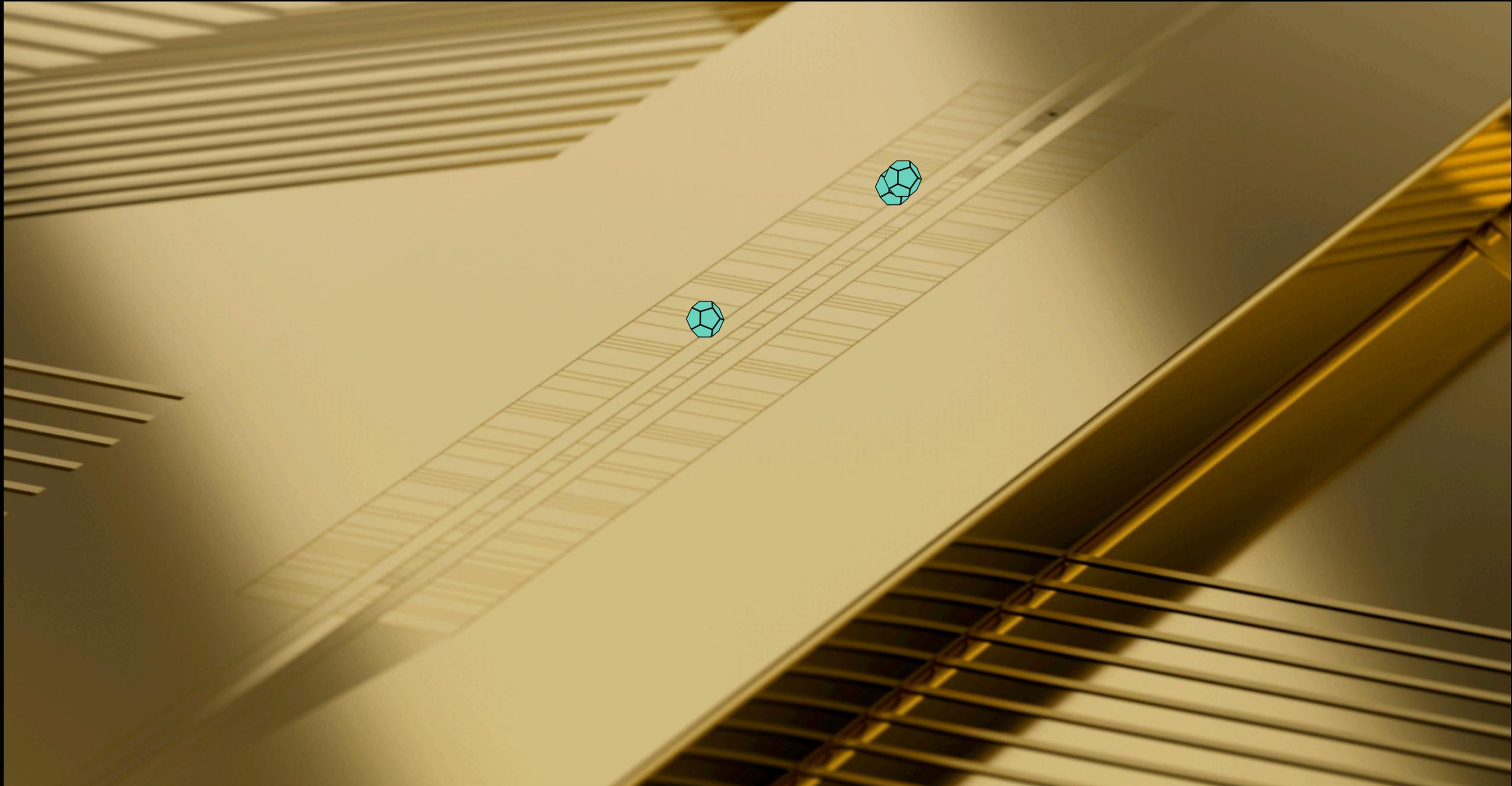
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With variational methods

- The ground state is prepared using the **variational quantum eigensolver (VQE)**
- A **trial state** is obtained using a parametrized quantum circuit $C(\theta)$ acting on some initial state

$$\cdot \quad |\Psi(\theta)\rangle = C(\theta) |\Psi_0\rangle$$

- The expectation value of the **Hamiltonian is measured**
- An optimizer updates the parameters towards the **minimum of the energy landscape**

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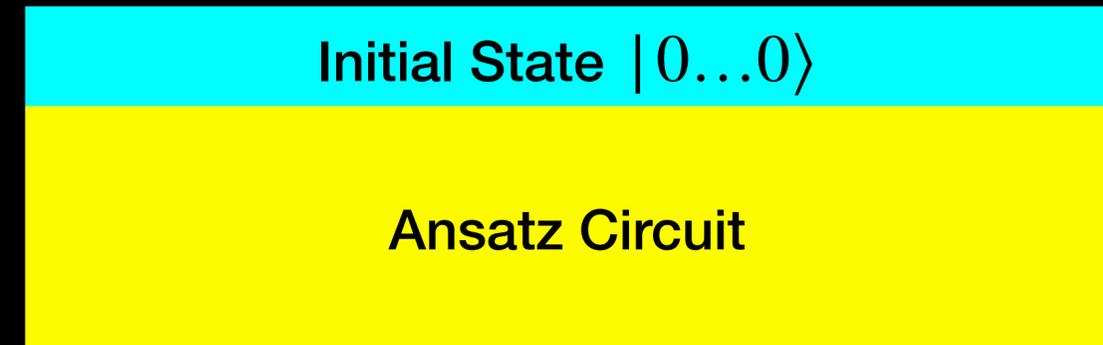
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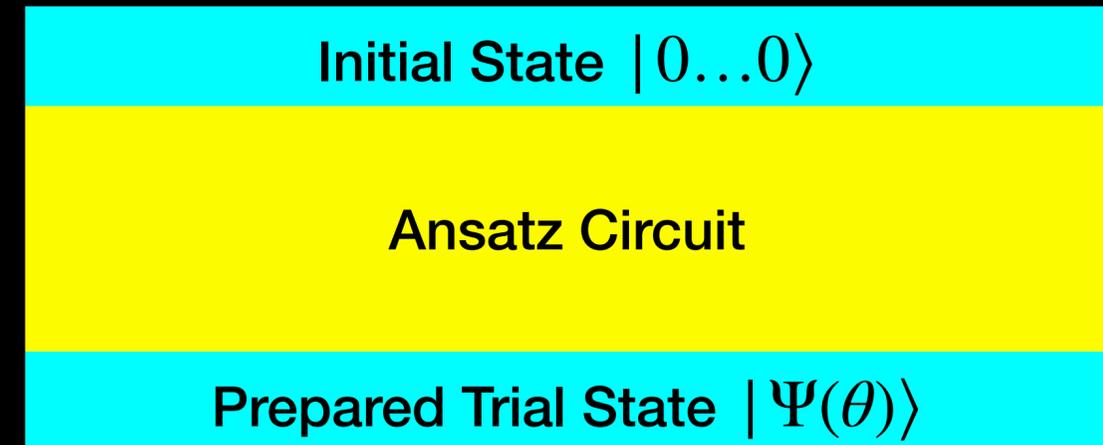
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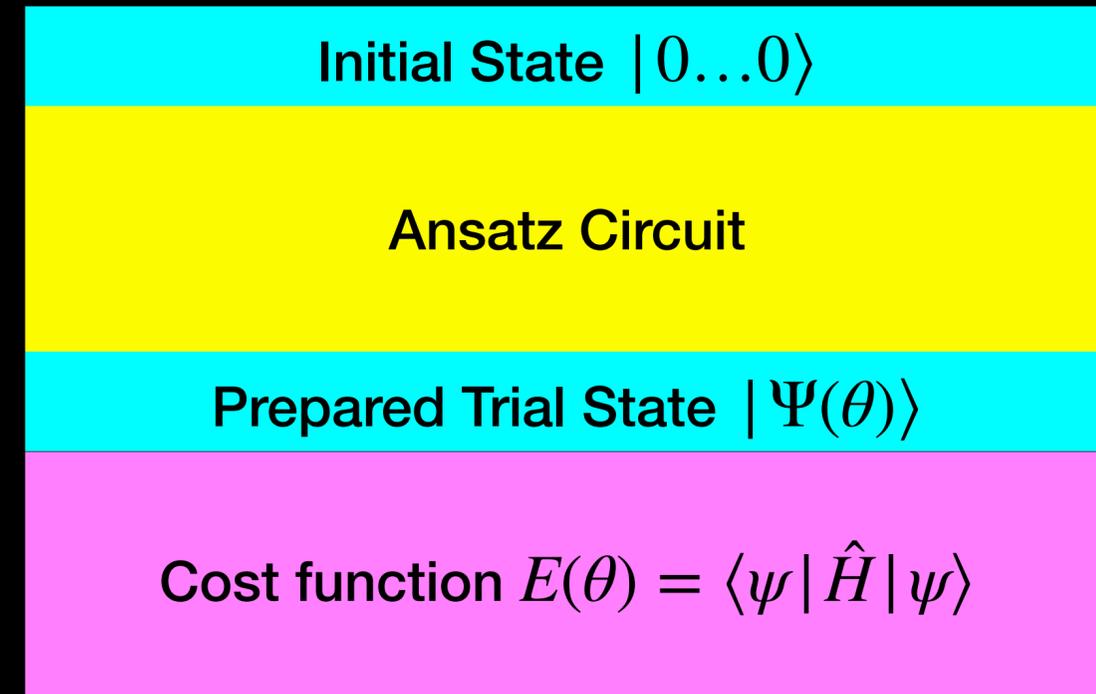
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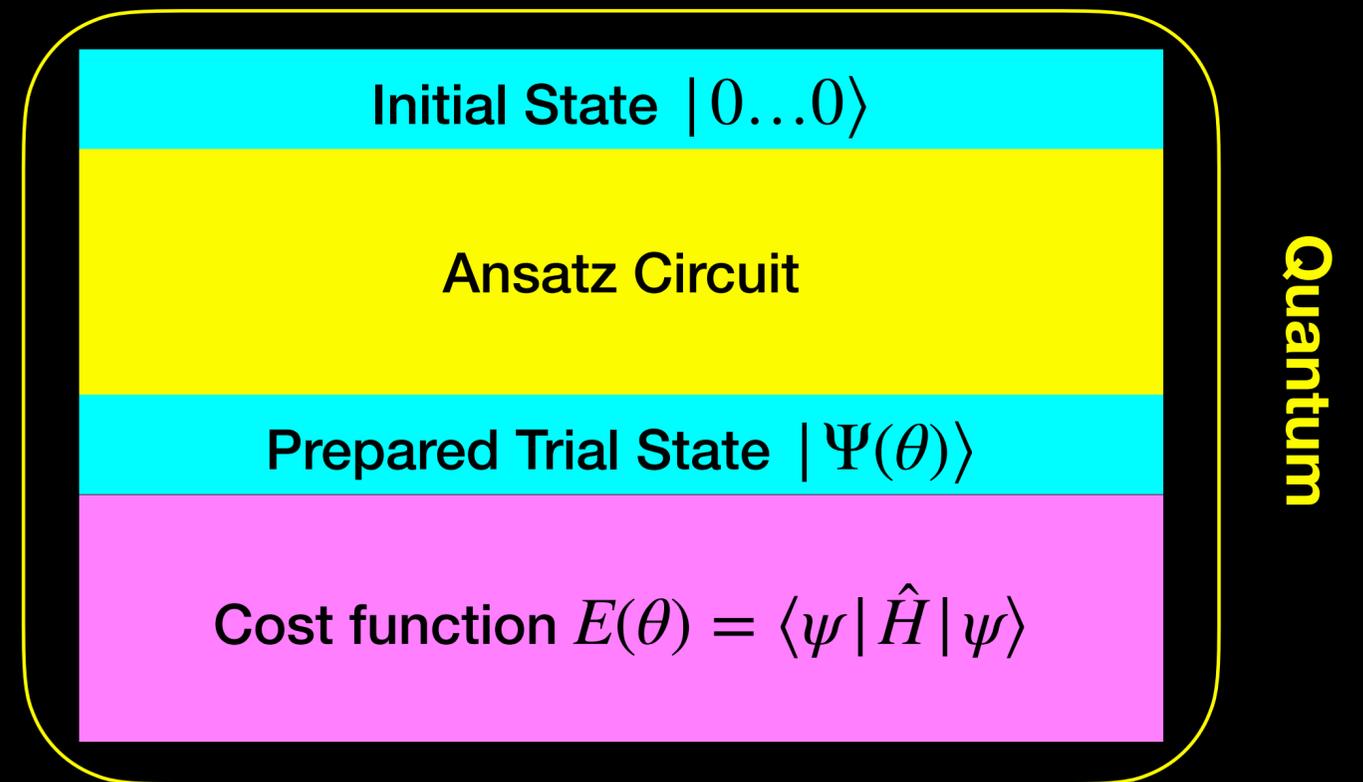
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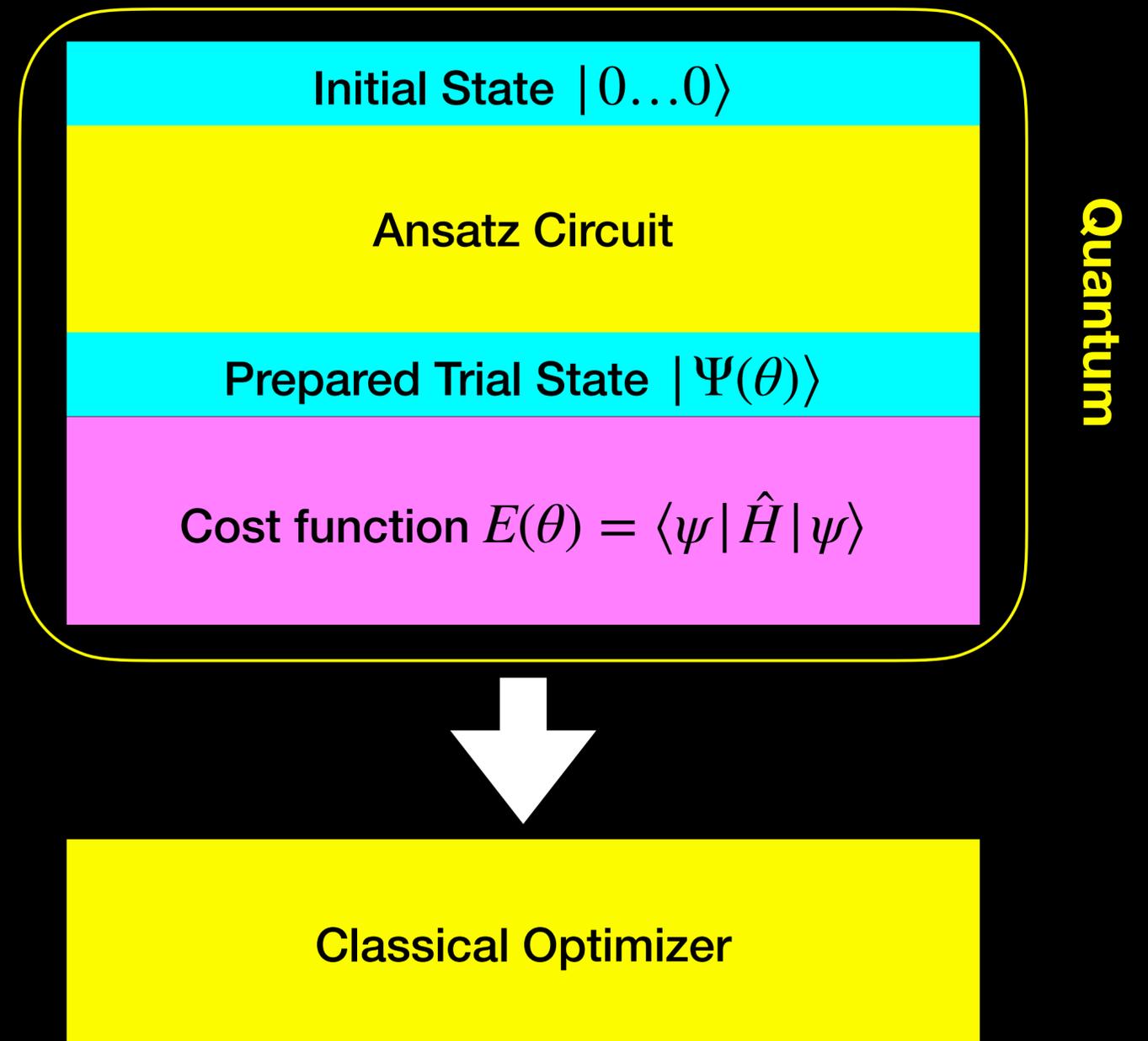
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- The ground state is prepared using the **variational quantum eigensolver (VQE)**
- A **trial state** is obtained using a parametrized quantum circuit $C(\theta)$ acting on some initial state

$$|\Psi(\theta)\rangle = C(\theta)|\Psi_0\rangle$$

- The expectation value of the **Hamiltonian is measured**
- An optimizer updates the parameters towards the **minimum of the energy landscape**



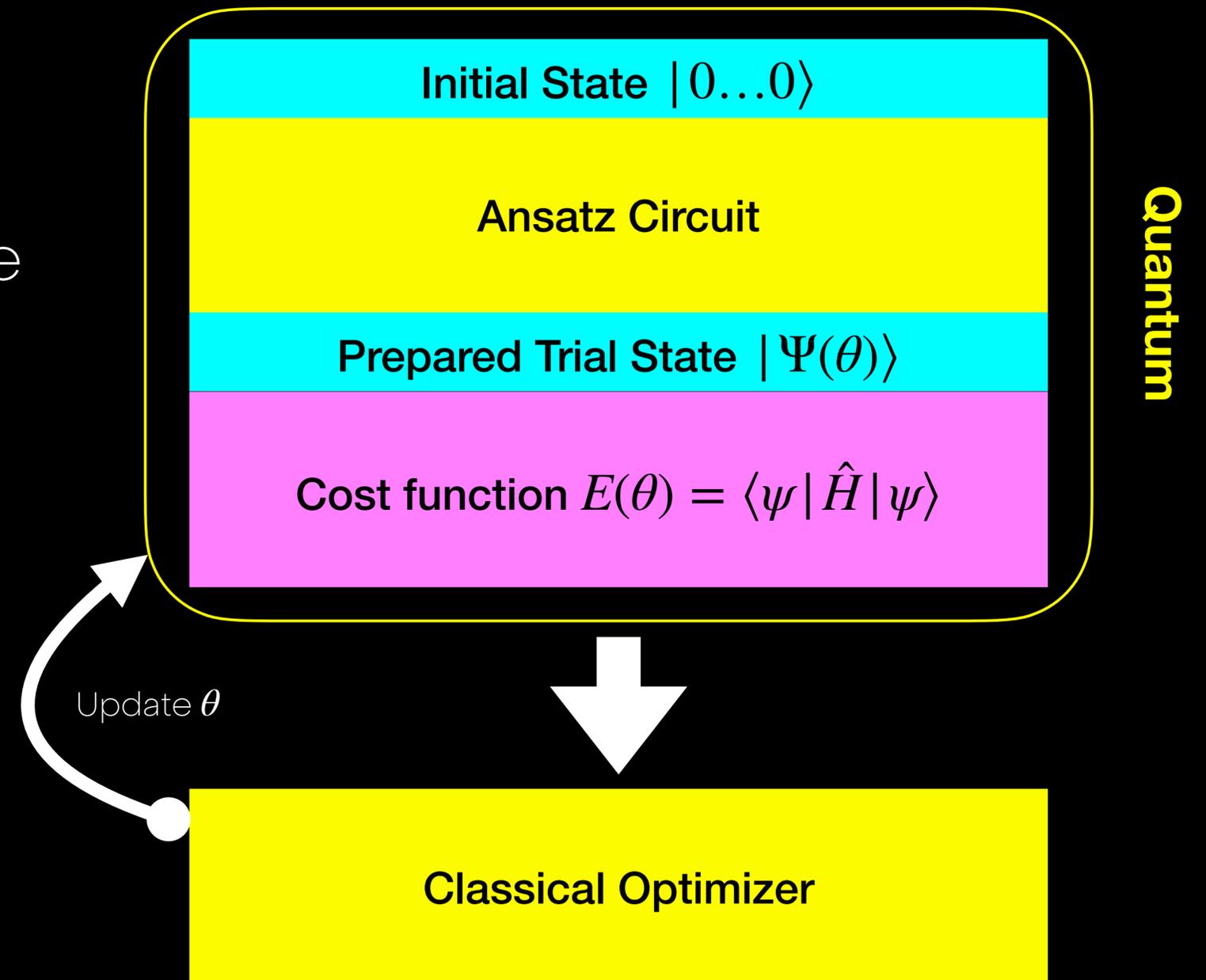
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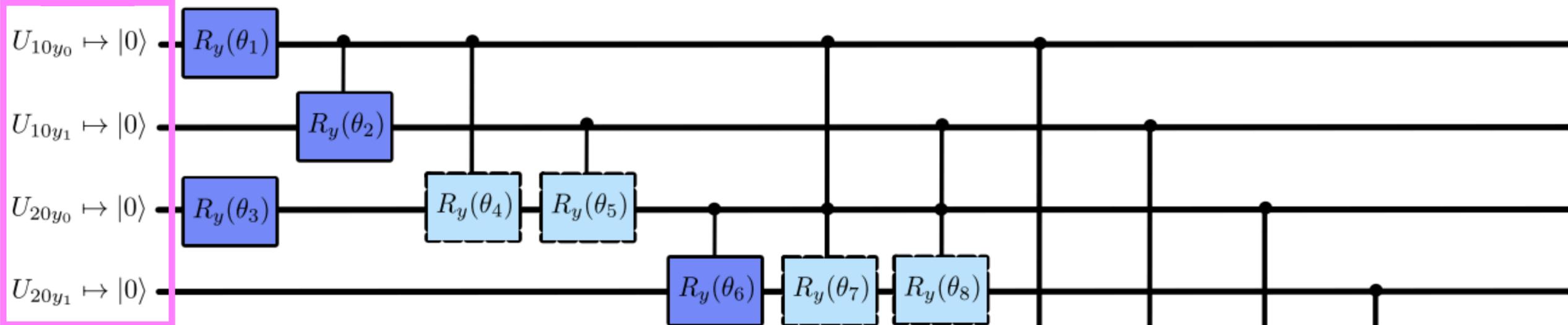
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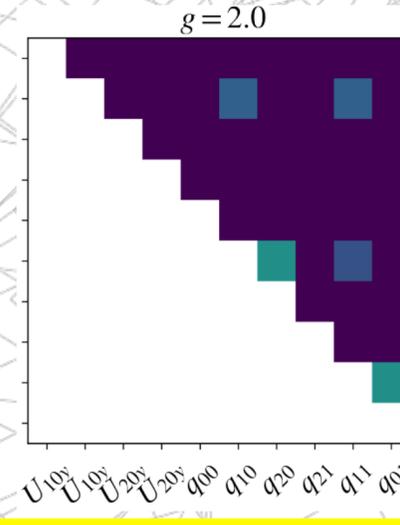
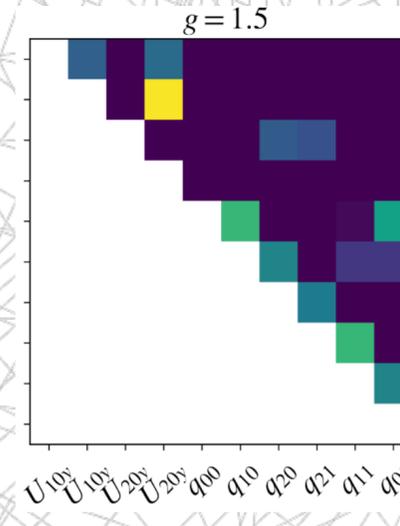
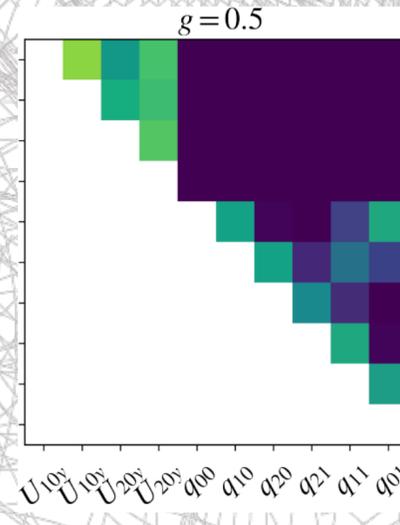
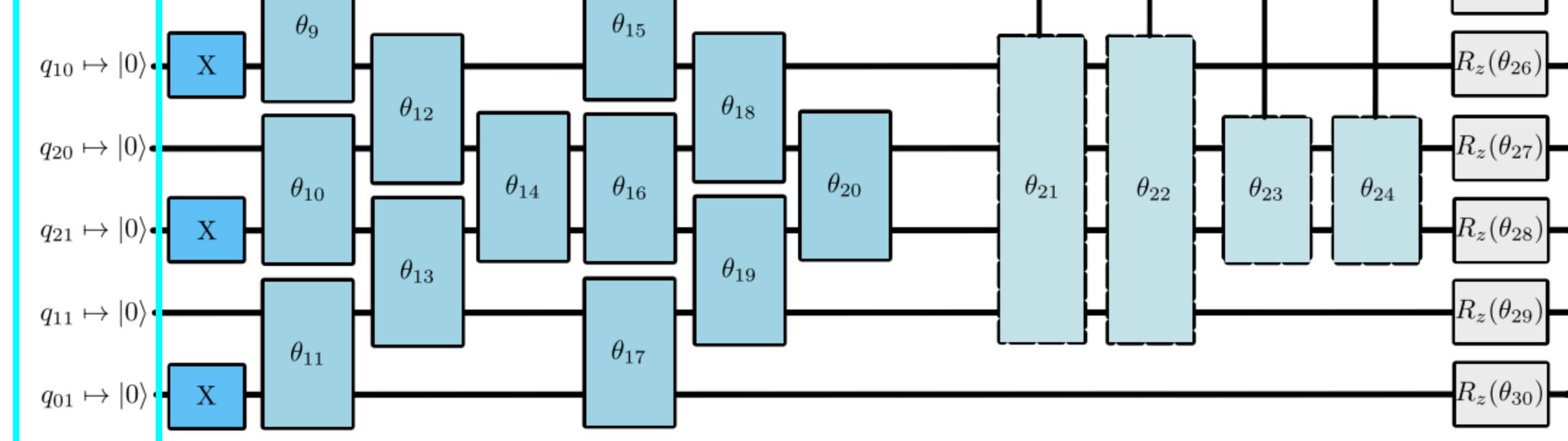


Mutual Information Ansatz

Links

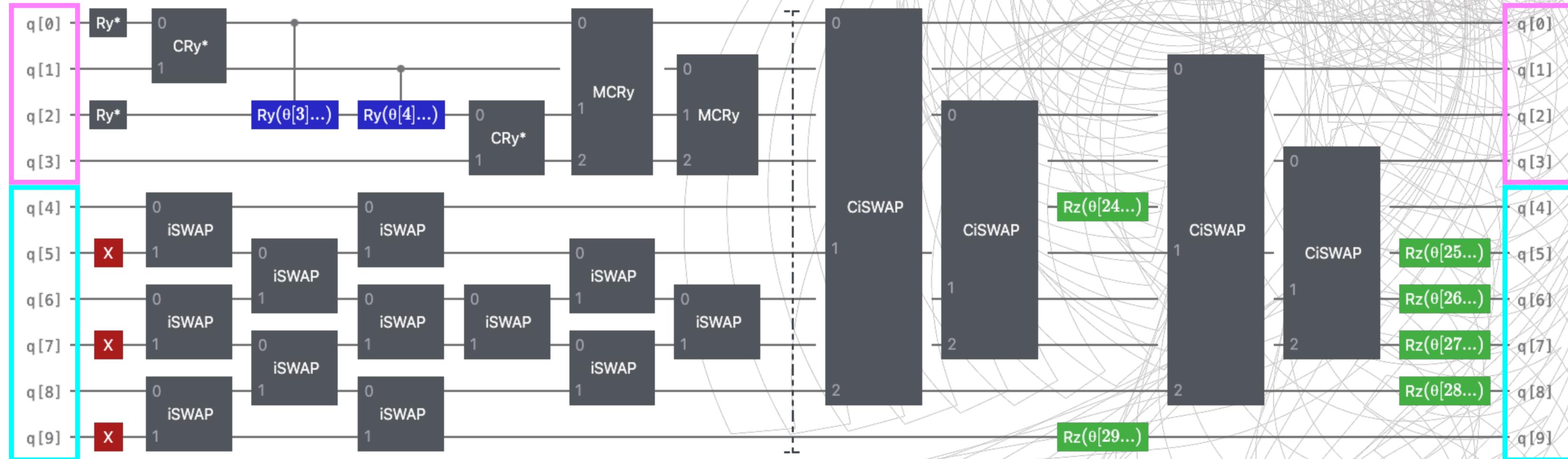


Sites



Example ansatz circuit

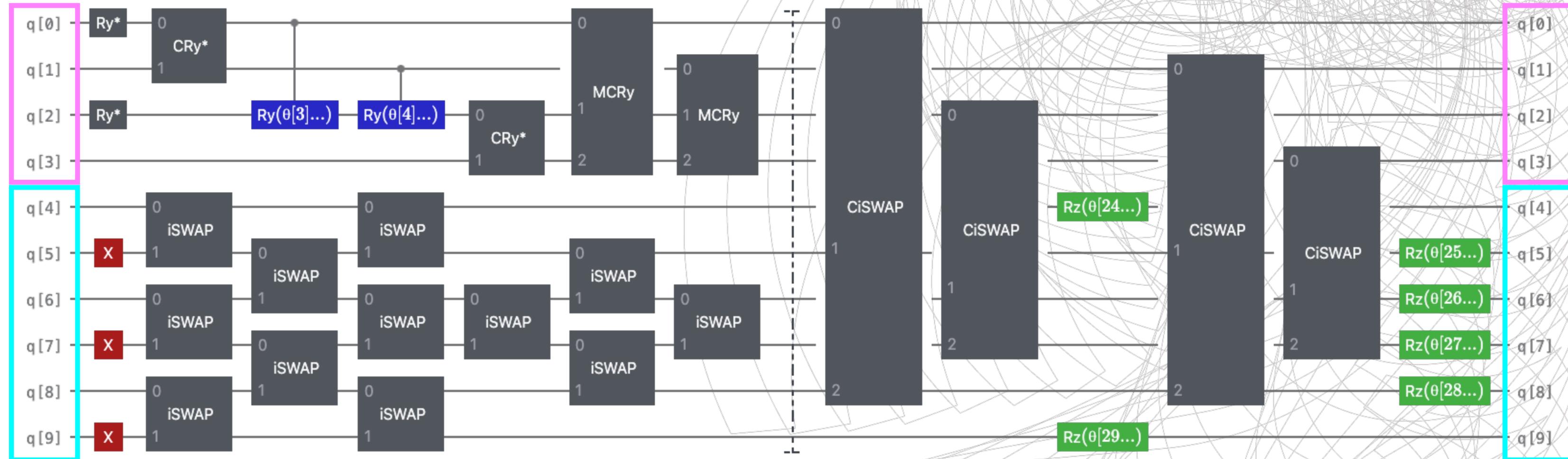
10 qubits, 30 parameters



$$|\Psi(\theta)\rangle = |\psi_{\text{sites}}\rangle \otimes |\psi_{\text{links}}\rangle = C(\theta) |\Psi_0\rangle$$

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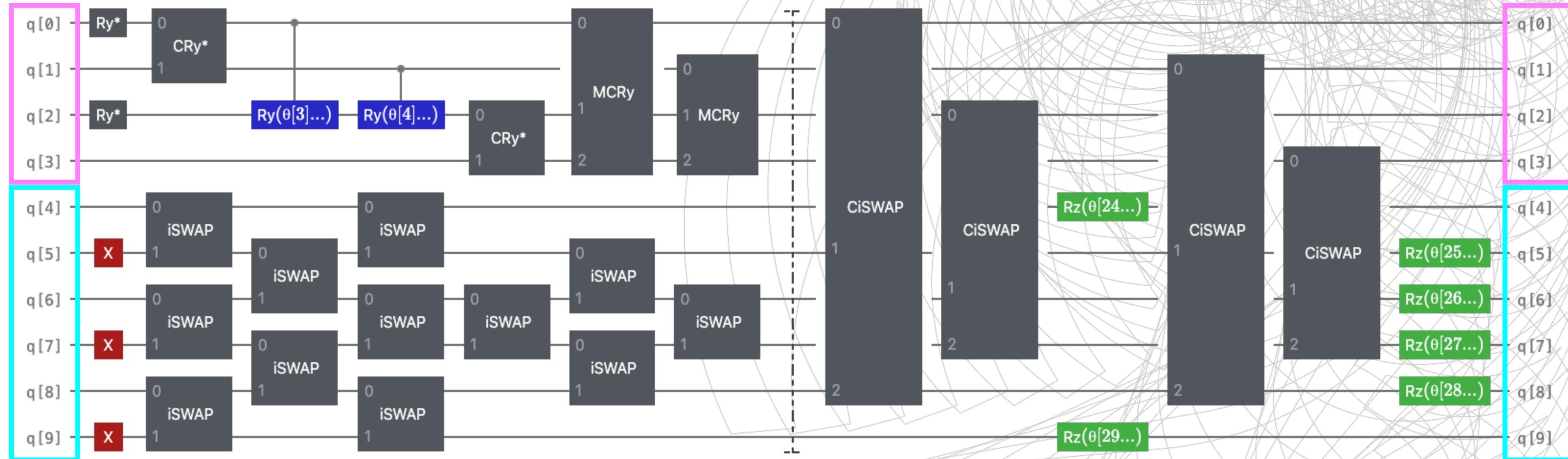
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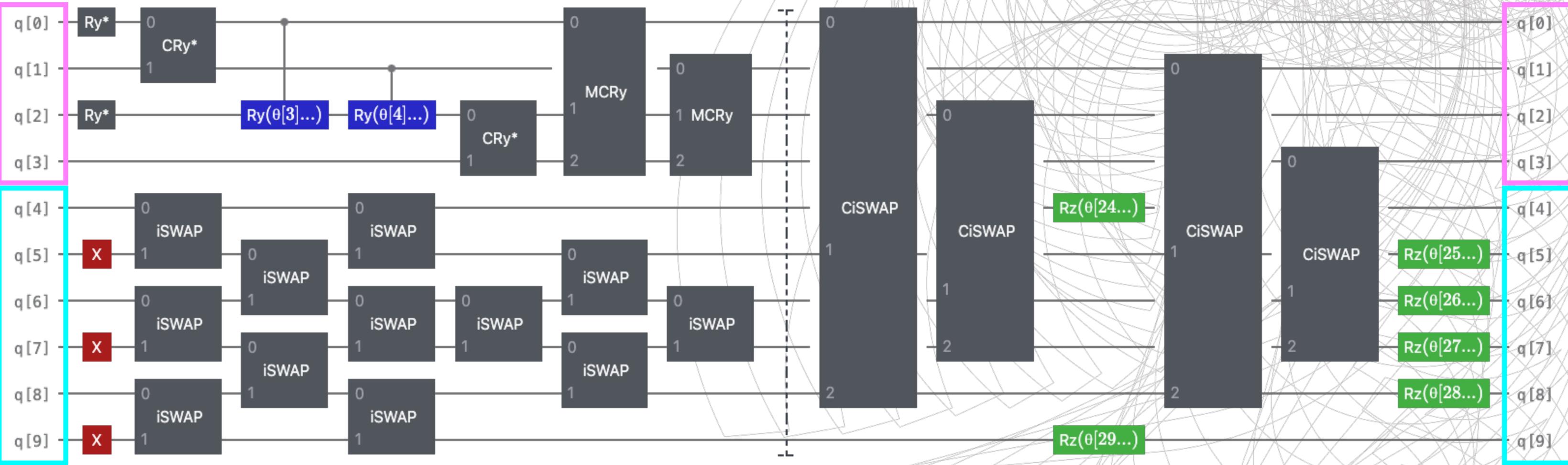
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Example ansatz circuit

10 qubits, 30 parameters



$$|\Psi(\theta)\rangle = |\psi_{\text{sites}}\rangle \otimes |\psi_{\text{links}}\rangle = C(\theta) |\Psi_0\rangle \Rightarrow$$

- |1010100101⟩
- |1010100100⟩
- |1000110100⟩
- |1000110111⟩
- ...

