



東北大学

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Study on hadron-hadron interaction  
with femtoscopic technique  
~Interaction of hyperons and Kaon~

Hadrons and Hadron Interactions in QCD  
2024 (HHIQCD 2024)  
@ Kyoto 21, Nov., 2024



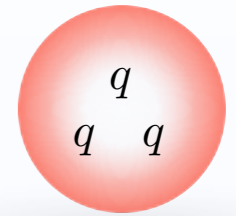
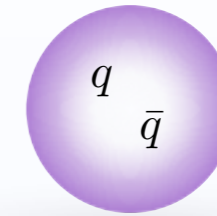
# Contents

- Introduction on Femtoscopy : hadron-hadron momentum correlation function in high-energy nuclear collisions
- $\bar{K}N$  correlation function and coupled-channel effect
- $\Lambda\alpha$  and  $\Xi\alpha$  correlation function
- Summary

# Study of hadron resonances and interaction

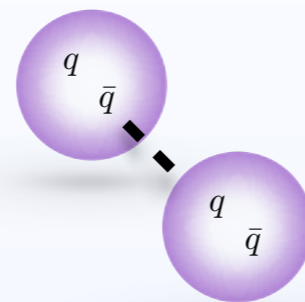
- Stable hadrons

- Described by simple quark configurations ( $q\bar{q}$ ,  $qqq$ )

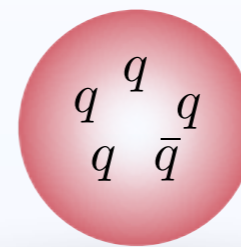


- Hadron resonances

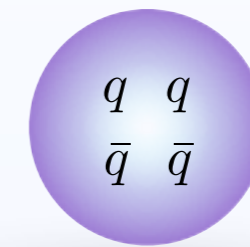
- Quark excitations
- Hadron molecules
- Exotic hadrons



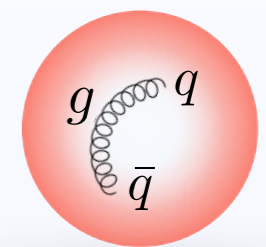
Hadron molecules



Pentaquark state



Tetraquark state



Hybrids

- Hadron resonances and hadron interaction

- Resonances : Eigenstates
- In scattering problem : Described as poles of scattering amplitude ( $\mathcal{F}$ )
- Detailed nature : Related with the detailed structure of amplitude

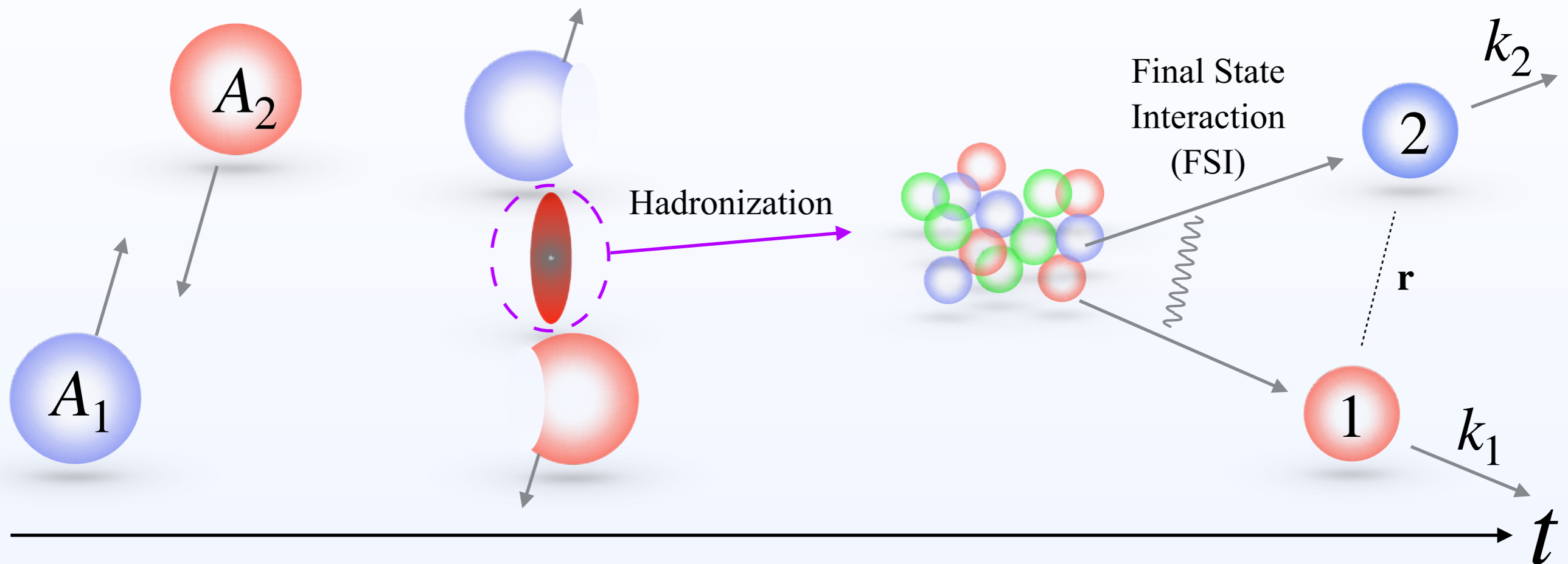
Understanding of resonance nature



Understanding of hadron interaction

# Femtoscscopy

- High energy nuclear collision and FSI

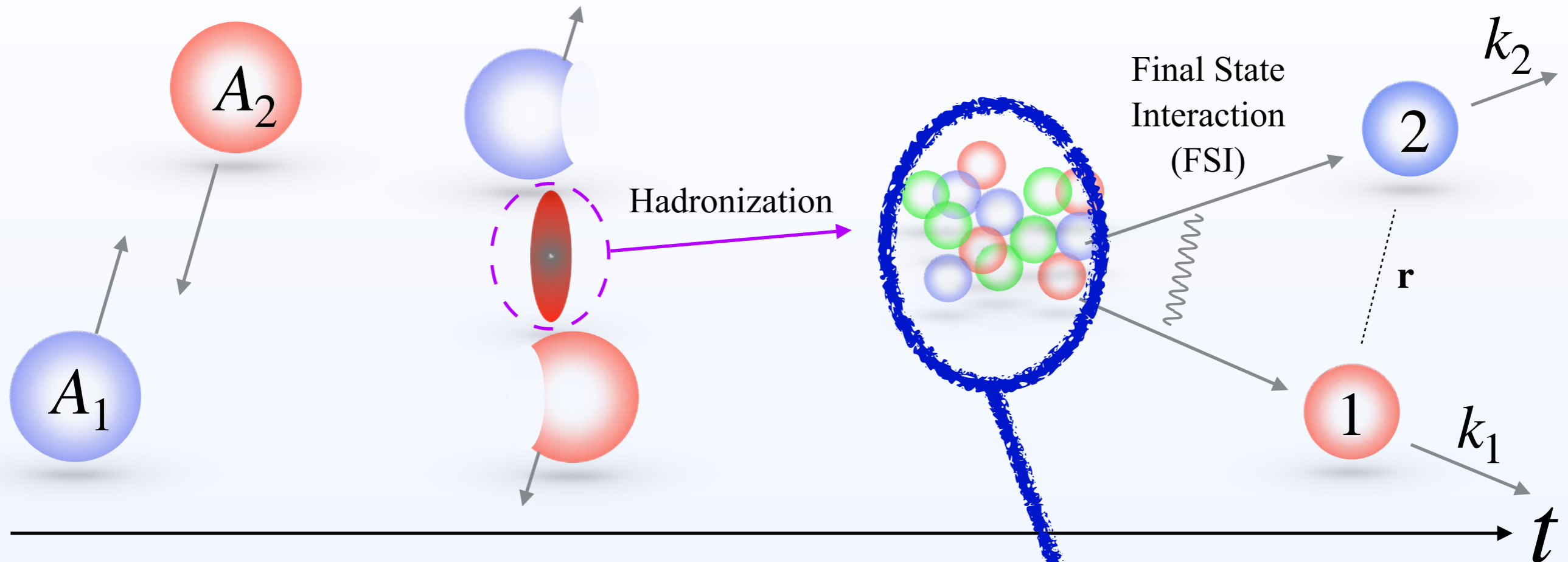


- Hadron-hadron correlation

$$C_{12}(k_1, k_2) = \frac{N_{12}(k_1, k_2)}{N_1(k_1)N_2(k_2)}$$
$$= \begin{cases} 1 & \text{(w/o correlation)} \\ \text{Others (w/ correlation)} \end{cases}$$

# Femtoscscopy

- High energy nuclear collision and FSI



- Hadron-hadron correlation

- Koonin-Pratt formula : S.E. Koonin, PLB 70 (1977)  
S. Pratt et. al. PRC 42 (1990)

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

$$\mathbf{q} = (m_2\mathbf{k}_1 - m_1\mathbf{k}_2)/(m_1 + m_2)$$

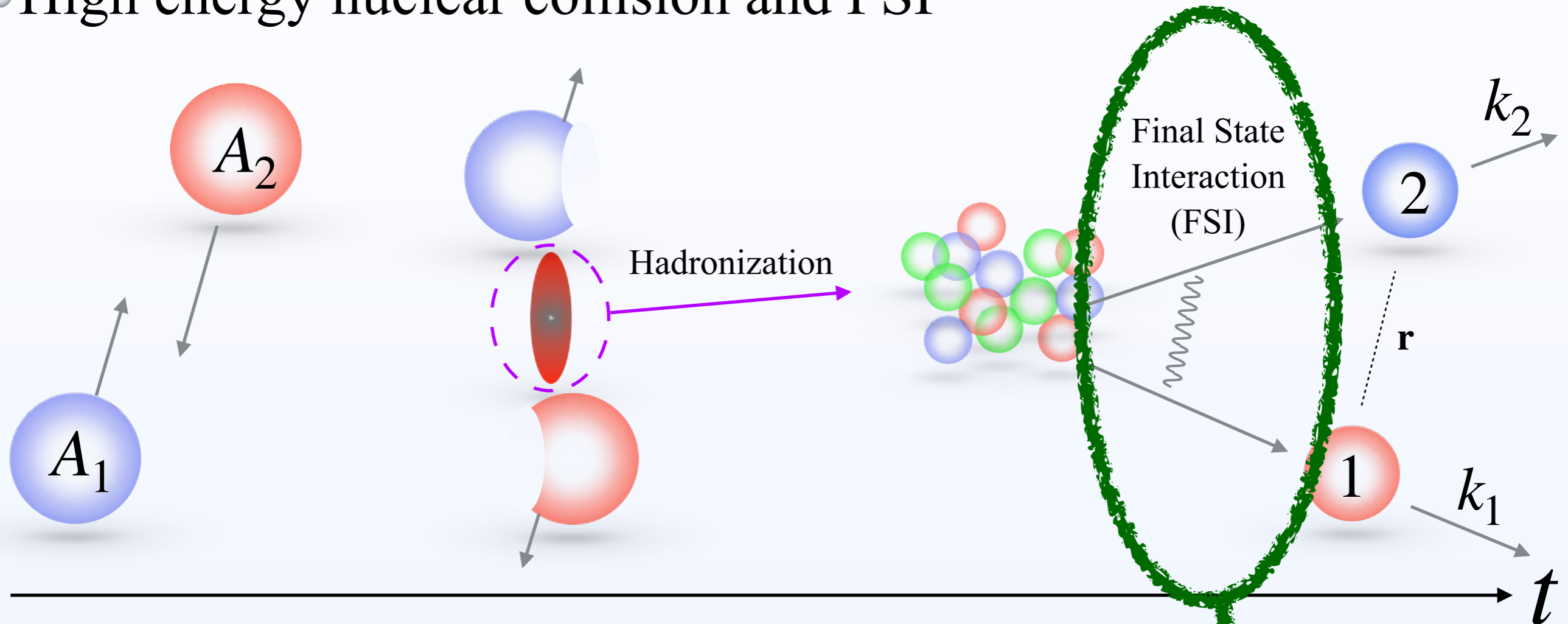
$S(\mathbf{r})$  : Source function

$\varphi^{(-)}(\mathbf{q}, \mathbf{r})$  : Relative wave function

- Depends on ...
- Collision detail ( $A_i$ , energy, centrality)
- Including information of...
  - size of hadron source,
  - momentum dependence, weight...

# Femtoscscopy

- High energy nuclear collision and FSI



- Hadron-hadron correlation

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S. Pratt et. al. PRC 42 (1990)

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$

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- Depends on ...

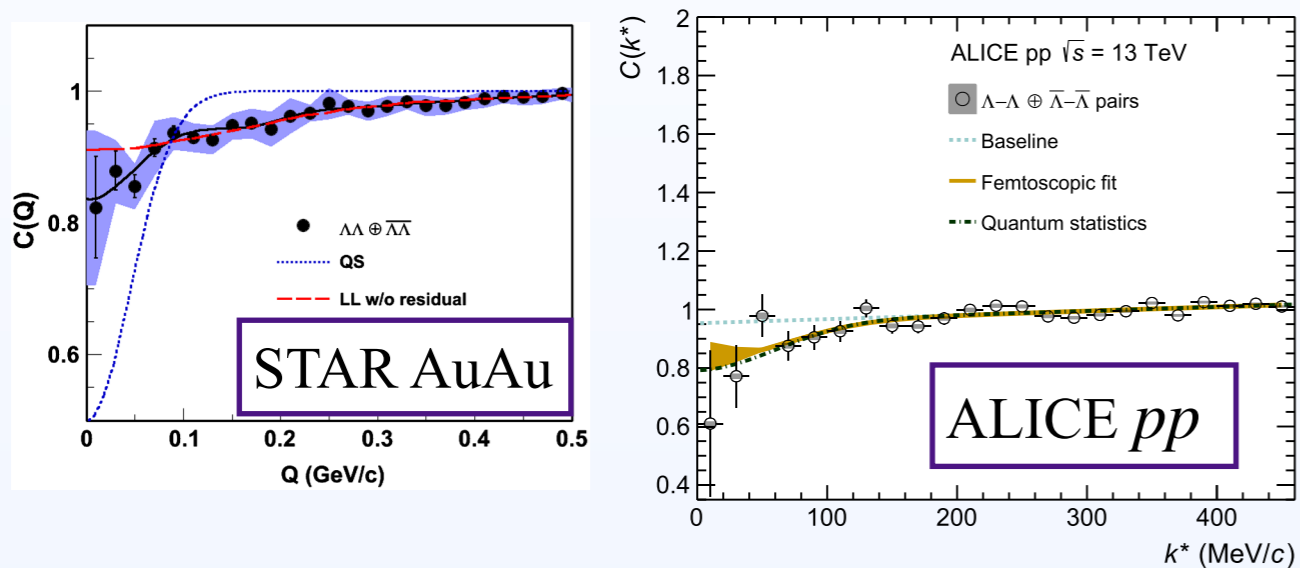
Interaction (strong and Coulomb)

quantum statistics (Fermion, boson)