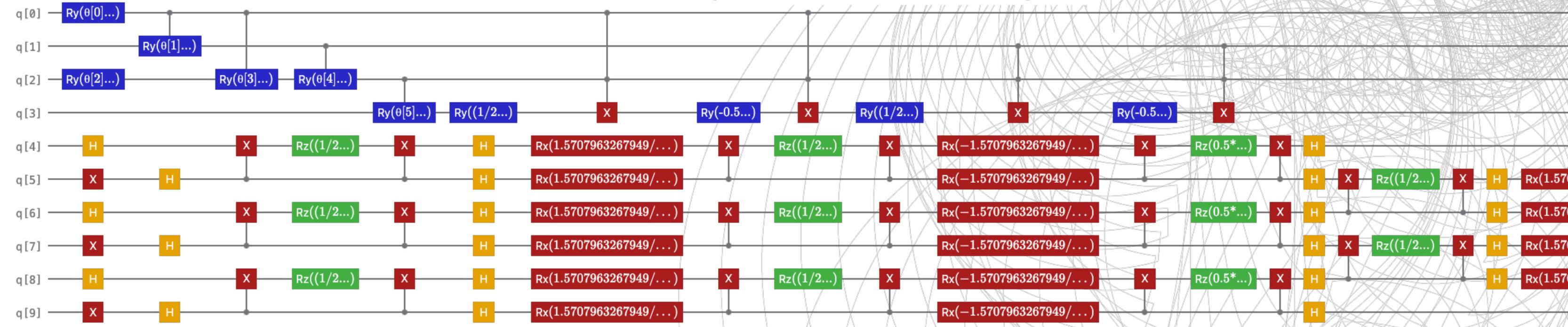
Example of gate decomposition

{H, X, Rz, Rx, Ry, CNOT}: ≈ 115 2-qubit gates

H-series Native Gates: ≈ 80 2-qubit gates

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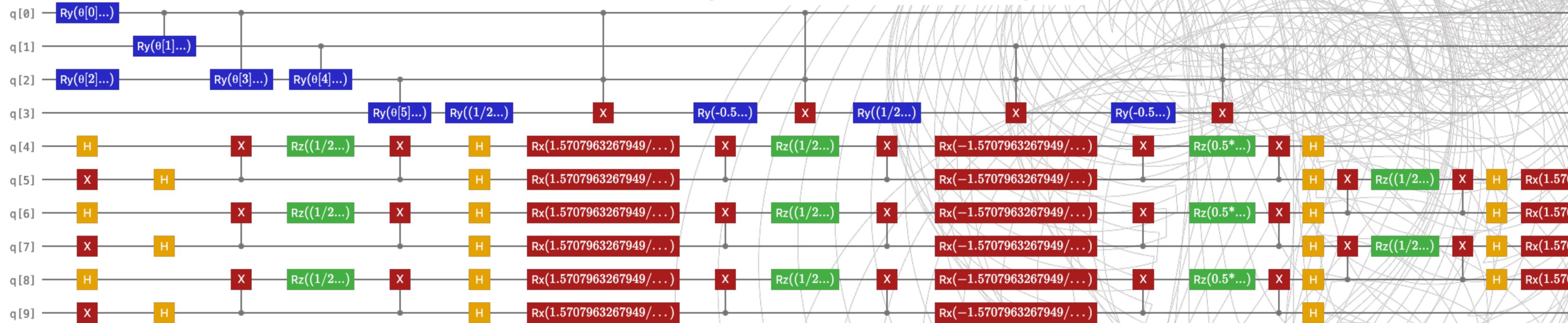
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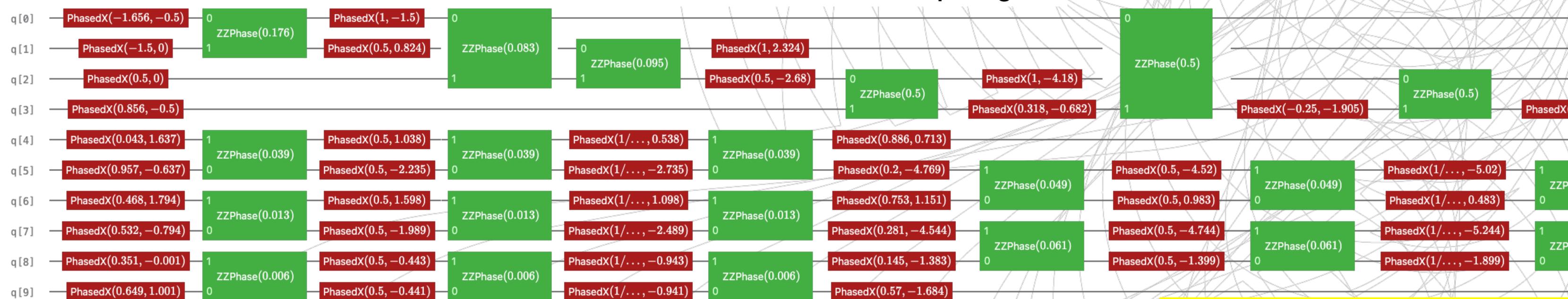
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Results

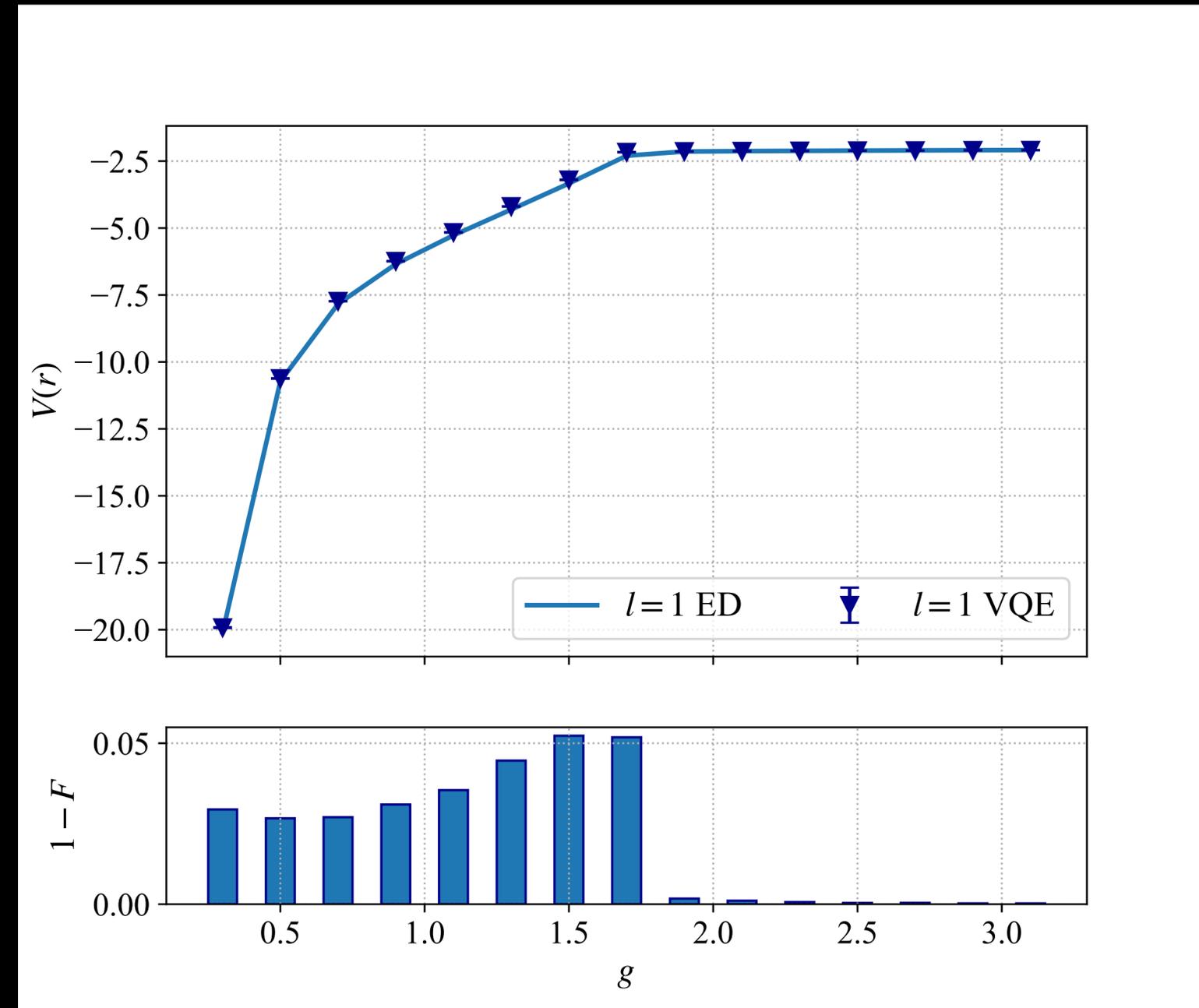


QUANTINUM

Consistency checks

Benchmark *quantum* results against *classical* methods

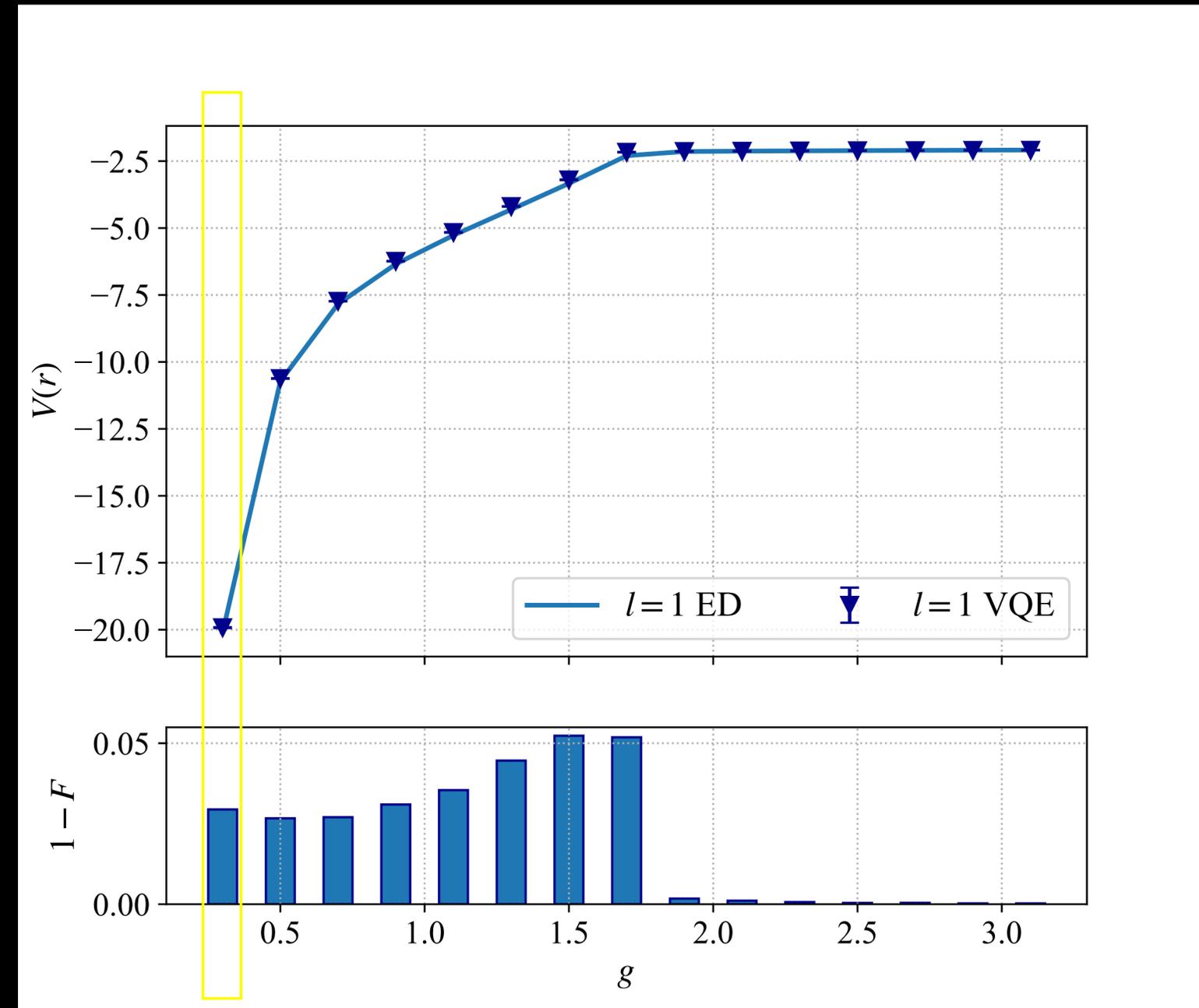
- A quantum state of 10 qubits can be represented with **classical memory on a laptop** and the Hamiltonian can be easily written as a matrix
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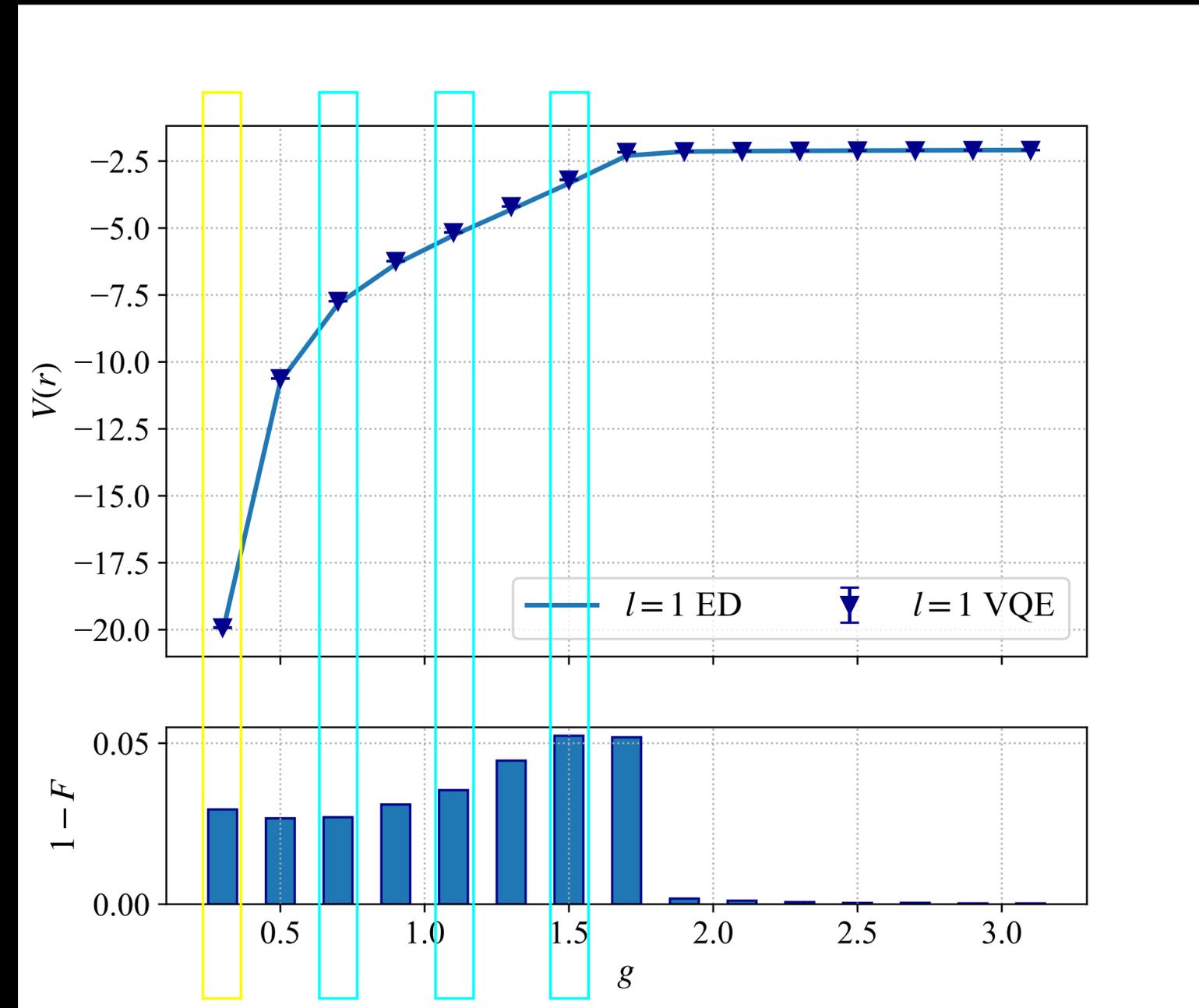
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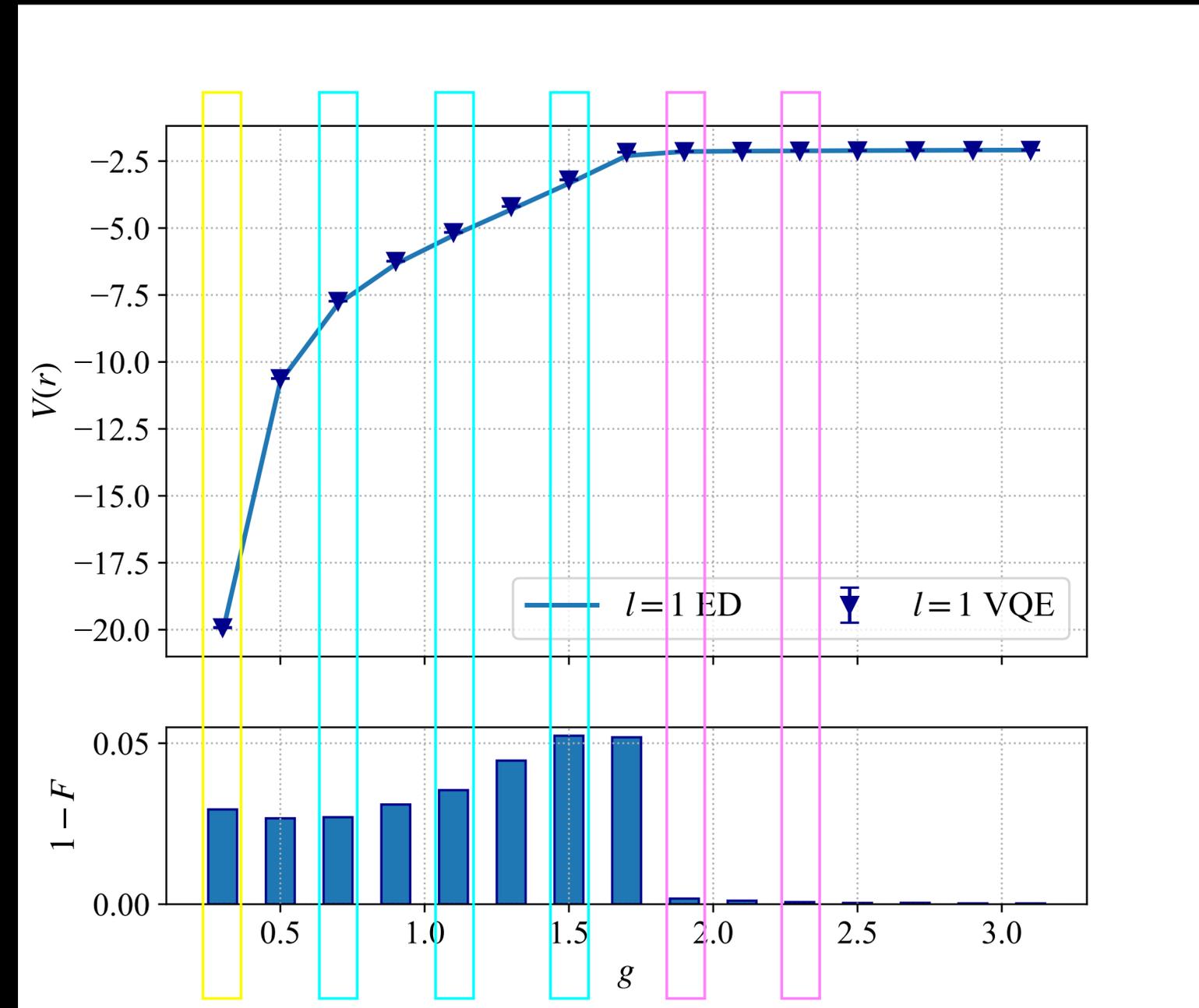
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Simulation and Emulation

And real hardware experiments

- Given the optimal parameters for the ansatz circuit at each coupling we can:
 - Simulate the circuit classically **without measuring**
 - Simulate the circuit classically **with measurements**
 - Simulate the circuit classically with measurements and **noisy operations**
 - Emulate the circuit on a trapped ion device
 - Run the circuit on a trapped ion device

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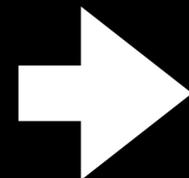
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$$|\Psi_{GS}\rangle = \sum_i^{2^N} c_i |i\rangle$$

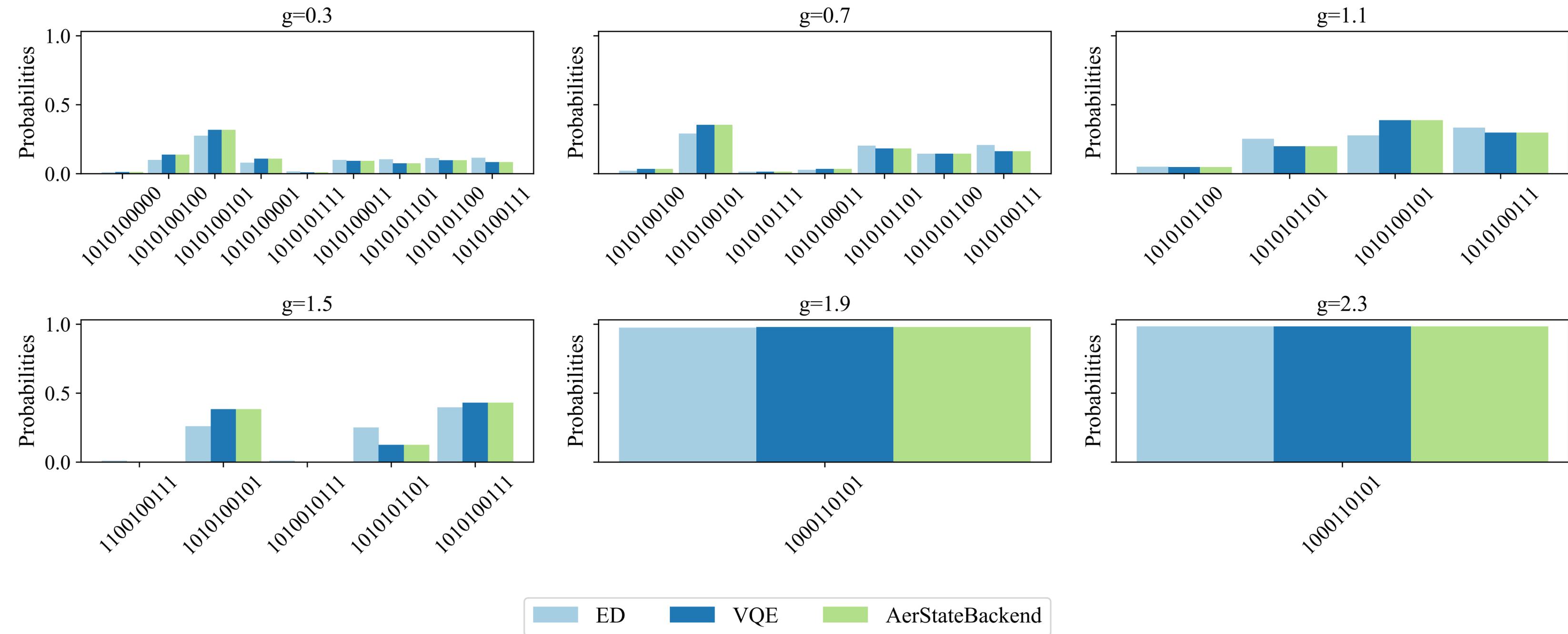


$$\text{prob}_i = |c_i|^2 = |\langle i | \Psi_{GS} \rangle|^2$$

Simulations **on laptops**

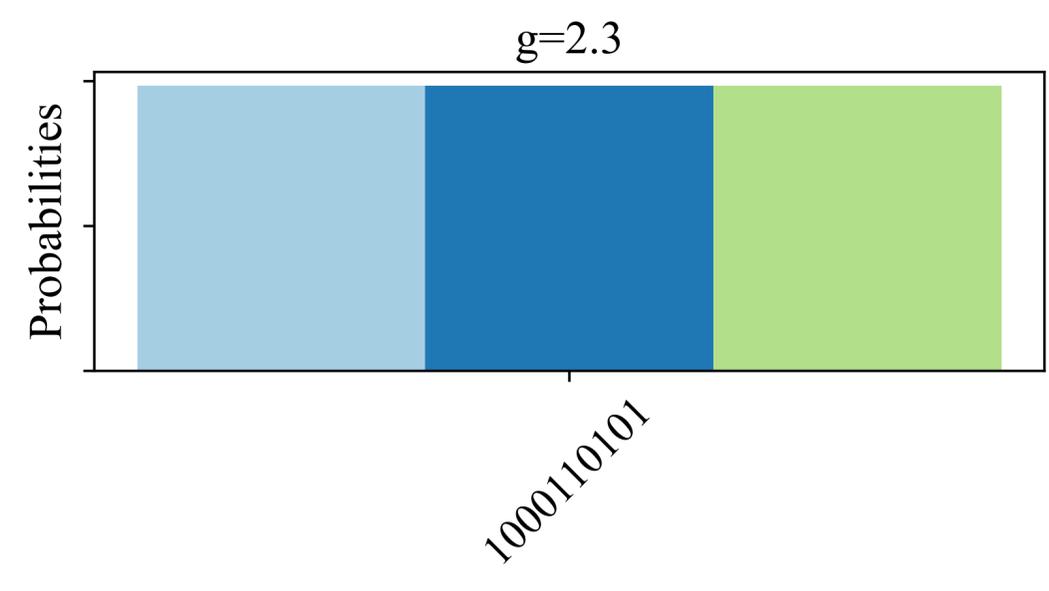
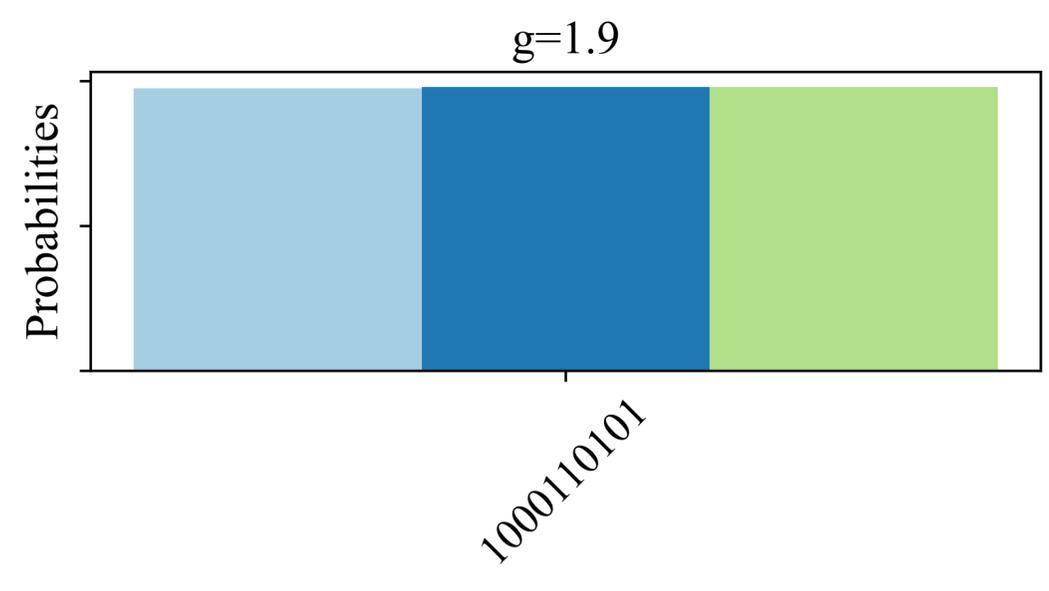
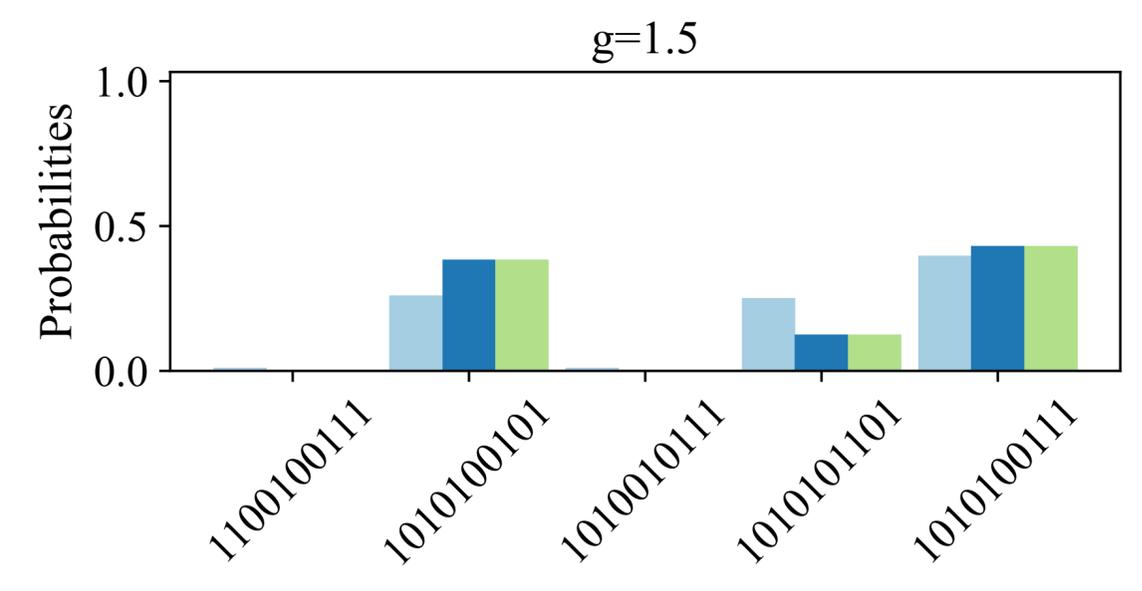
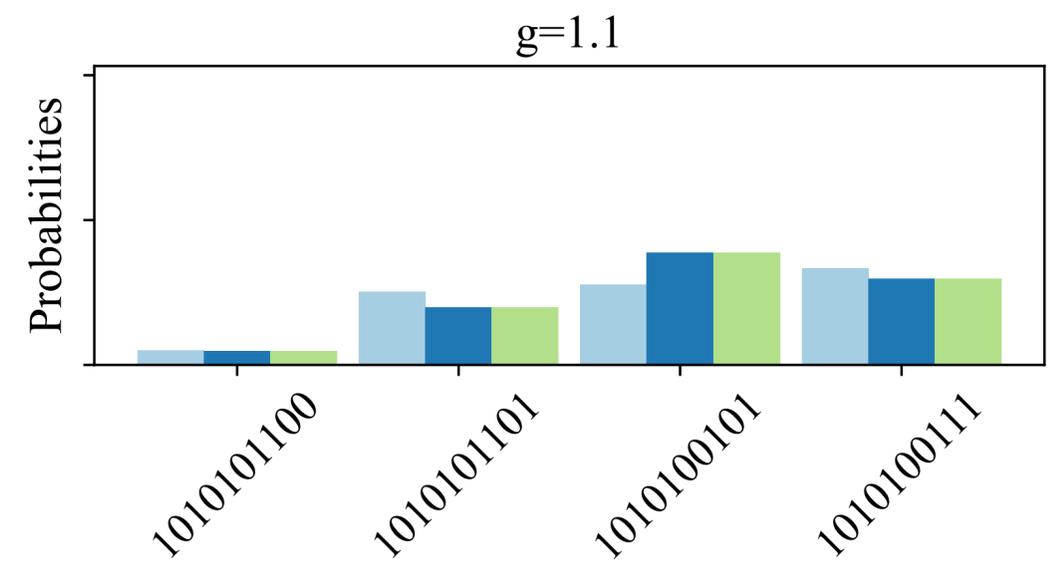
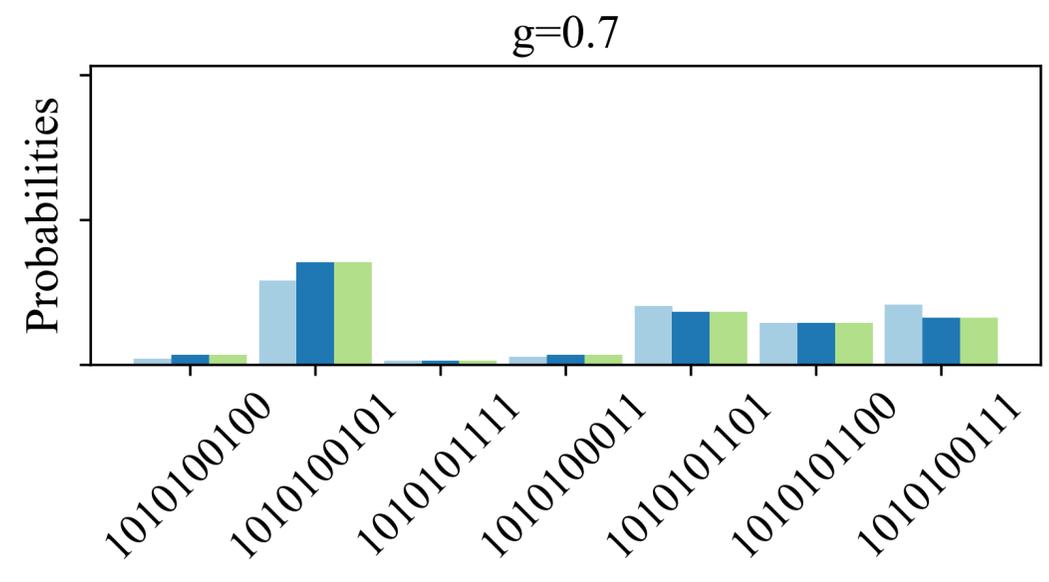
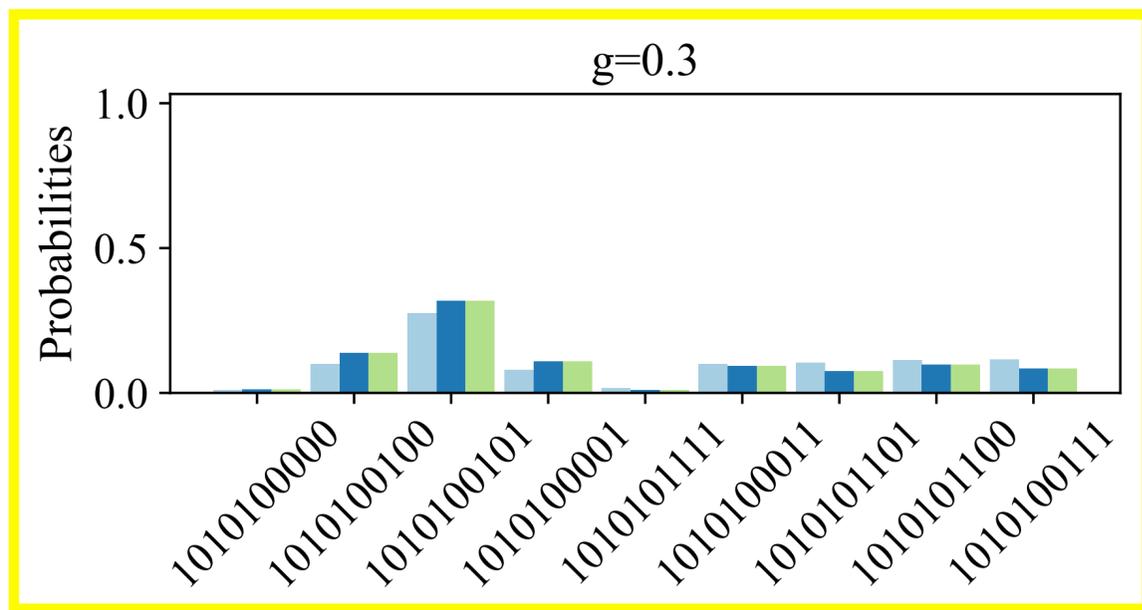
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Probabilities of the different states for each g value: AerStateBackend shots=None