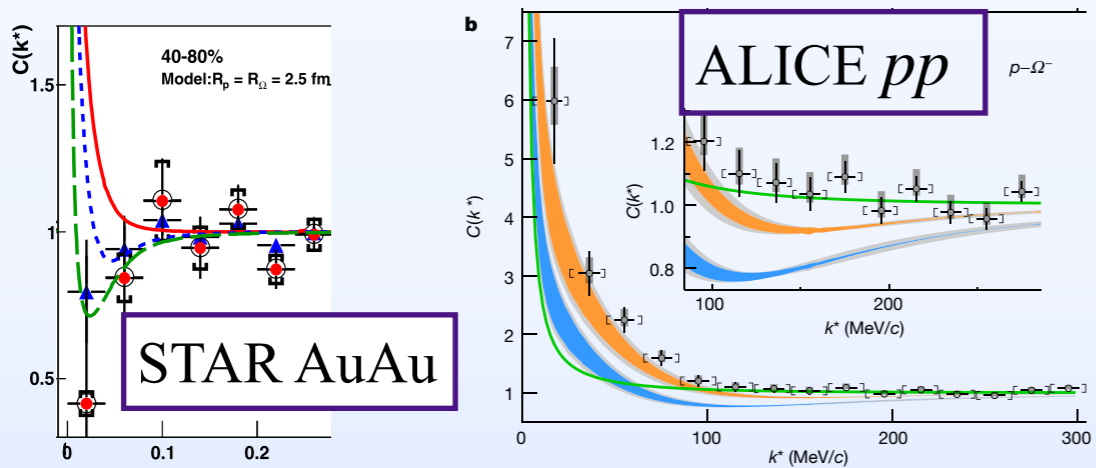
# Femtoscopic data

- Experimental data in various sectors

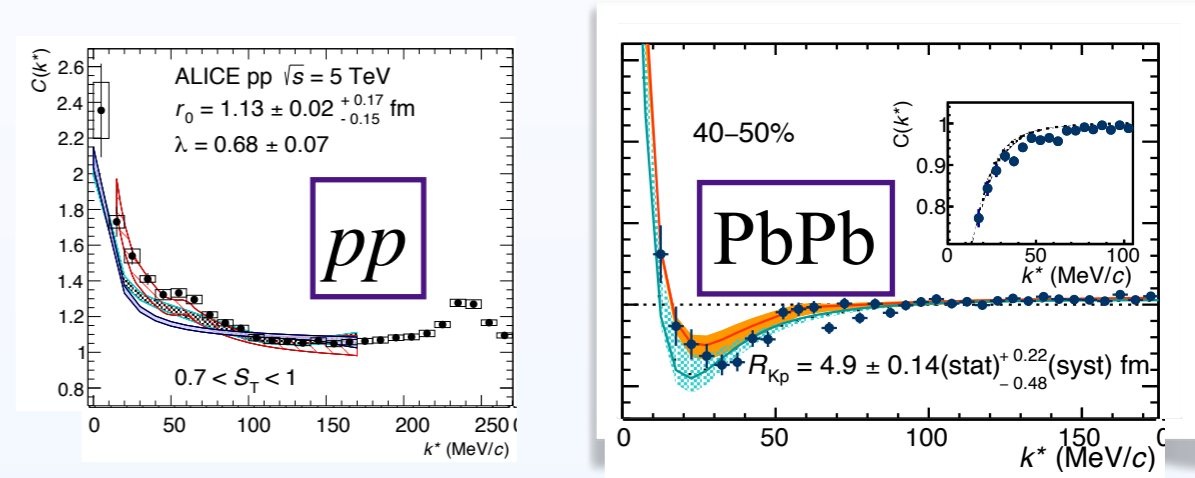
- $\Lambda\Lambda$  STAR AuAu: PRL 114,022301(2015)  
ALICE pp: PLB 797 (2019) 134822  
PbPb: PRC99, 024001 (2019)



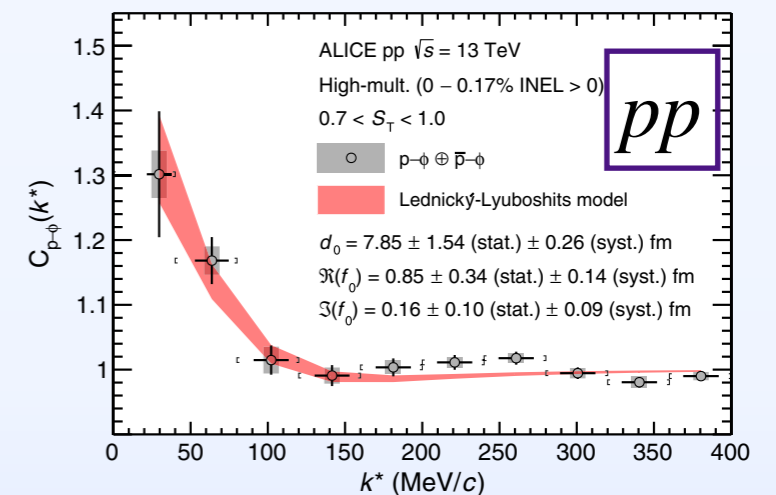
- $p\Omega$  STAR AuAu: PLB 790, 490 (2019)  
ALICE pp: Nature 588 (2020) 232



- $K^\pm p$  ALICE pp: PRL 124 (2020) 9, 092301  
PbPb: PLB 822 (2021) 136708  
STAR AuAu: NPA 982 (2019) 359



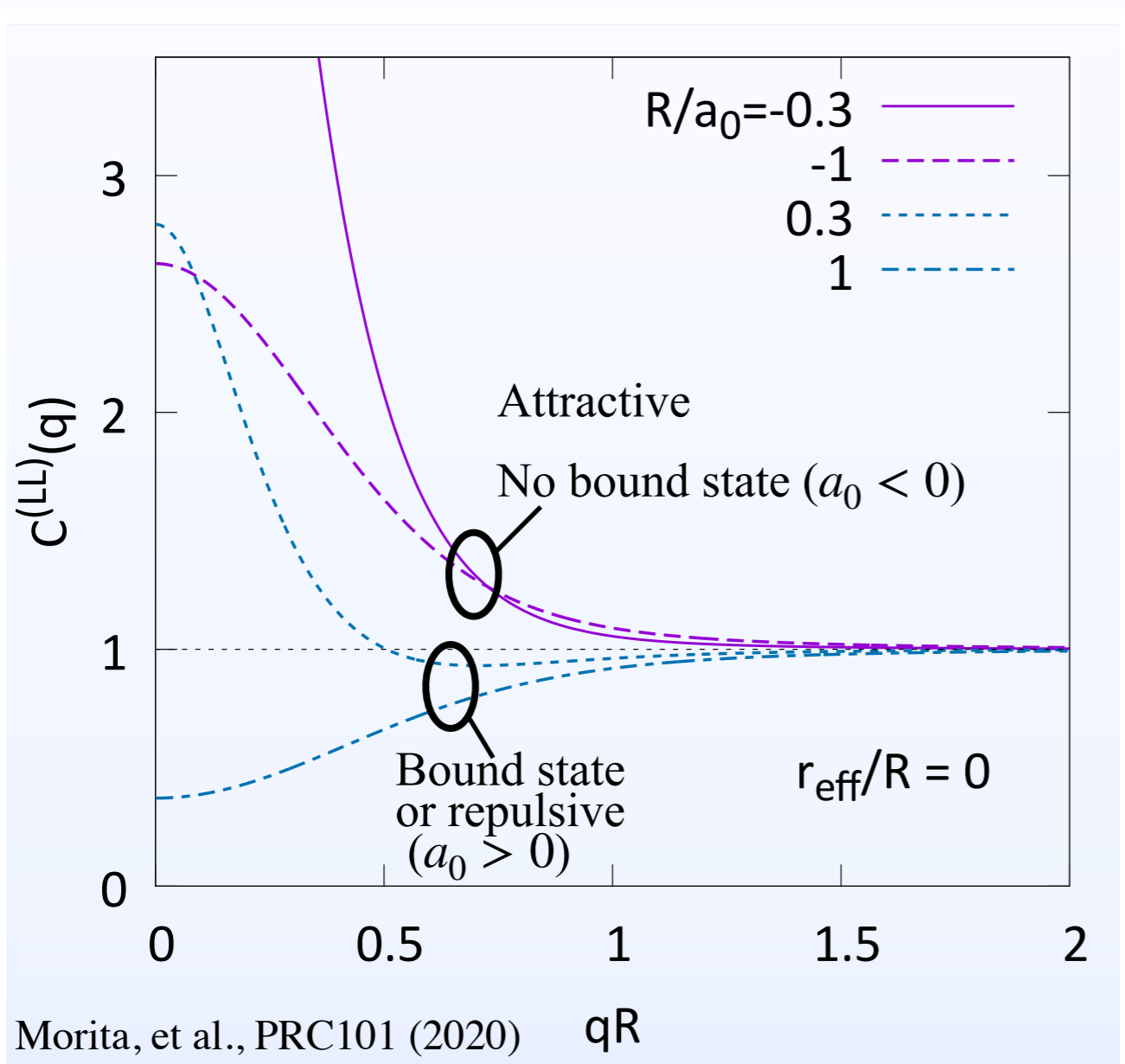
- $p\phi$  ALICE pp: PRL 127 (2021) 17, 172301



# Source size dependence

- Line shapes of  $C(q)$ : relation to interaction

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2$$



- Scattering length  $a_0$  and source size  $R$  determines the suppression/enhancement of line shape  
 $* a_0 = -\mathcal{F}(q=0)$

- Repulsive int. ( $a_0 > 0$ , small  $|a_0|$ )

➔ Suppressed  $C(q)$

- Attractive int. w/ bound state ( $a_0 > 0$ , large  $|a_0|$ )

➔ Suppressed  $C(q)$  for Large  $R$   
 Enhanced  $C(q)$  for small  $R$

- Attractive int. w/o bound state ( $a_0 < 0$ )

➔ Enhanced  $C(q)$



# Source size dependence

- Line shapes of  $C(q)$ : relation to interaction

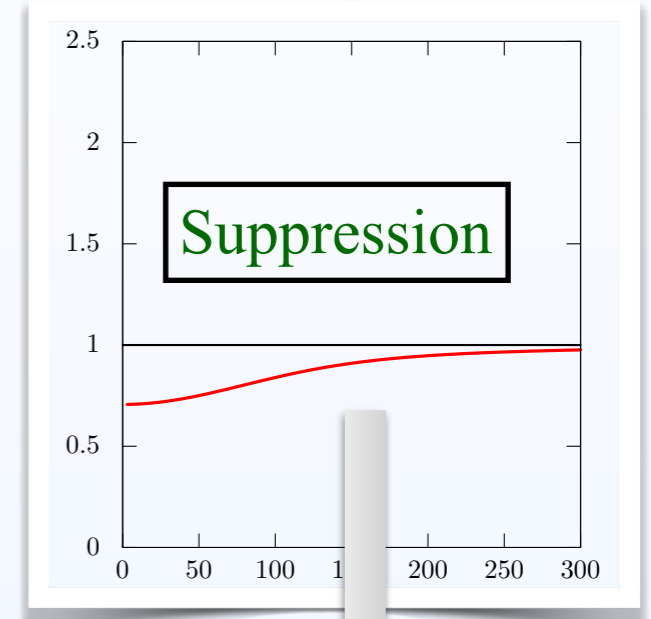
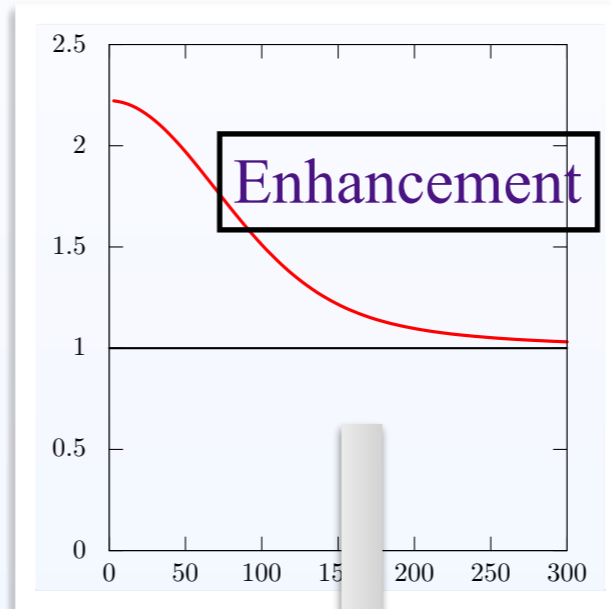
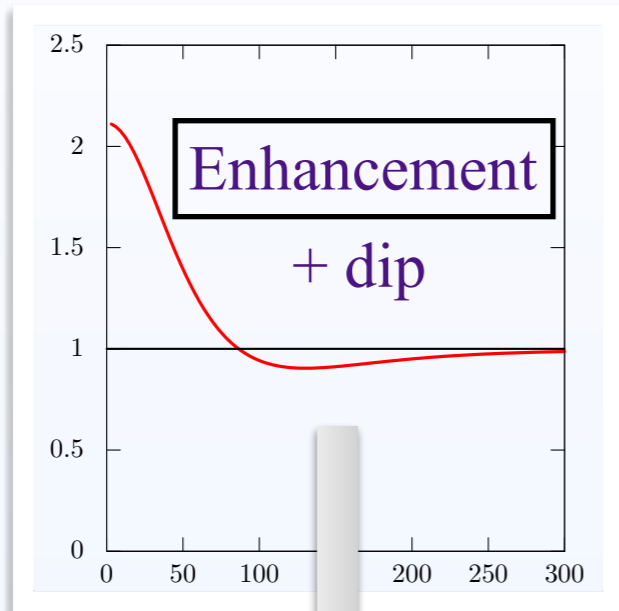
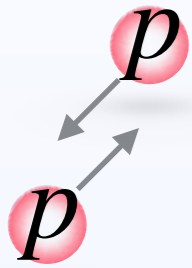
- Attractive interaction