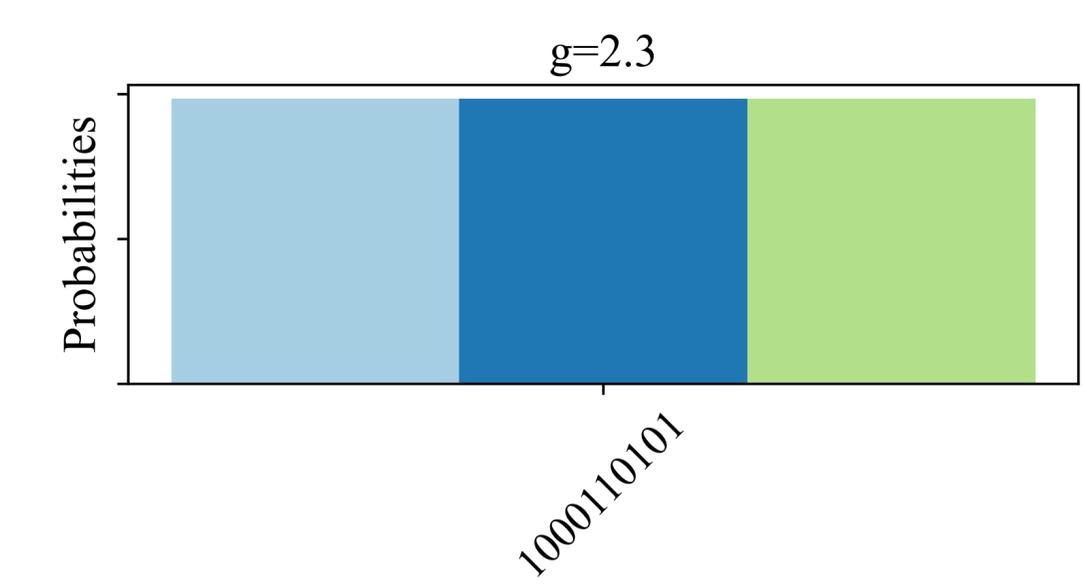
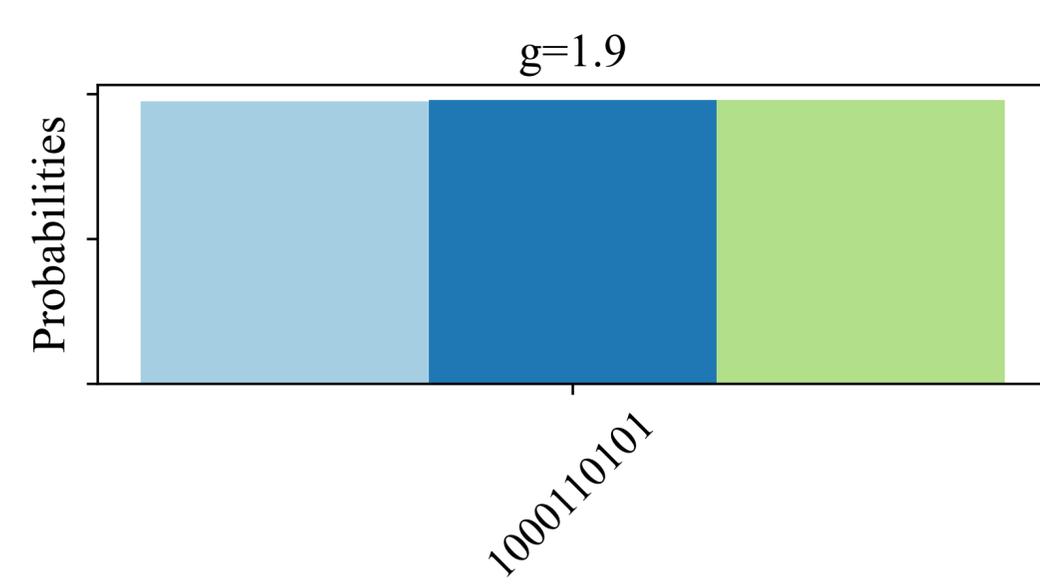
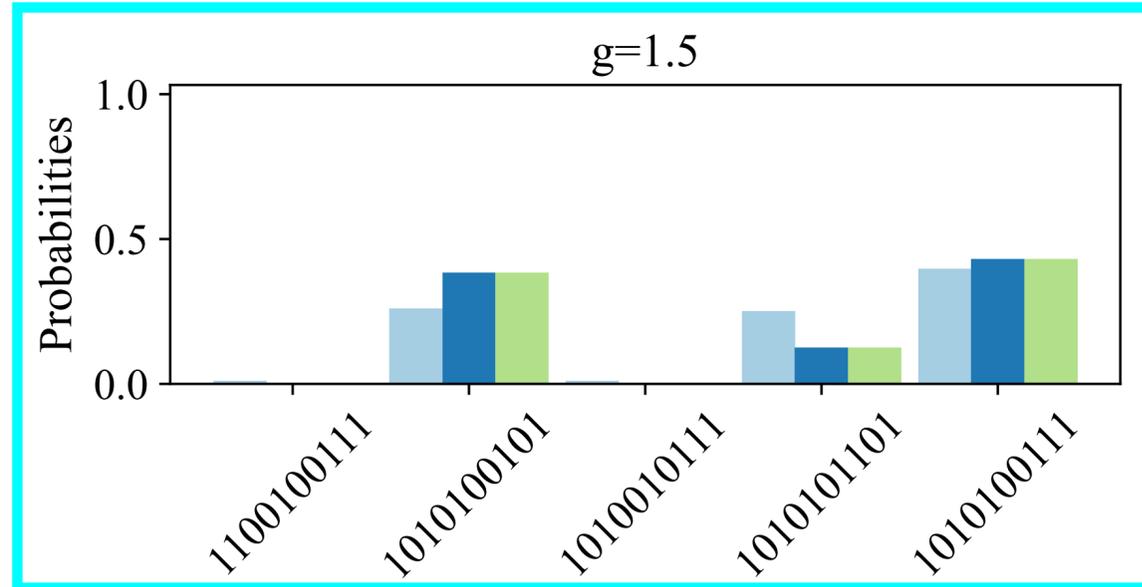
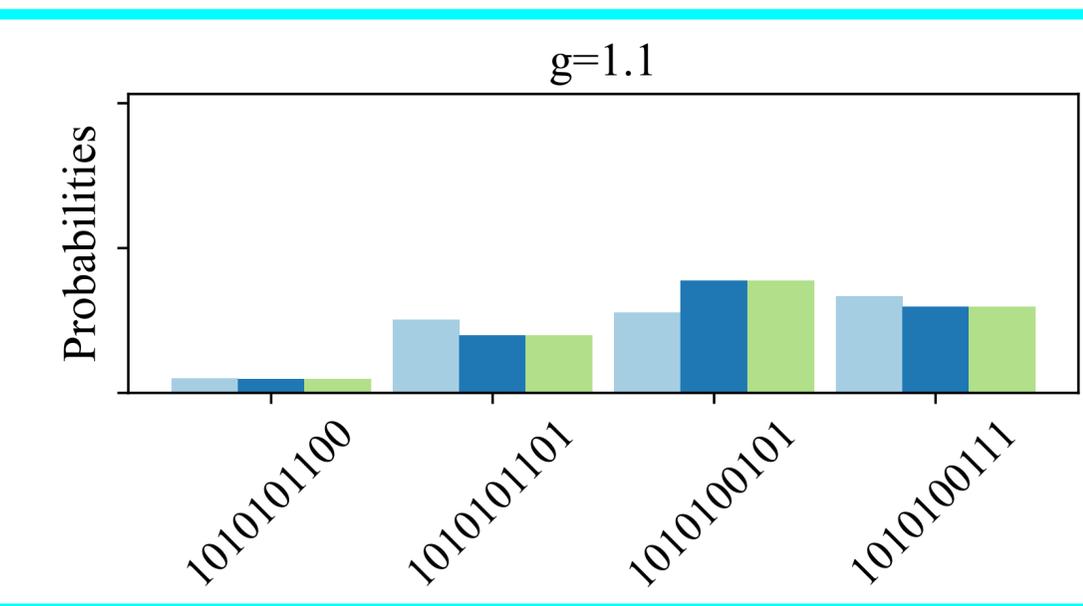
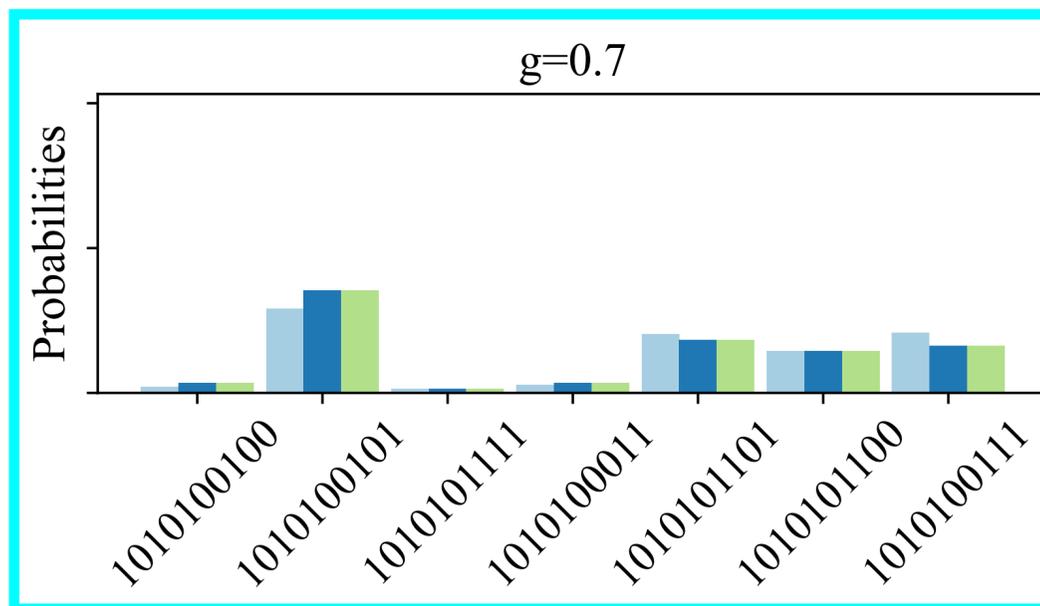
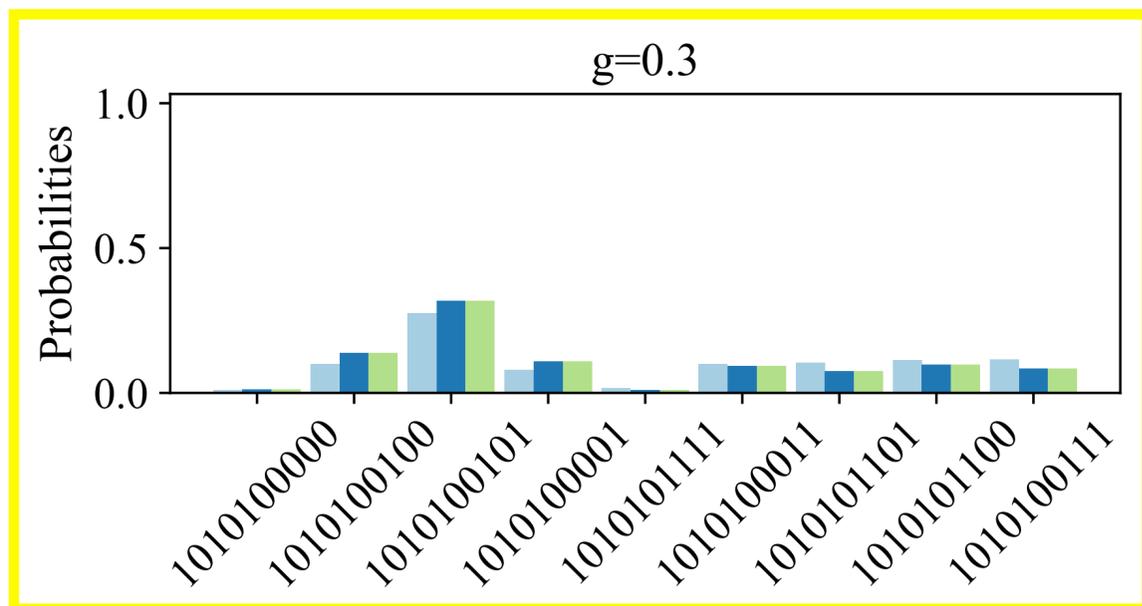
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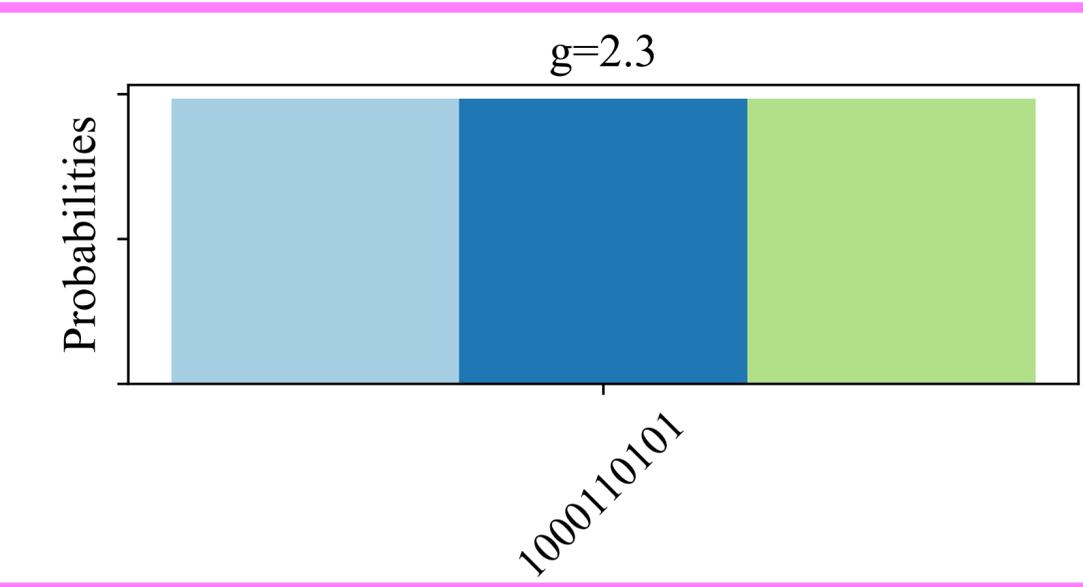
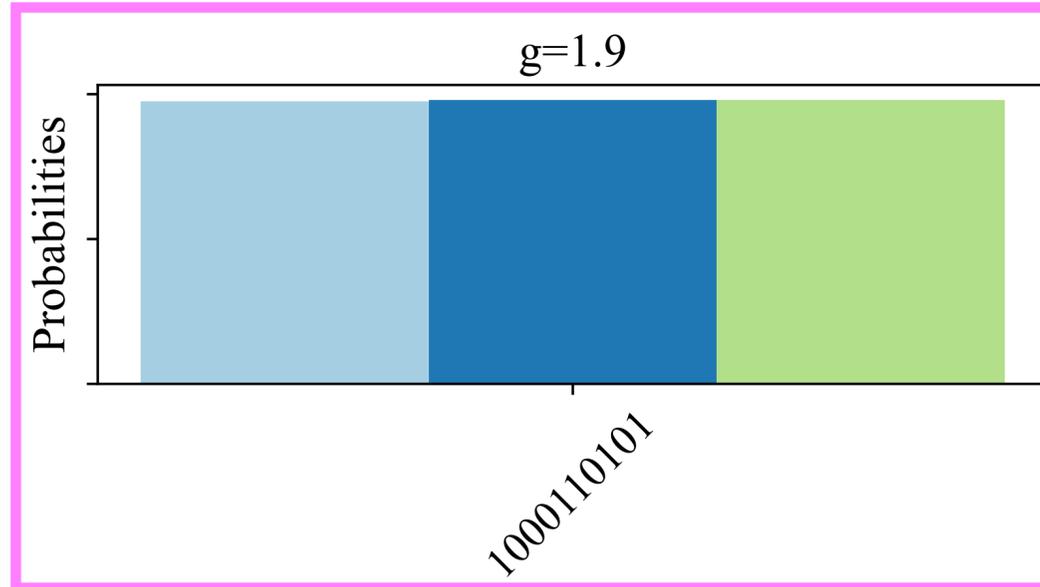
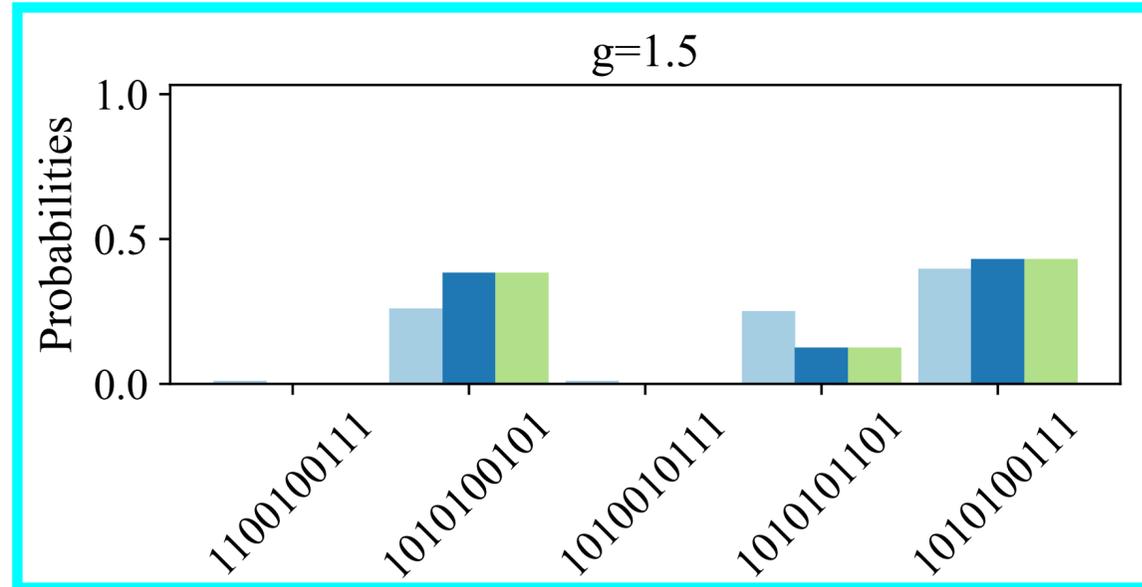
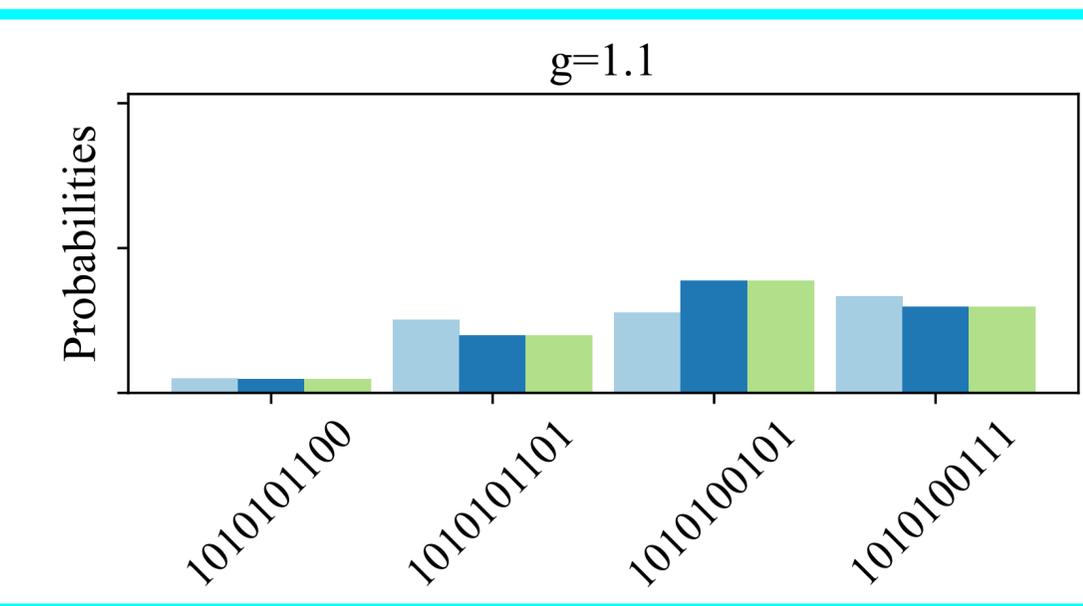
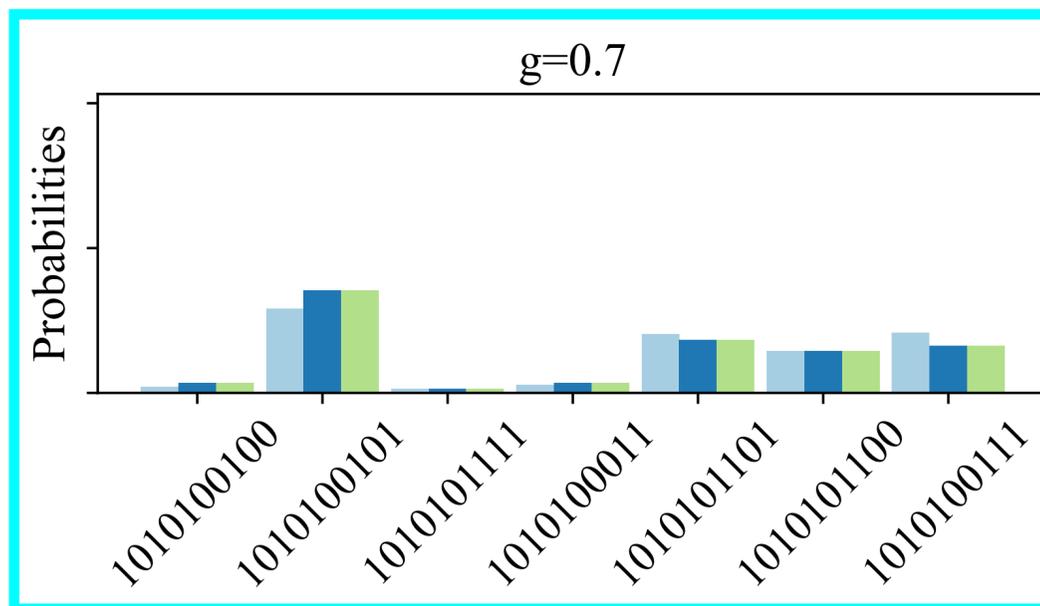
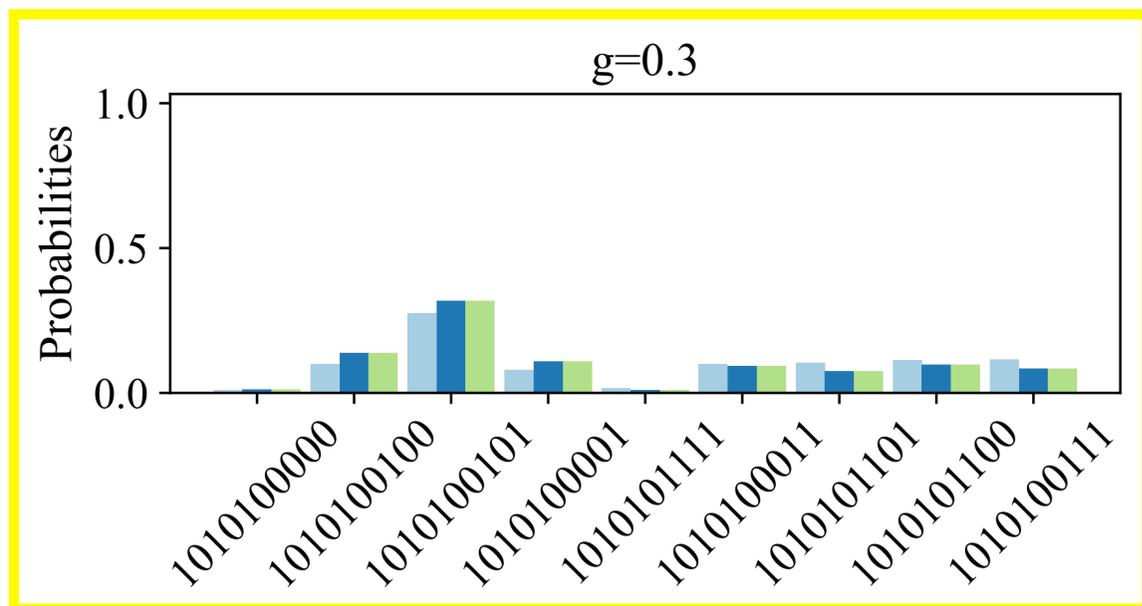
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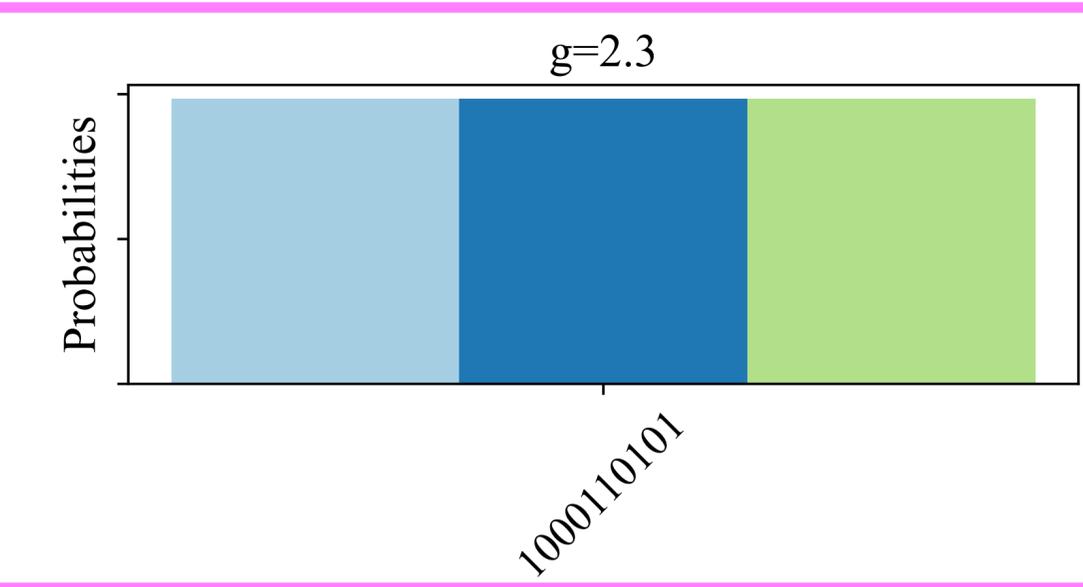
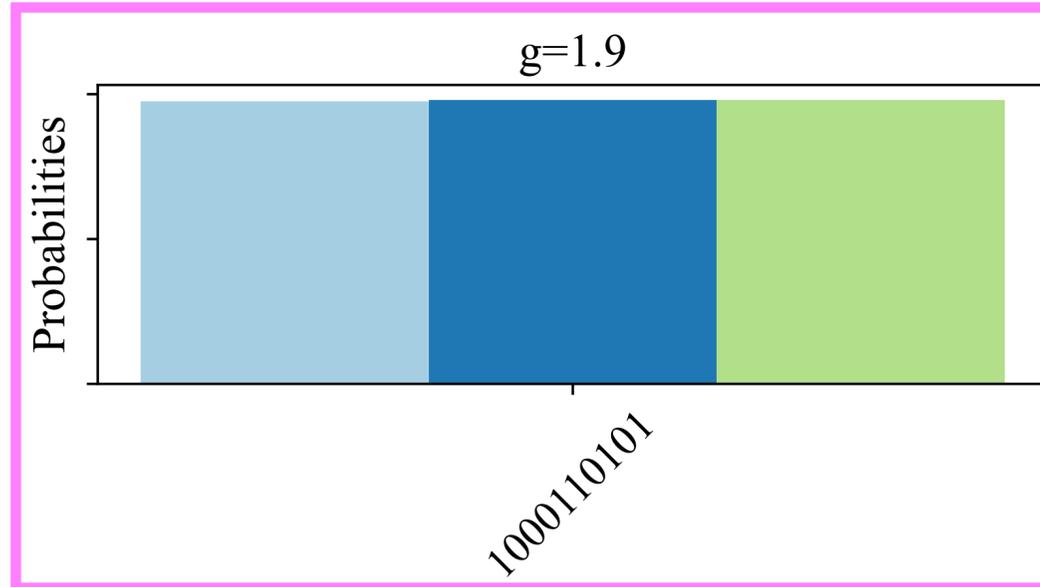
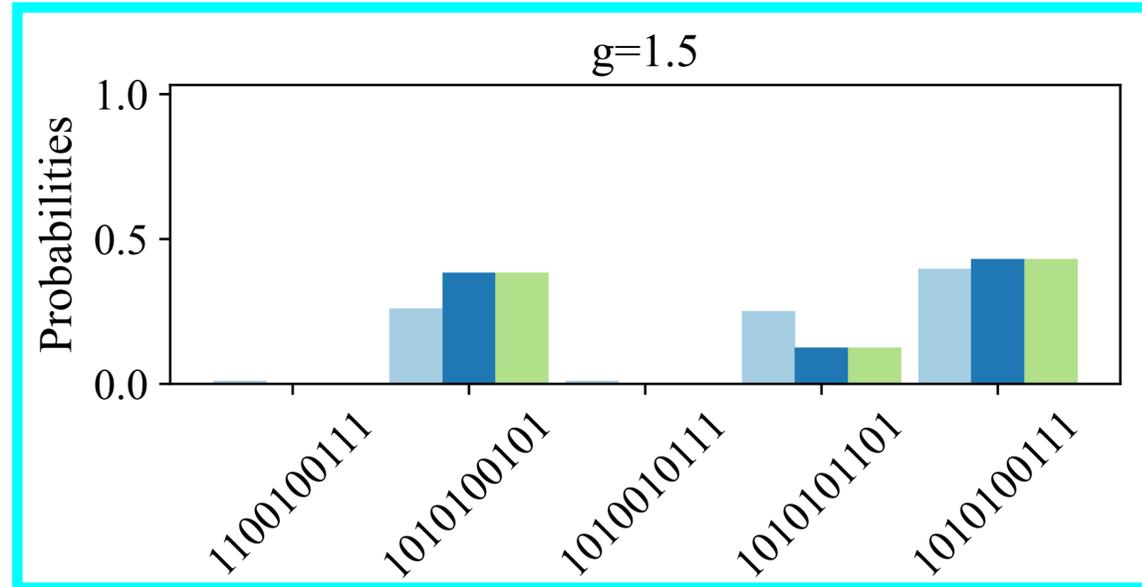
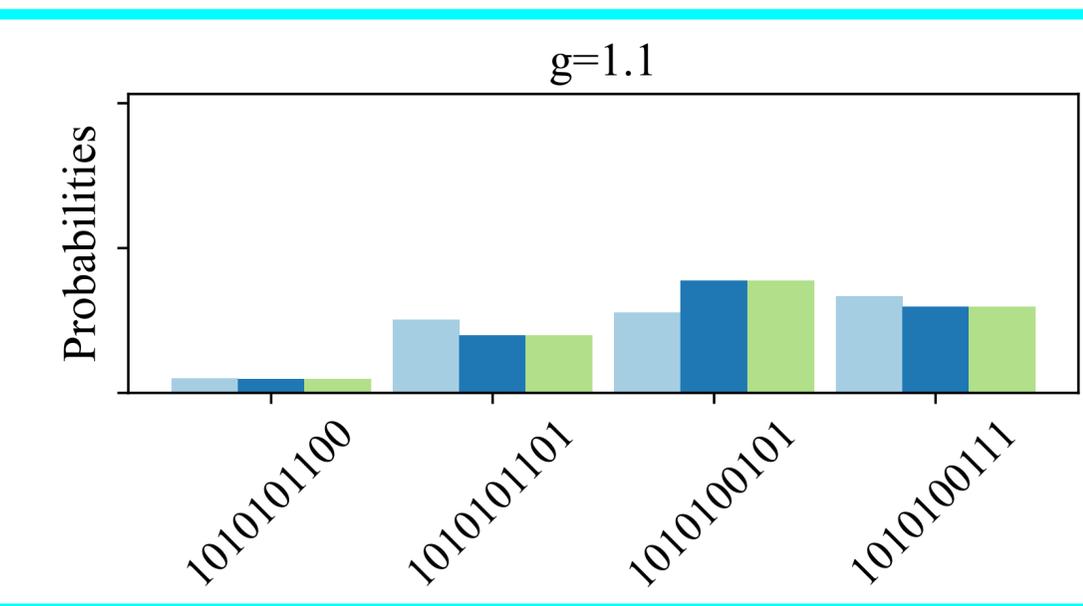
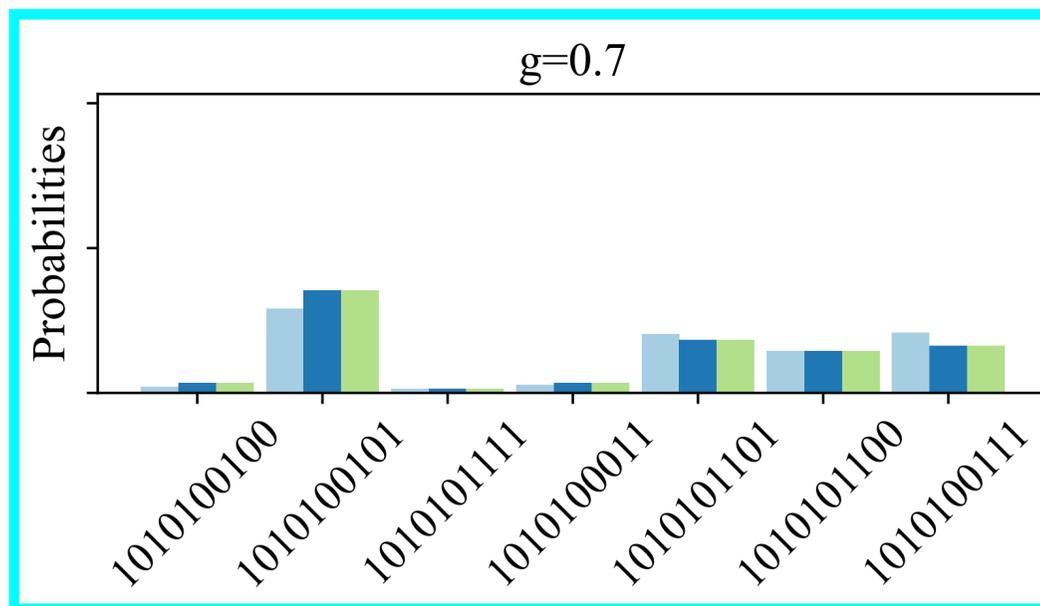
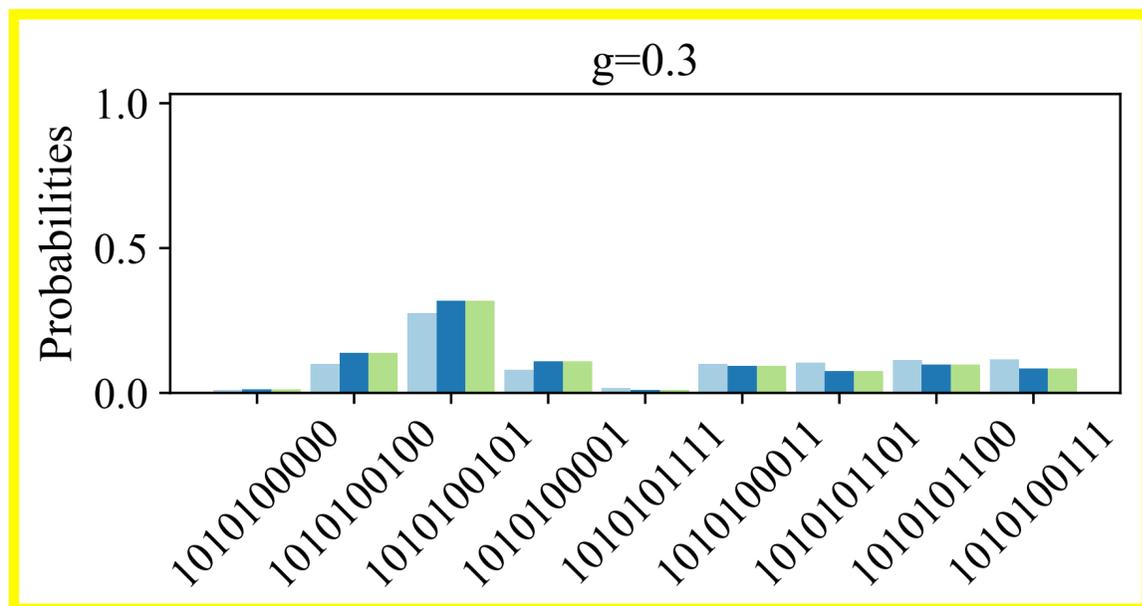
ED VQE AerStateBackend

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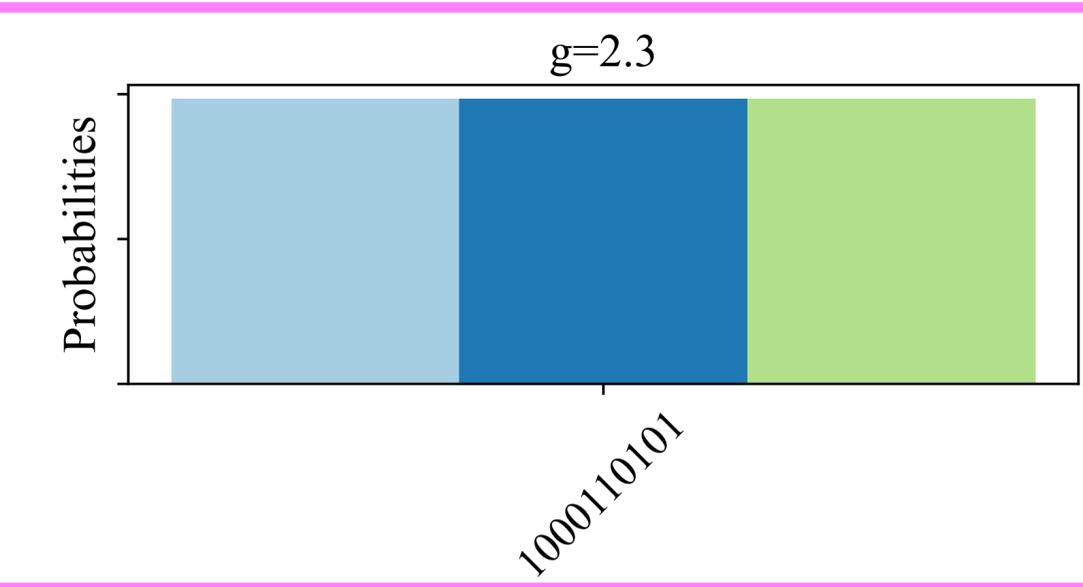
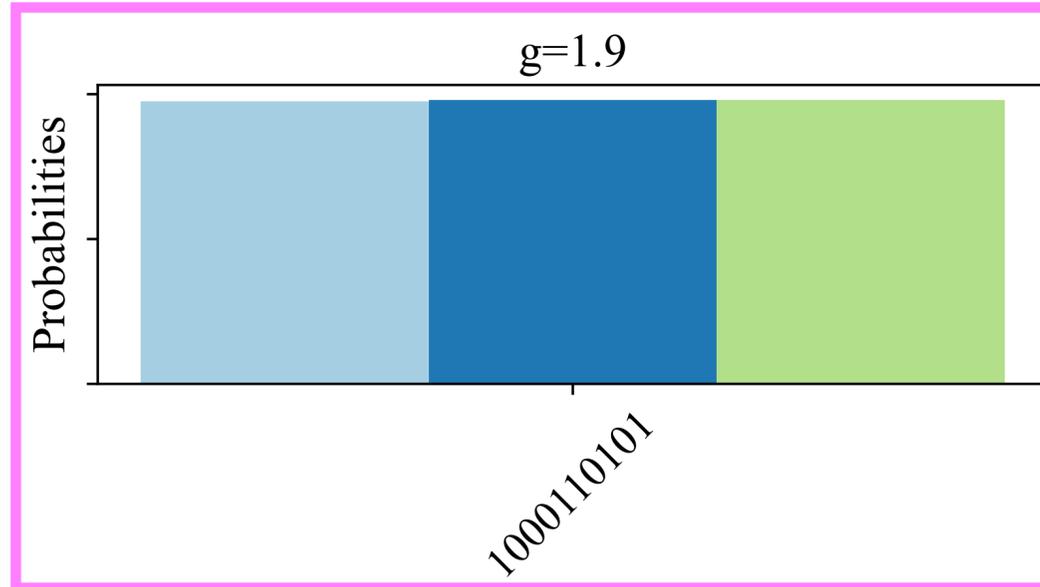
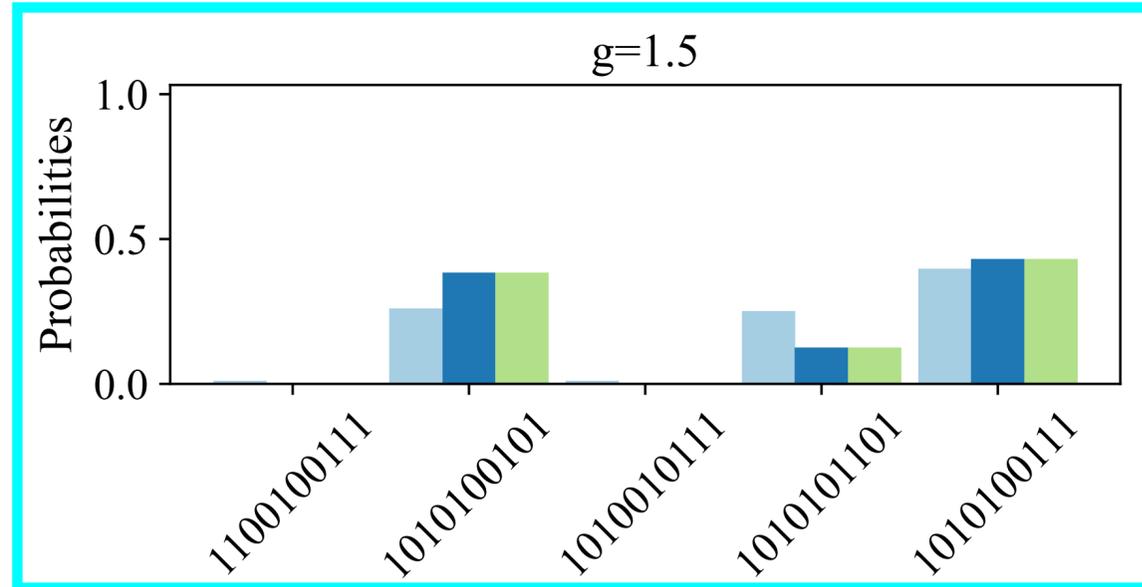
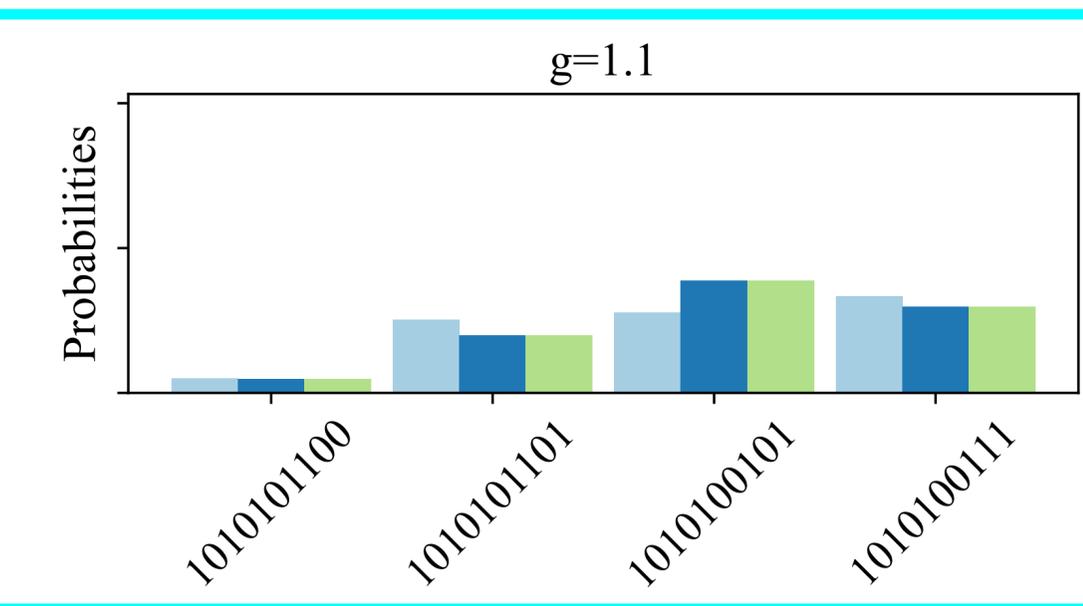
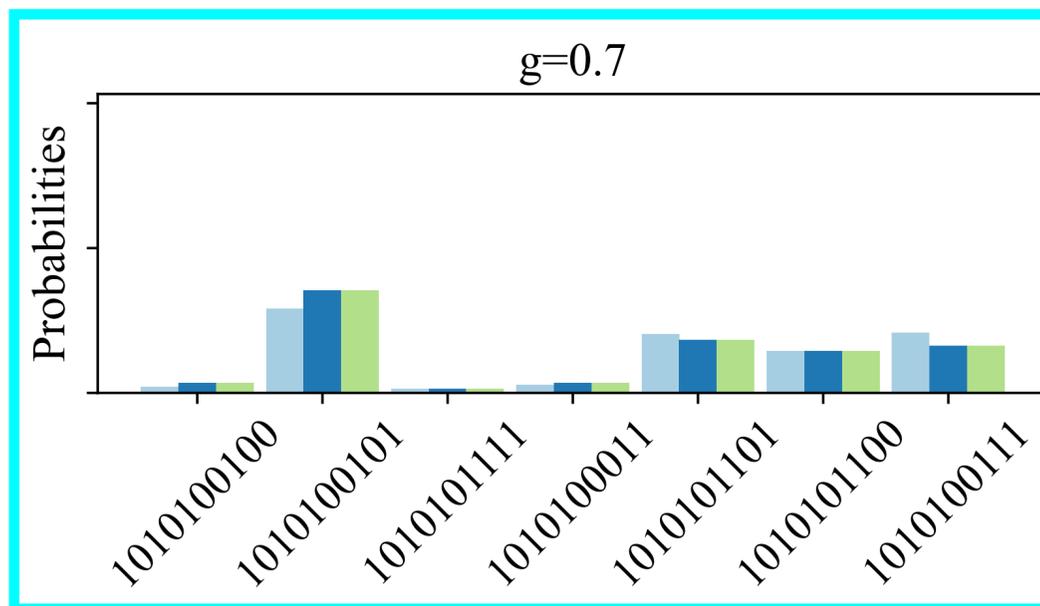
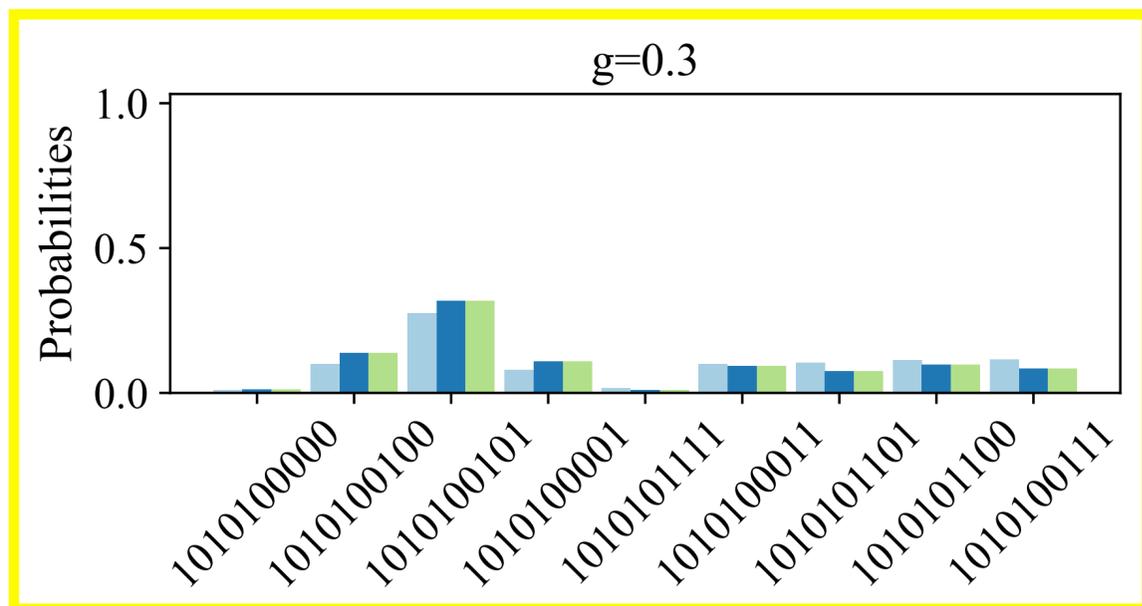
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ED VQE AerStateBackend