w/ bound state

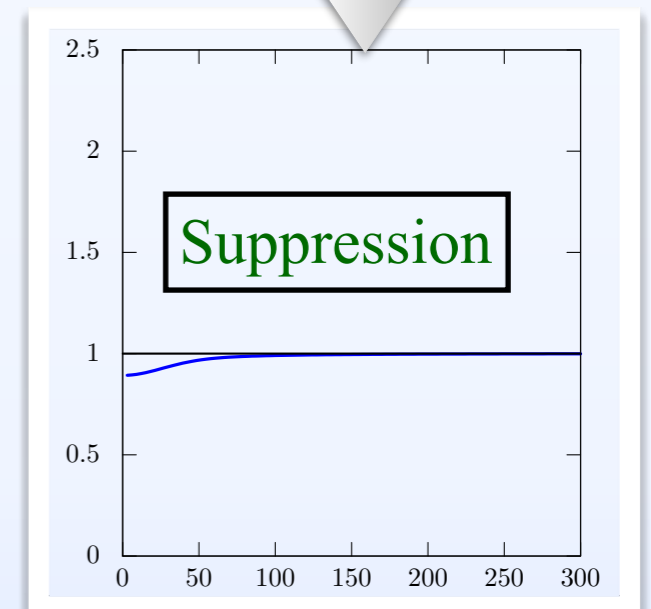
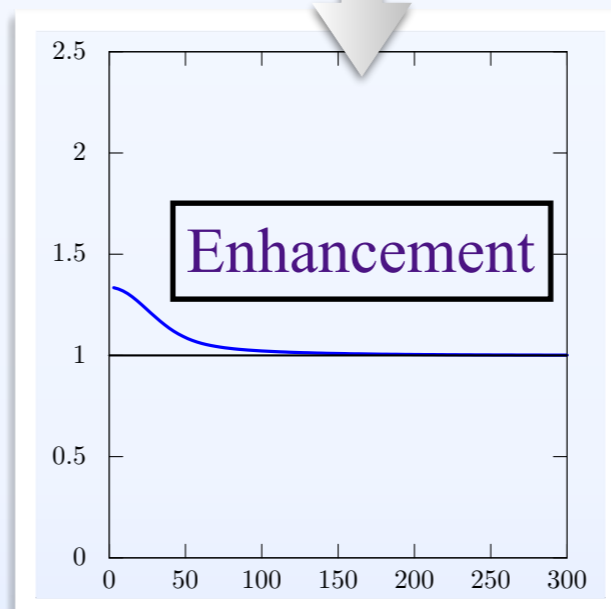
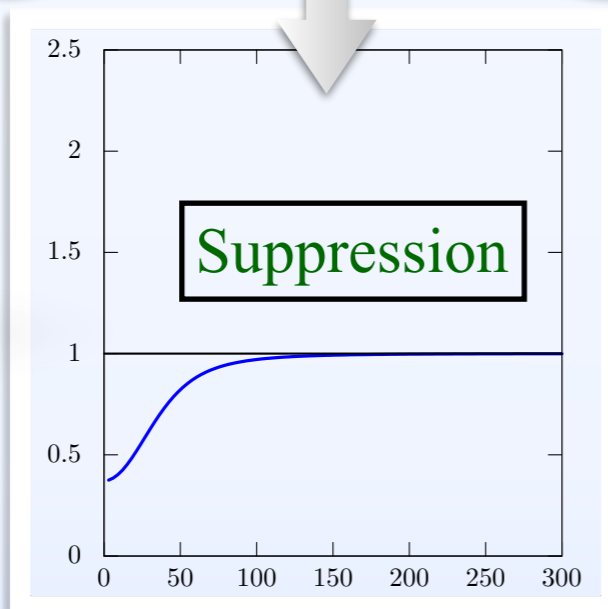
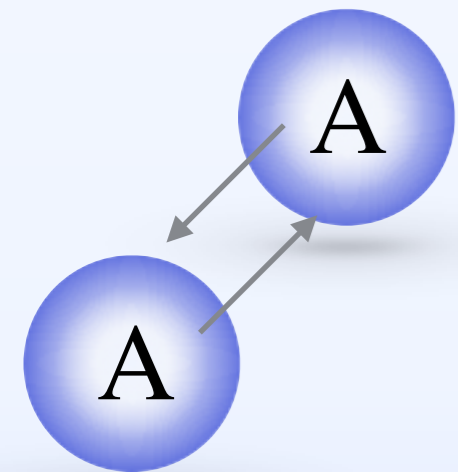
w/o bound state

- Repulsive interaction

- Small source



- Large source



Source size dependence for typical for bound state cases!

## • How to control source size $R$

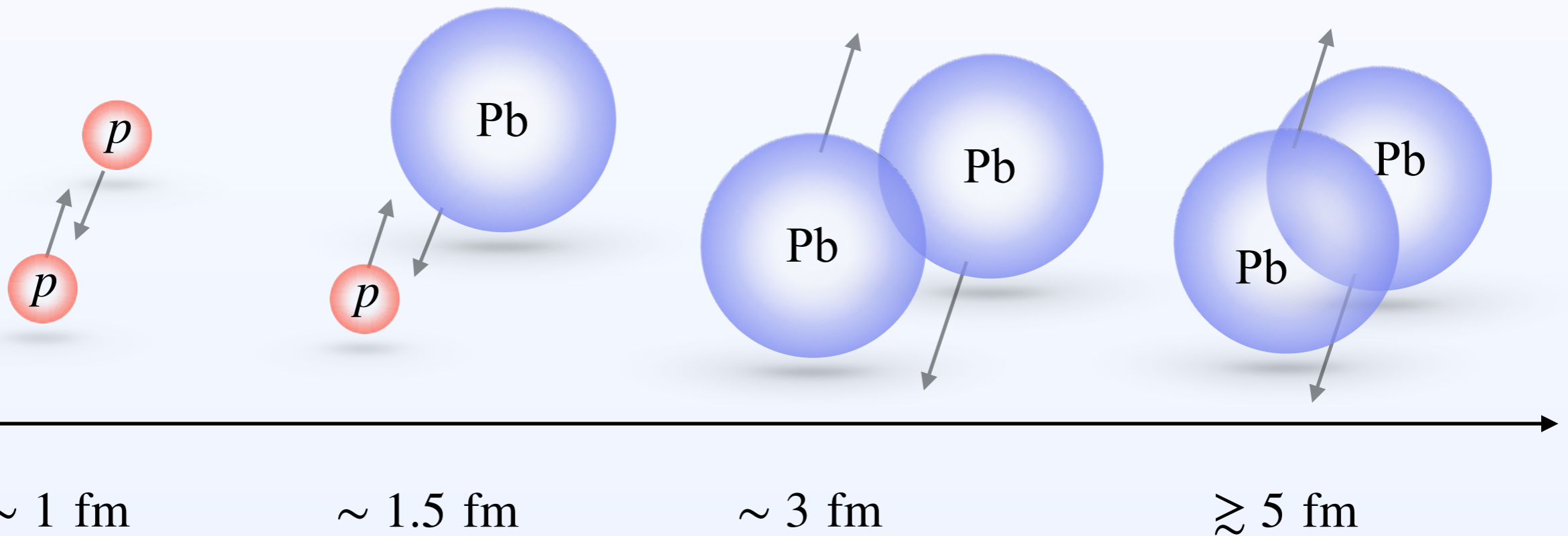
•  $pp$  collisions

•  $p\text{Pb}$  collisions

•  $\text{PbPb}$  collisions

• peripheral

• central





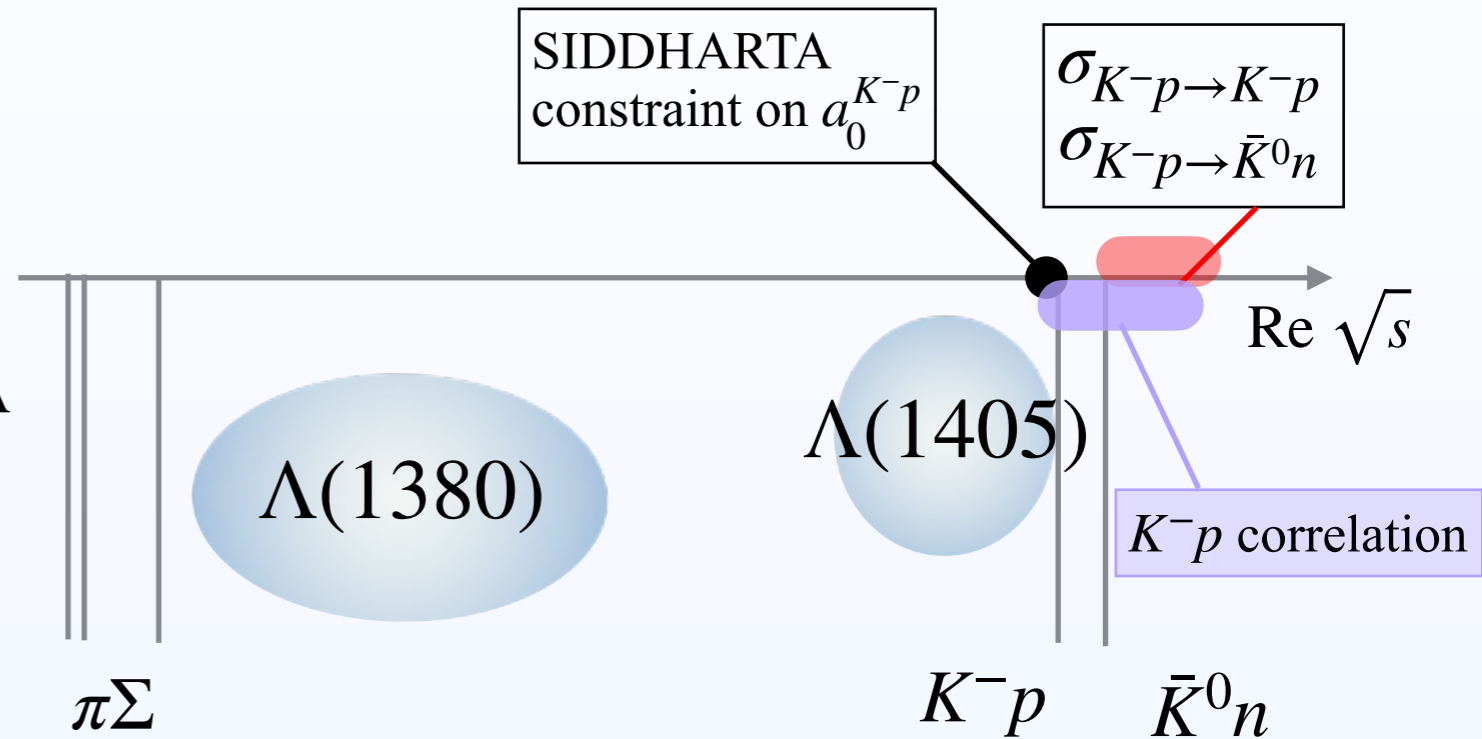
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# $\bar{K}N$ interaction and $K^-p$ correlation

## • $\bar{K}N$ interaction and $\Lambda(1405)$

- Coupled-channel system of  $\pi\Sigma$ - $\pi\Lambda$ - $\bar{K}N$
- Strong attraction reproducing quasi-bound state  $\Lambda(1405)$
- Strong constraint on  $a_0^{K^-p}$  by SIDDHARTA experiment of Kaonic hydrogen  
M. Bazzi, et al., PLB 704 (2011)
- Structure of  $\Lambda(1405)$  and  $\Lambda(1380)$ 
  - two pole structure  
J. A. Oller and U. G. Meißner, PLB500, 263 (2001)
  - $\bar{K}N$  molecular picture  
R.H. Dalitz, S.F. Tuan, PRL 425 (1959).



## • Chiral SU(3) based $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ potential Miyahara, Hyodo, Weise, PRC 98 (2018)

- Constructed based on the amplitude with NLO chiral SU(3) dynamics  
Ikeda, Hyodo, Weise, NPA881 (2012)
- Coupled-channel, energy dependent as

$$V_{ij}^{\text{strong}}(r, E) = e^{-(b_i/2 + b_j/2)r^2} \sum_{\alpha=0}^{\alpha_{\text{max}}} K_{\alpha,ij} (E/100 \text{ MeV})^\alpha$$

- Constructed to reproduce the chiral SU(3) amplitude around the  $\bar{K}N$  sub-threshold region 12

# $\bar{K}N$ interaction and $K^-p$ correlation

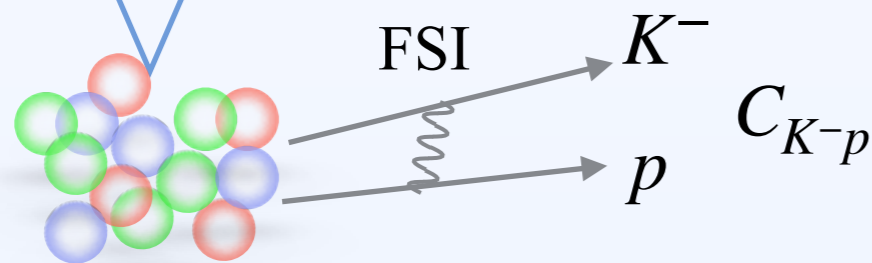
- Koonin-Pratt-Lednicky-Lyuboshits-Lyuboshits (KPLLL) formula

$$C(\mathbf{q}) = \int d^3\mathbf{r} S(\mathbf{r}) |\psi^{(-)}(q; r)|^2 + \sum_{j \neq K^-p} \omega_j \int d^3\mathbf{r} S_j(\mathbf{r}) |\psi_j^{(-)}(q; r)|^2$$

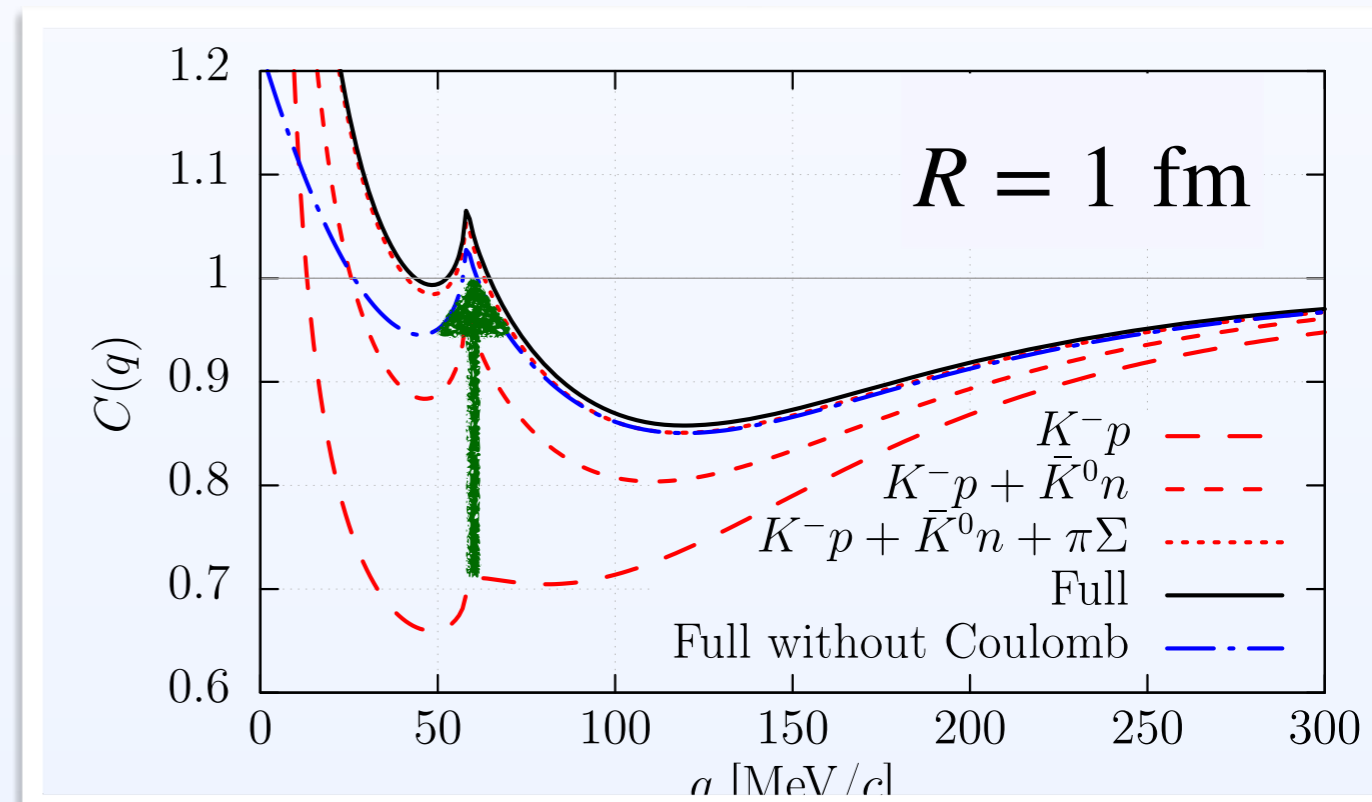
S.E. Koonin, PLB 70 (1977)  
 S. Pratt et. al. PRC 42 (1990)  
 R. Lednicky, et.al. Phys. At. Nucl. 61(1998)

- Contribution from coupled-channel source

$K^-p, \bar{K}^0n, \pi^0\Sigma^0, \pi^+\Sigma^-, \pi^-\Sigma^+, \pi^0\Lambda$

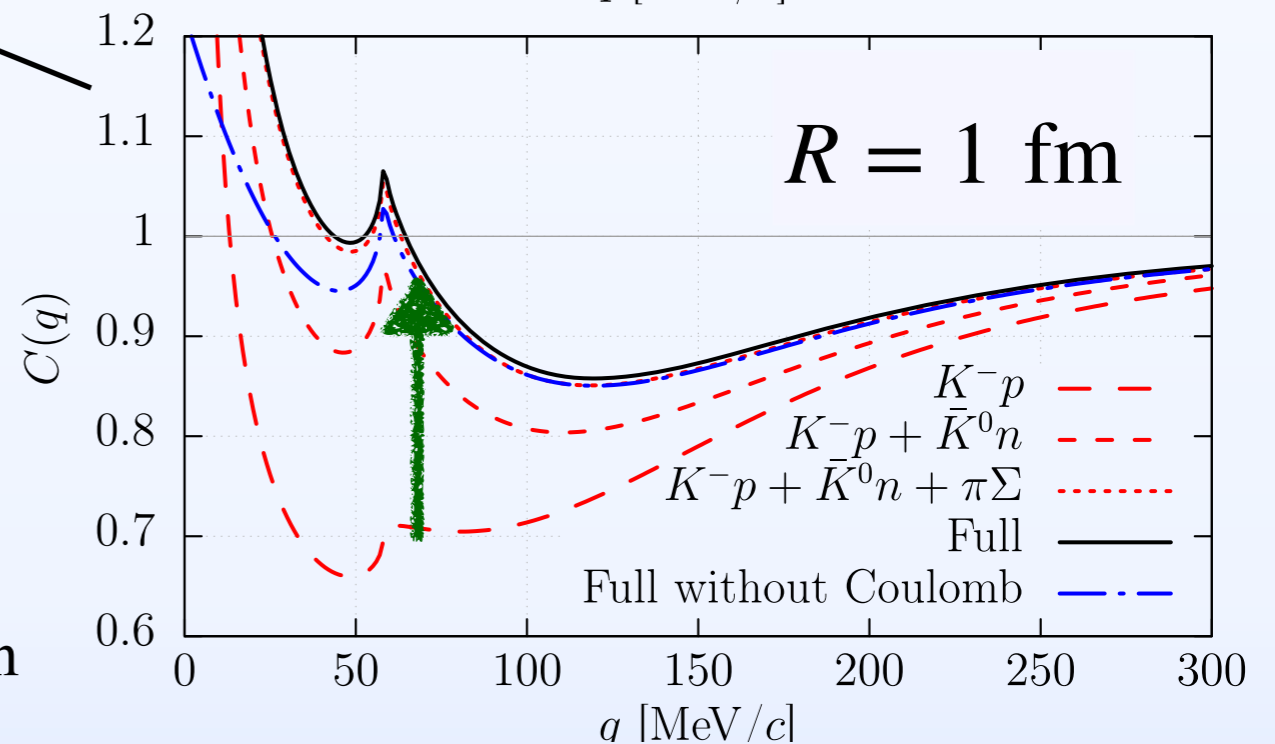
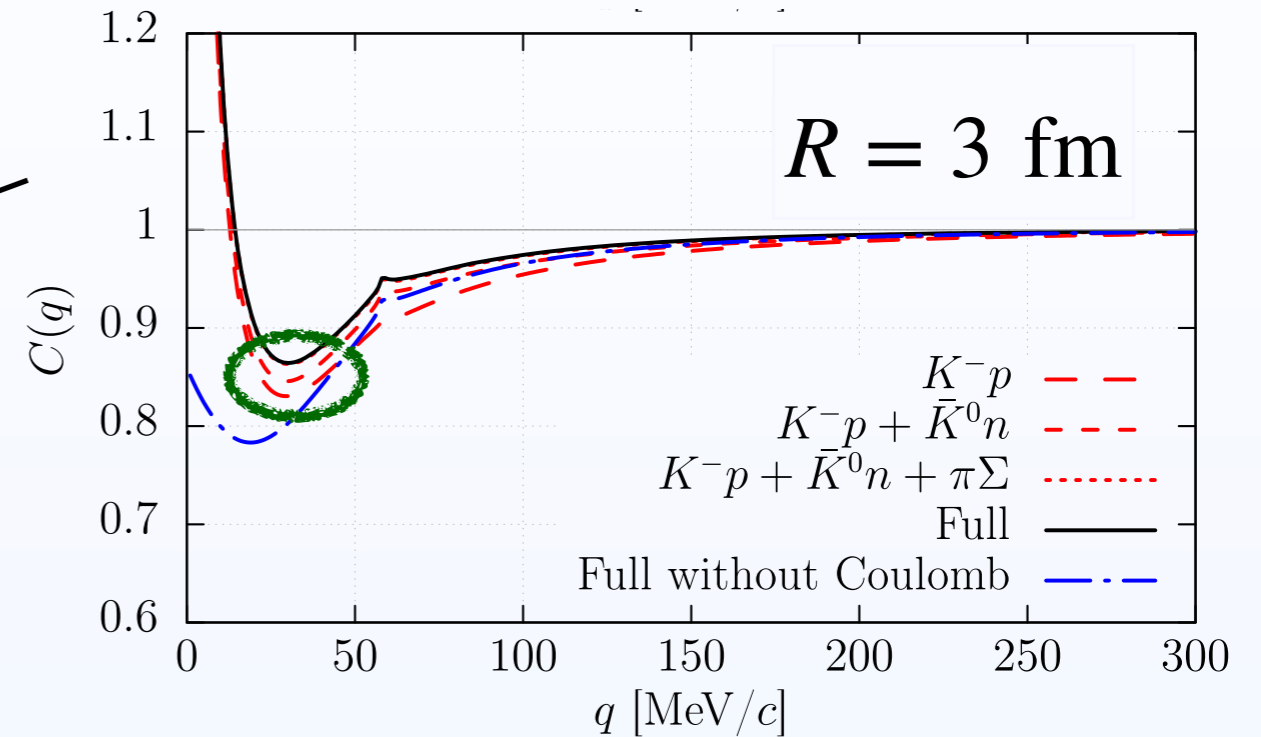
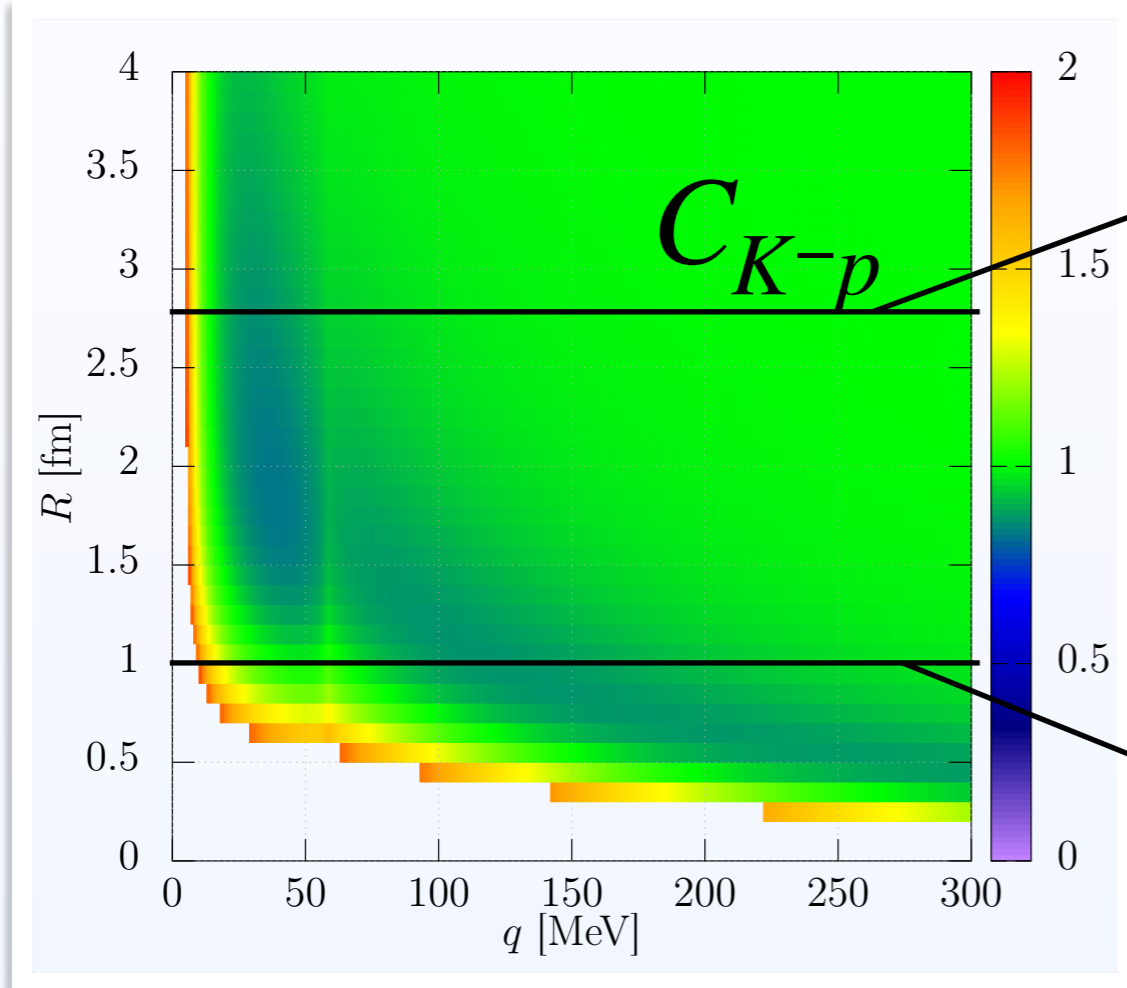


- Enhance  $C(q)$
- Enhance cusp structure
- $\omega_i$  : production rate  
(compared to measured channel)



# $\bar{K}N$ interaction and $K^-p$ correlation

## Source size dependence of coupled-channel effect



- Strong source size dependence  
 $\Leftarrow$  Due to the near-threshold  $\Lambda(1405)$  pole
- $C(q)$  with large source
  - Less prominent cusp structure
  - Weaker coupled-channel source contribution
- Large source :  $K^-p$  scattering
- Small source : detailed coupled-channel effect

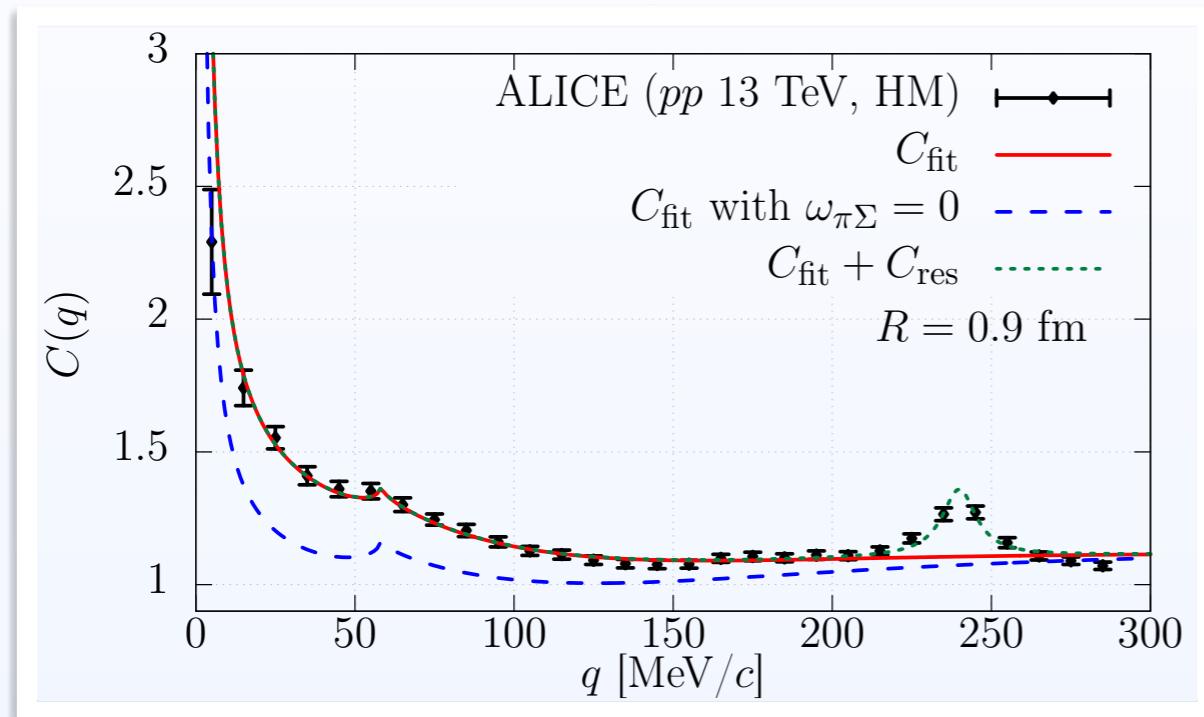


# $\bar{K}N$ interaction and $K^-p$ correlation

## Source size dependence with $K^-p$ data

- ALICE  $pp$  collision data

ALICE PRL 124, 092301 (2020)