$|\Psi_{GS}\rangle = c_0 |1000110101\rangle + \dots$

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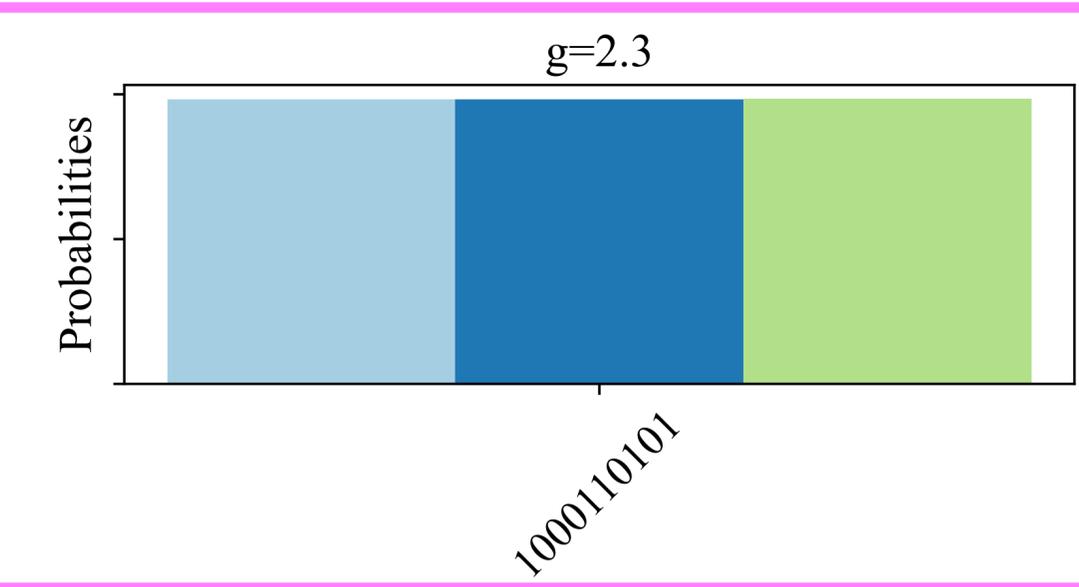
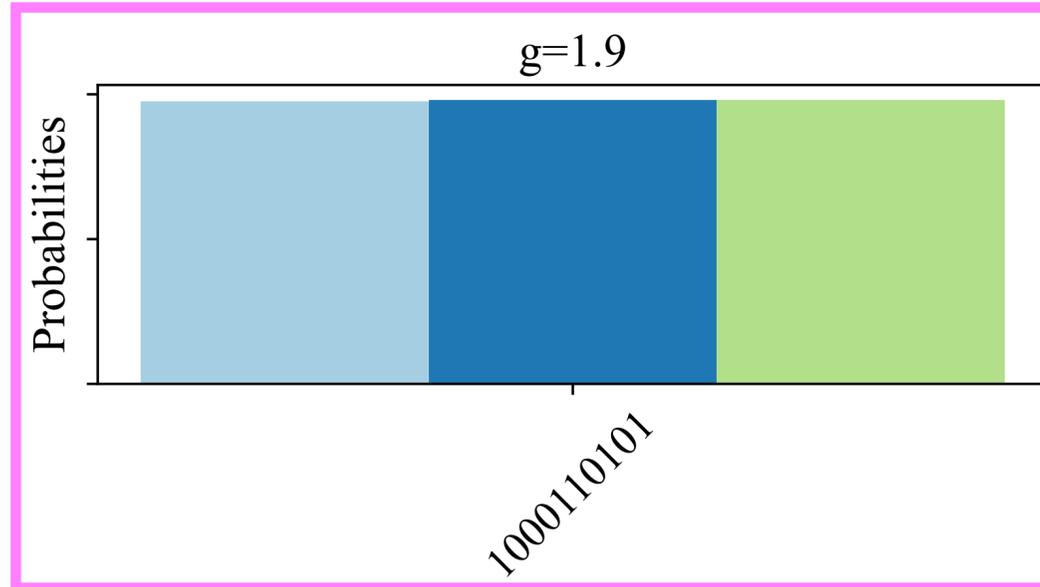
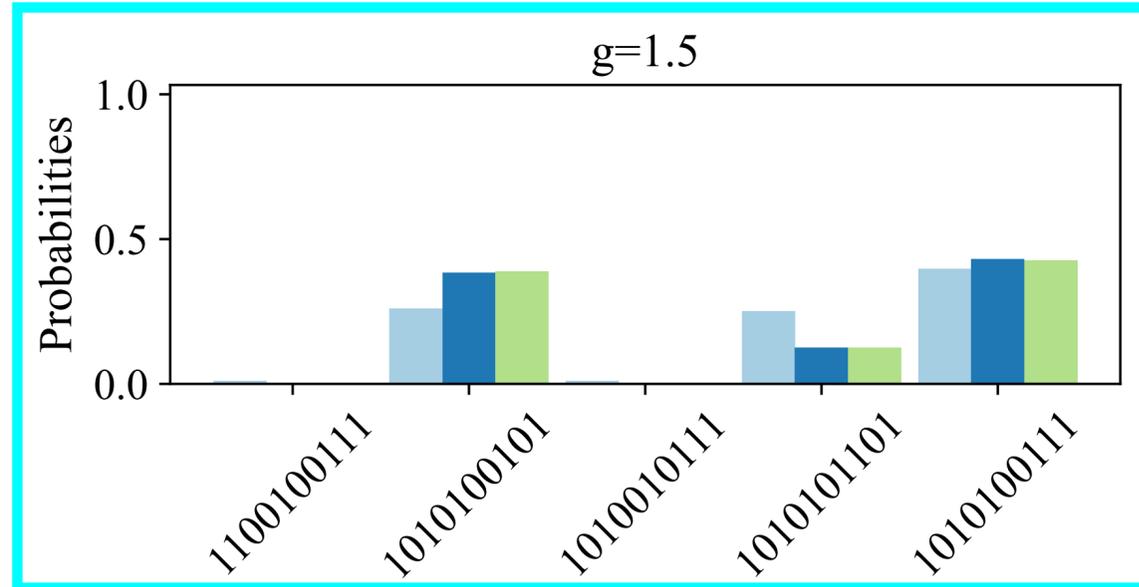
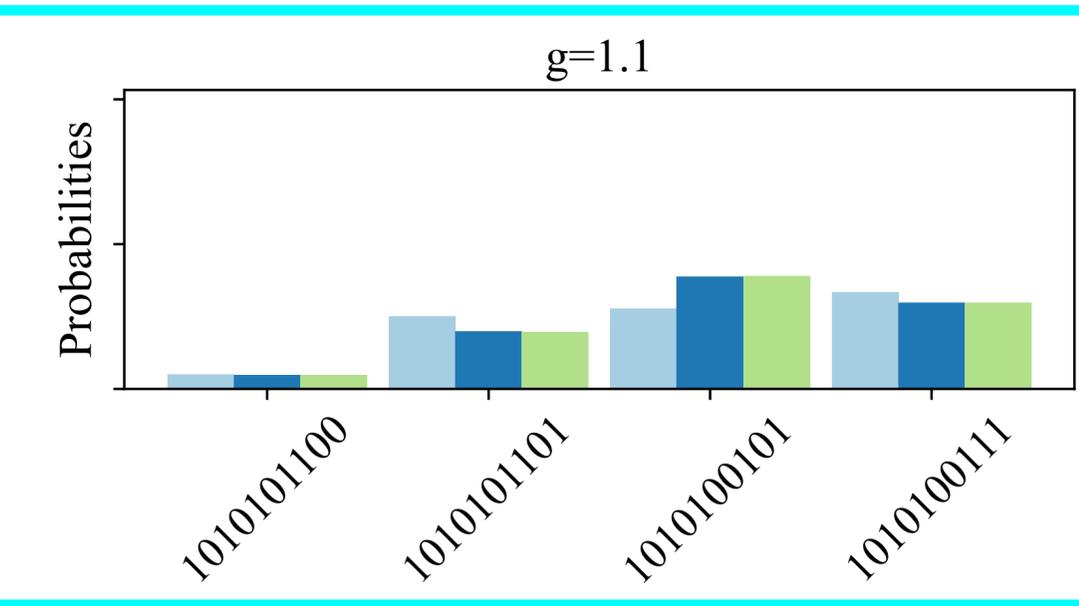
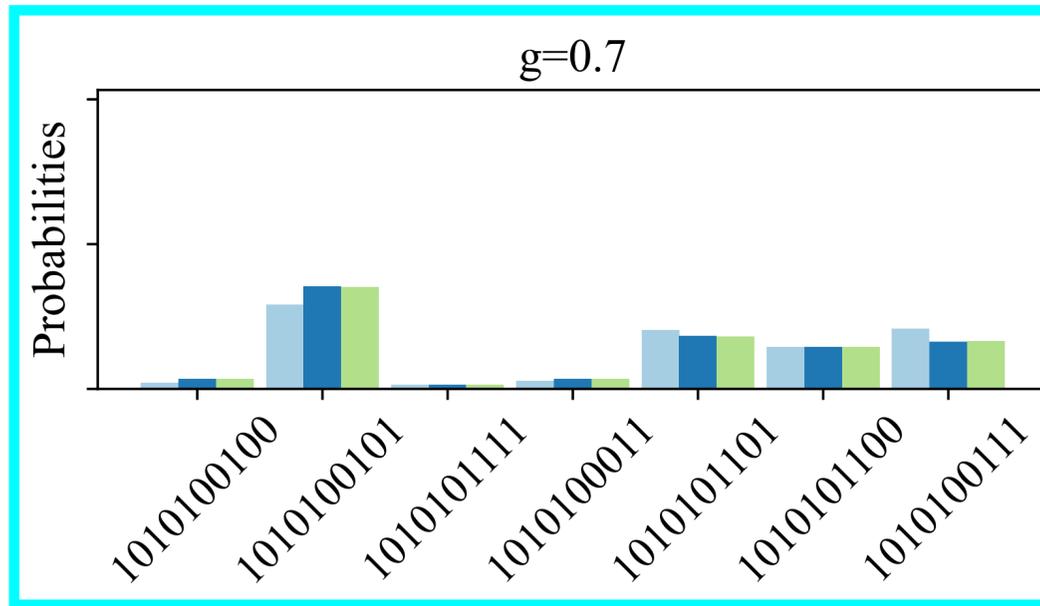
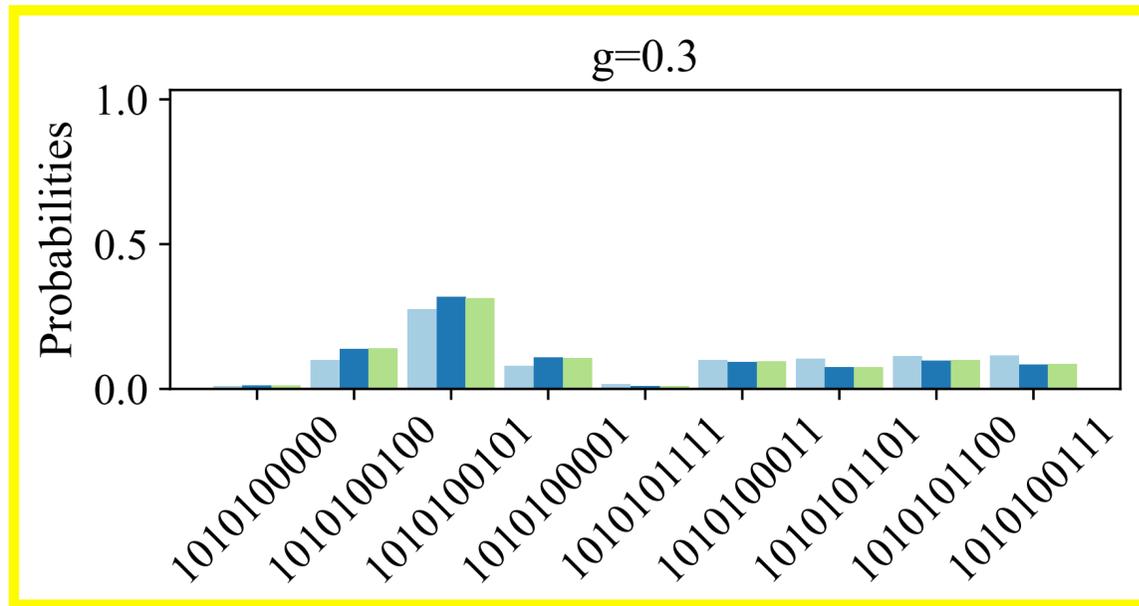


$|\Psi_{GS}\rangle = c_0 |1010100101\rangle + c_1 |1010101101\rangle + c_2 |1010100111\rangle$

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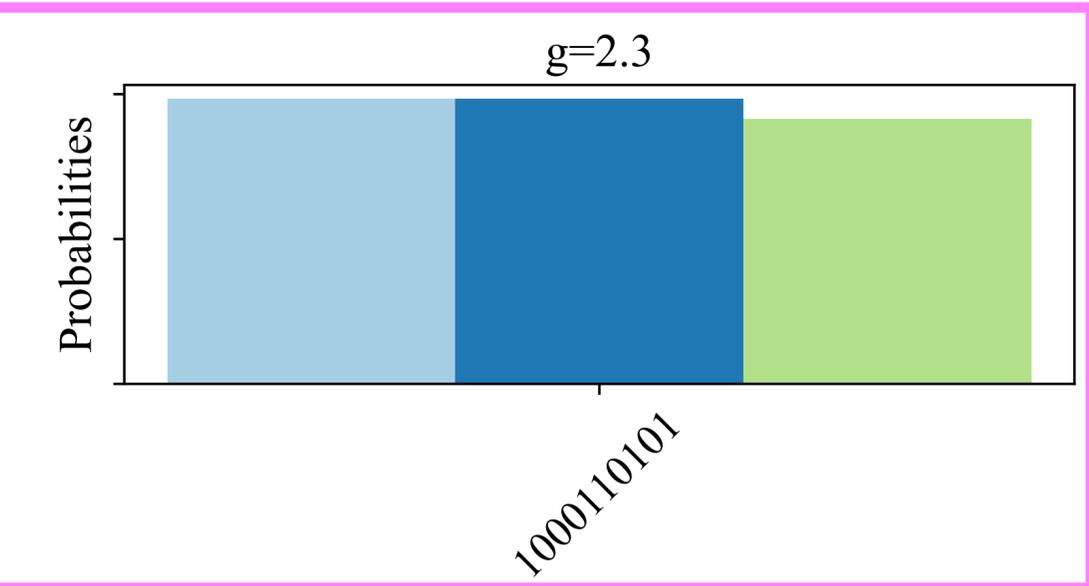
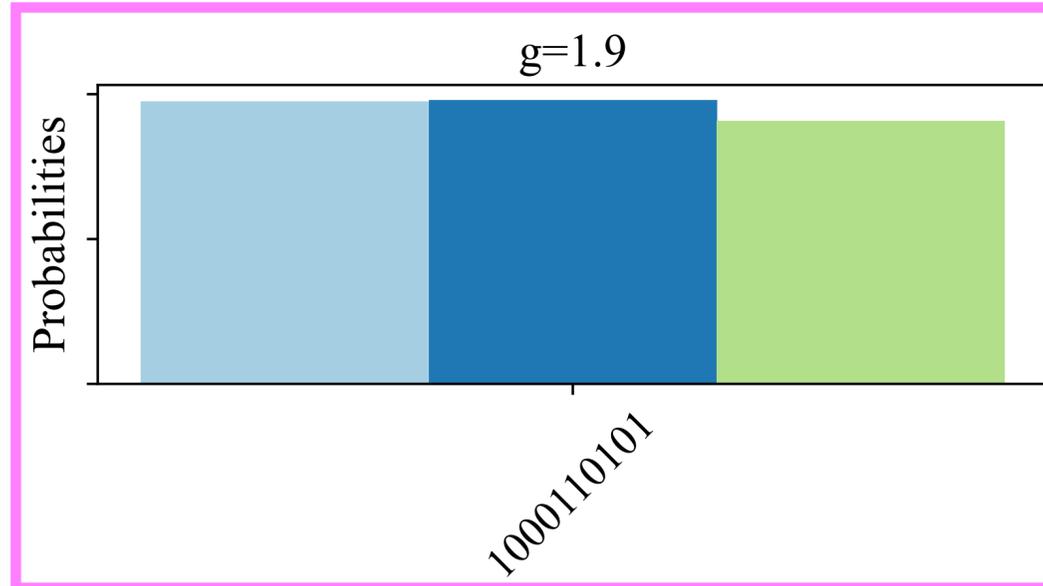
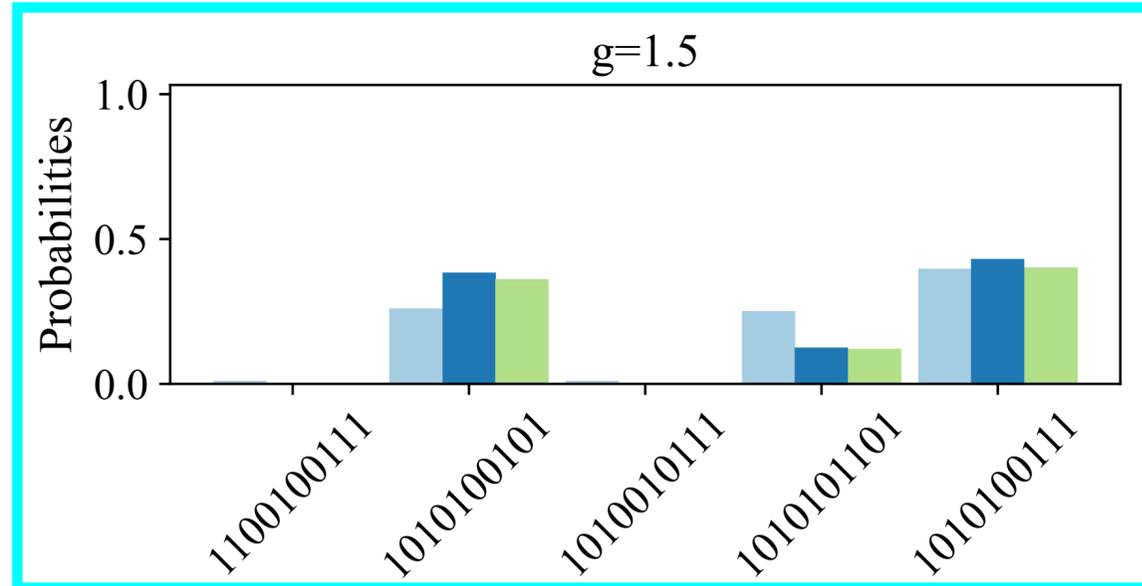
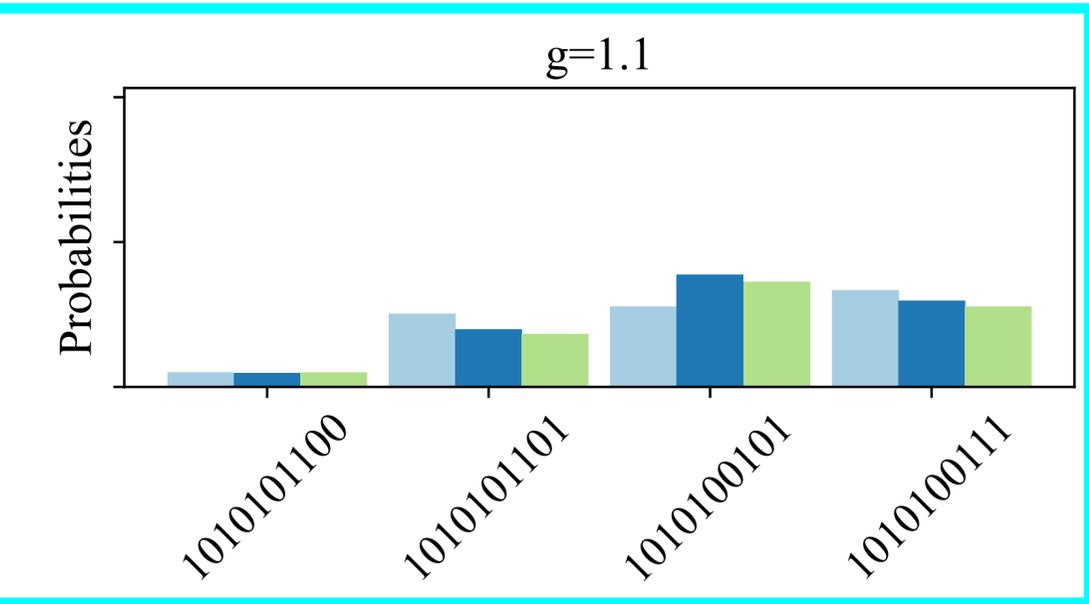
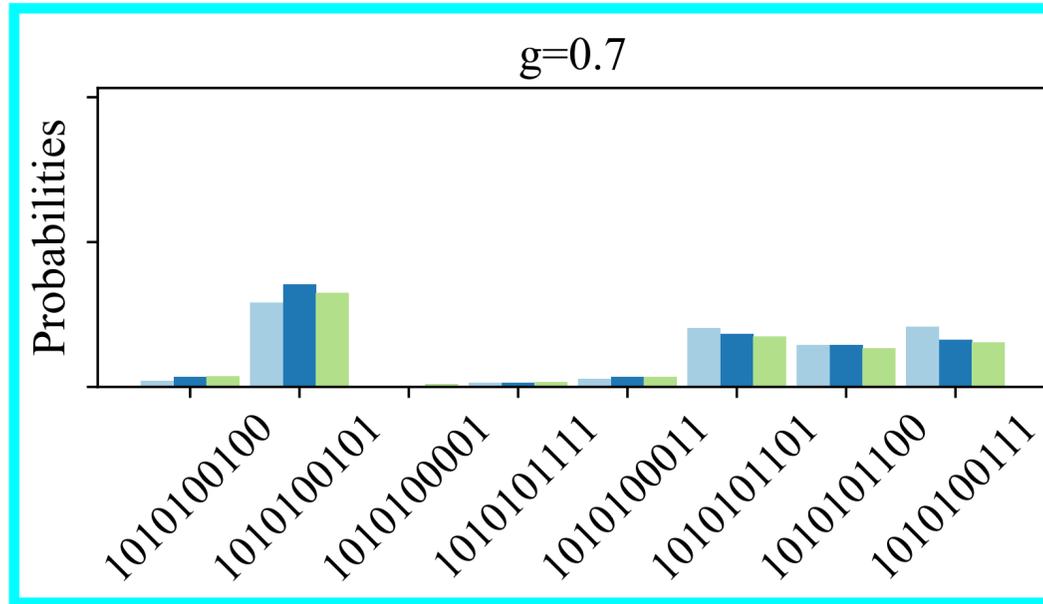
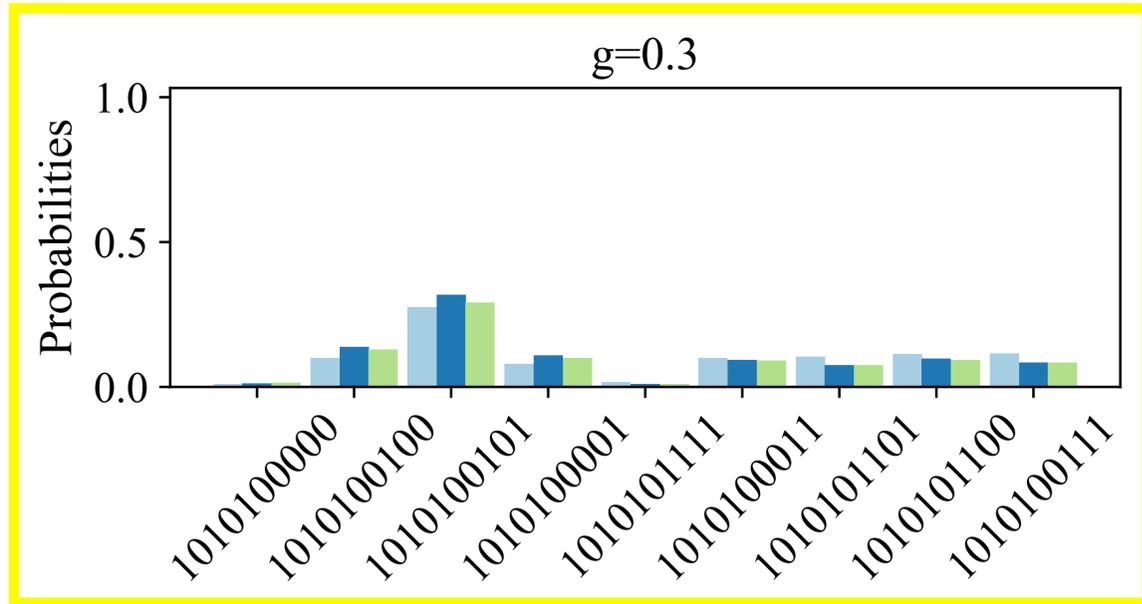
Probabilities of the different states for each g value: AerBackend shots=100000



ED VQE AerBackend

Small Noise Levels

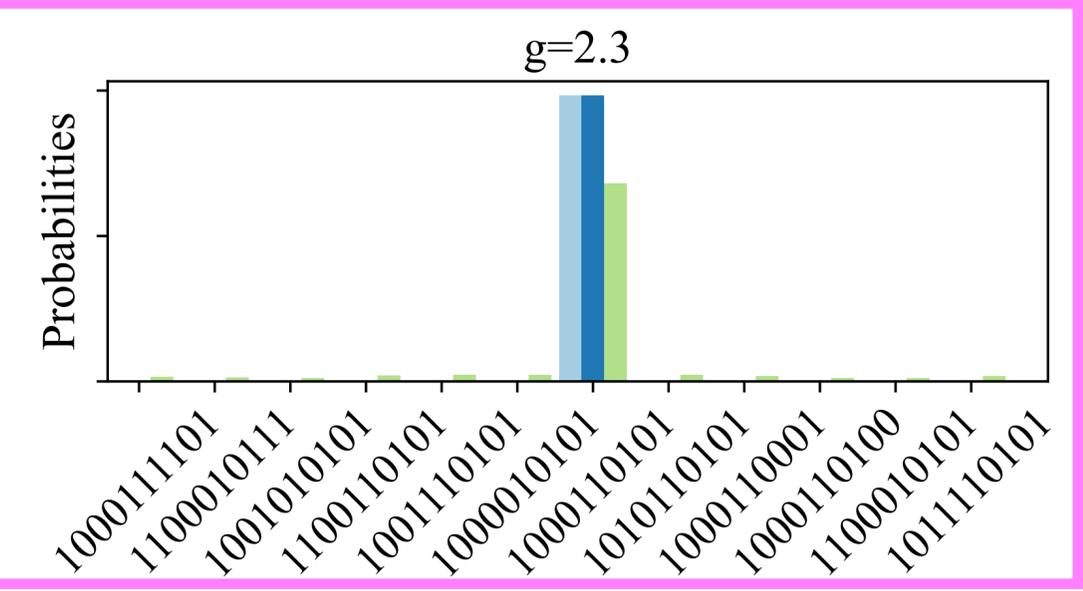
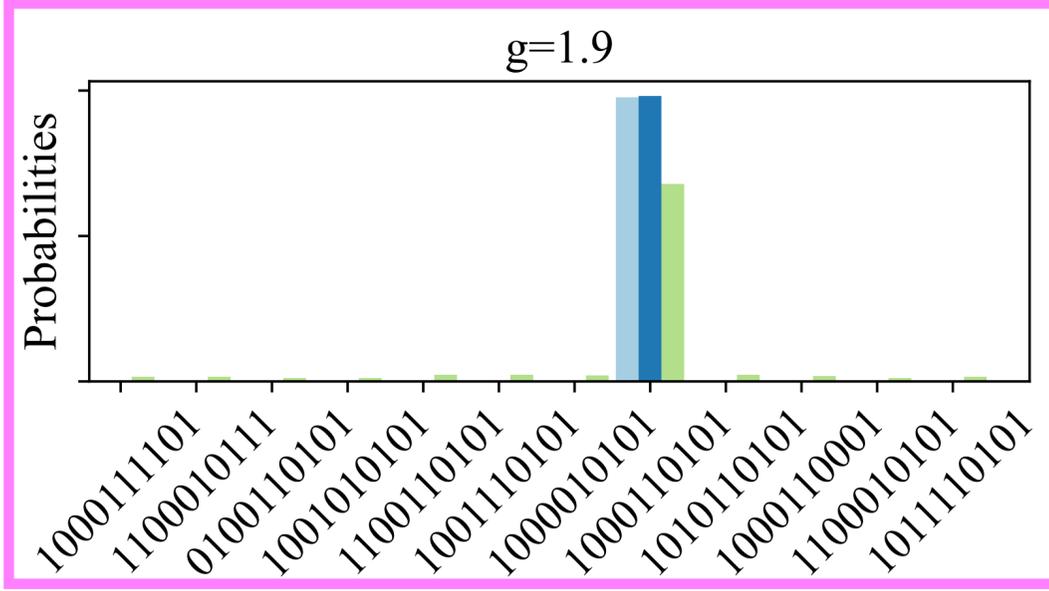
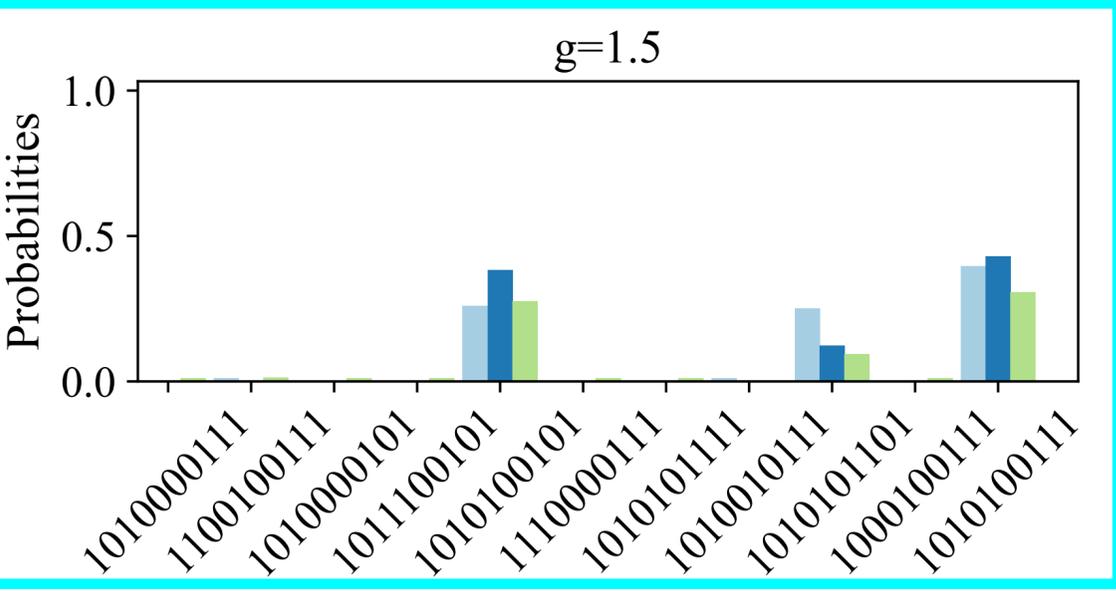
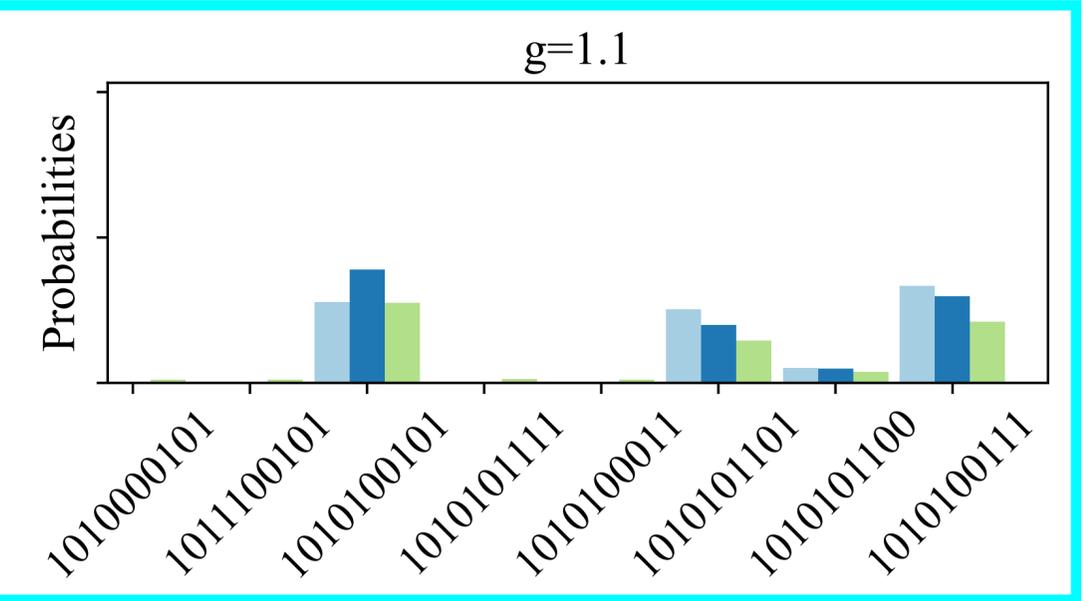
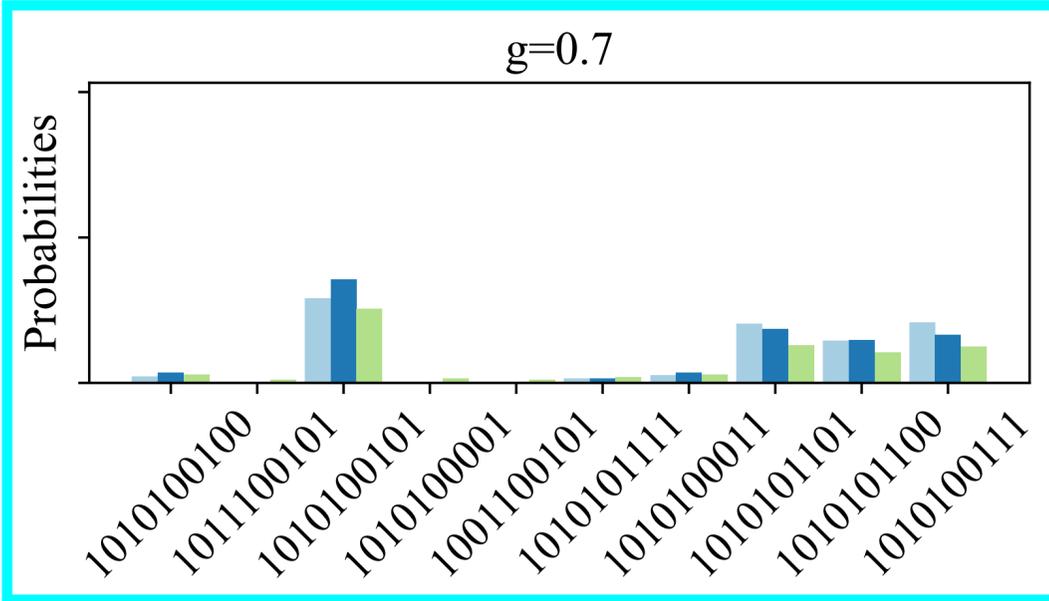
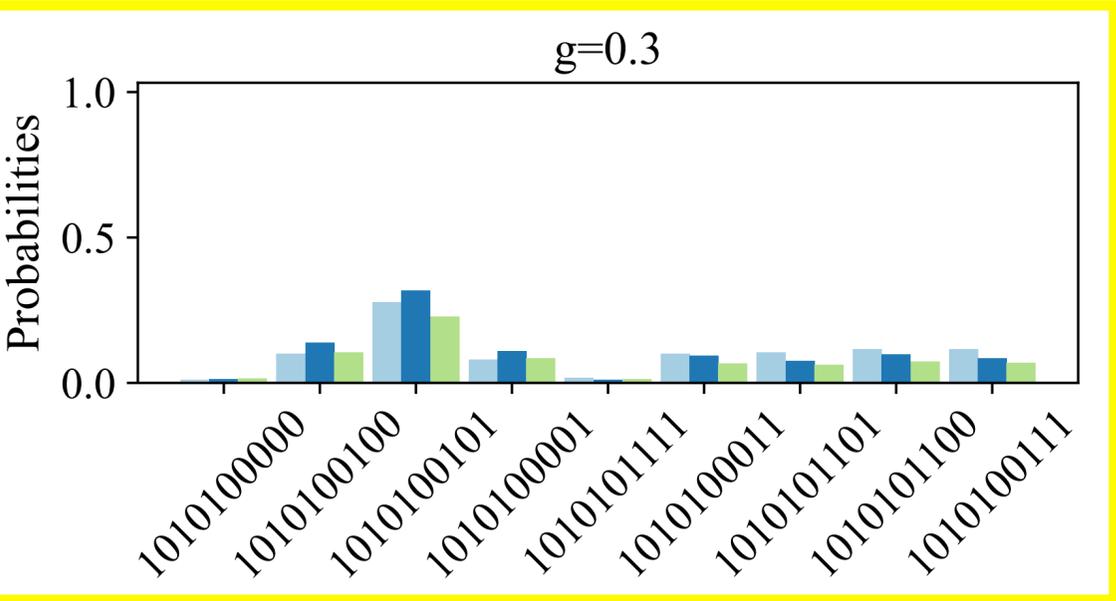
Probabilities of the different states for each g value: AerBackend shots=10000



ED VQE AerBackend(noise)