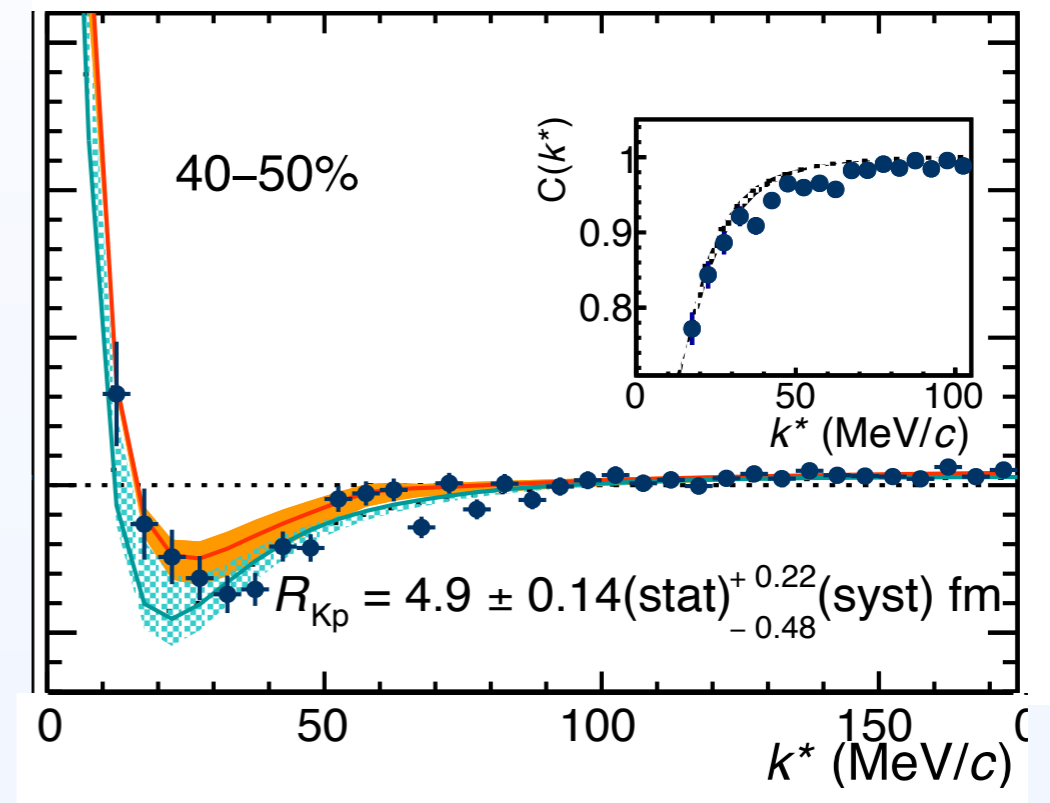


Kamiya, Hyodo, Morita, Ohnishi, Weise, PRL 124 (2020) 13, 132501

- Small source
- Clear  $\bar{K}^0n$  cusp structure
- Sizable contribution from coupled-channel source required to reproduce data

- ALICE PbPb collision data

ALICE PLB 822 (2021) 136708



- Large source
- Weaker cusp
- Consistent with analysis only with  $K^-p$  source



- Chiral SU(3) dynamics describes the both correlation data well.

# $\bar{K}N$ interaction and $K^-p$ correlation

- Latest  $K^-p$  correlation results

ALICE, EPJC 83 (2023)

- $pPb$  : 0-20%, 20-40% 40-100%
- $PbPb$  : 60-70%, 70-80% 80-90%

- **Discrepancy** around  $\bar{K}^0n$  threshold between **chiral SU(3) model** and exp. data for small source data

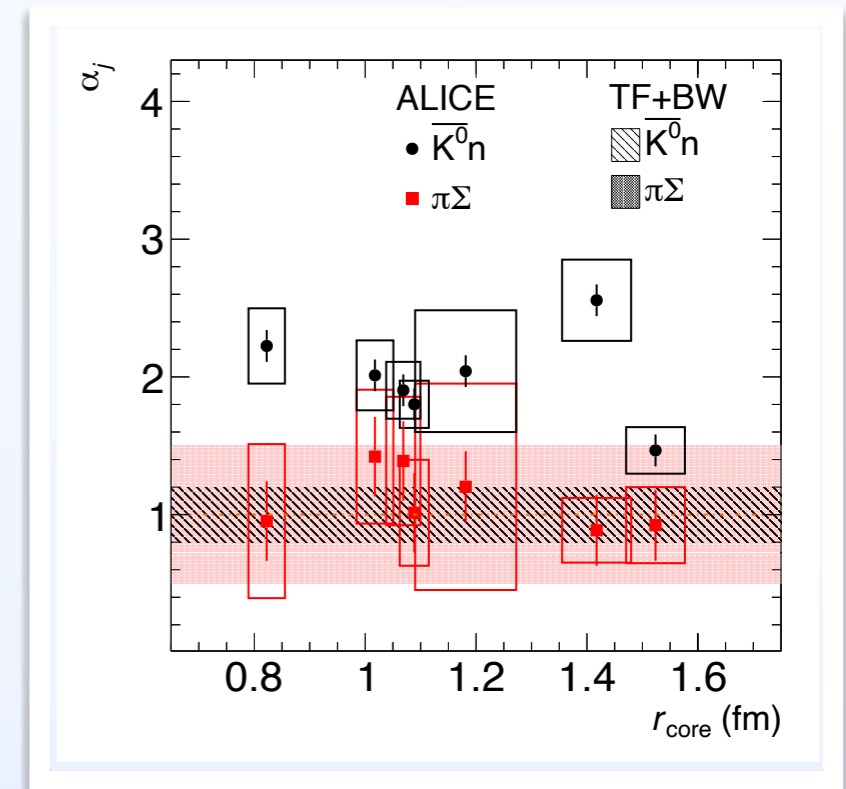
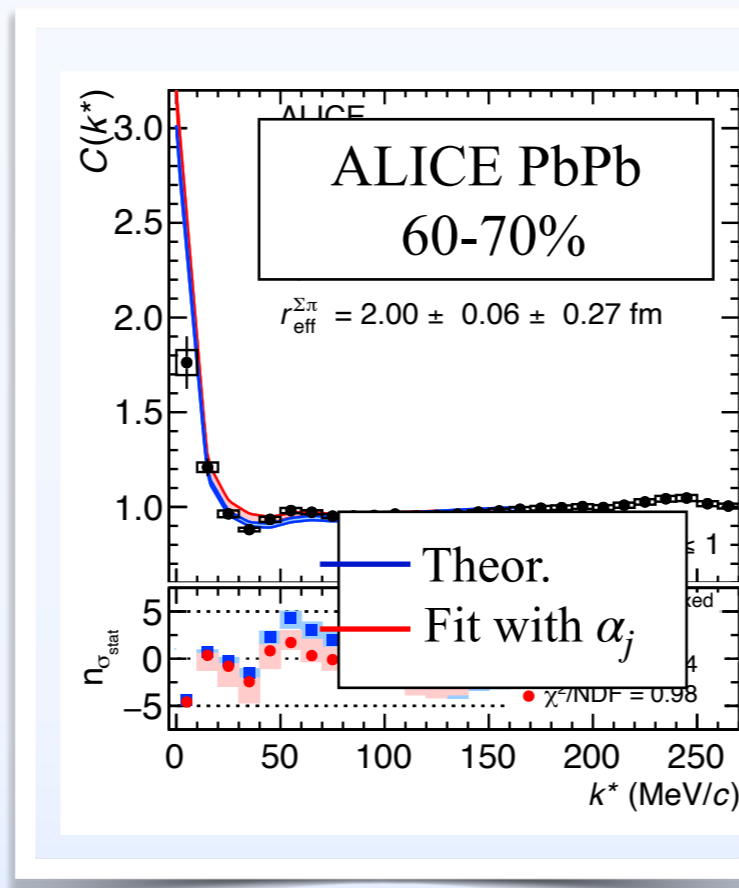
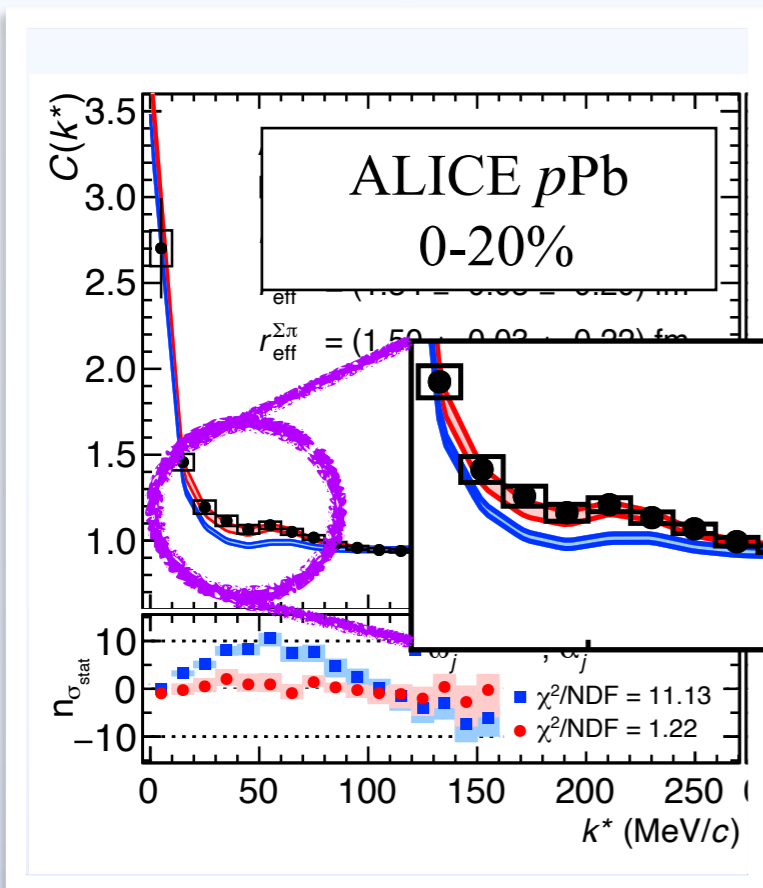
- Analysis with scale factor  $\alpha_j$

- Scale the coupled-channel source contribution by scaling factor

$$C_{K^-p} = C_{K^-p}^{\text{el}} + \sum_j \alpha_j C_j^{\text{inel}}$$

- $\alpha_{\bar{K}^0n} \sim 2$  gives better agreement

→ implying the stronger coupling







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# $N\Xi$ interaction and $H$ -dibaryon state

- $\Lambda\Lambda$ - $N\Xi$  interaction ( $S = -2$ ) and  $H$ -dibaryon

- $J = 0$ : Unique sector in flavor Octet-Octet baryon int.

$$8 \otimes 8 = \mathbf{1} \oplus 8_A \oplus 8_S \oplus 10 \oplus \bar{10} \oplus 27$$

- Pauli arrowed
- Attractive color-magnetic int.

- Flavor-singlet dihyperon “H” R. L. Jaffe, PRL 38 (1977), 195.

Predicted as “single hadron” below  $\Lambda\Lambda$

- Binding energy of double  $\Lambda$  hypernucleus  
Takahashi et al., PRL87 (2001) 212502

→  $\Lambda\Lambda$  does NOT form (deep) bound state

- HAL QCD  $\Lambda\Lambda$ - $N\Xi$  coupled-channel potential

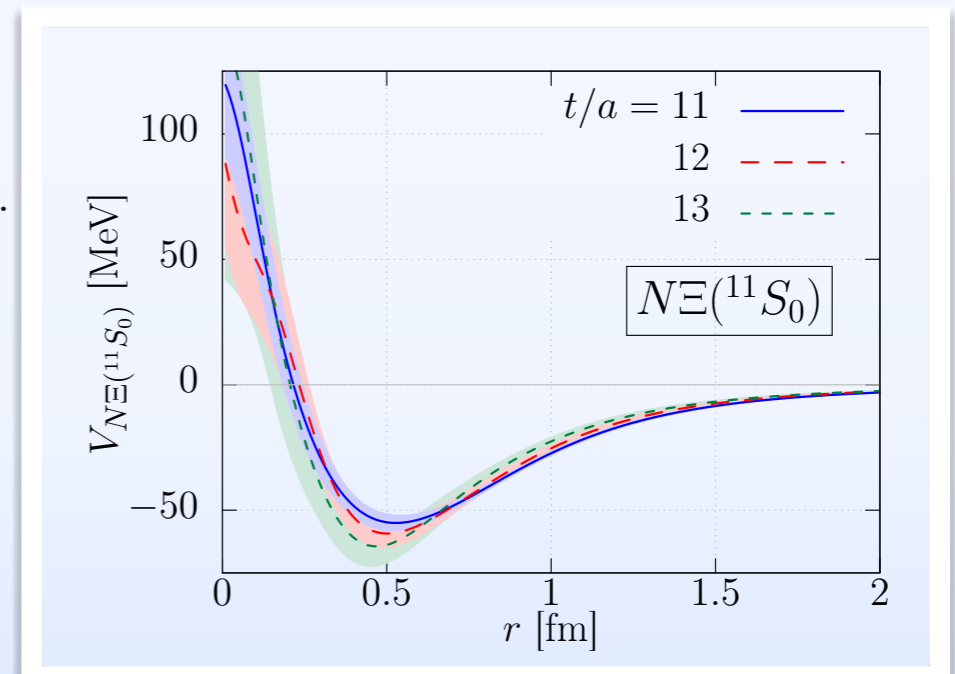
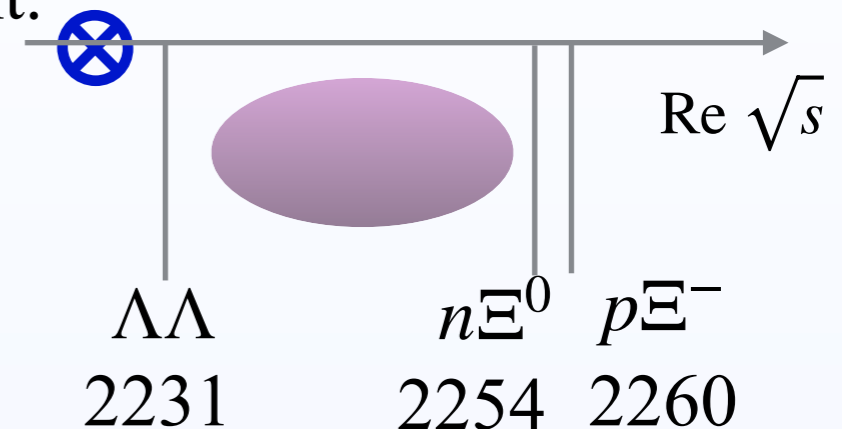
K. Sasaki et al. [HAL QCD], NPA 998 (2020), 121737.

- Strong attraction in  $J = 0, I = 0$   $N\Xi$  channel

$$a_0^{p\Xi^-(J=0)} = -1.21 - i1.52$$

$H$  dibaryon state is just barely unbound.

Fate of  $H$ -dibaryon?



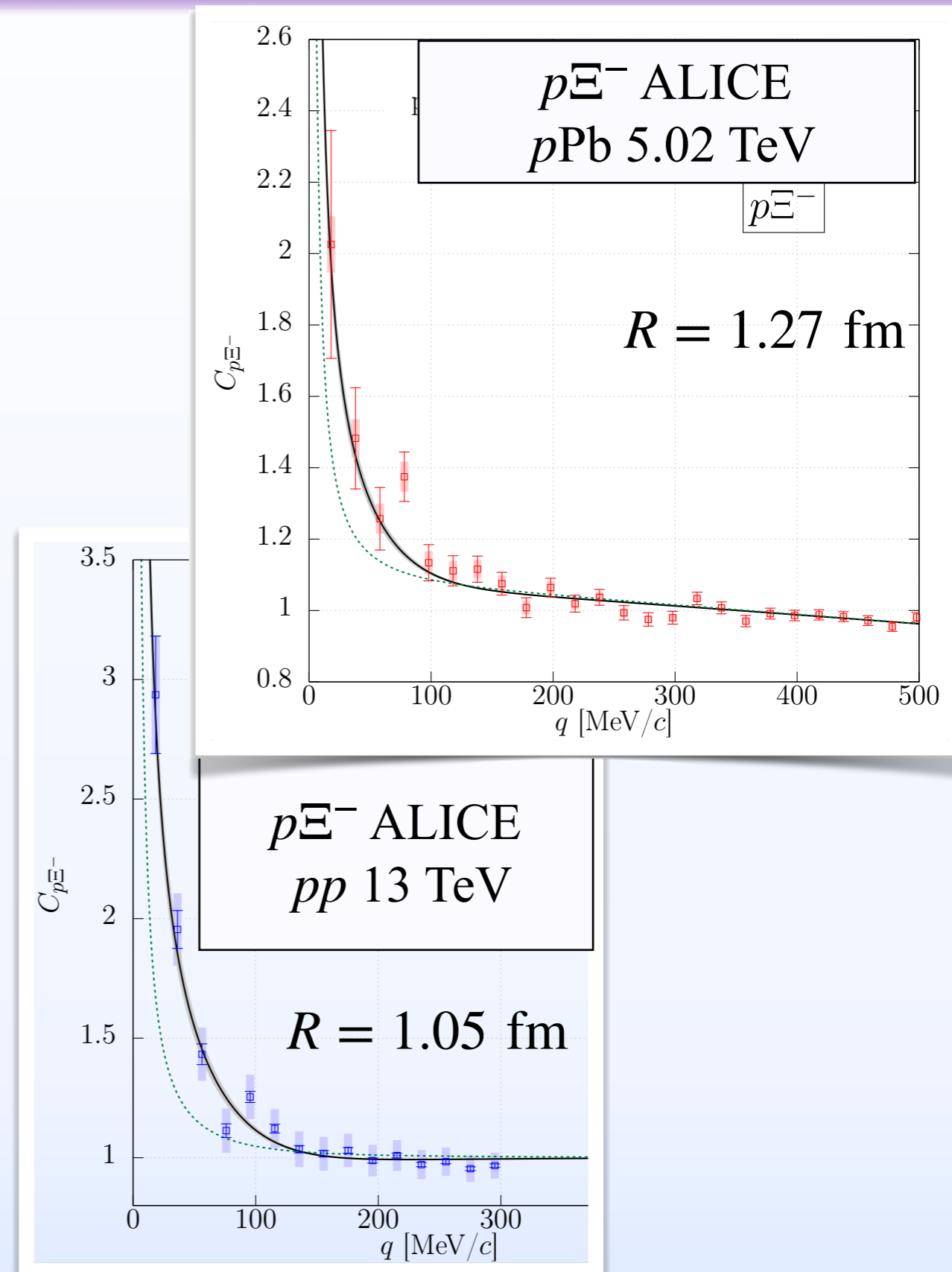
# $p\Xi^-$ correlation function

- $p\Xi^-$  correlation function

$$C_{p\Xi^-} = \frac{1}{4} C_{p\Xi^-, \text{singlet}} + \frac{3}{4} C_{p\Xi^-, \text{triplet}}$$

Couples to  $\Lambda\Lambda$   
(H-dibaryon channel)

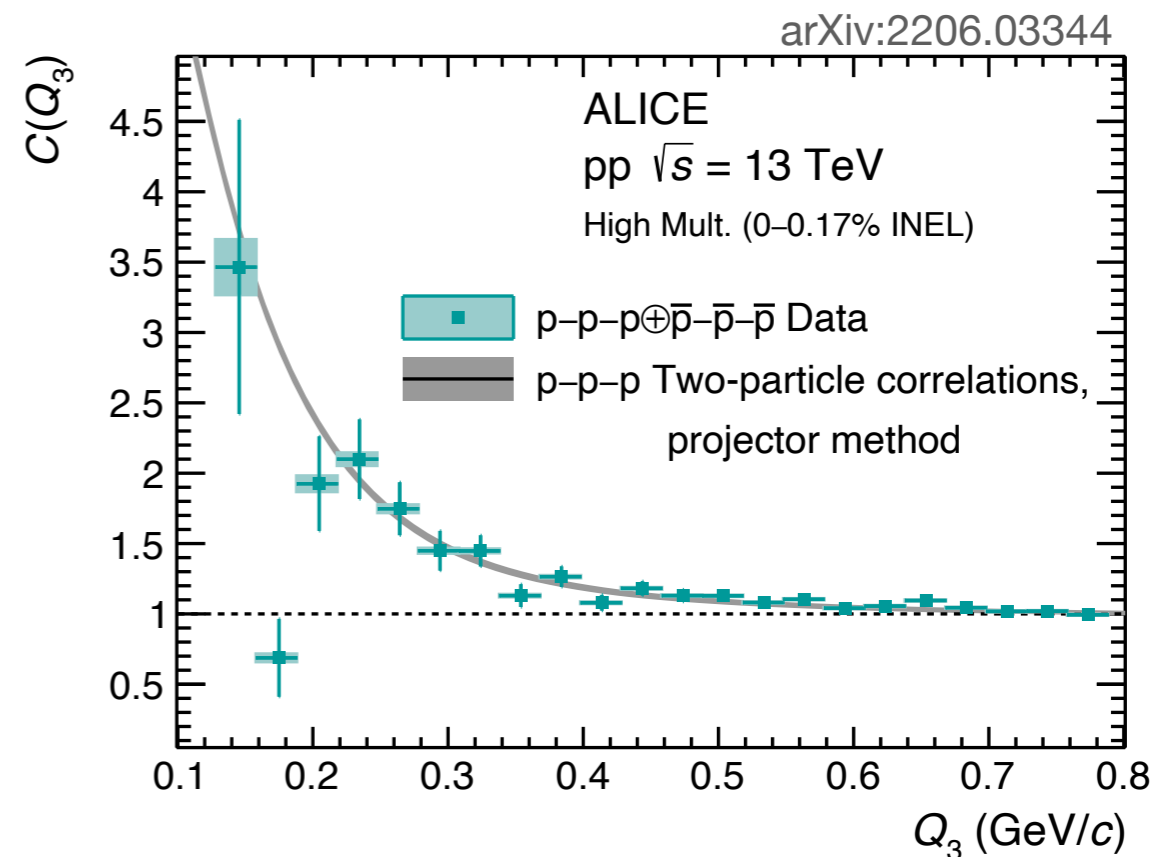
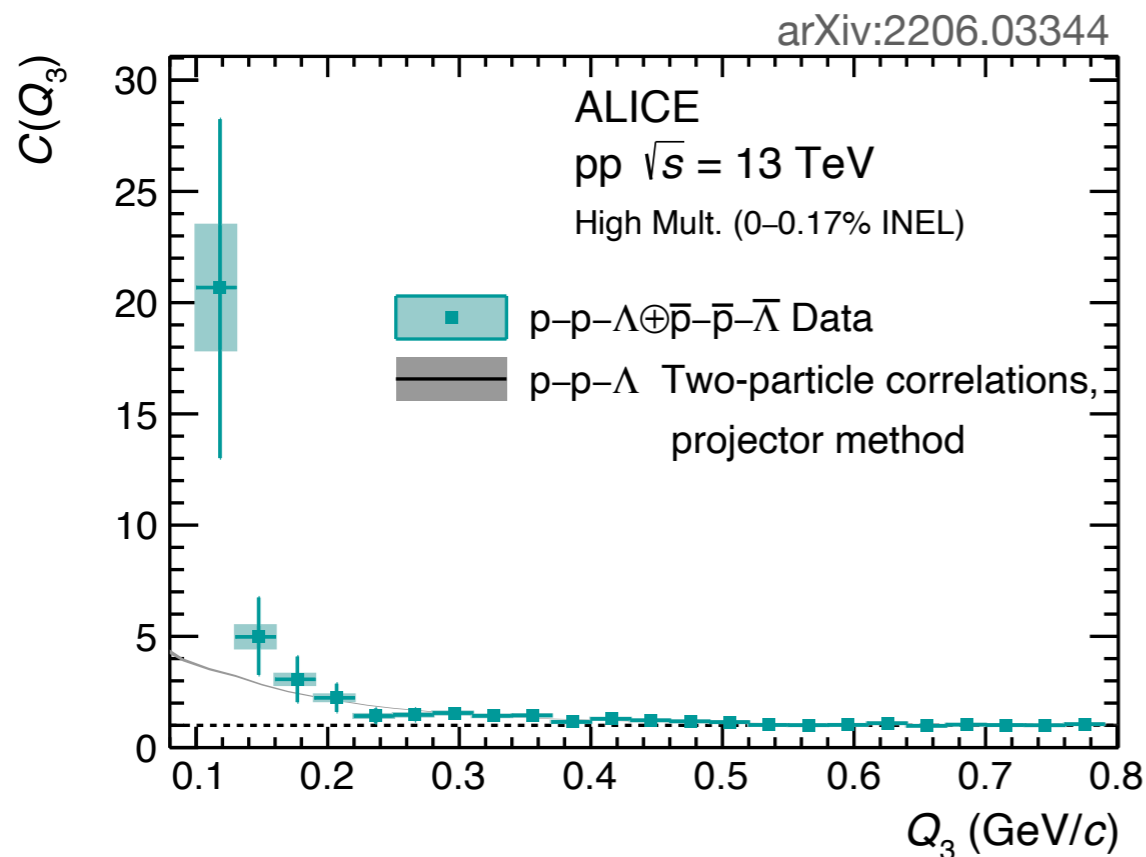
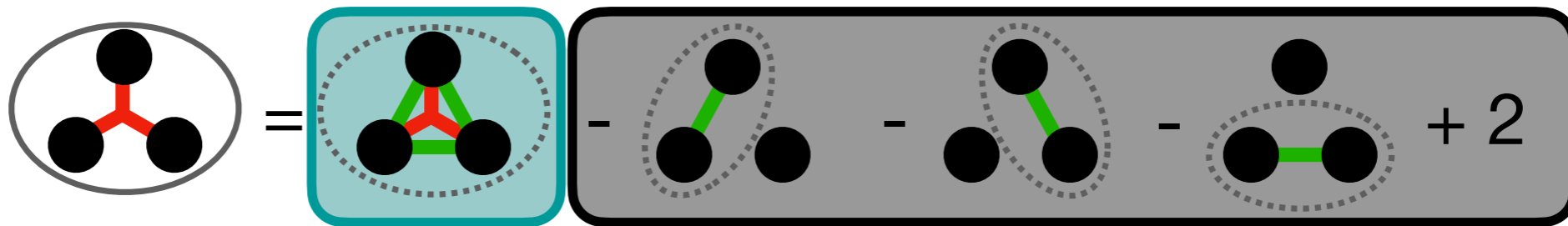
- Enhancement from **pure Coulomb** case
- Comparison with ALICE data  
pPb 5.02 TeV, pp 13 TeV collisions :  
S. Acharya et al. [ALICE], PLB 797 (2019).
- Spin channel reduction  
Singlet : Stronger enhancement  
Triplet : Weaker enhancement



# Correlation with few body systems

Three body correlation : 
$$C(\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3) \equiv \frac{P(\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3)}{P(\mathbf{p}_1)P(\mathbf{p}_2)P(\mathbf{p}_3)}$$