Large Noise Levels

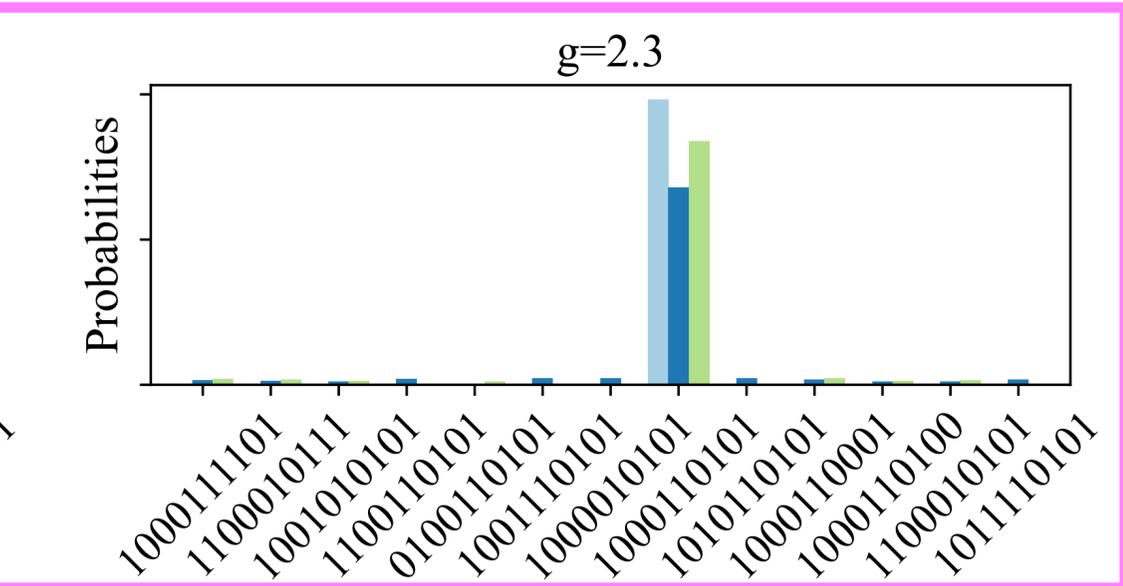
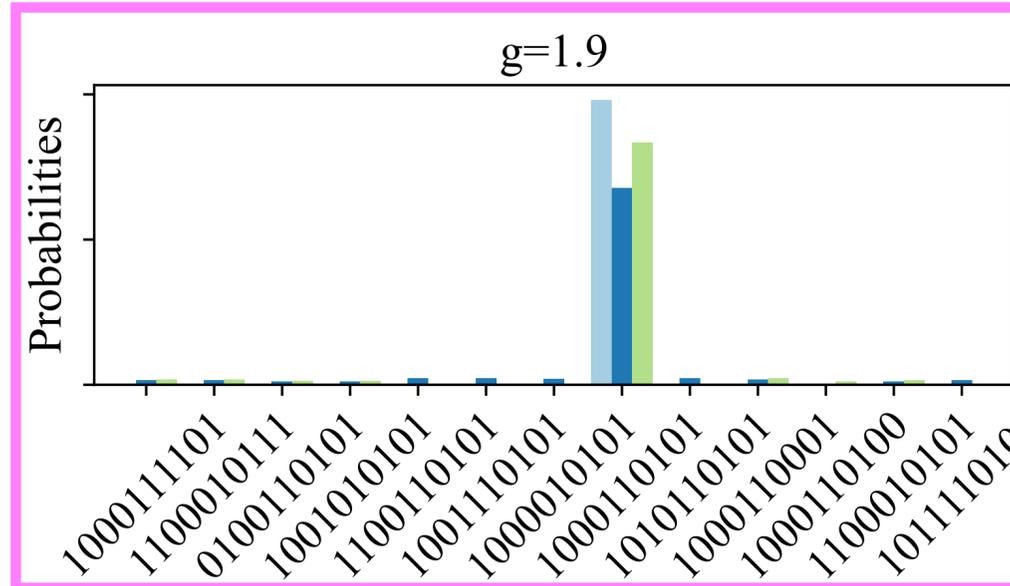
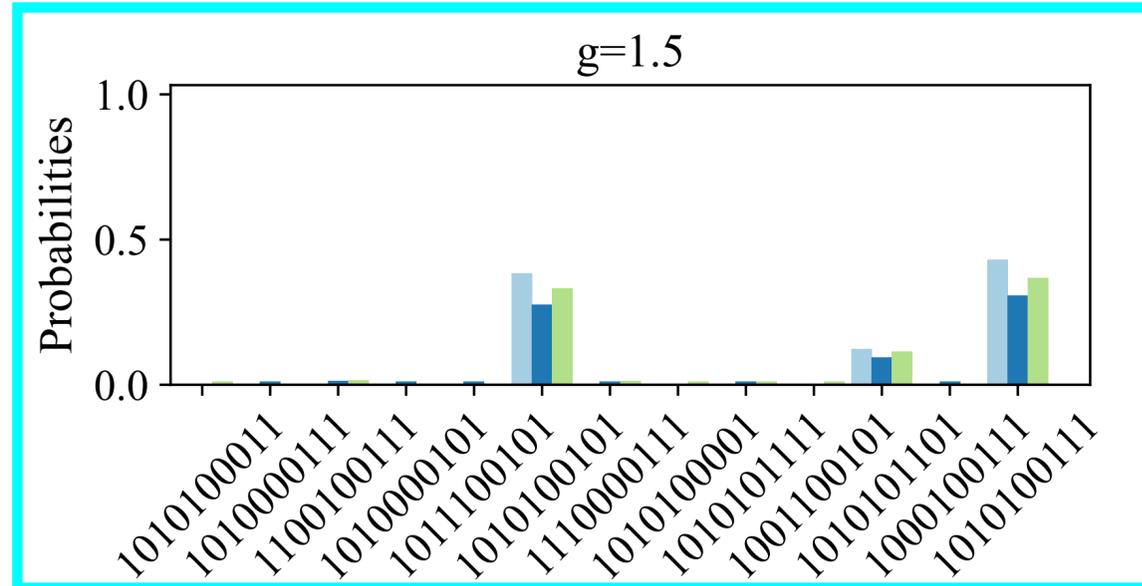
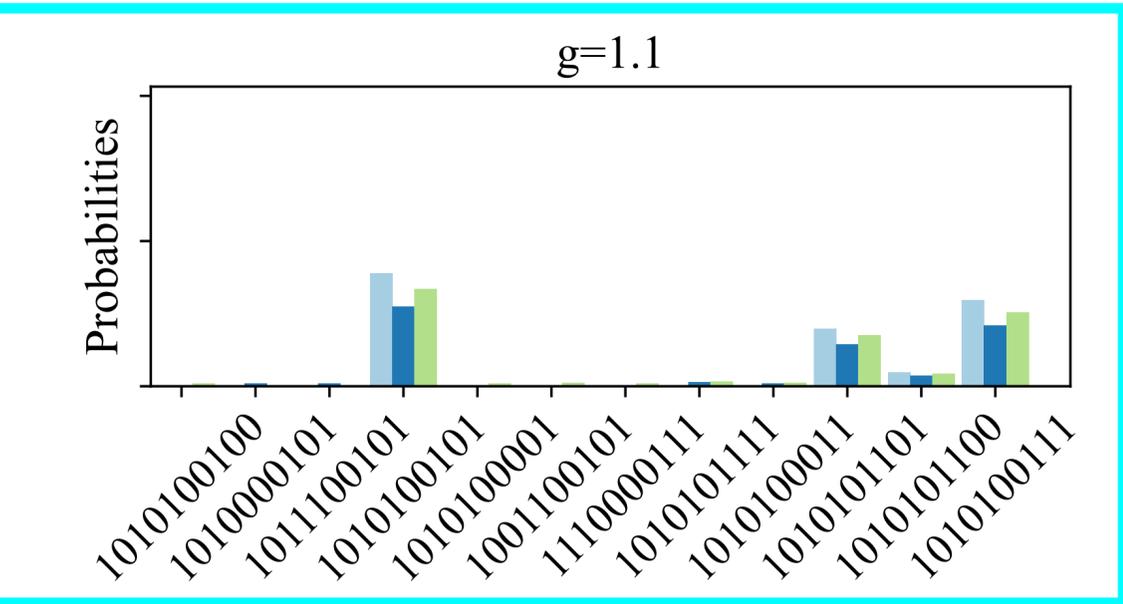
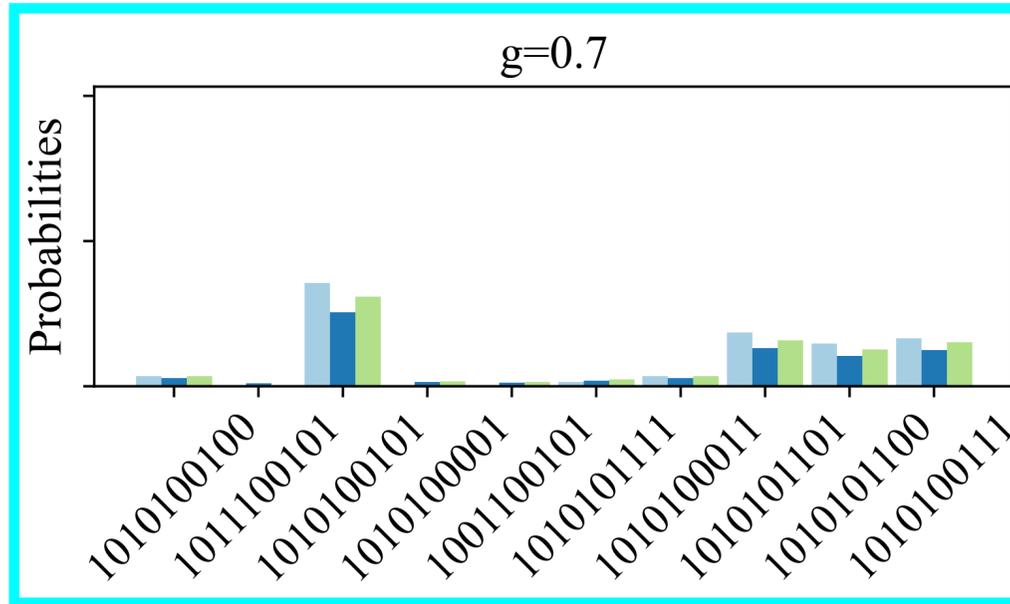
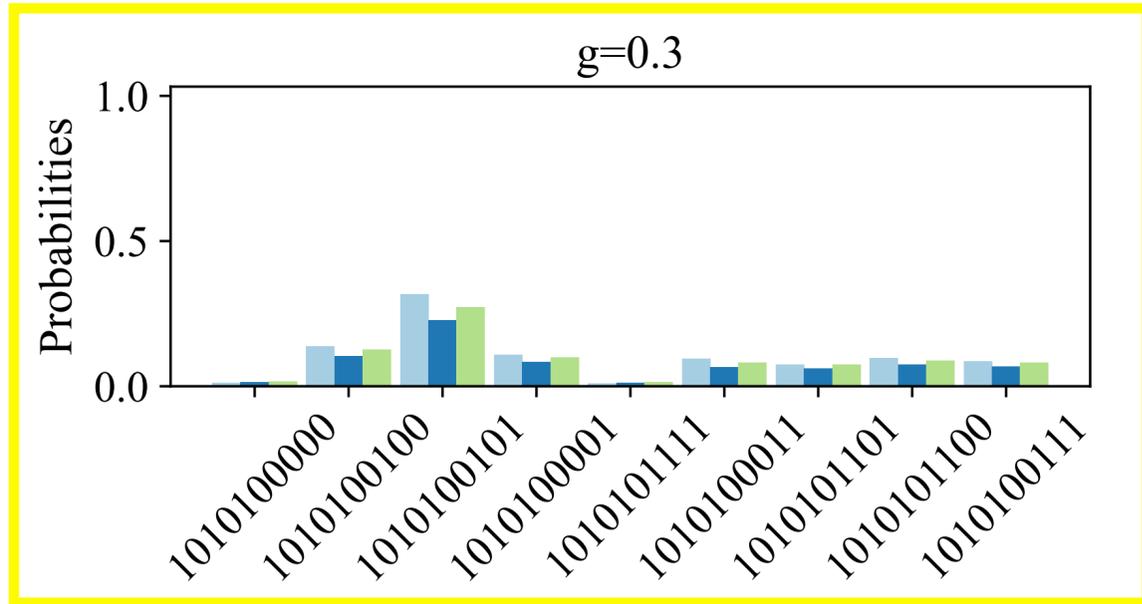
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ED VQE AerBackend(noise)

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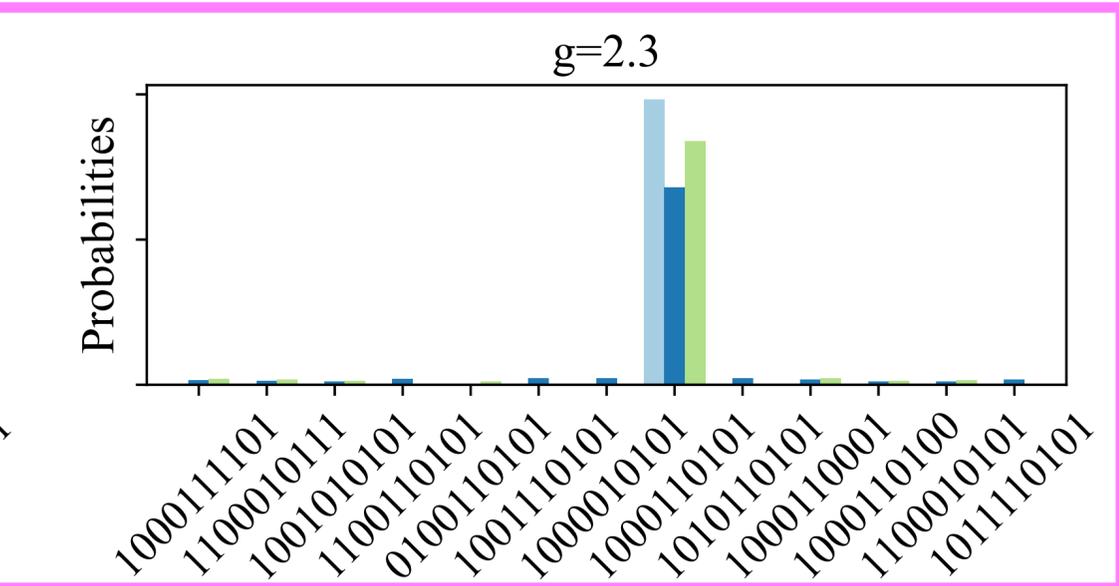
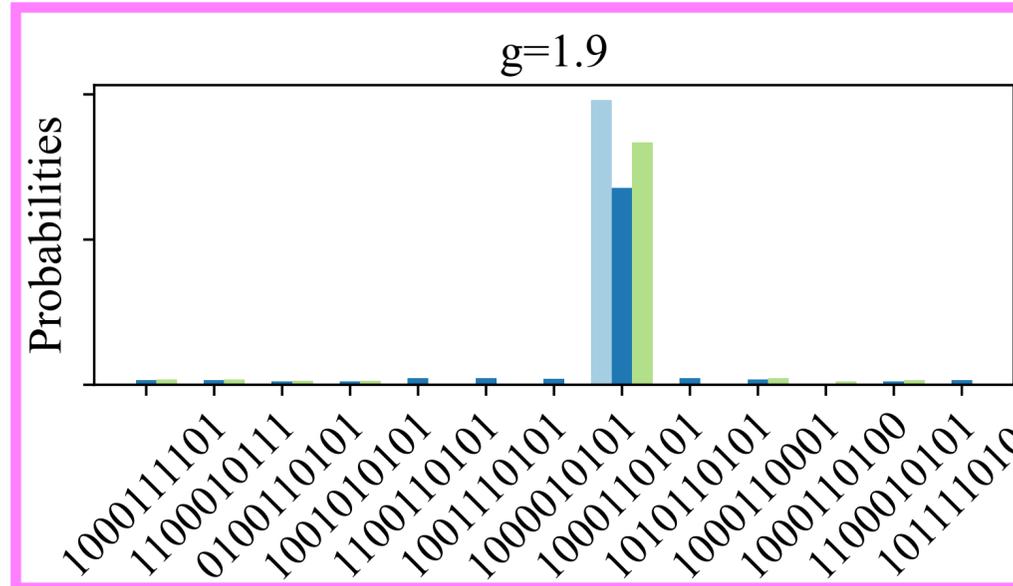
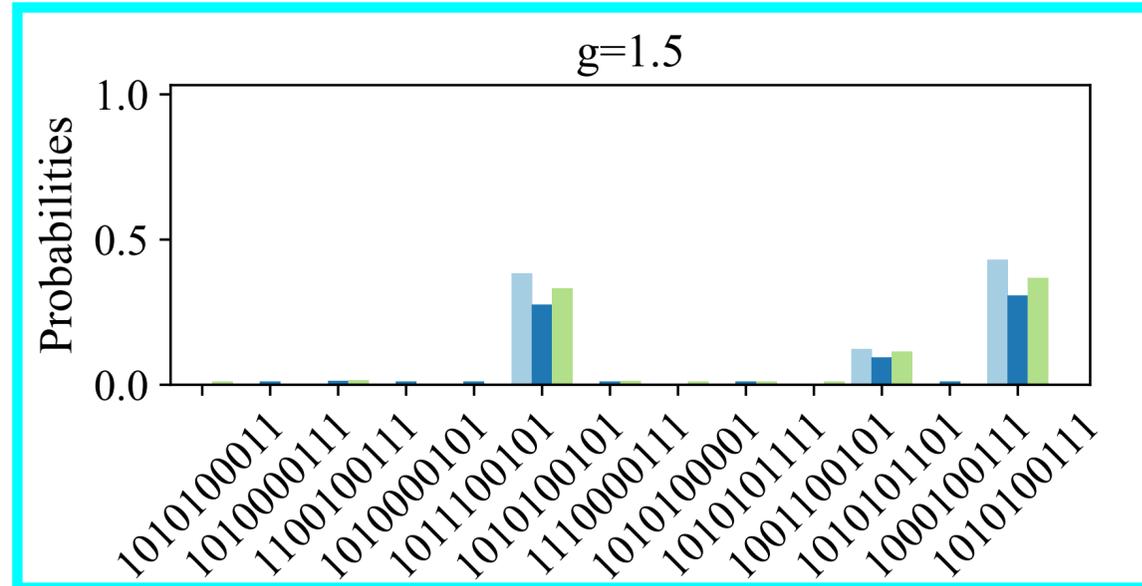
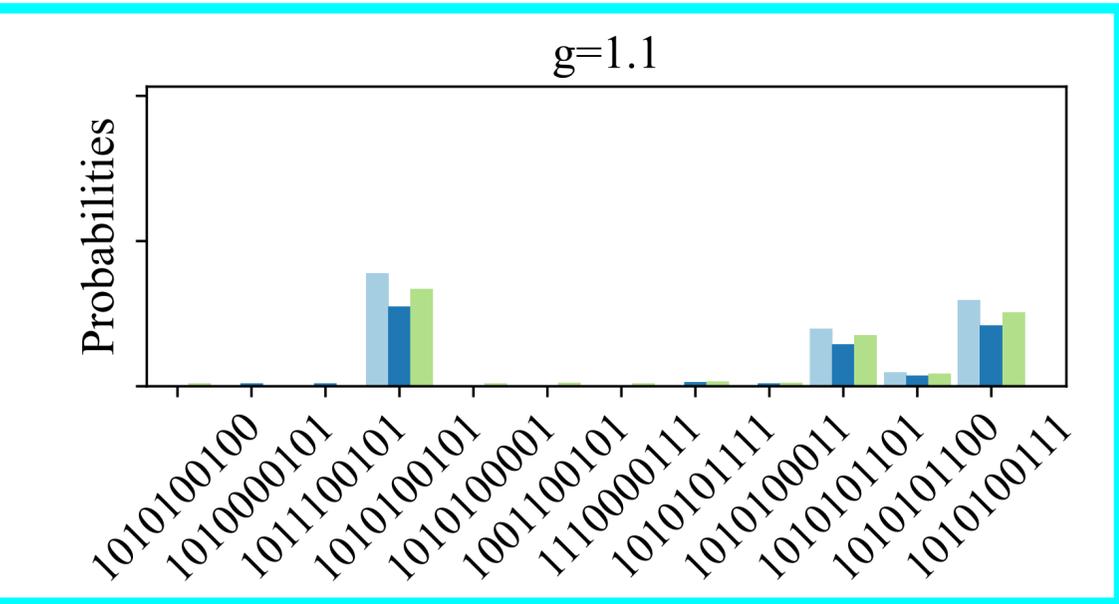
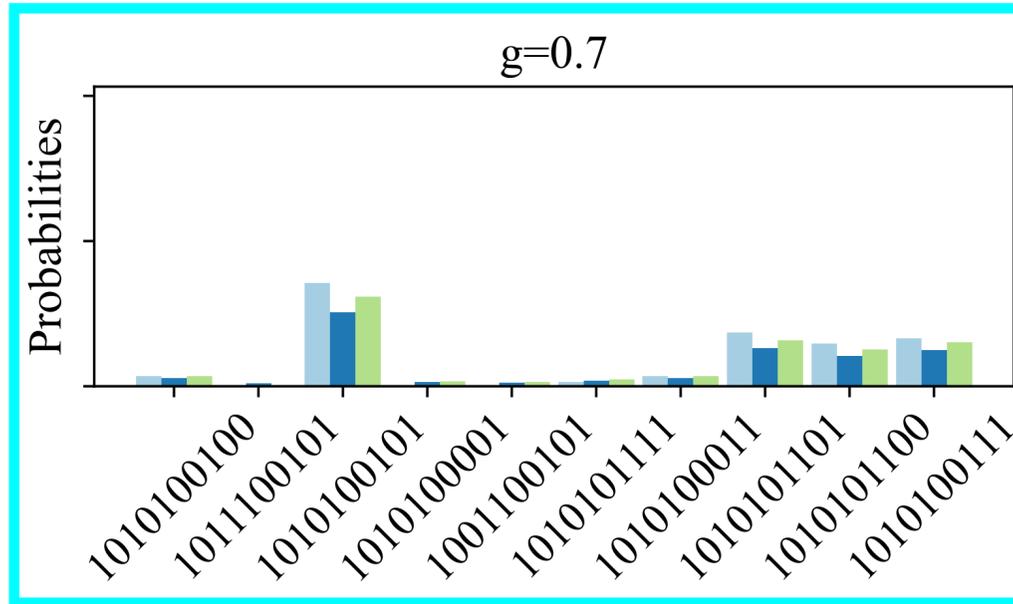
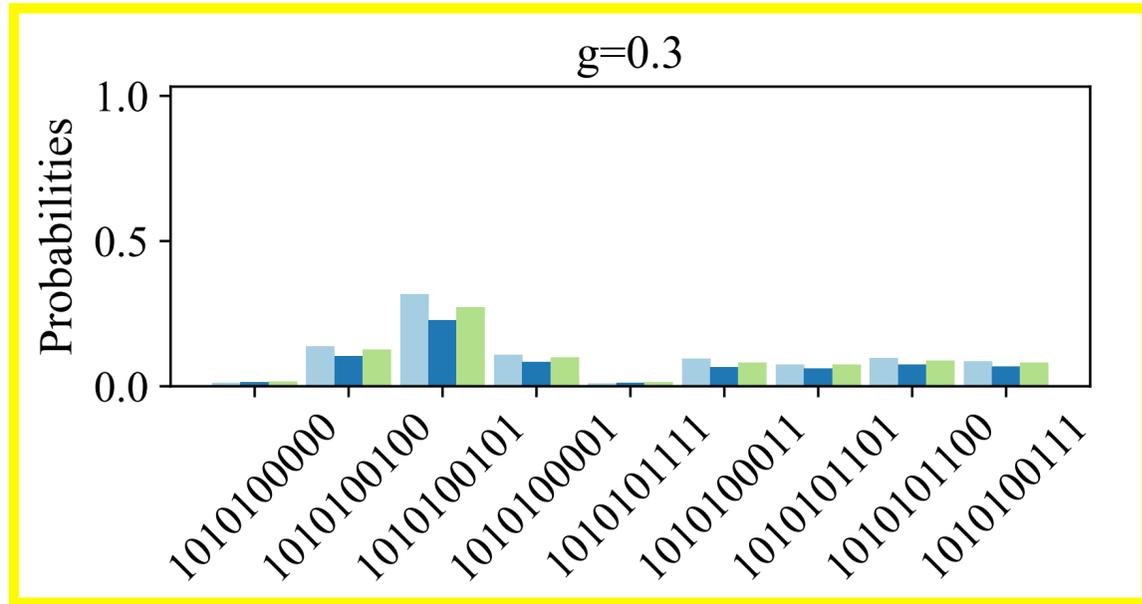
Probabilities of the different states for each g value: AerBackend shots=10000



VQE
 AerBackend(noise)
 Mitigated

Large Noise Levels

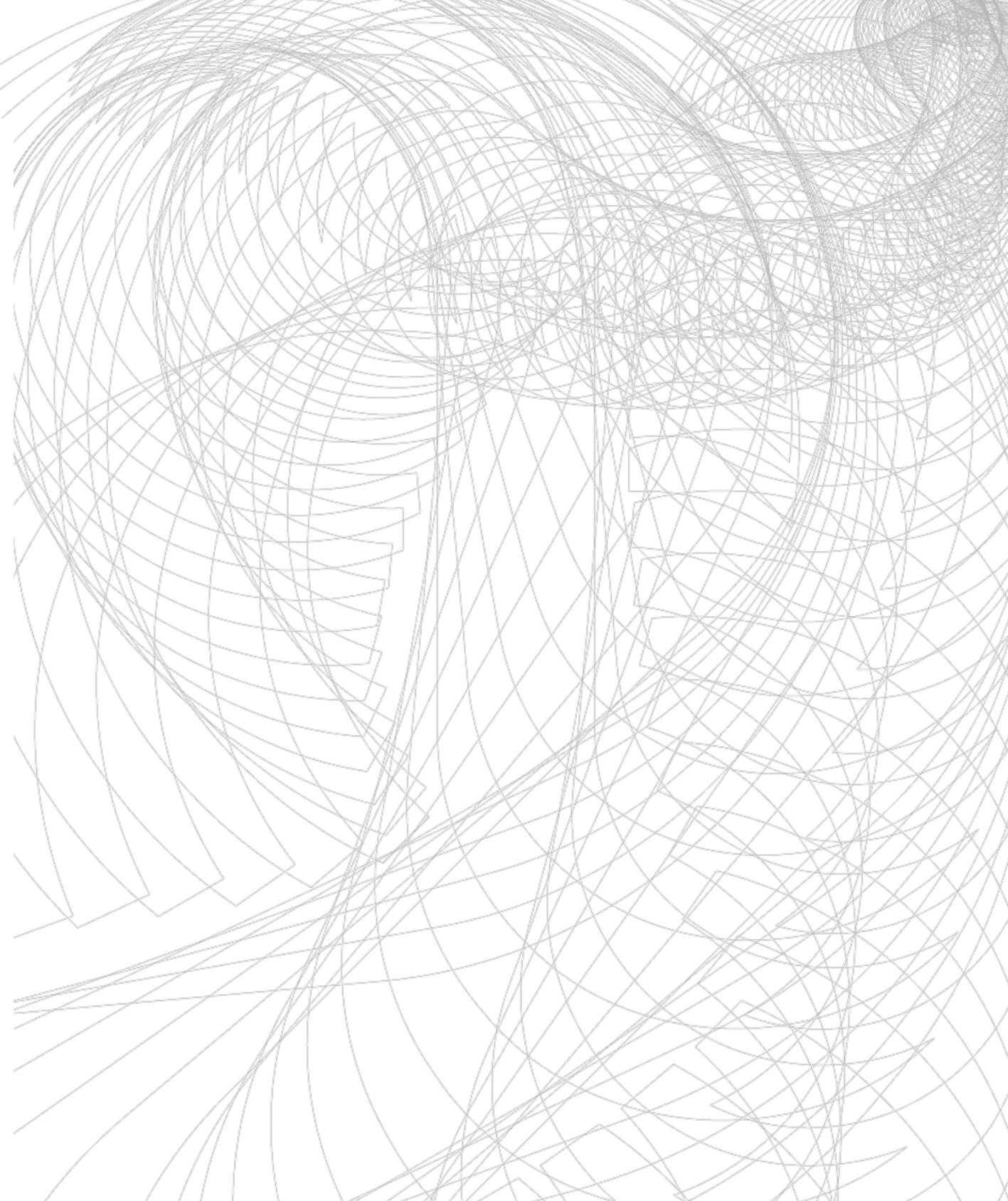
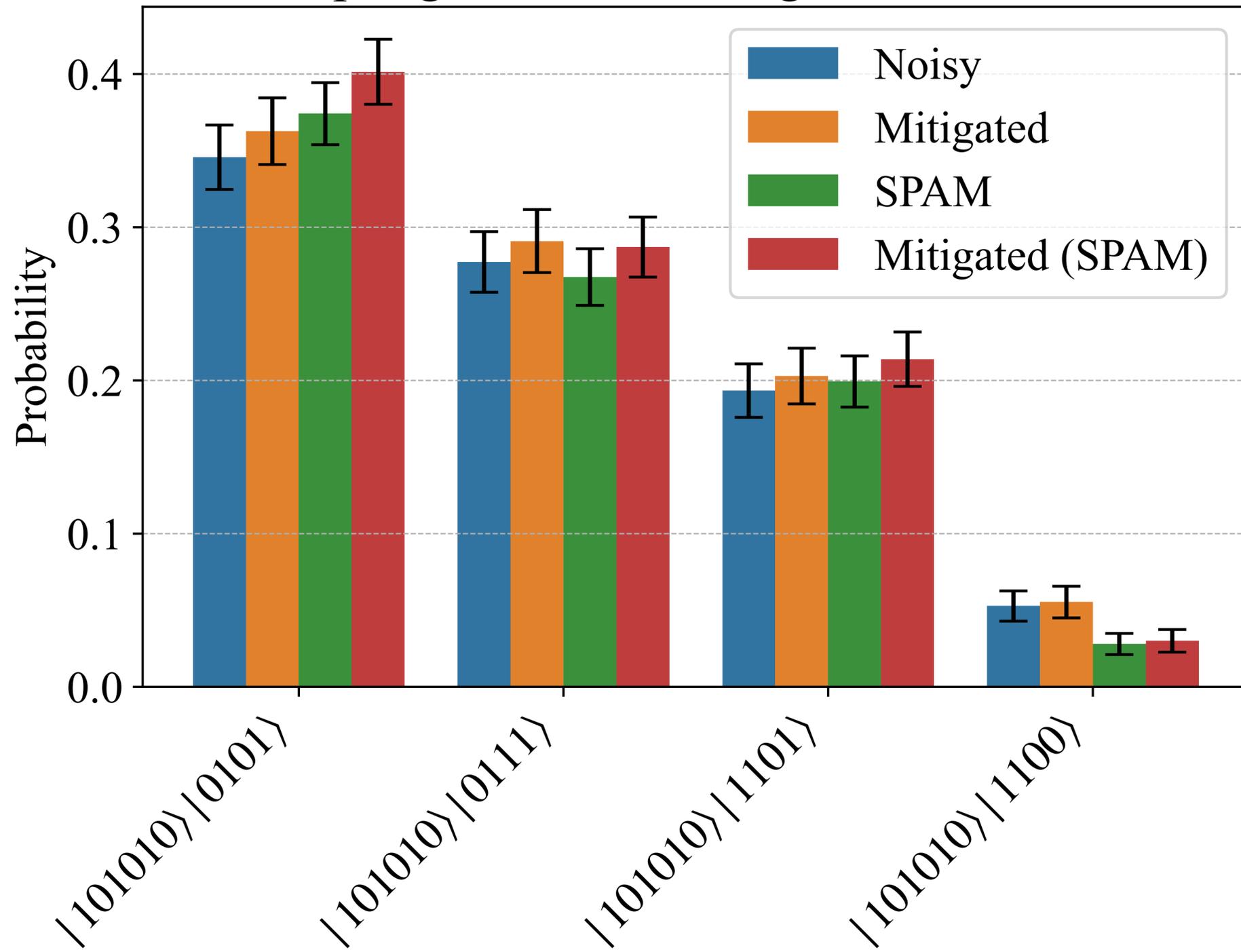
Probabilities of the different states for each g value: AerBackend shots=10000



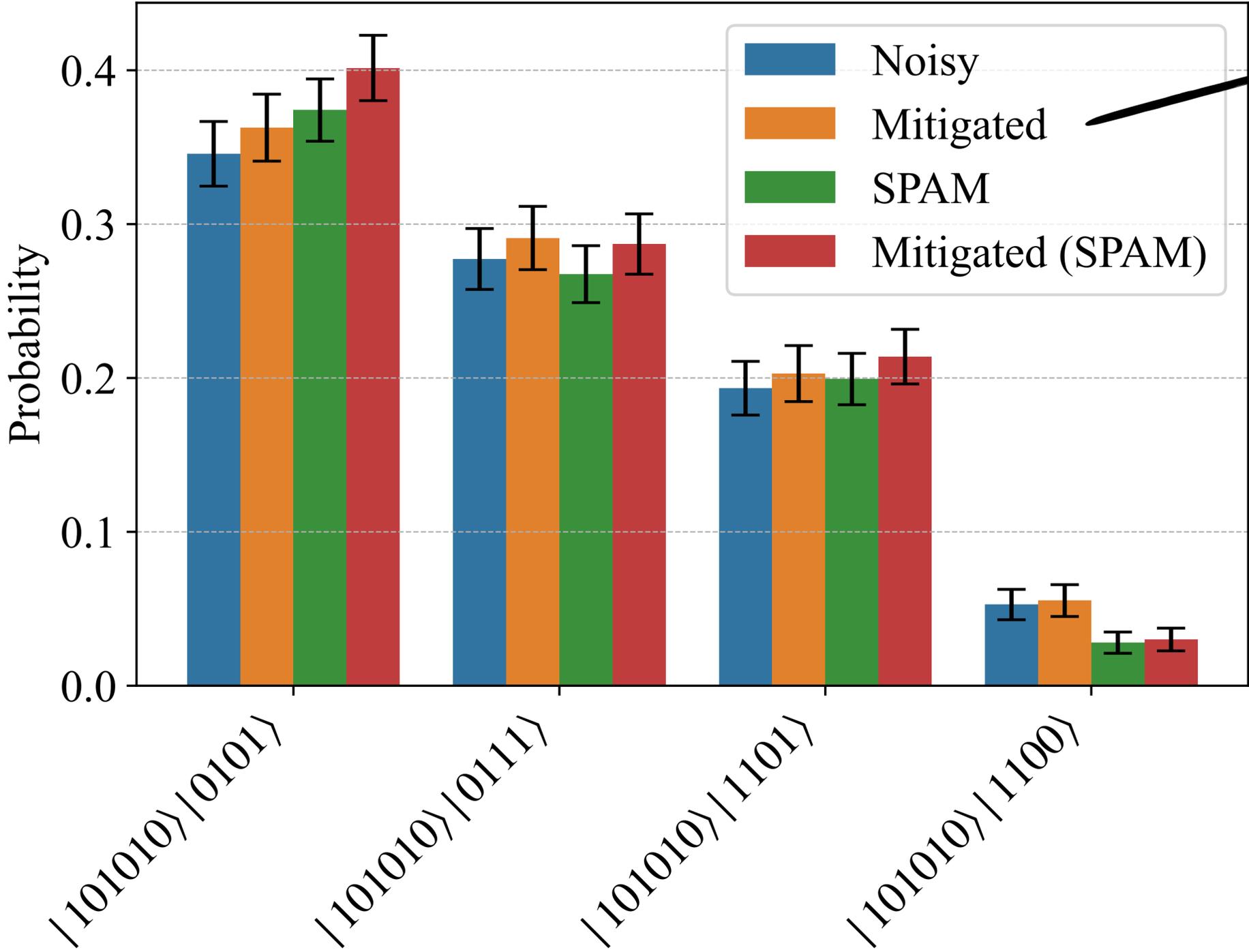
Remove unphysical states/bitstrings

Emulations on H1-1E

Sampling with H1-1E at $g=1.1$, shots=512

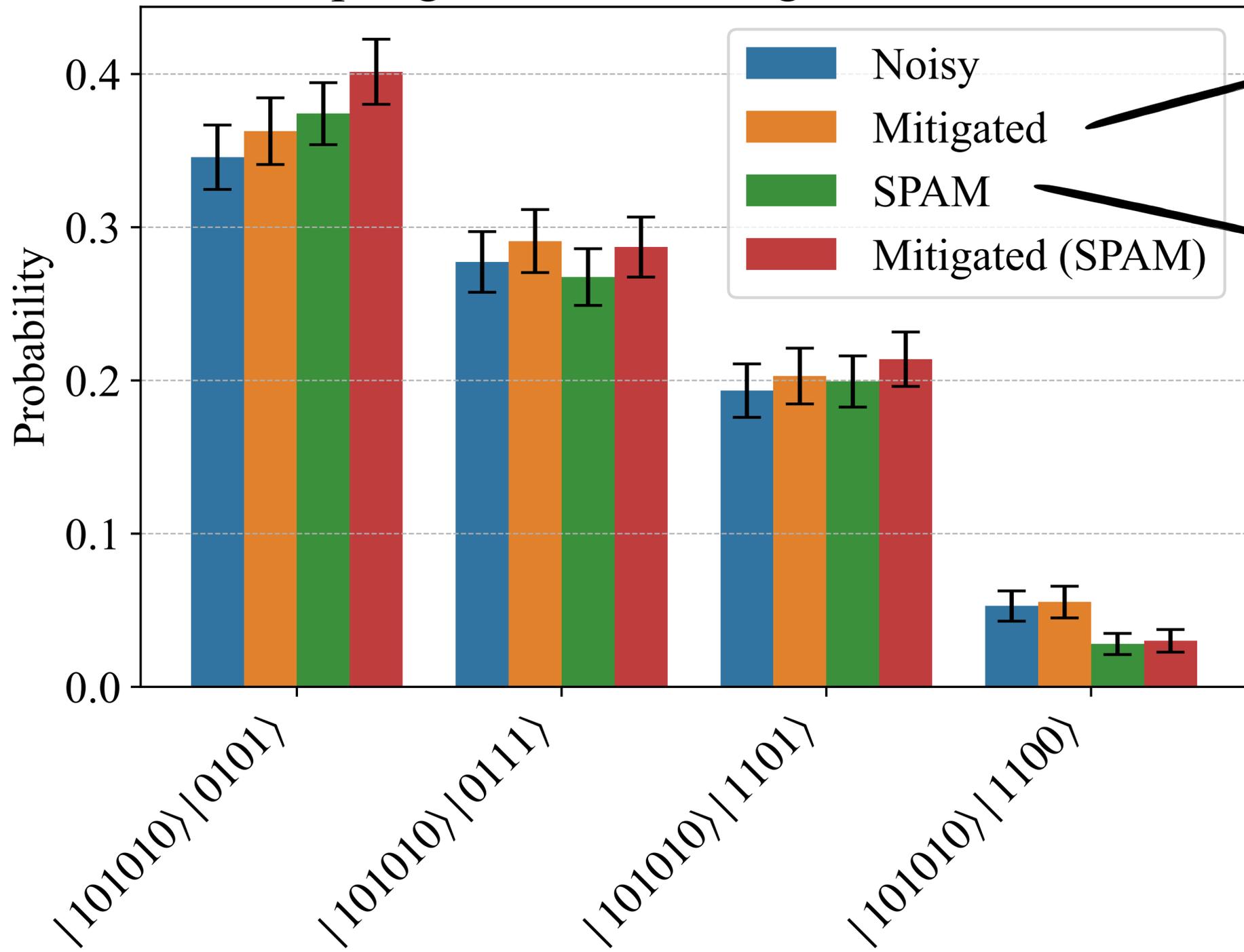


Sampling with H1-1E at $g=1.1$, shots=512



Remove bit strings that do not correspond to physical configurations

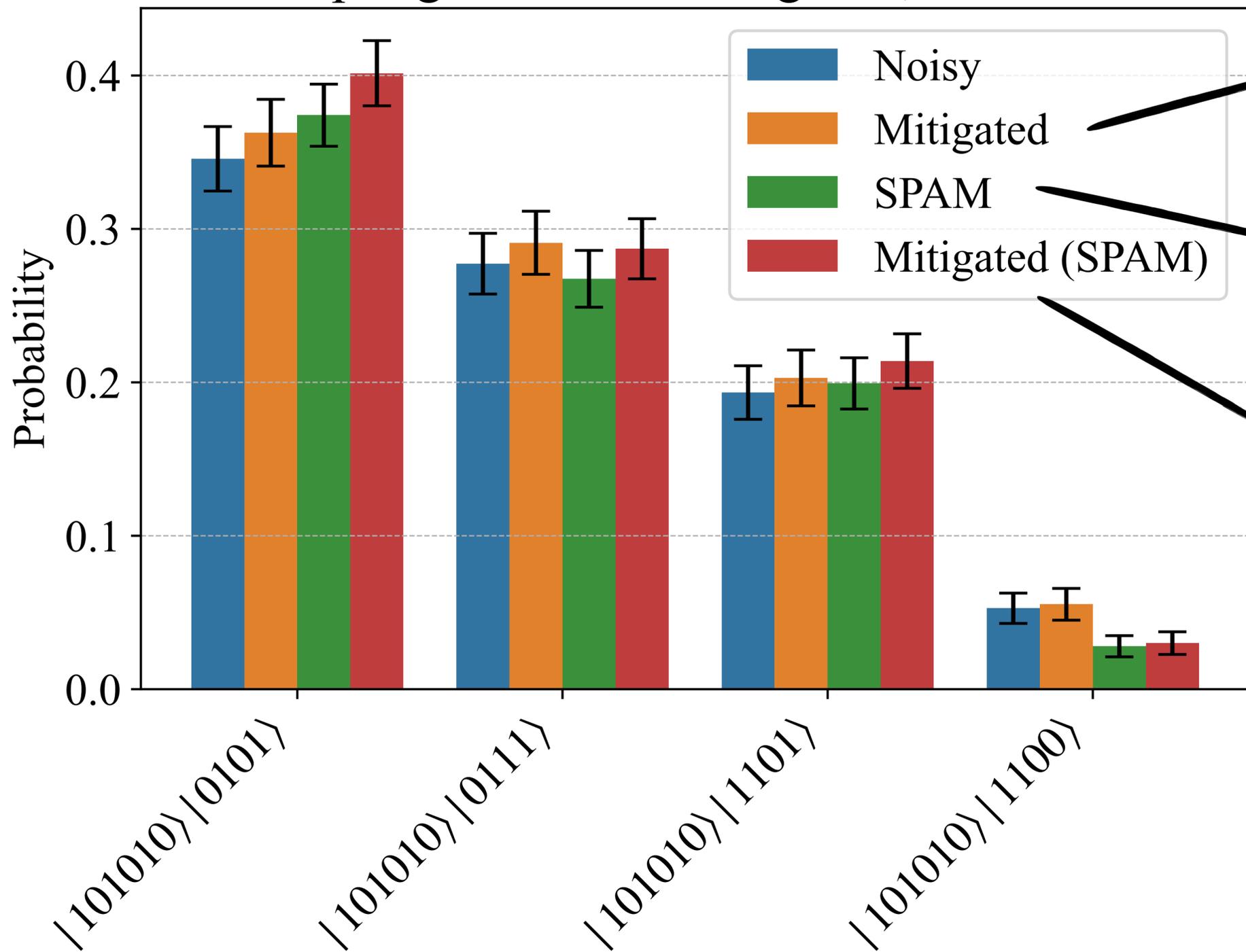
Sampling with H1-1E at $g=1.1$, shots=512