## p-p- $\Lambda$ and p-p-p correlations

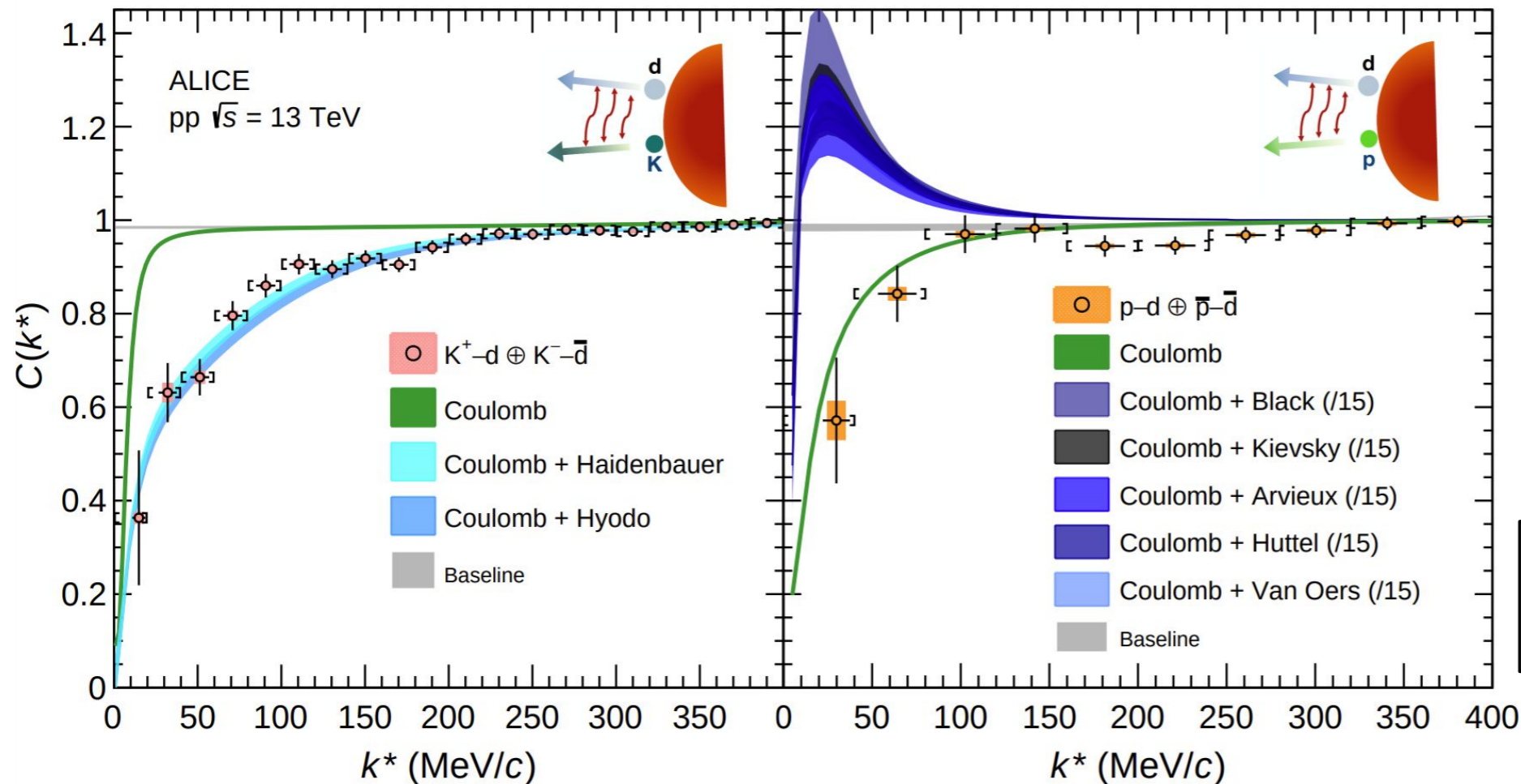


# Correlation with few body systems

## Lednicky model vs ALICE data



ALICE



$$r_{K+d} = 1.41 \pm 0.04 \text{ fm}$$

$$r_{pd} = 1.08 \pm 0.06 \text{ fm}$$

pd model calculations scaled by 1/15

Kd data well reproduced

**$\Rightarrow$  fully formed deuterons present assuming small source**

pd data not described

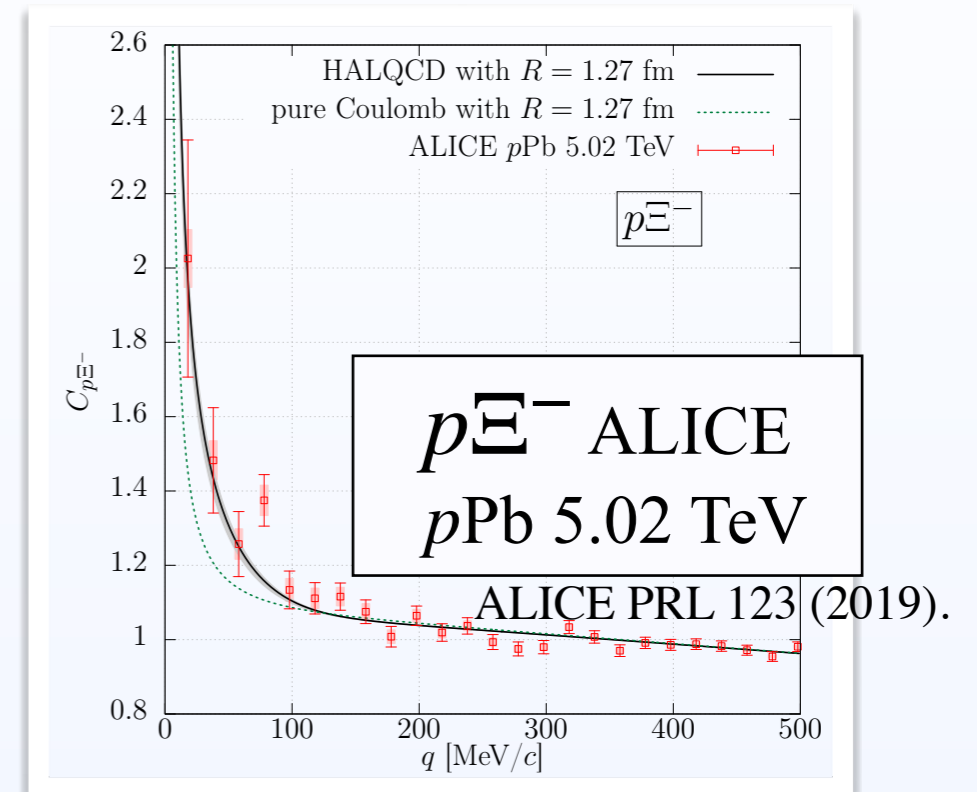
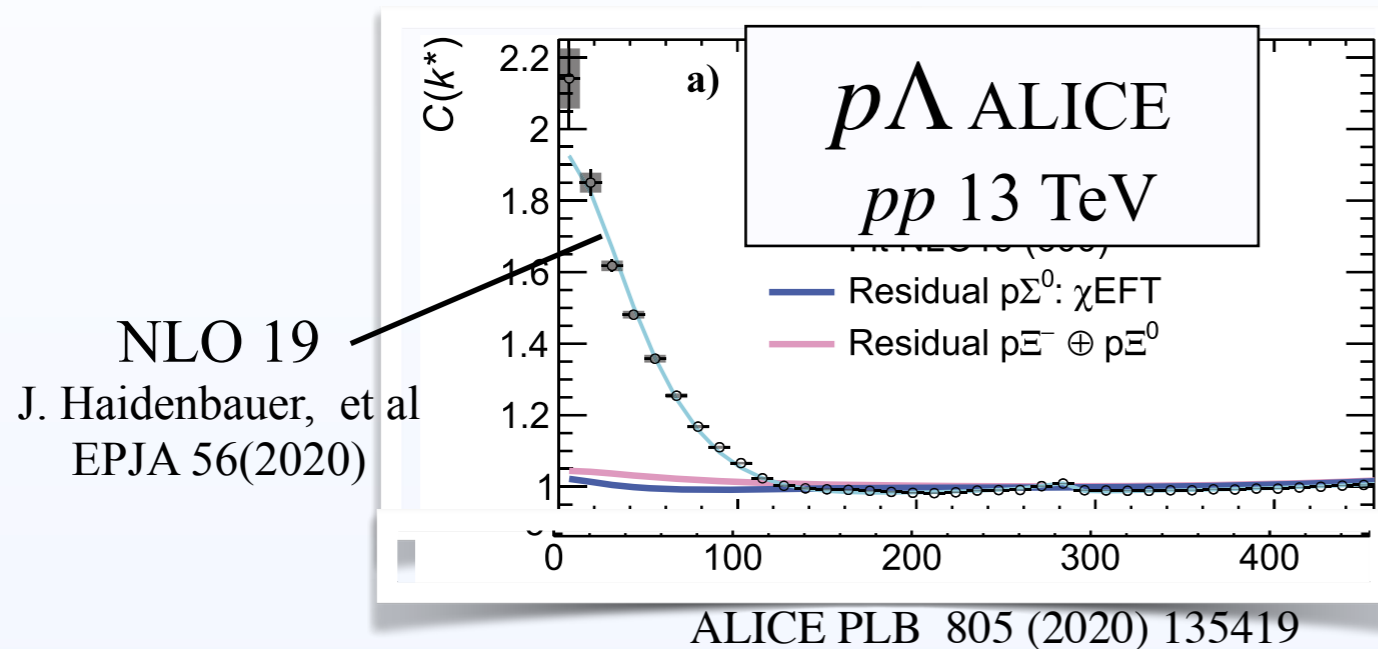
**$\Rightarrow$  pd can't be treated as effective two-body system**

Considering protons, deuterons as distinguishable point-like particles leads to huge discrepancy

**arXiv:2206.03344 (ALICE)**

# $Y$ - $\alpha$ correlation

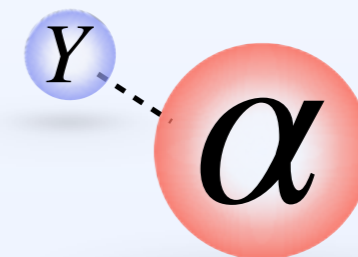
- Good agreement of  $Y$ - $N$  correlation function



Y. Kamiya, et al. PRC 105, 014915 (2022)

## $Y$ - $\alpha(^4\text{He})$ correlation

- Large binding energy of  $\alpha$ 
  - > • Good description by two body treatment
- $Y$ - $\alpha$  potential: smeared potential range
  - > • Detailed potential shape may be investigated



Further constraint on the  $YN(YY)$  int?

# $\Lambda\alpha$ correlation

## • $N\Lambda$ interaction at finite density

- Chiral EFT with NLO D. Gerstung, N. Kaiser, W. Weise, EPJA 55 (2020)  
 —>  $\Lambda NN$  three body interaction gives the additional repulsion A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, PRC 110 (2024), 014001  
 —> stiffer EOS

- **Chi3**: Skyrme type  $\Lambda$  potential based on Chiral EFT with three body  
 ▲ Jinno, K. Murase, Y. Nara, and A. Ohnishi, PRC 108 (2023) 6, 065803

$$U_{\Lambda}^{\text{local}} = a_1^{\Lambda} \rho_N + a_2^{\Lambda} \tau_N - a_3^{\Lambda} \Delta \rho_N + a_4^{\Lambda} \rho_N^{4/3} + a_5^{\Lambda} \rho_N^{5/3}$$

- Well reproduces the binding energy of  $\Lambda$  in hypernuclei

- $N\Lambda$  potential model with different density dependence

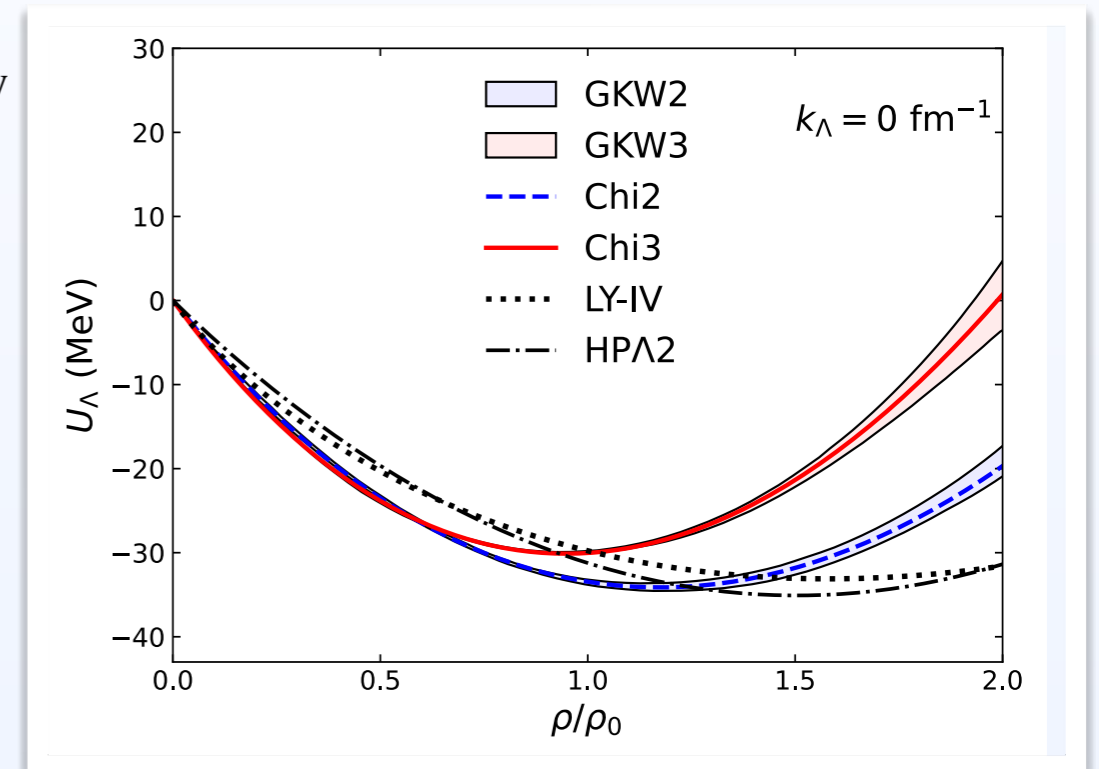
D. E. Lancsok and Y. Yamamoto, PRC 55, 2330 (1997)

N. Guleria, S. K. Dhiman, and R. Shyam, Nucl. Phys. A 886, 71 (2012)

- HPA2

- LY-IV

Weaker density dependence



# $\Lambda\alpha$ correlation

## $\Lambda\alpha$ potential w/ Skyrme type pot.

- Nucleon density with Gaussian form:

$$\rho(r) = A(2\nu_c/\pi)^{3/2} e^{-2\nu_c r^2}$$

- high central density  $\sim 2\rho_0$

- Unknown  $a_3^\Lambda$ : fit to reproduce the  ${}^5_\Lambda\text{He}$  experimental  $E_B = 3.12$  MeV

## Simple potential models

Kumagai-Fuse, S. Okabe, Y. Akaishi, PLB 345 (1995)

- Isle potential

$$V(r) = V_1 e^{-r^2/b_1^2} + V_2 e^{-r^2/b_2^2}$$

repulsive core

attractive part

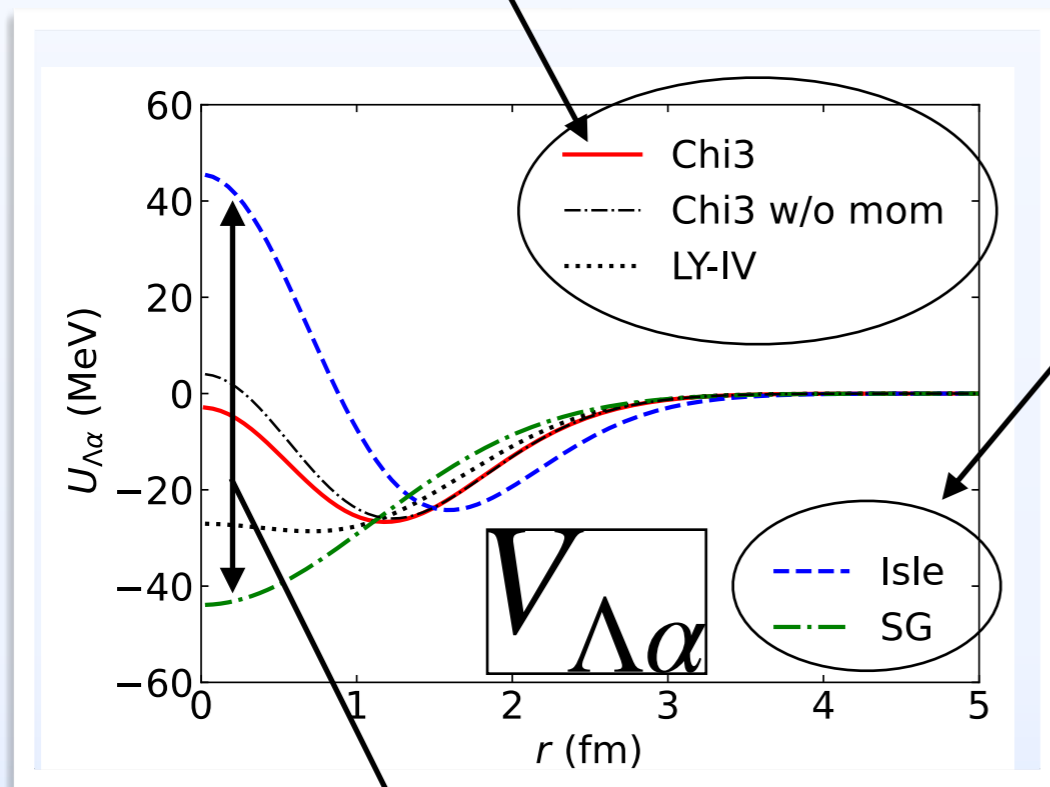
(short range)

(long range)

- Single Gaussian (Isle)

$$V(r) = V e^{-r^2/b^2}$$

- parameters are chosen to reproduce  $E_B$



Large difference in strength of repulsive core

- Difference of  $\rho$  dependence of Skyrme pot. appear in the strength of repulsive core.