Remove bit strings that do not correspond to physical configurations

State Preparation and Measurement correction

Sampling with H1-1E at $g=1.1$, shots=512

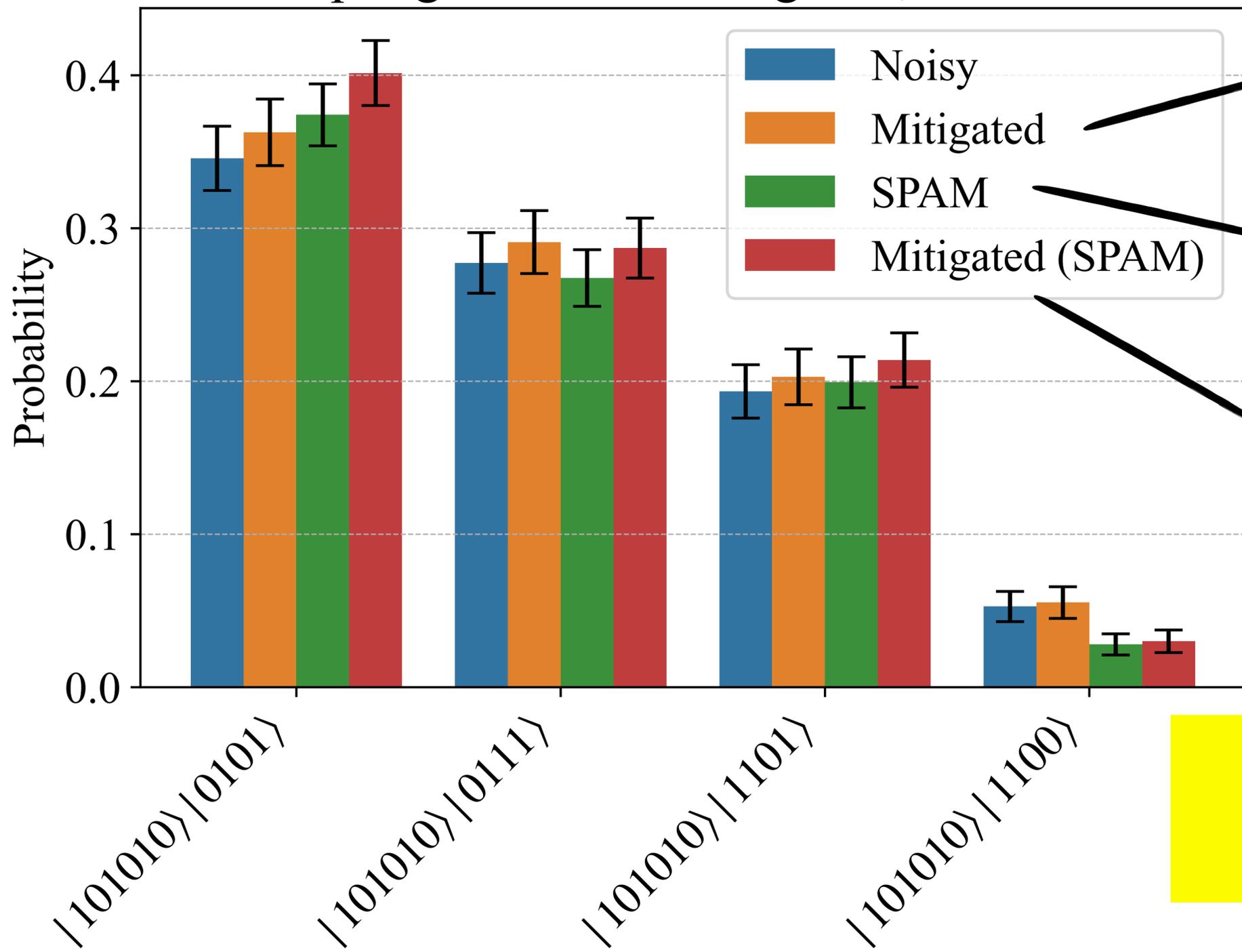


Remove bit strings that do not correspond to physical configurations

State Preparation and Measurement correction

Filter out unphysical bit strings after correcting for SPAM

Sampling with H1-1E at g=1.1, shots=512

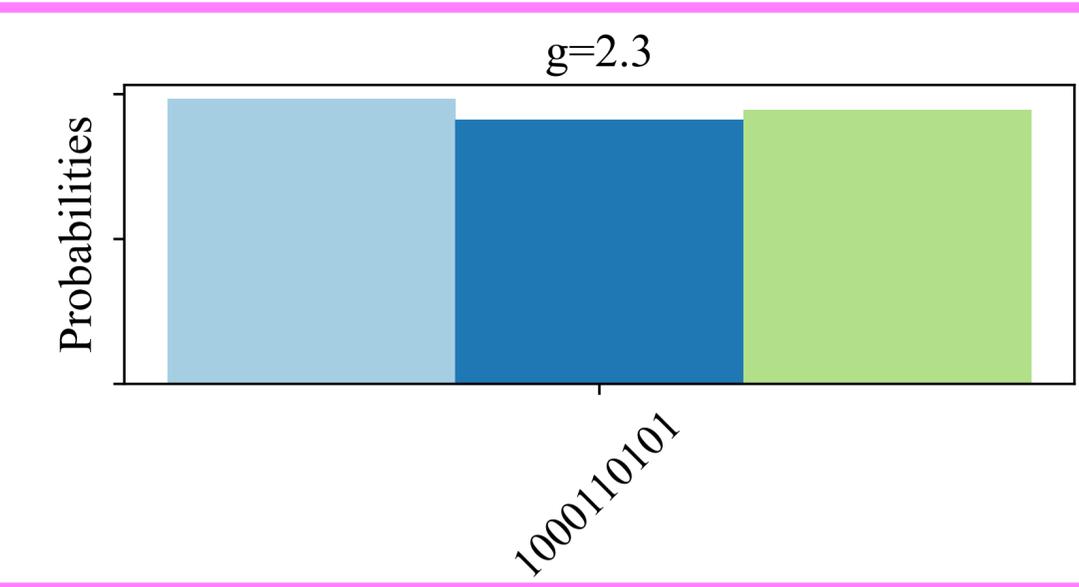
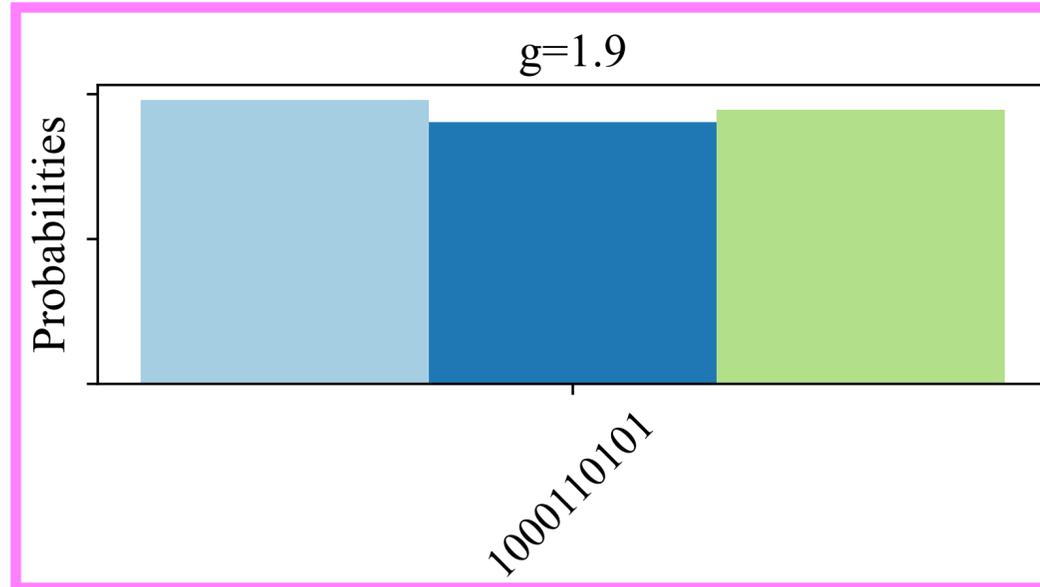
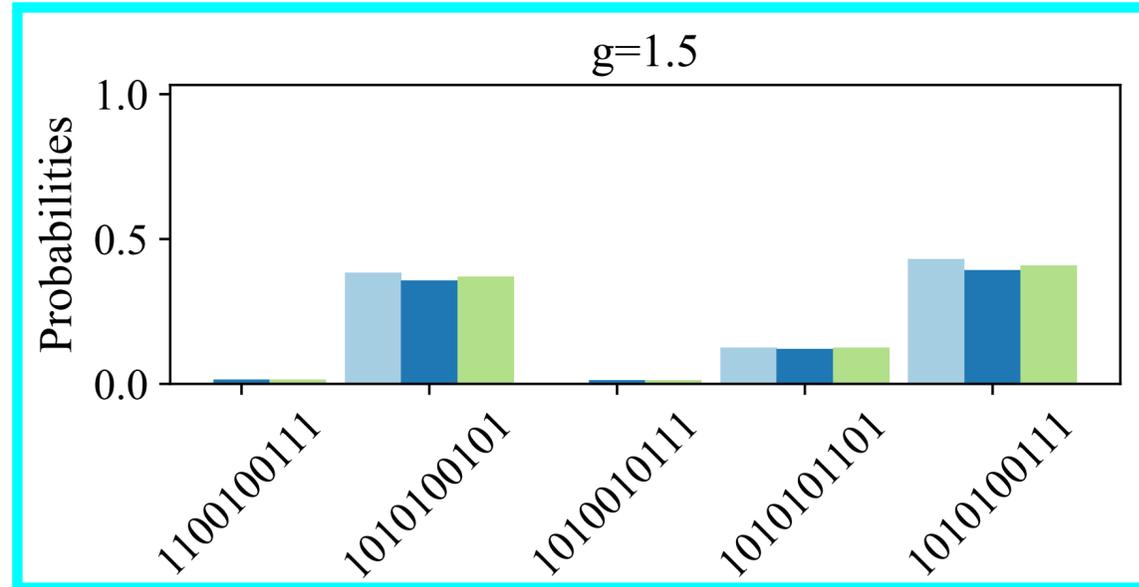
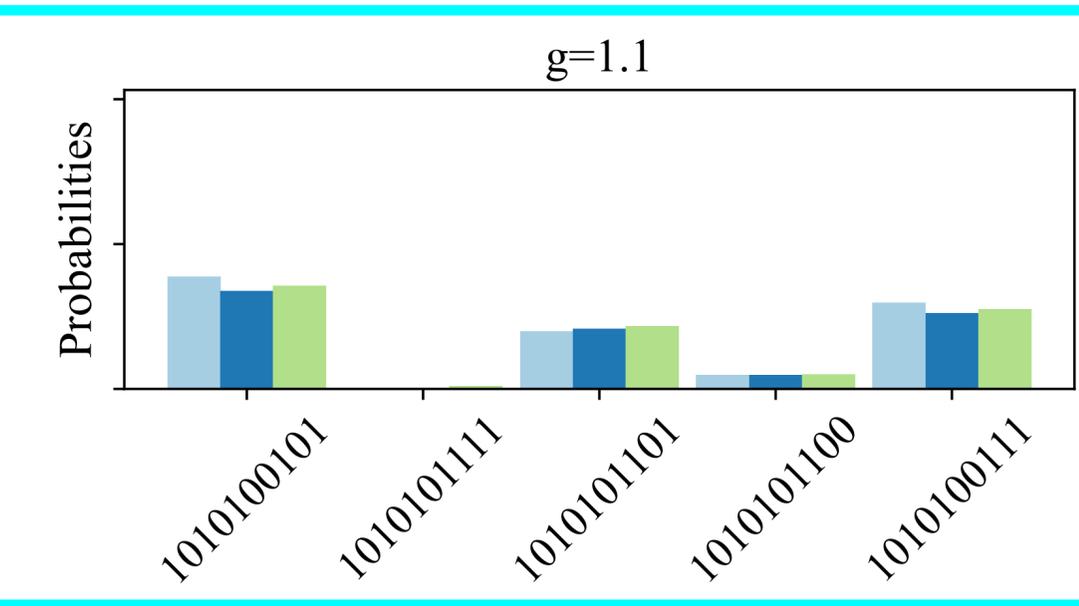
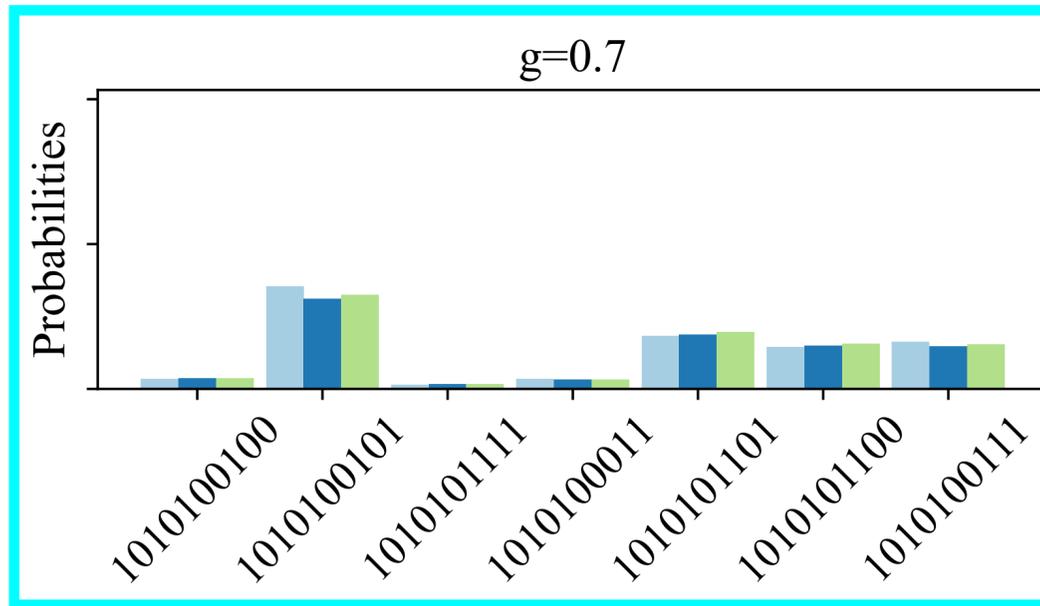
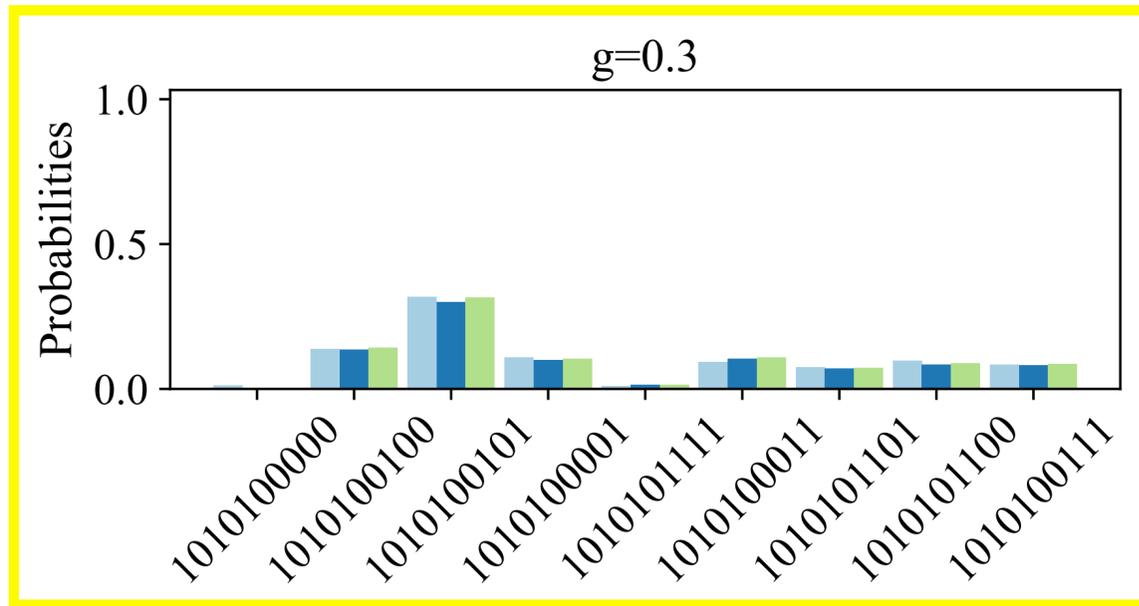


Remove bit strings that do not correspond to physical configurations

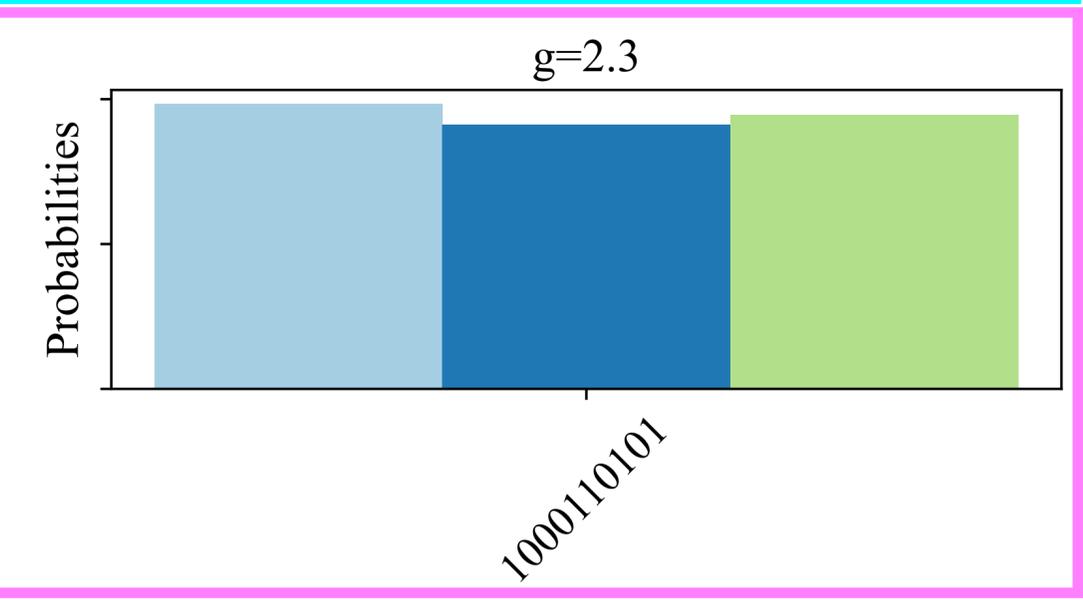
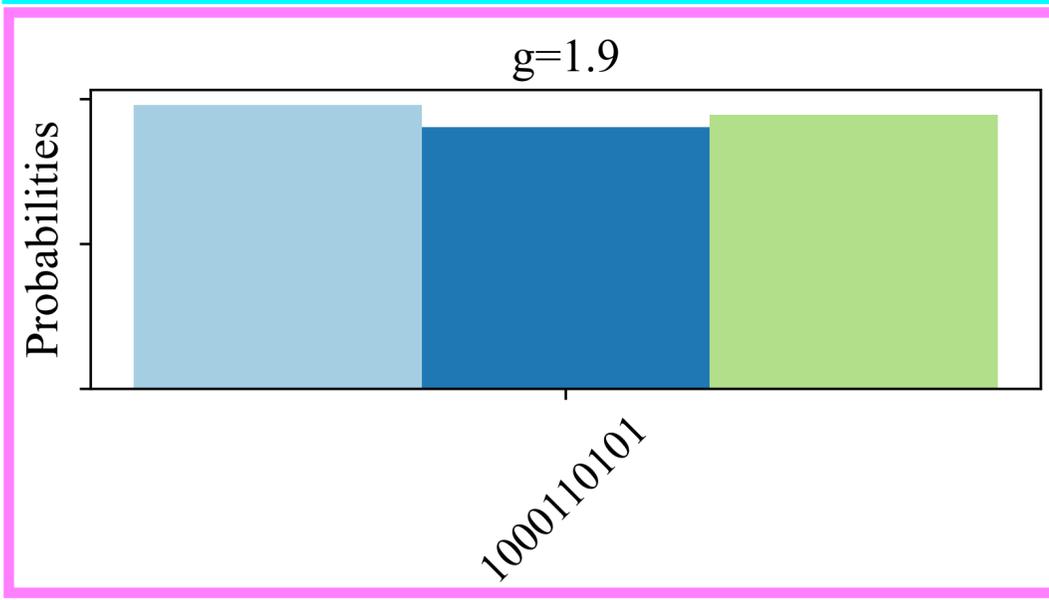
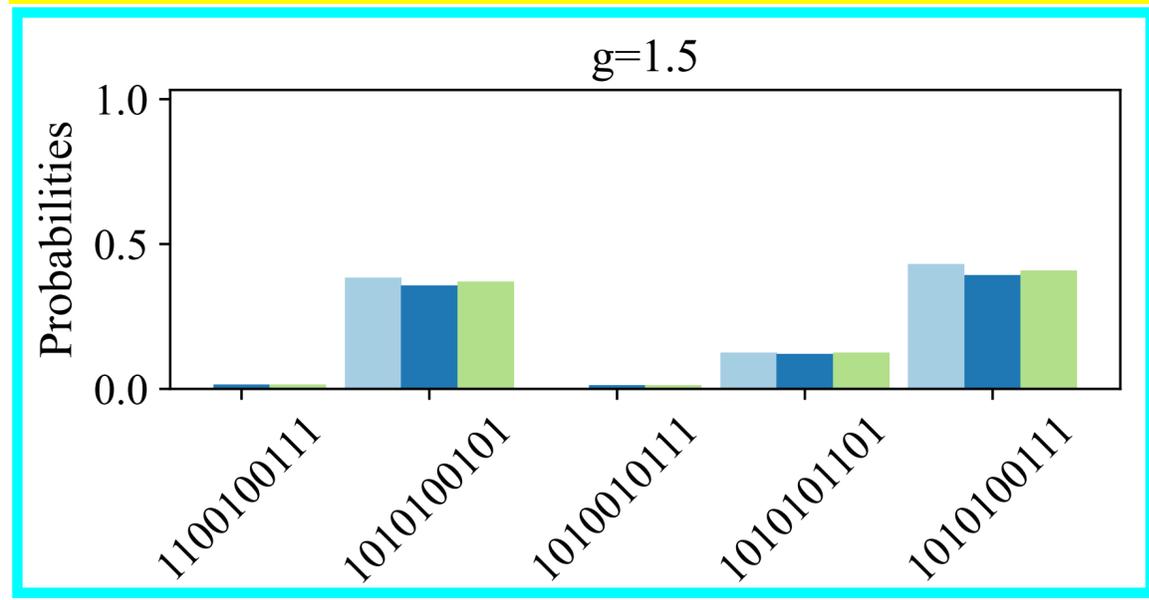
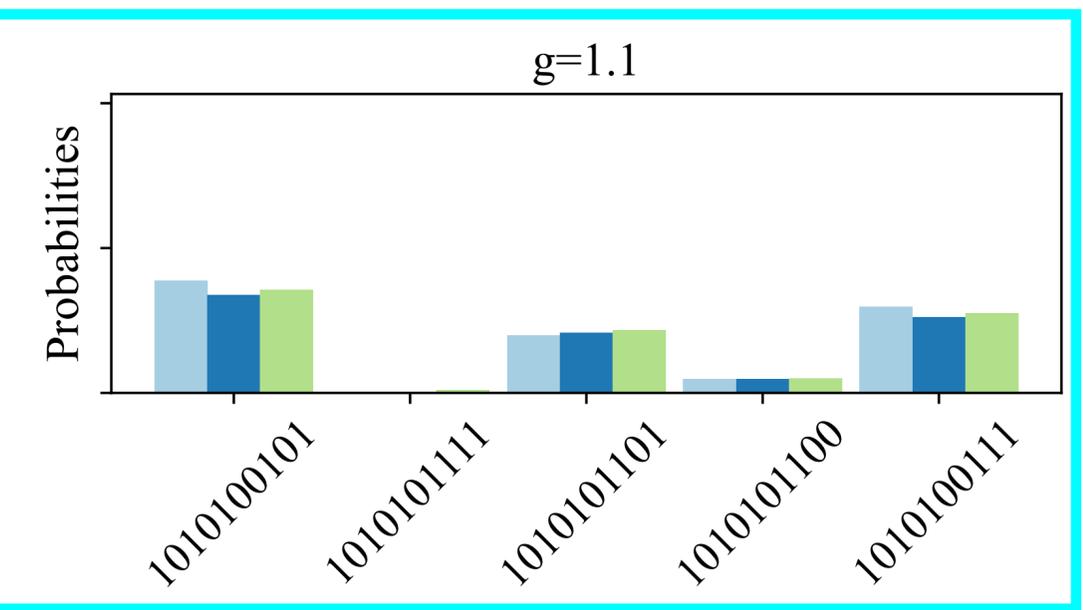
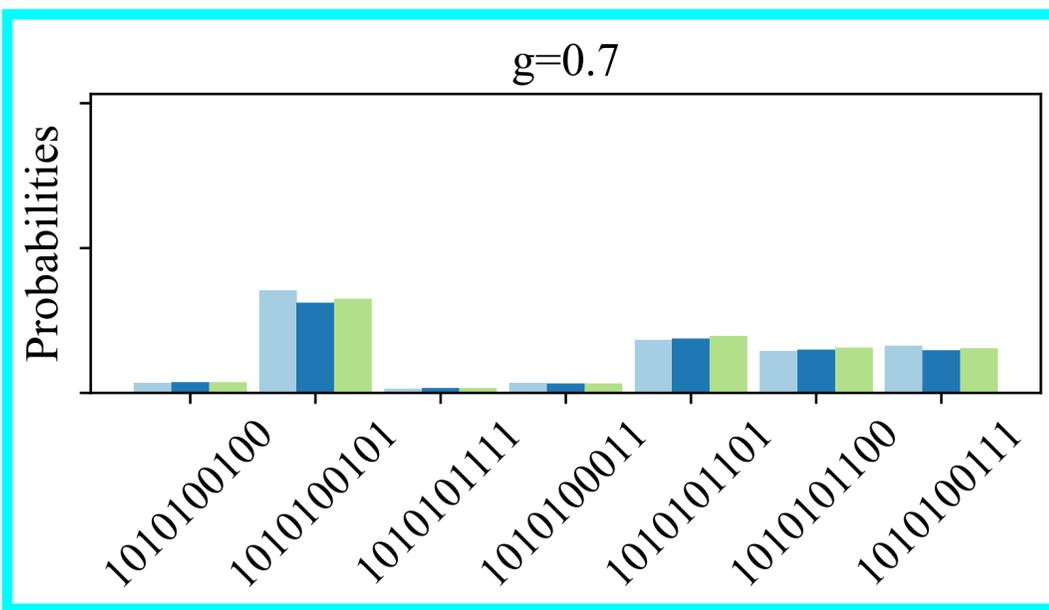
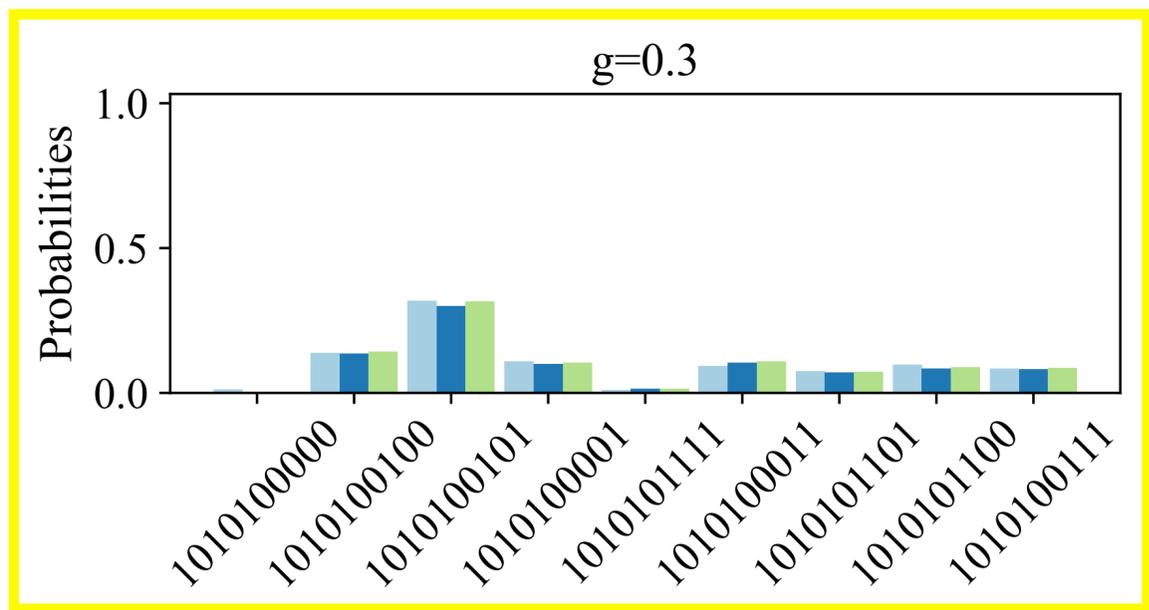
State Preparation and Measurement correction

Filter out unphysical bit strings after correcting for SPAM

Very little difference raw results and mitigated ones



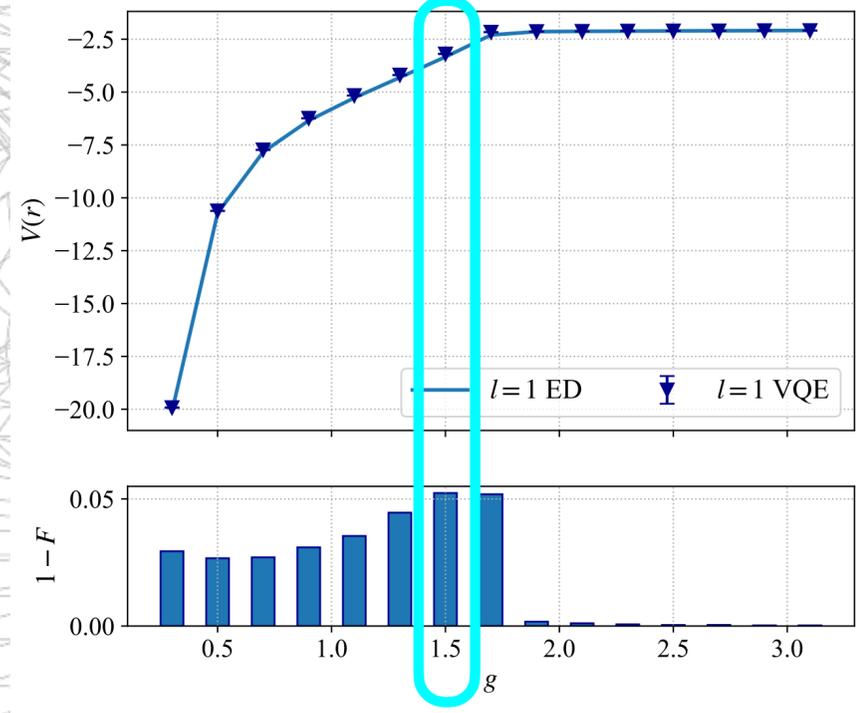
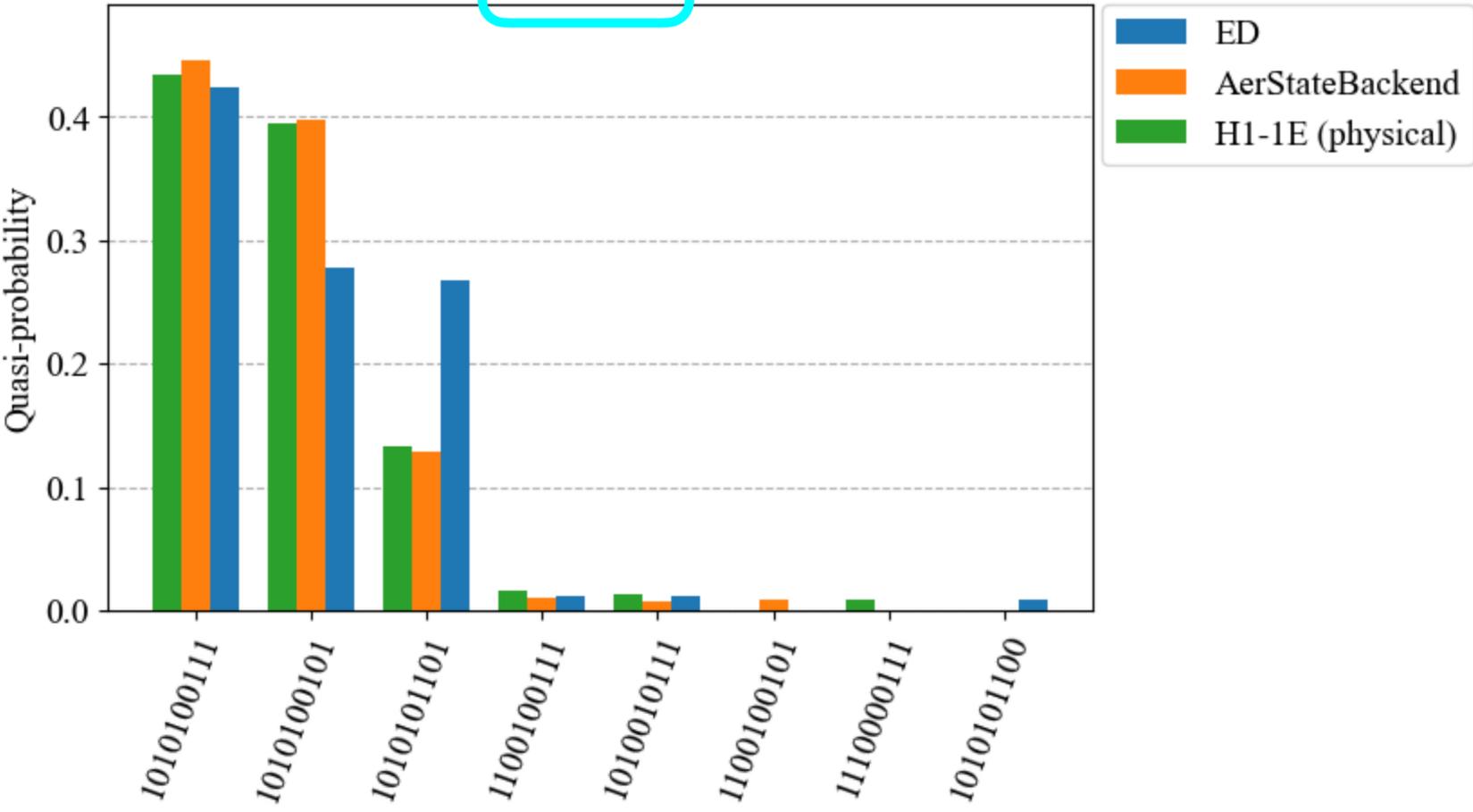
VQE
 H1-1E (raw)
 H1-1E (physical)



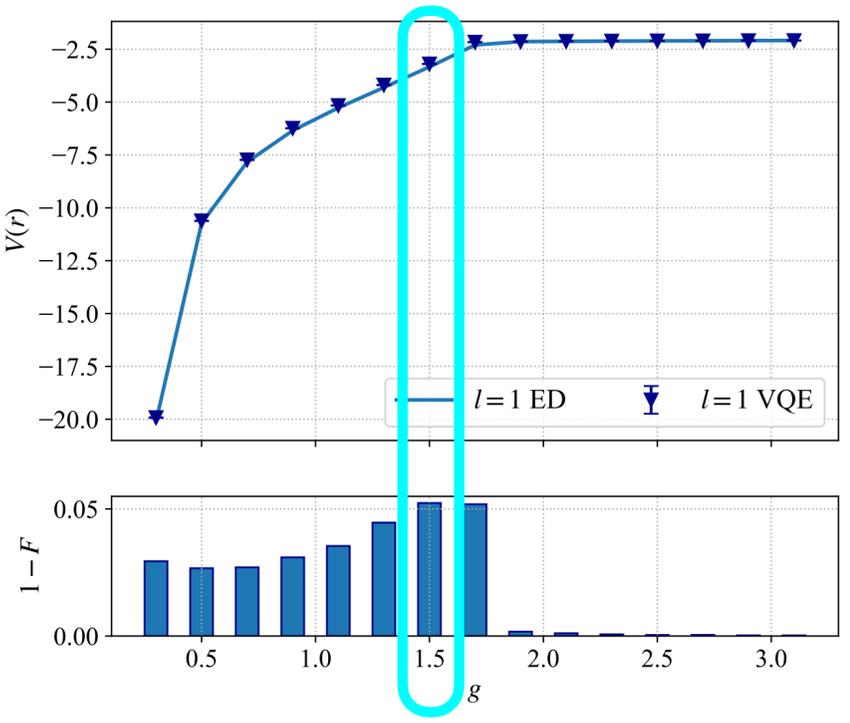
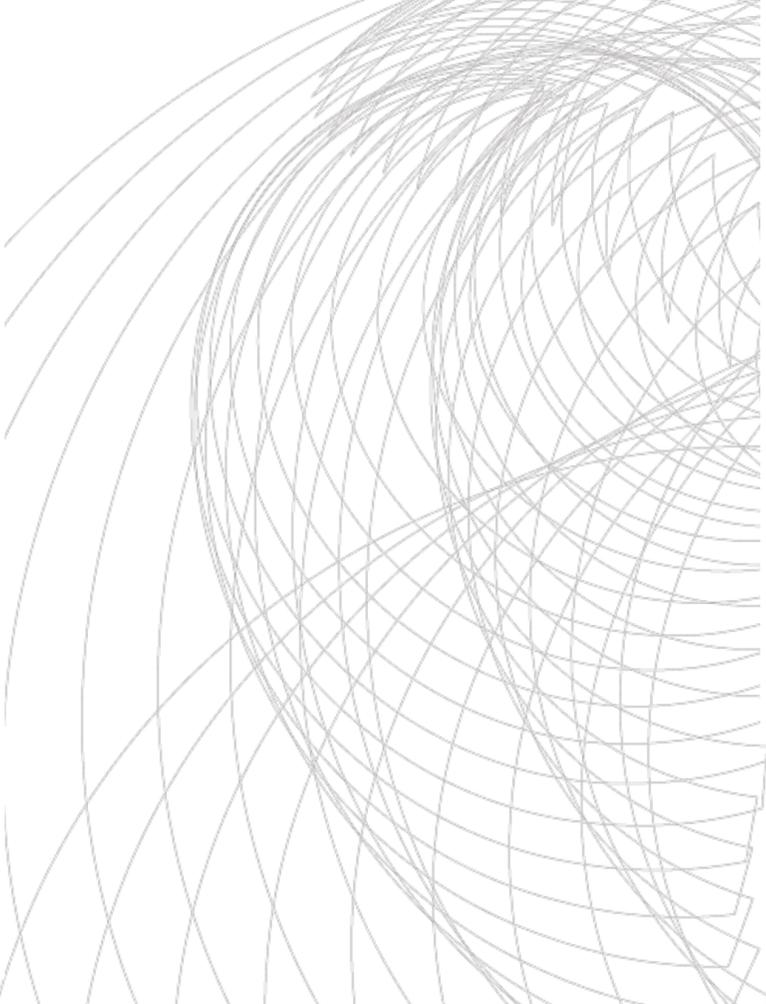
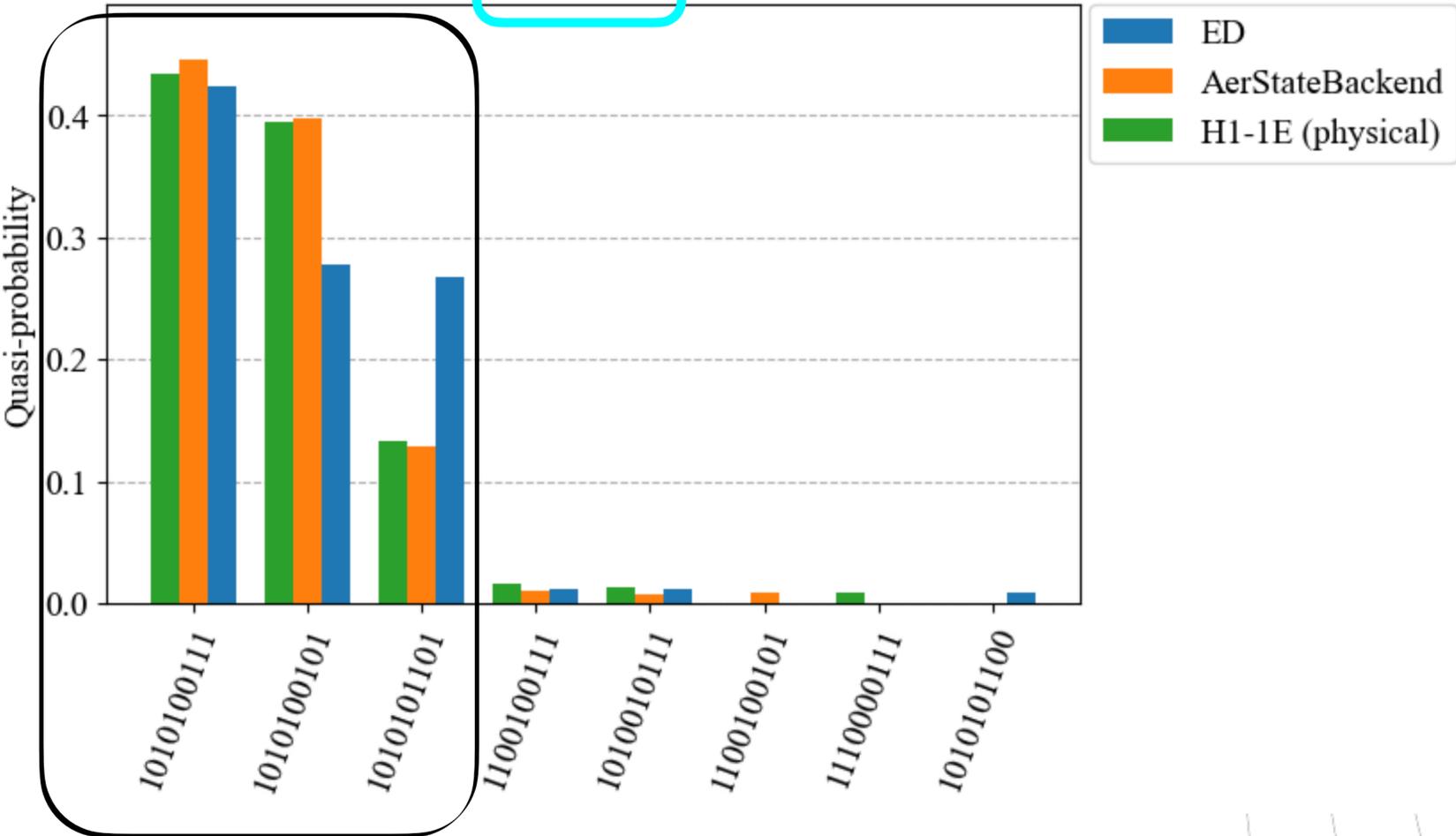
VQE H1-1E (raw) H1-1E (physical)

Remove unphysical states/bitstrings

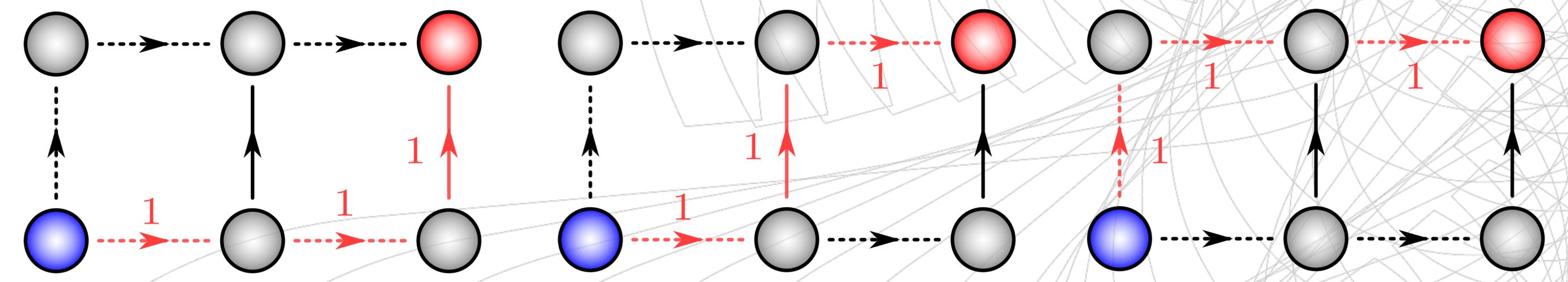
$g=1.5$ $l=1$



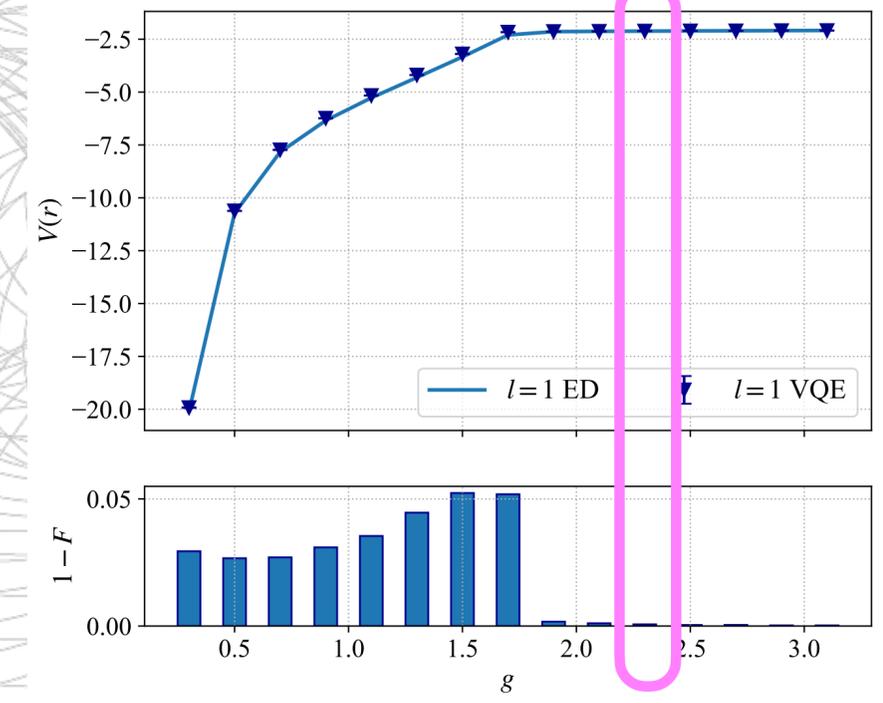
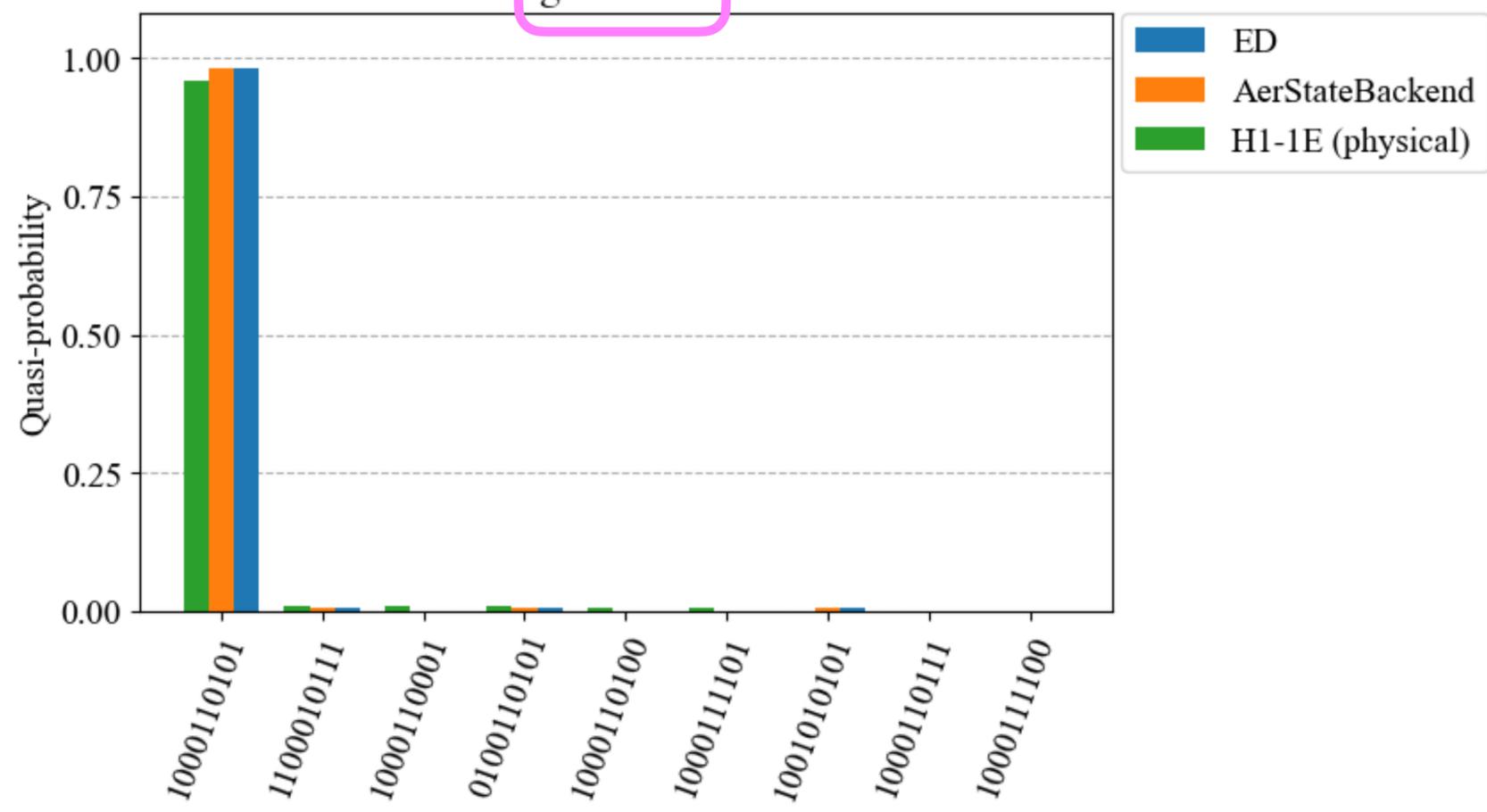
$g=1.5 \quad l=1$



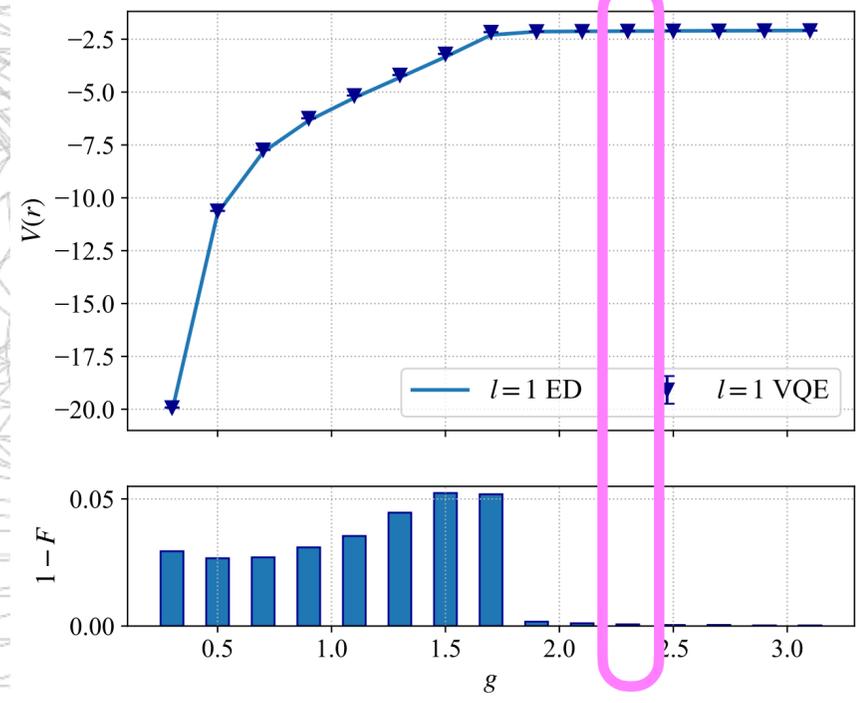
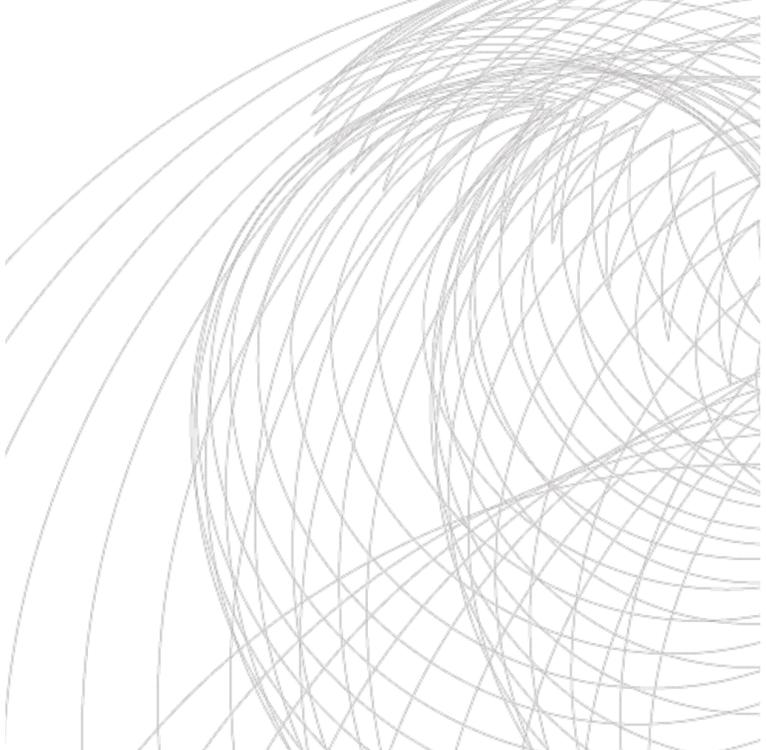
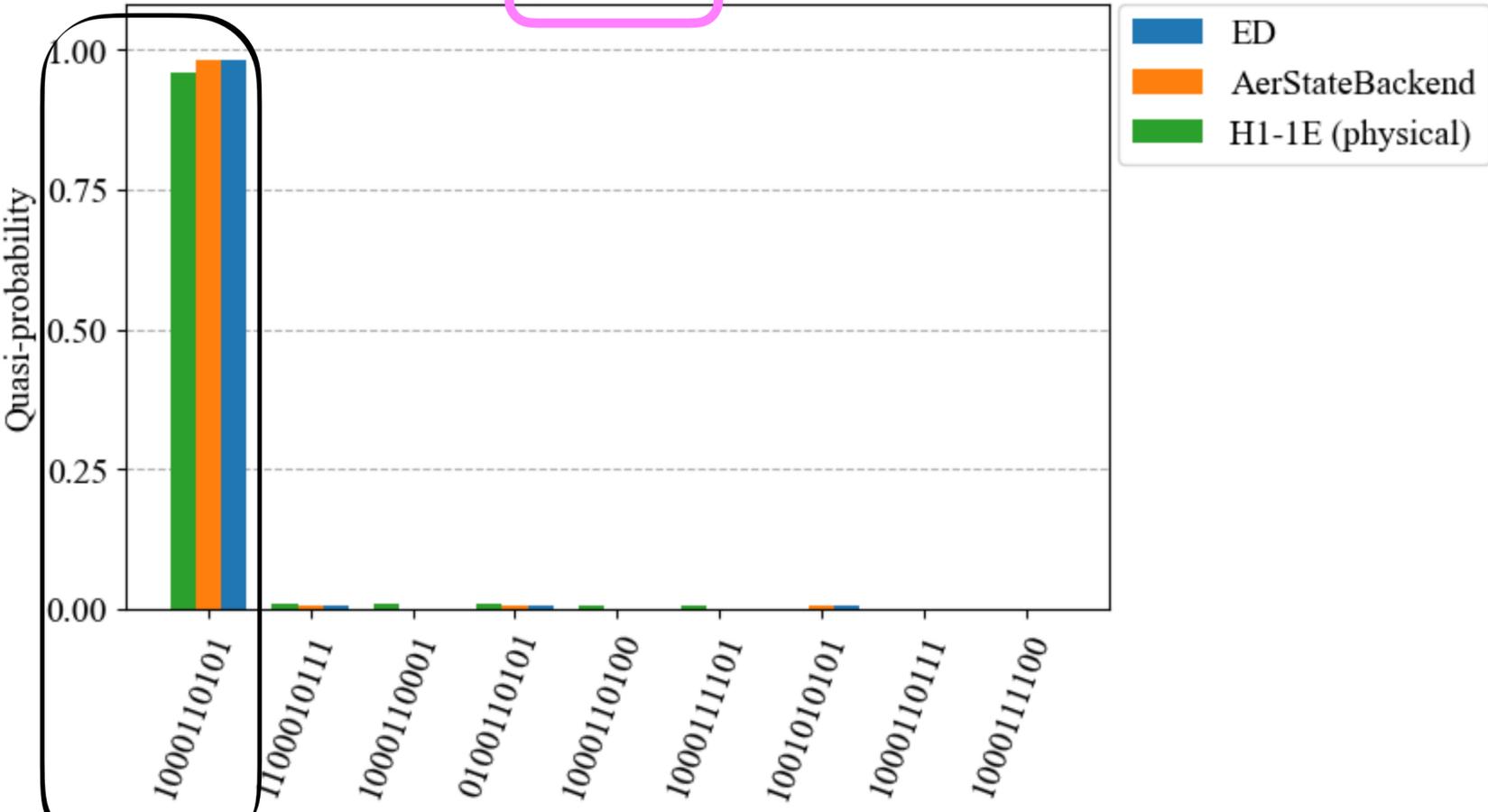
Configurations with electric fluxes between static charges



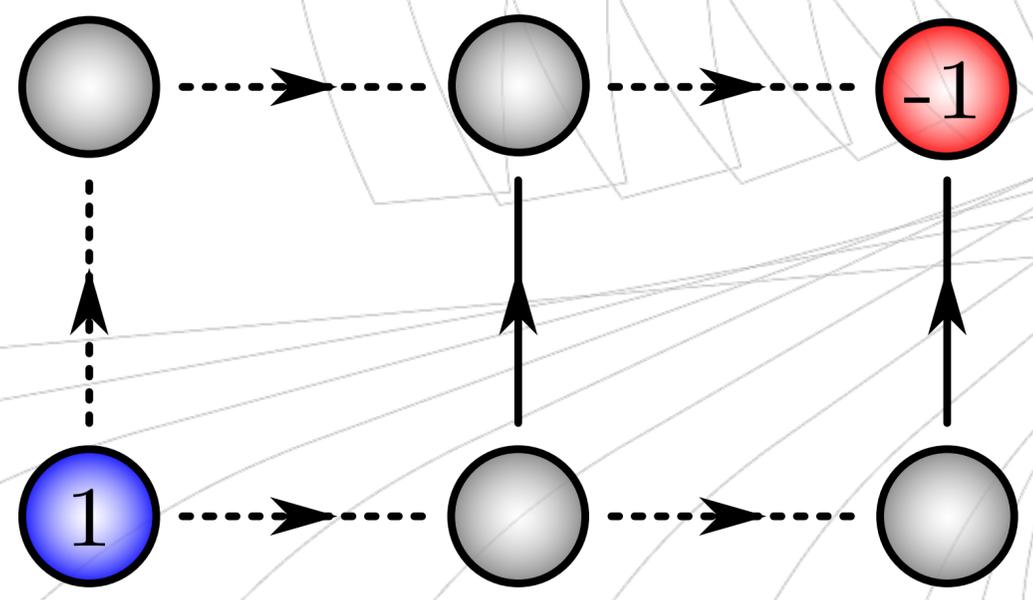
$g=2.3$ $l=1$



$g=2.3$ $l=1$



Configurations with string breaking: no flux and mesons

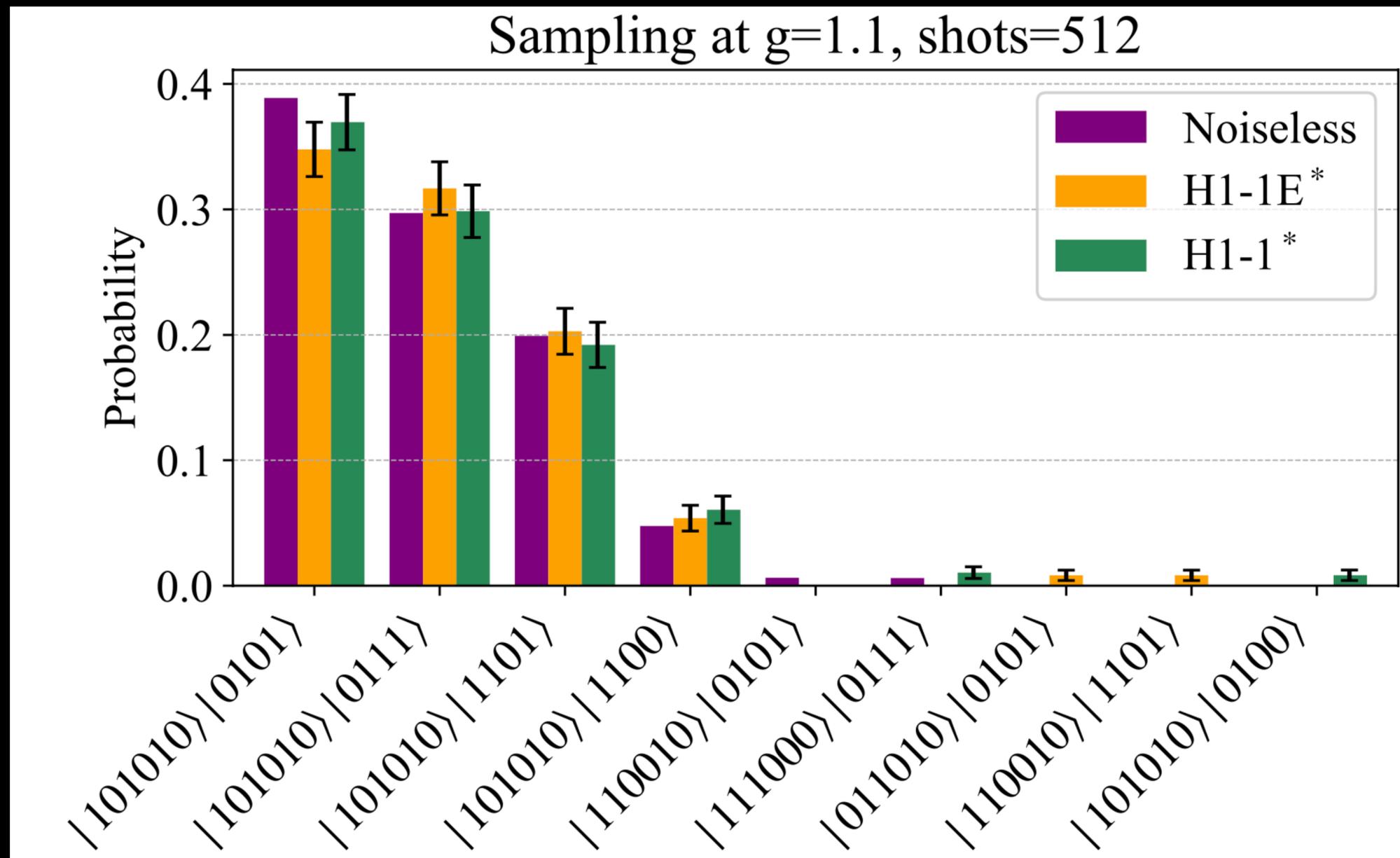


Experiments on H1-1

Real Hardware

Results compared to Emulator

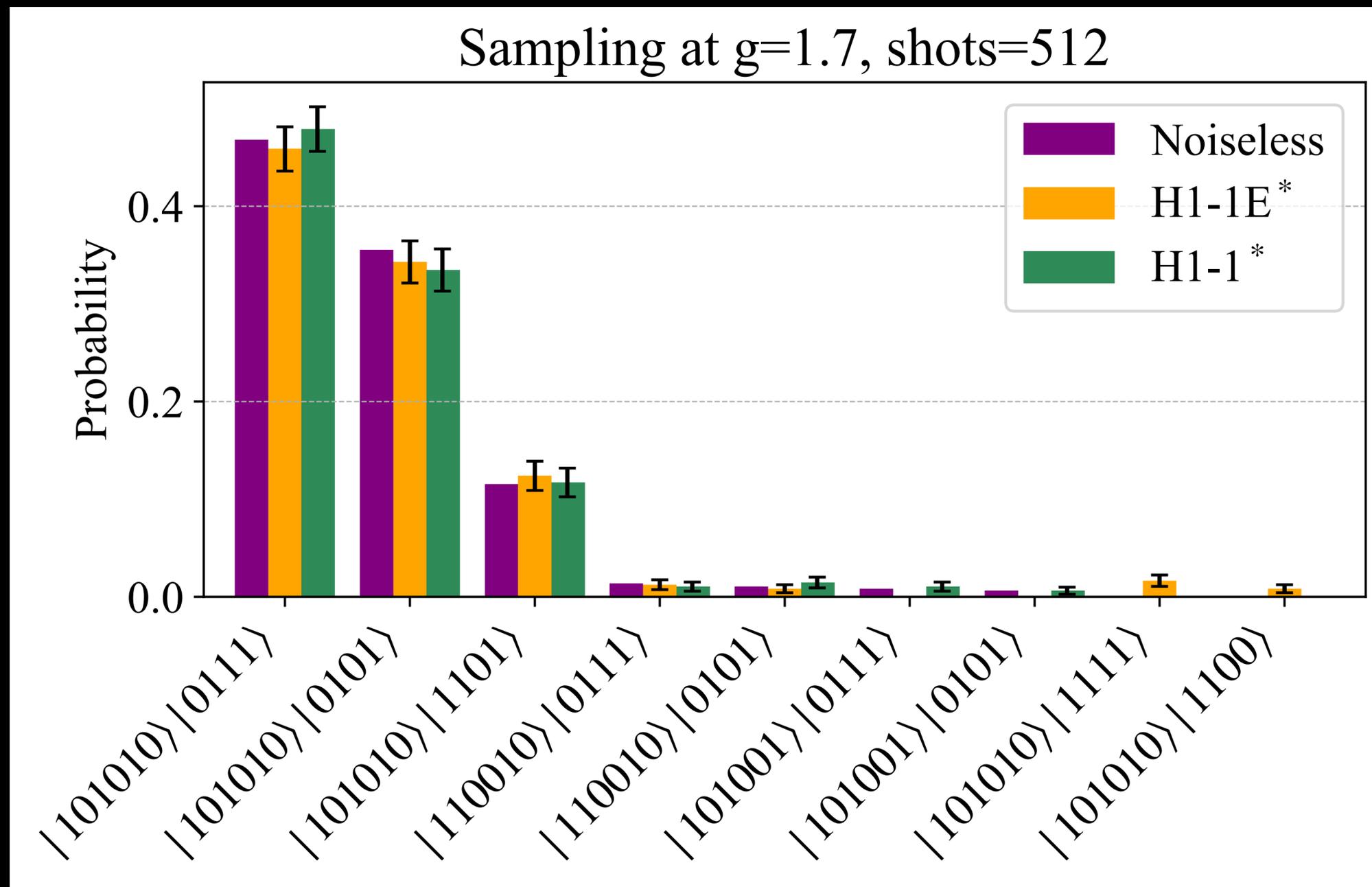
- 512 shots on H1-1
- No error mitigation!
(post processing)



Real Hardware

Results compared to Emulator

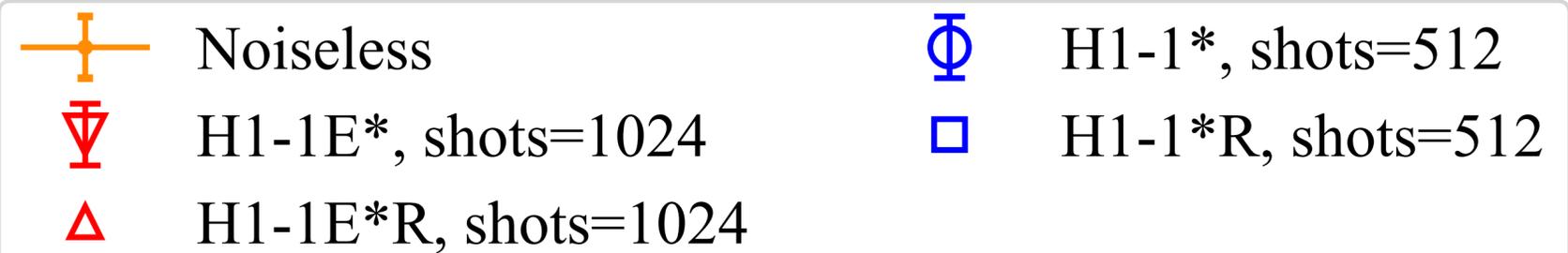
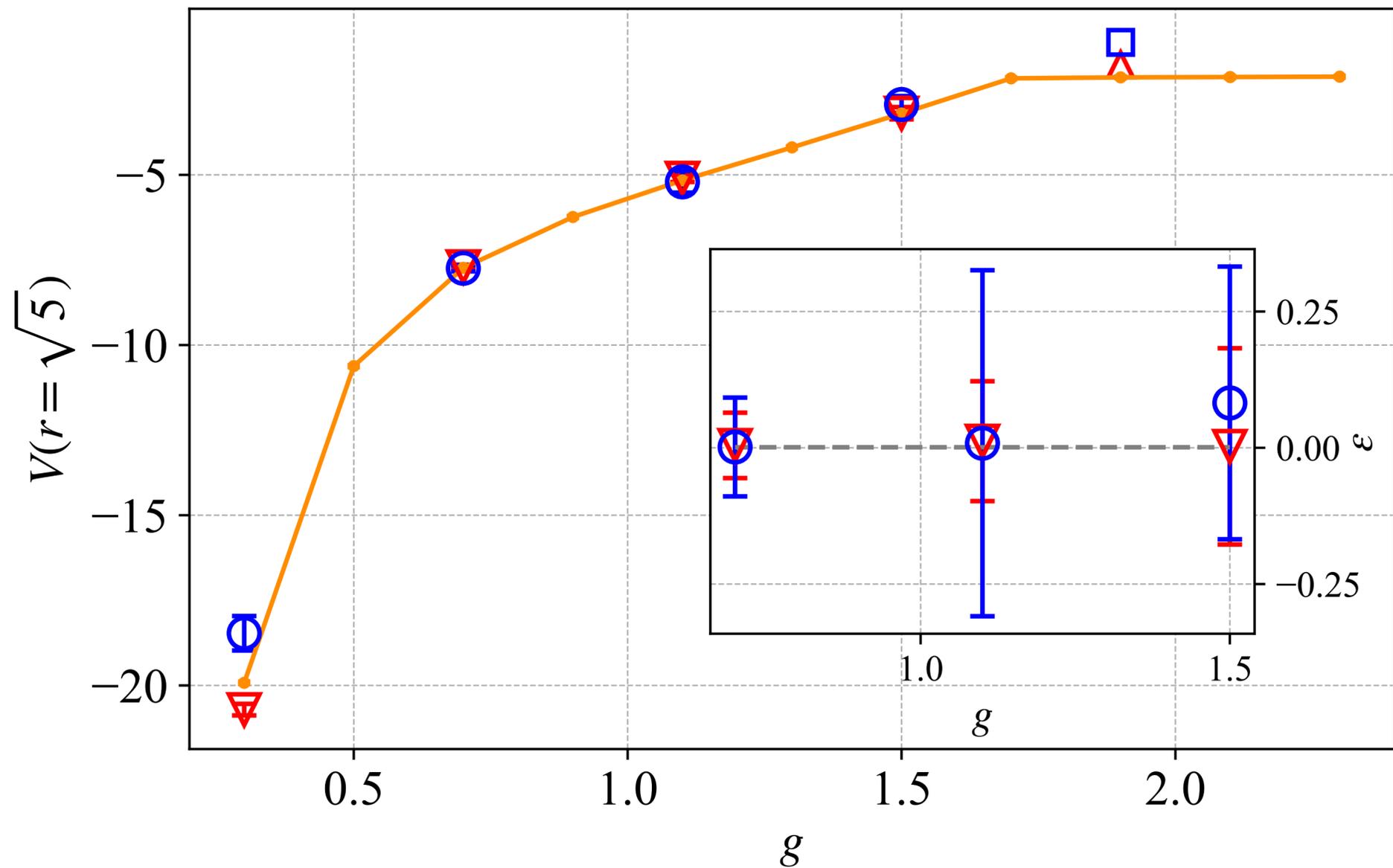
- 512 shots on H1-1
- No error mitigation!
(post processing)



Real Hardware

Results compared to Emulator

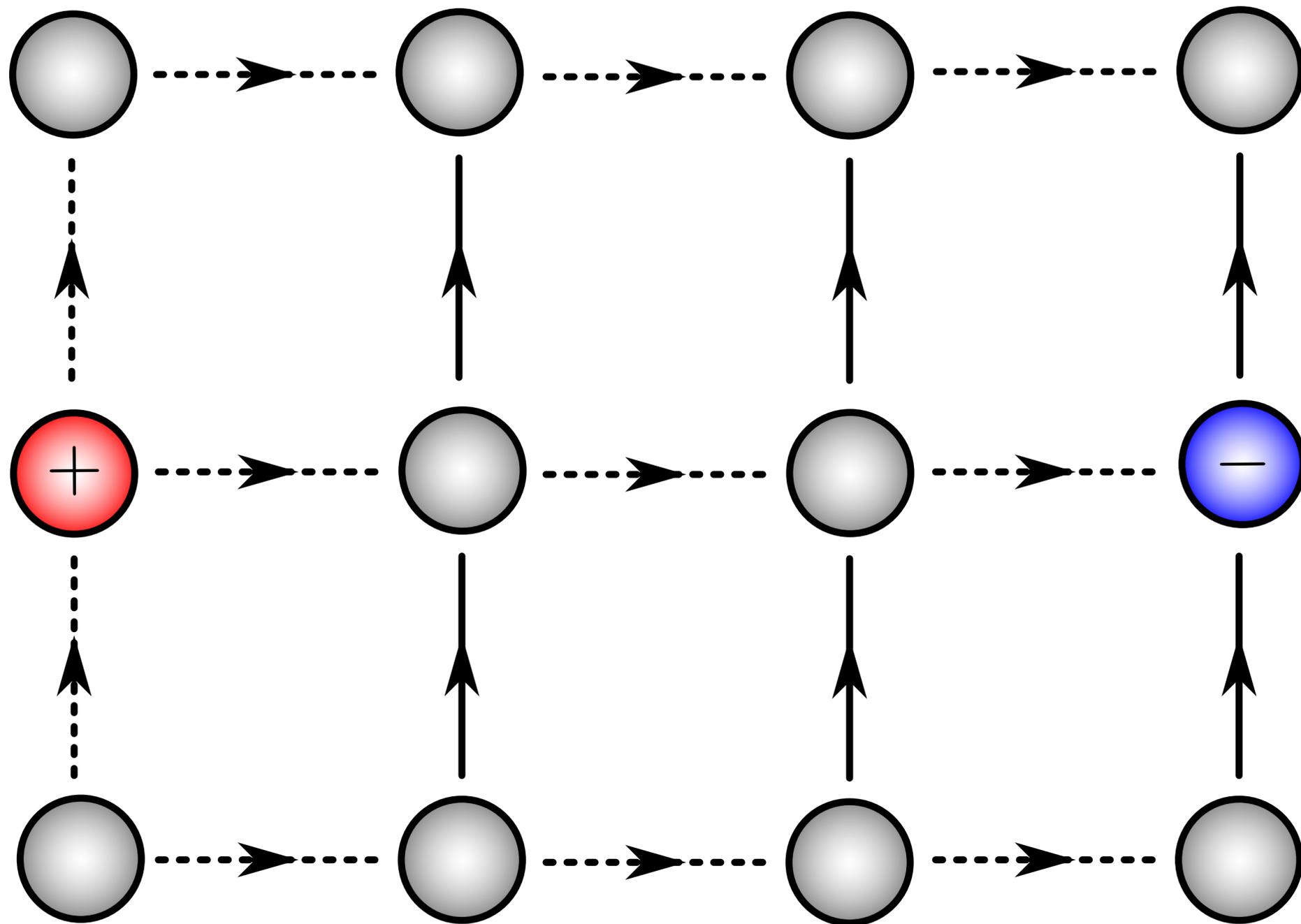
- 512 shots on H1-1
- Simple Error Mitigation: (PMSV, SPAM)



Conclusions

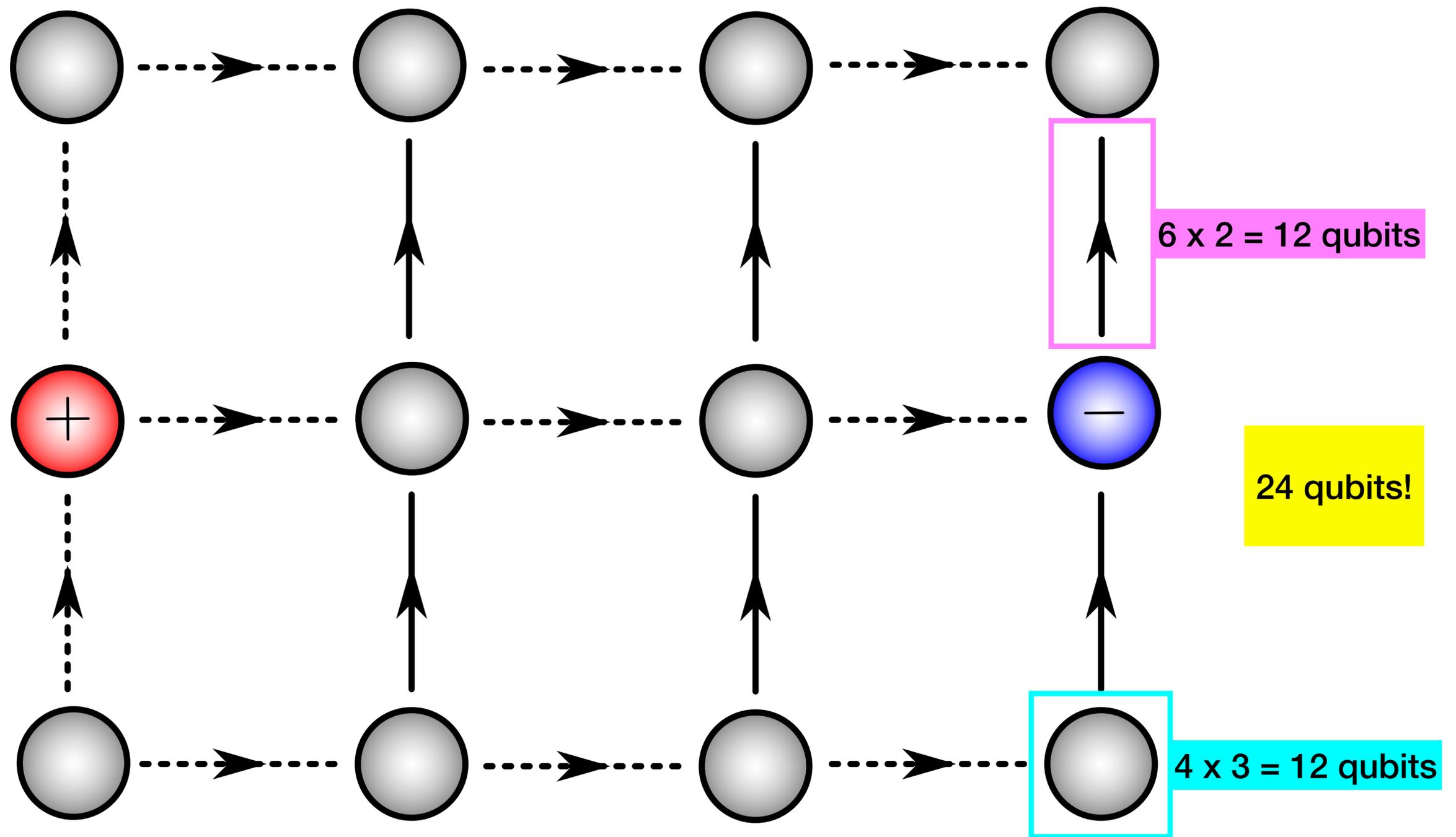
And future directions

- We demonstrated on real quantum hardware a calculation of **the confining potential of (2+1)D QED** in the Hamiltonian formulation
- Access to the ground state, even in a variational sense, allows us to **visualize the confining fluxes** between static charges
- **String breaking and the formation of “mesons”** is observed
- Scaling up the **quantum state preparation** step is important to fully leverage the computational power of quantum hardware
- Formulations of Hamiltonian lattice gauge theories **that are different from Kogut–Susskind might help scaling to non-Abelian theories in (3+1)D**



Can we do larger lattices?

Even if we have 20 qubits?

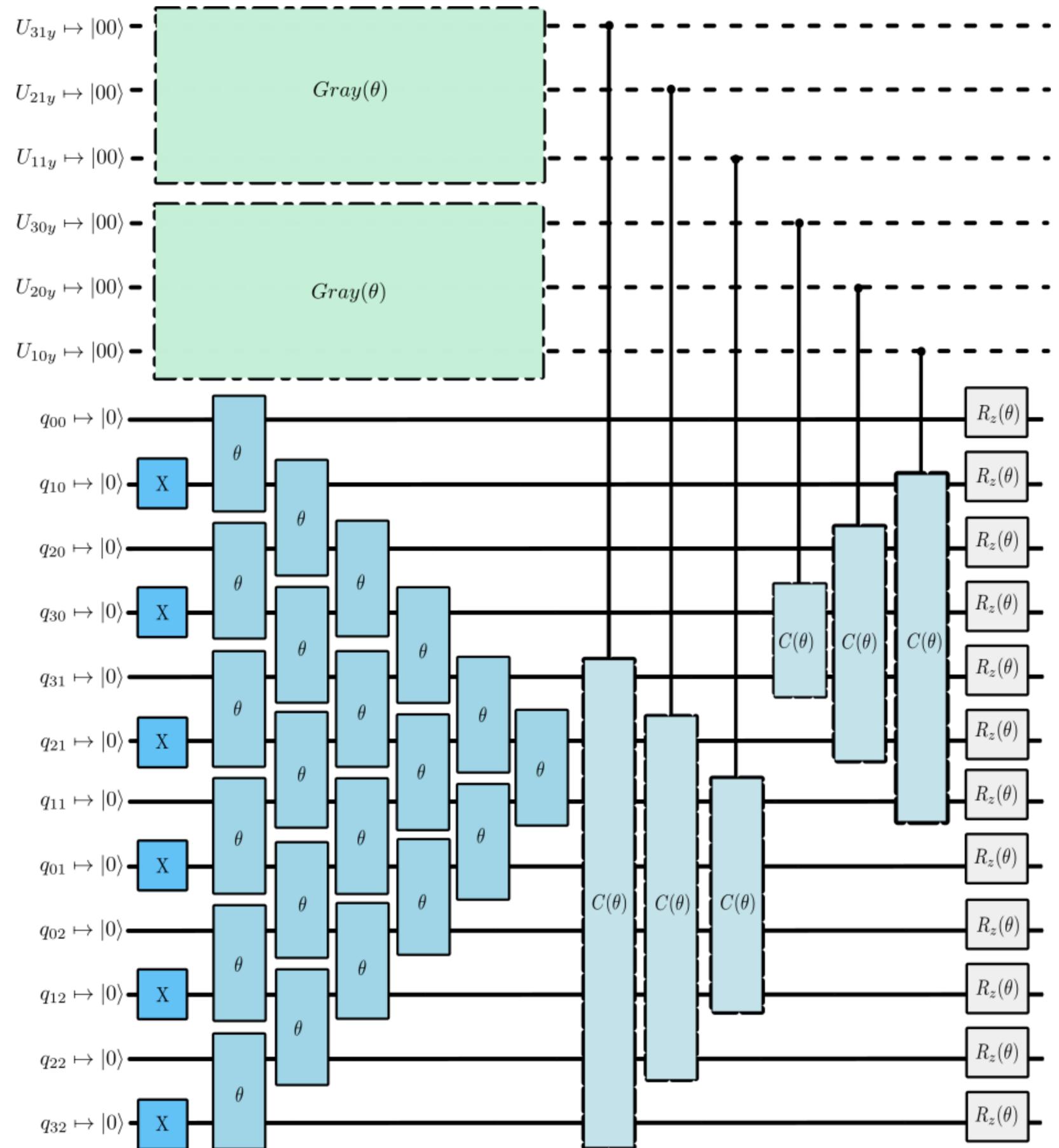


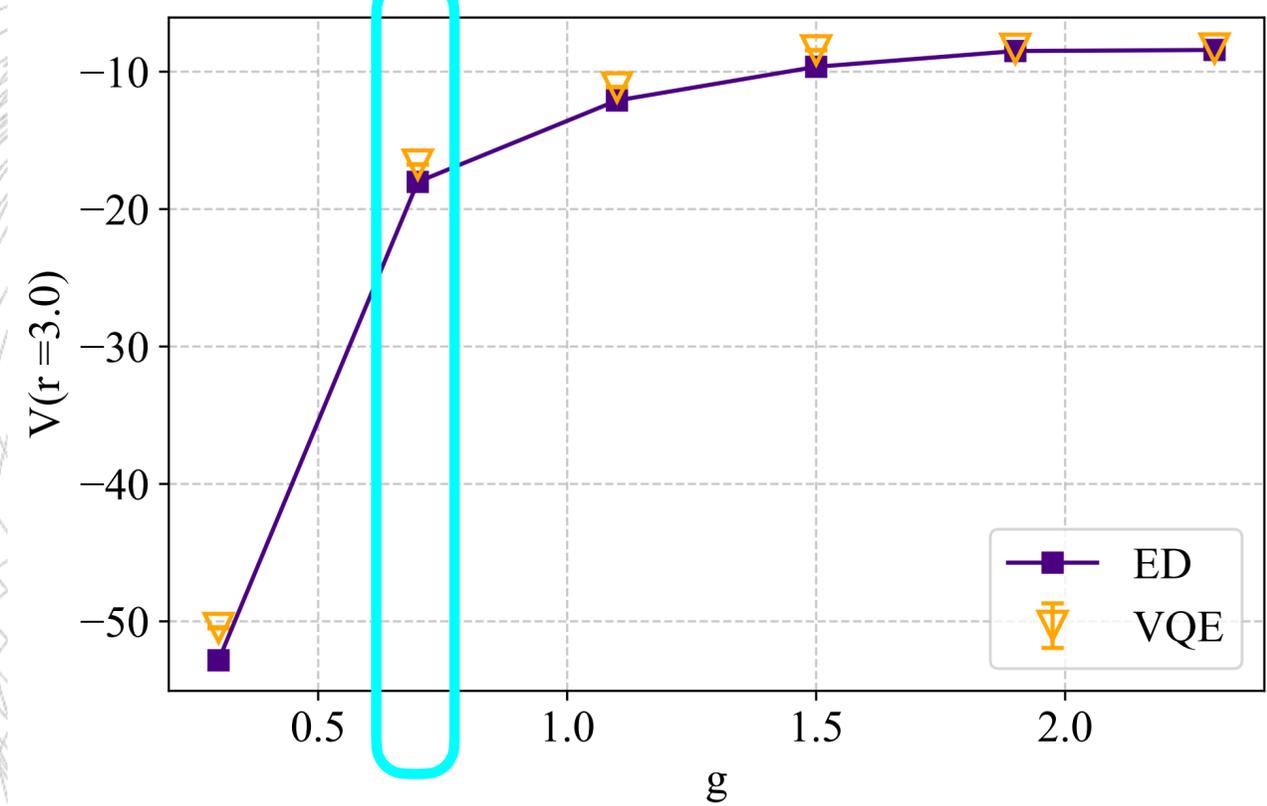
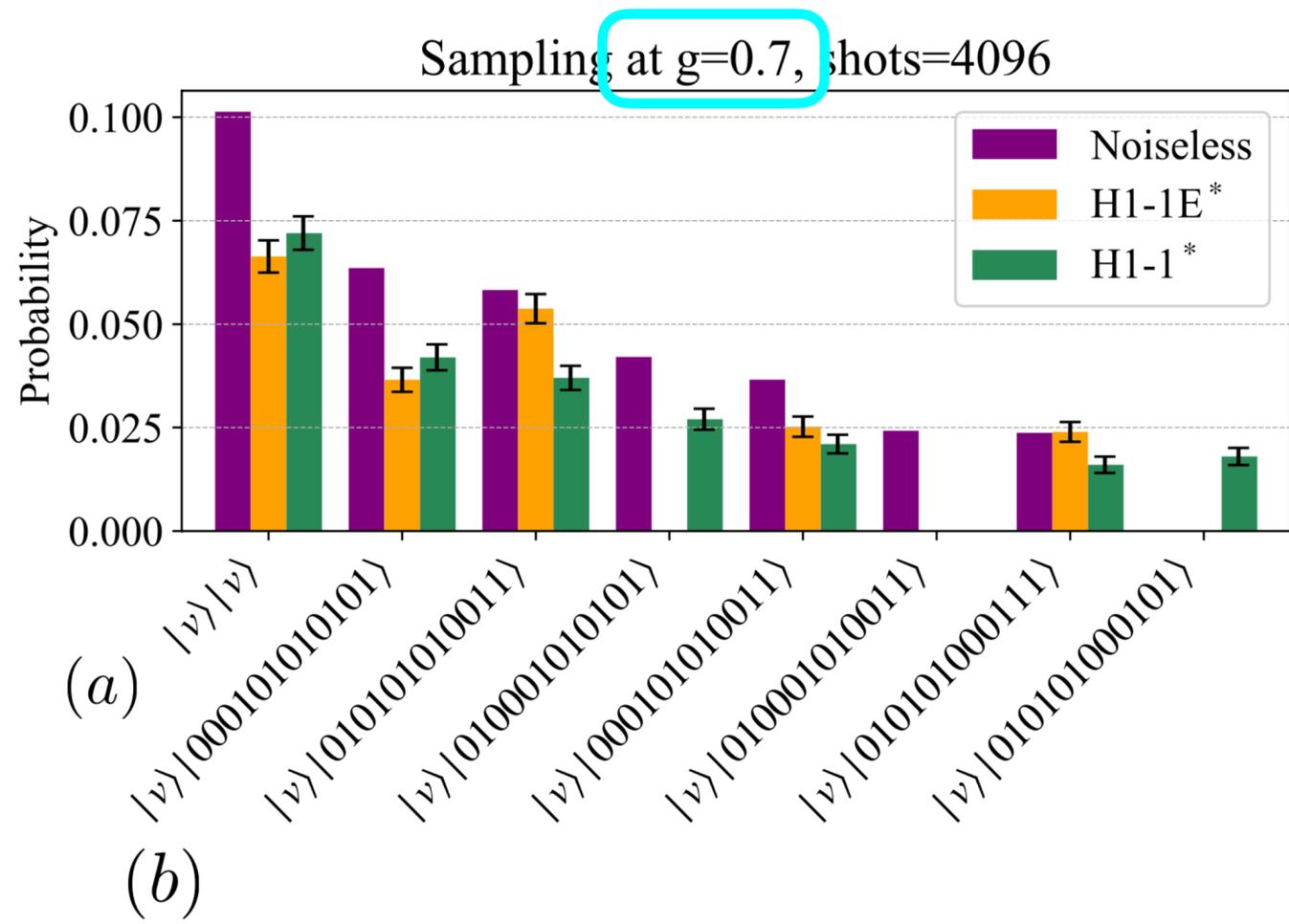
Can we do larger lattices?

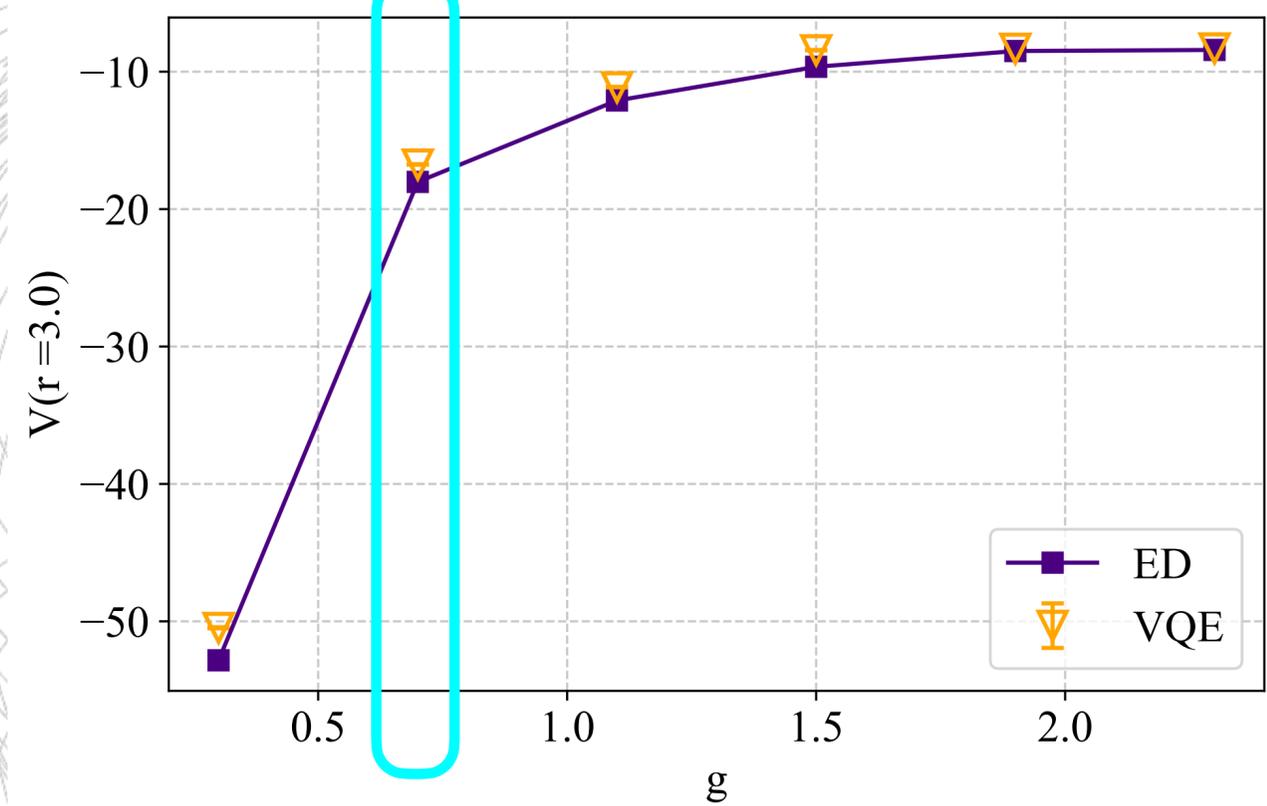
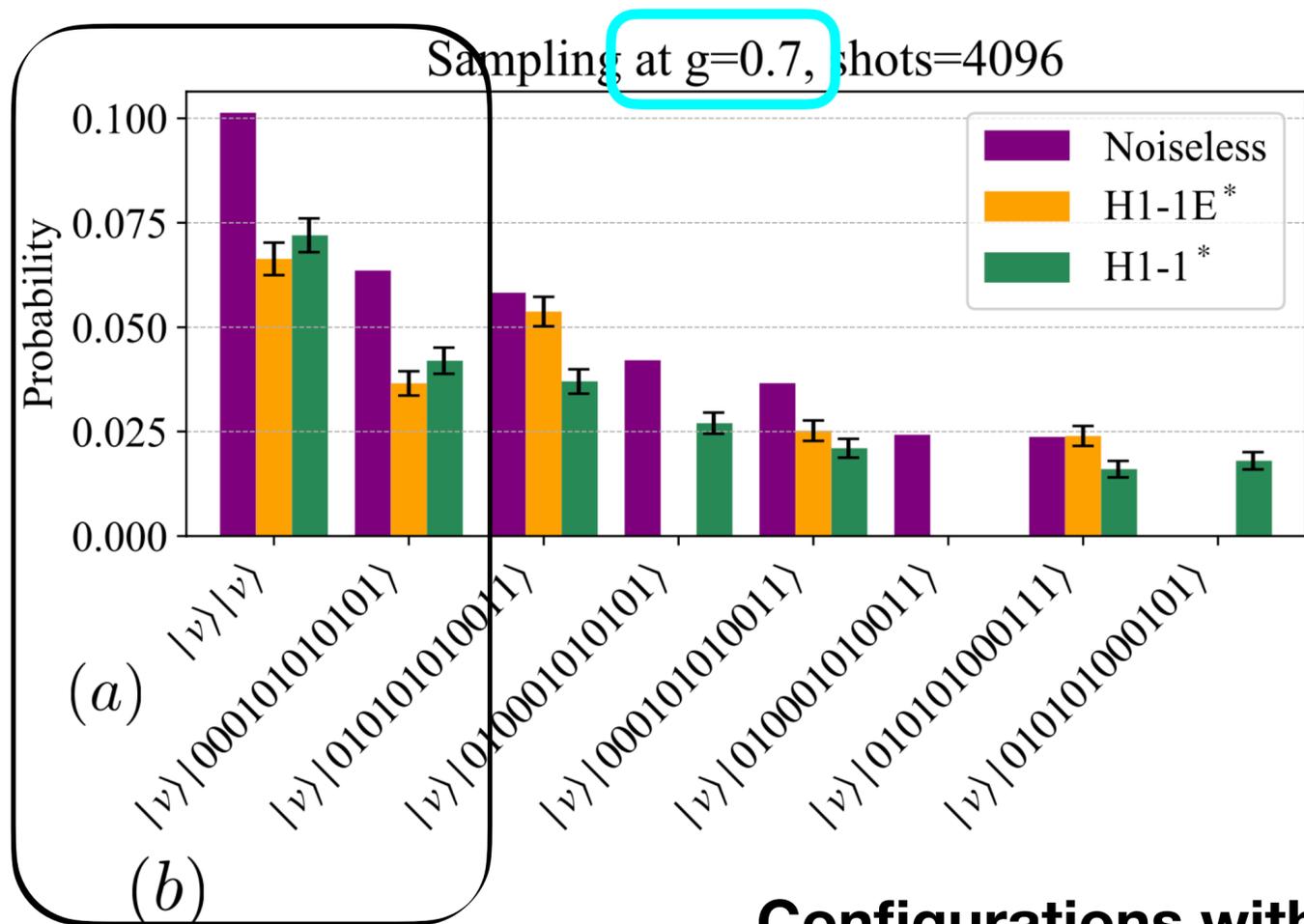
Even if we have 20 qubits?

Mutual Information Ansatz

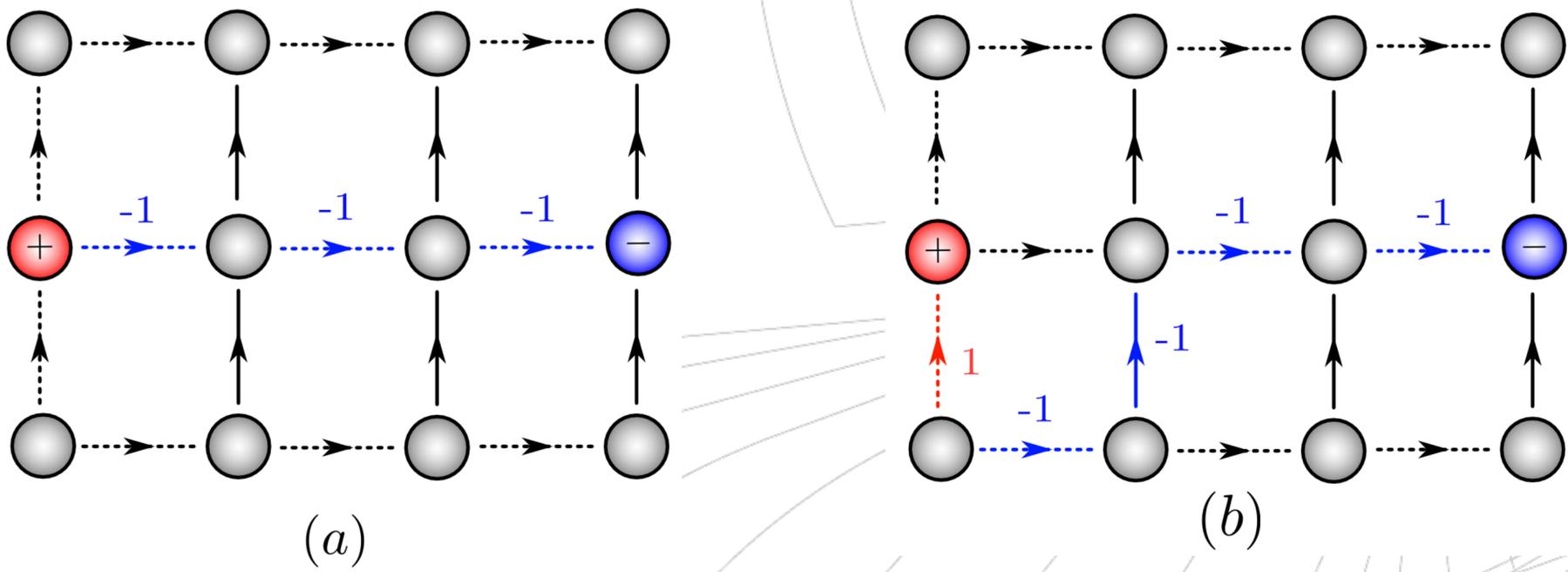
- More qubits: 10 \rightarrow 24
- More parameters: 30 \rightarrow 81
- More 2qb gates: 152 \rightarrow 450

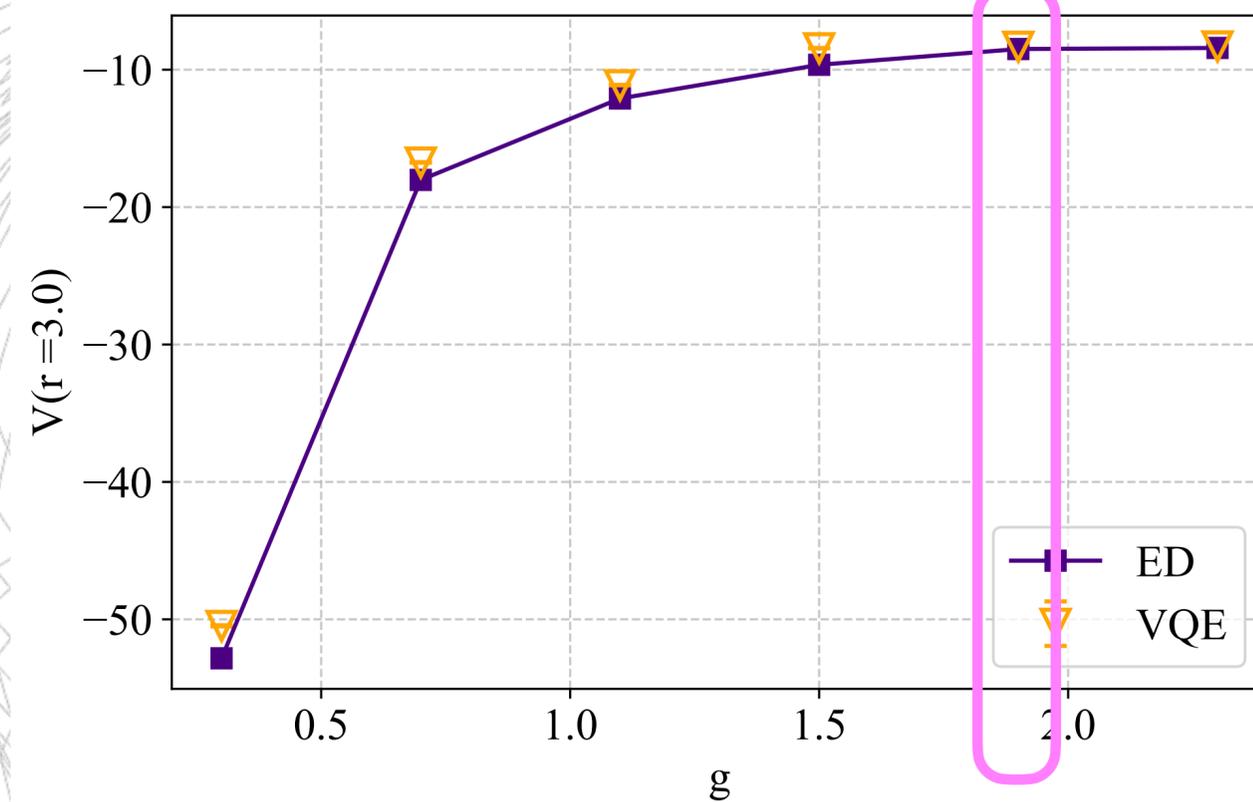
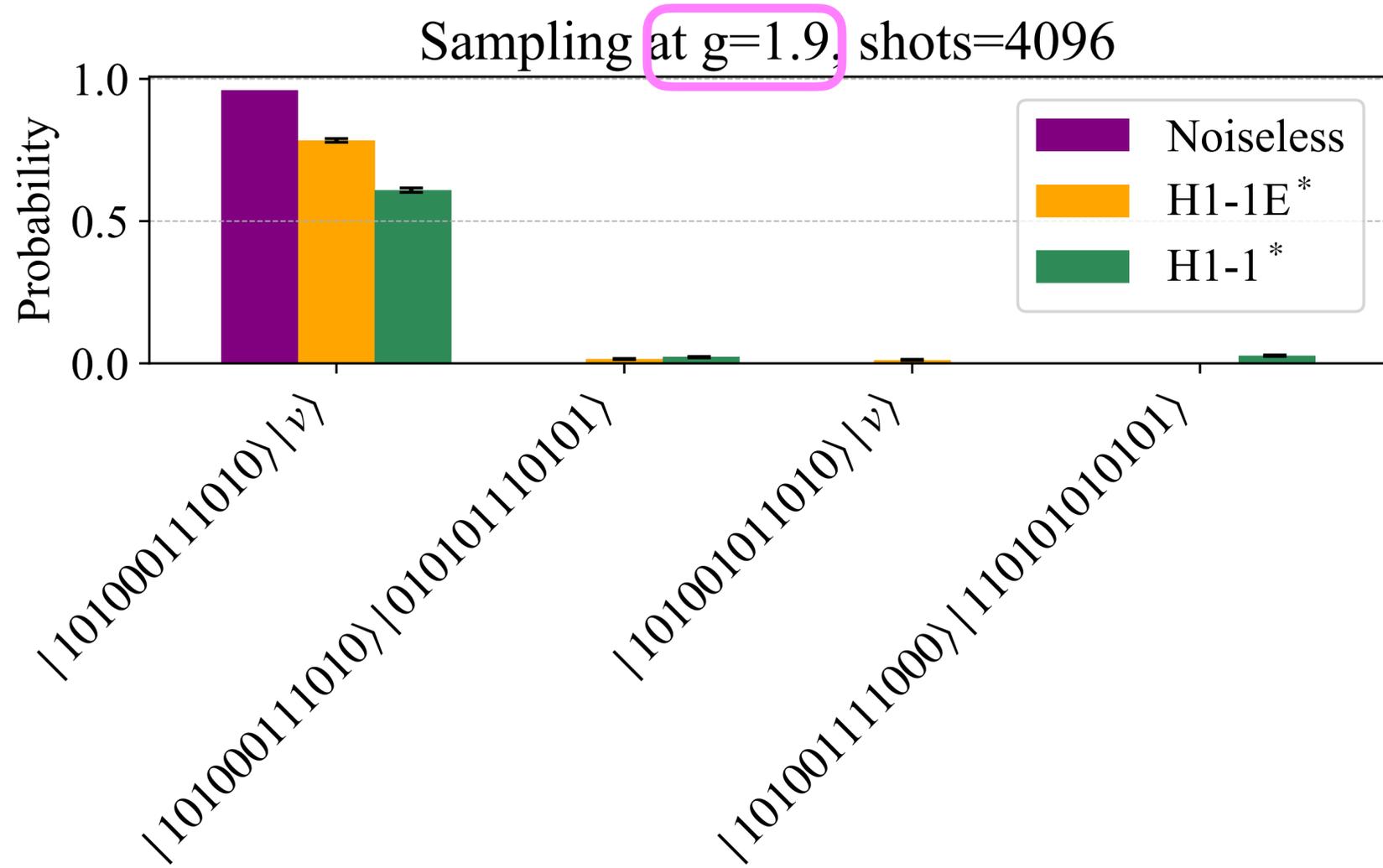




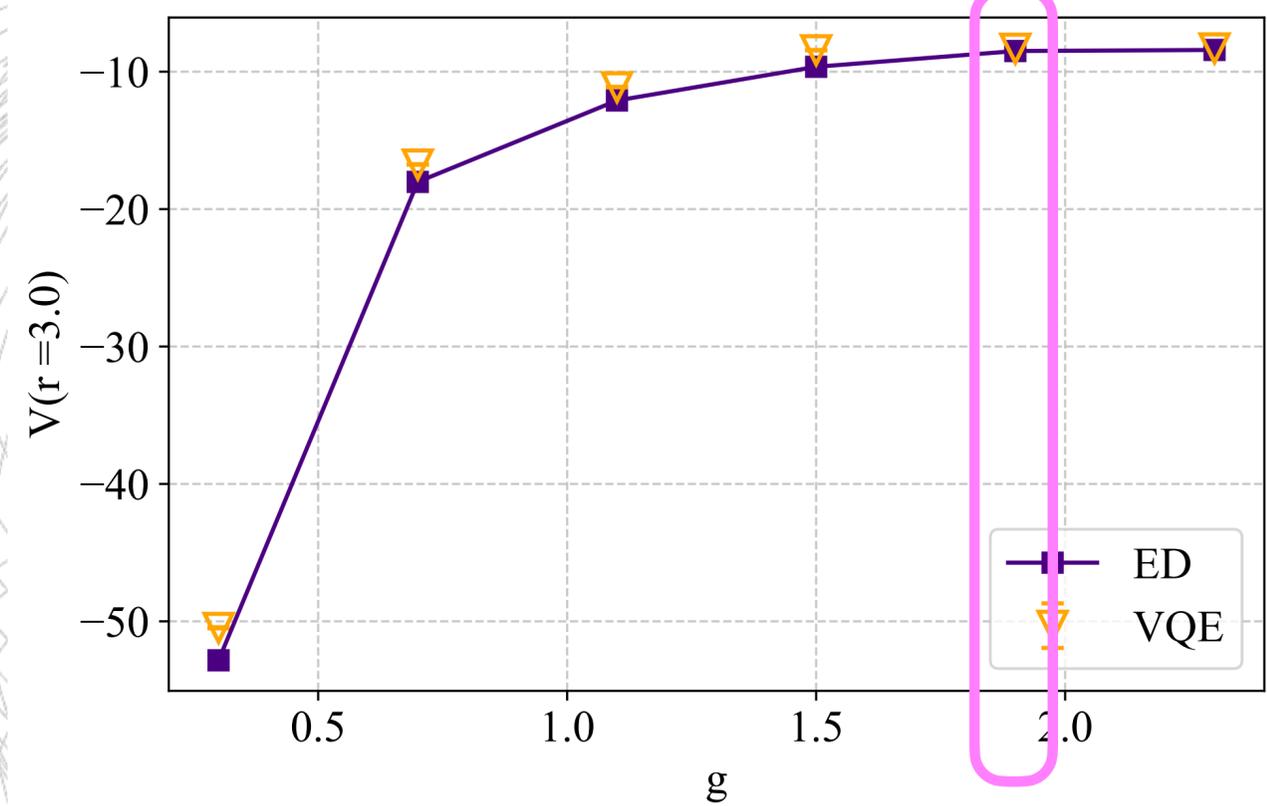
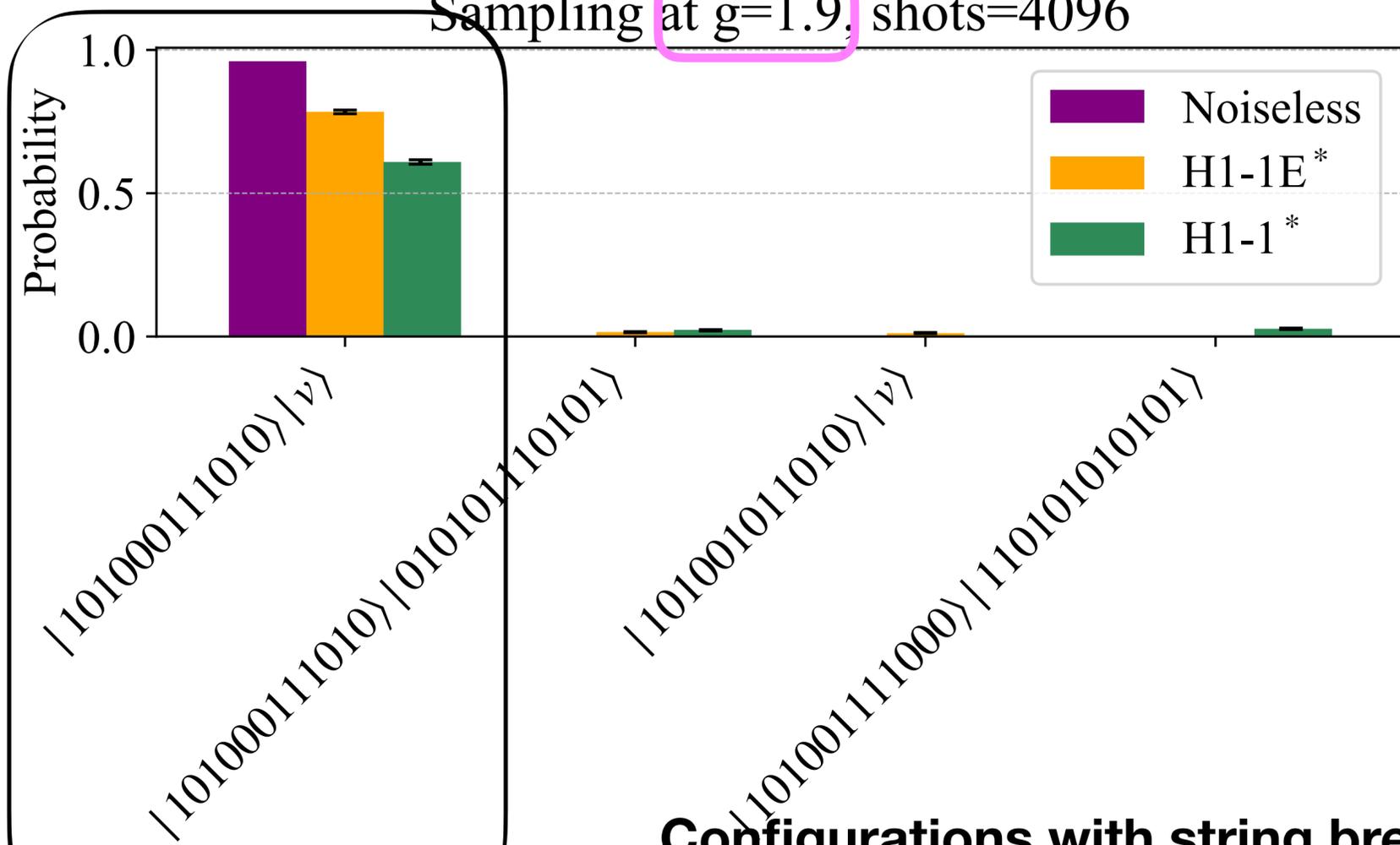


Configurations with electric fluxes between static charges





Sampling at $g=1.9$, shots=4096



Configurations with string breaking: no flux and mesons