- Strength of repulsive core

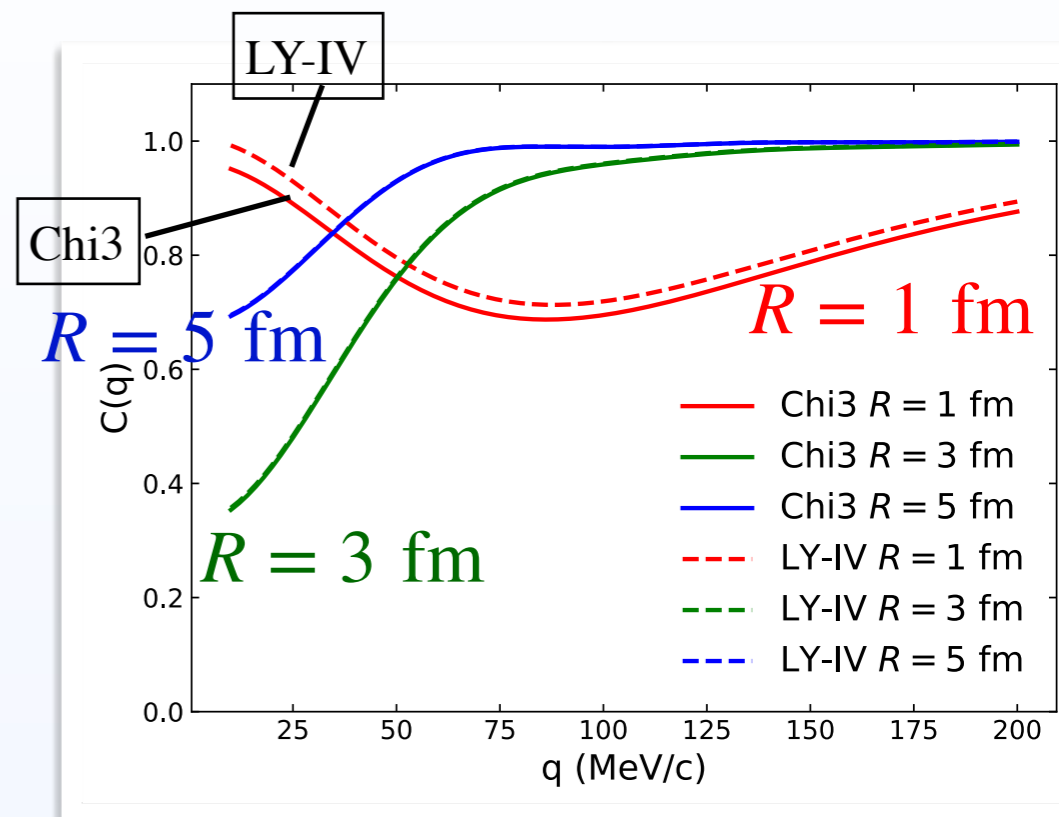
Isle > Chi3 > LY-IV > SG (No core)

- Potential shape dependence of  $C_{\Lambda\alpha}$ ?



# $\Lambda\alpha$ correlation

- Source size dependence of  $C_{\Lambda\alpha}$ 
    - Characteristic lineshapes for weak binding system ( ${}^5_{\Lambda}\text{He}$ )
      - Dip for small source
      - Suppression for large source
    - Potential difference appear only in small source results
- ➔ Large source results are useful to check  $E_B$  of  ${}^5_{\Lambda}\text{He}$



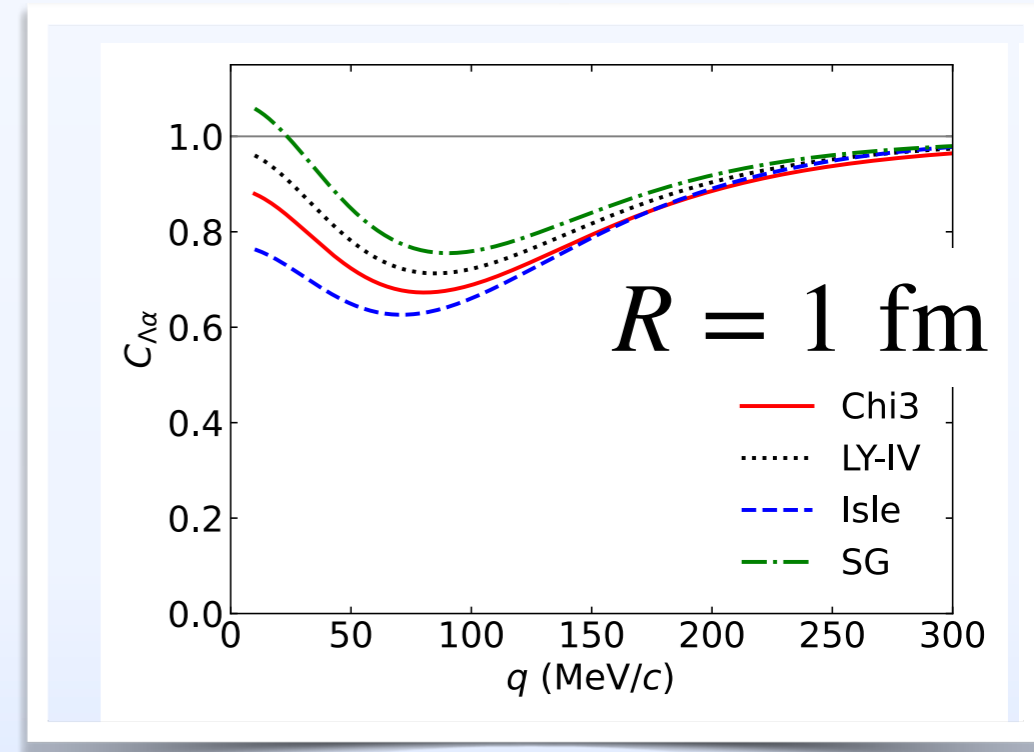
## Effect of repulsive core

- $C(q)$  are ordered from bottom to top as  
Isle -> Chi3 -> LY-IV -> SG (No core)

Same ordering with the strength of repulsive core

—> Stronger core causes Stronger suppression

➔ Strength of the repulsive core can be tested with  $C_{\Lambda\alpha}(q)$  from small source.



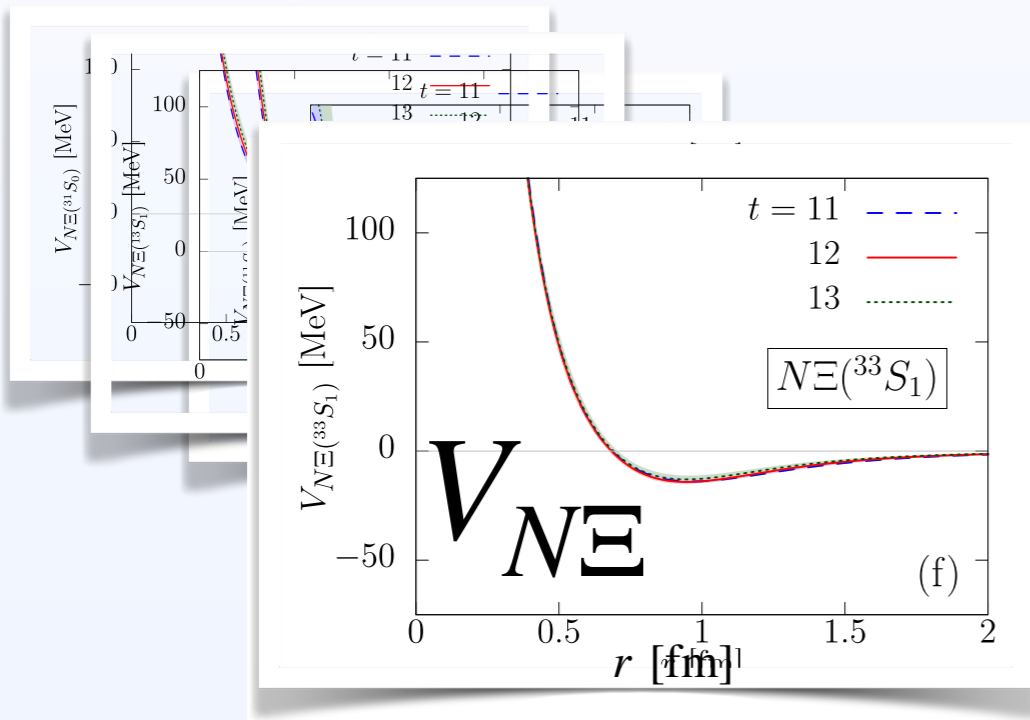
# $\Xi\alpha$ correlation

- $N\Xi$  potential and  $\Xi\alpha$  potential

- HAL QCD  $N\Xi$  potential

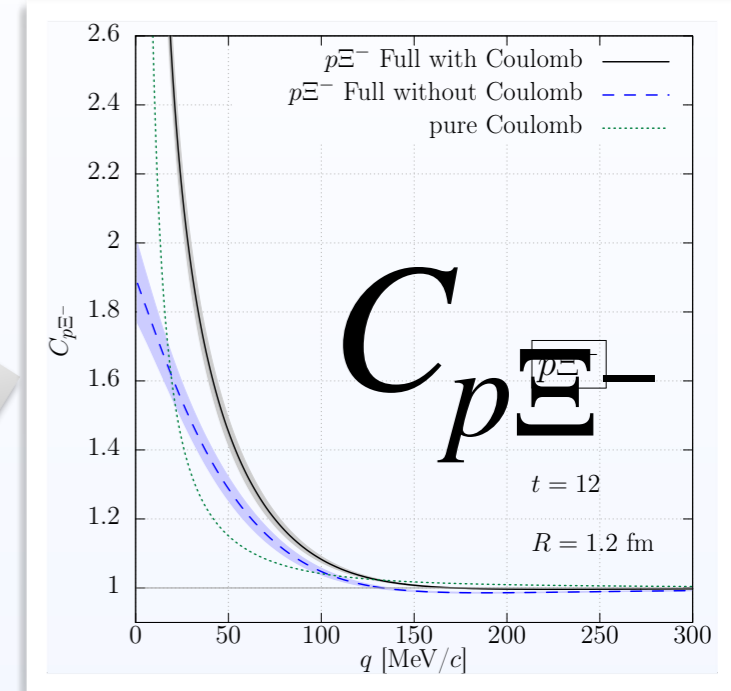
K. Sasaki et al., NPA, 121737 (2019).

4 components for  $s$ -wave:  $^{11}S_0$ ,  $^{13}S_1$ ,  $^{31}S_0$ ,  $^{33}S_1$



- Repulsive core
- Long tail attraction by  $\pi$  exchange

- $p\Xi^-$  correlation  
 $\sim [C(^{11}S_0) + 3C(^{13}S_1) + C(^{31}S_0) + 3C(^{33}S_1)]/8$



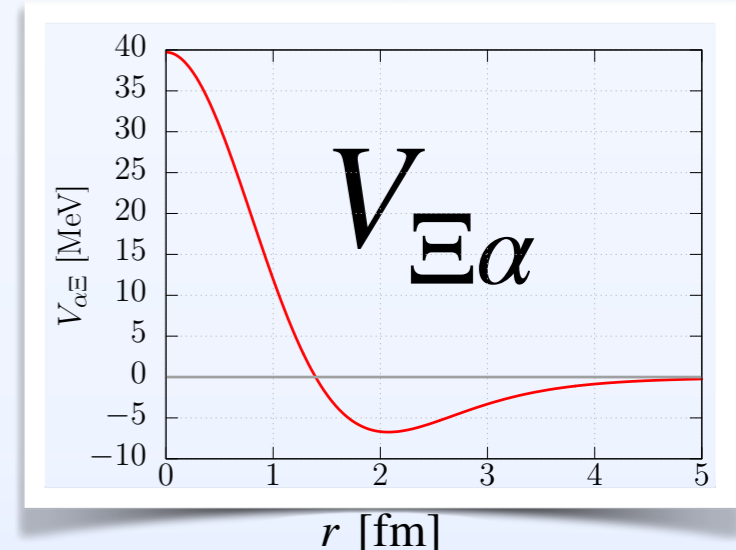
- Large enhancement from  $^{11}S_0$

- Folding  $\Xi\alpha$  potential

E. Hiyama, M. Isaka, T. Doi, and T. Hatsuda, PRC 106, 064318 (2022).

$$\frac{[V(^{11}S_0) + 3V(^{13}S_1) + 3V(^{31}S_0) + 9V(^{33}S_1)]}{16}$$

- Large weight of  $^{11}S_0$



- Different channel weight

- Effect of smeared repulsive core/attraction?

# $\Xi\alpha$ correlation

## Predictions for $\Xi\alpha$ bound state: ${}^5_{\Xi}\text{H}$

- Coulomb assisted bound state  $\leftarrow$  HAL QCD pot.

- Bound state found only for Coulomb attractive pair

E. Hiyama, et al PRC 106, 064318 (2022).

K. Sasaki et al., NPA, 121737 (2019).

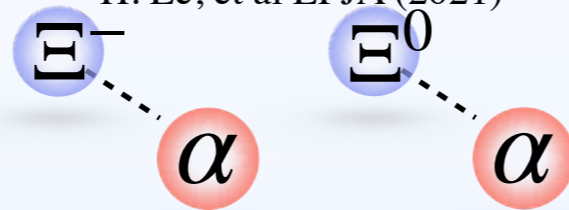
$$E_B = 0.47 \text{ MeV}$$



- Deeper bound state  $\leftarrow$  chiral effective SU(3) pot.

H. Le, et al EPJA (2021)

$$E_B = 2.16 \text{ MeV}$$



Large difference comes from  ${}^{33}S_1$   
H. Le, et al EPJA (2021)

- Behavior for Coulomb assisted bound state?
- Can we distinguish  ${}^5_{\Xi}\text{H}$  with  $C_{\Xi\alpha}$ ?

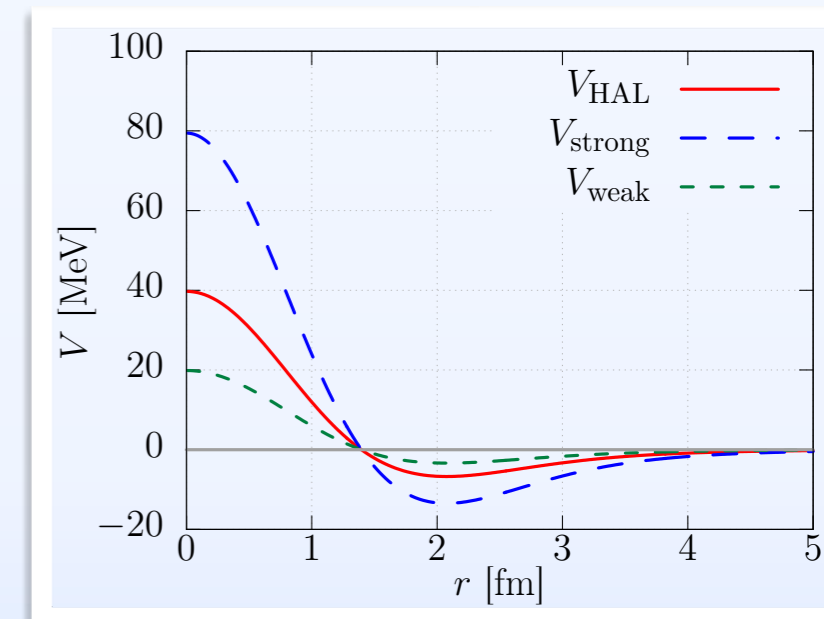
## Folding potential and variations

- $V_{\text{HAL}}$ : Folding potential based on  $S = -2$  HAL QCD potential

E. Hiyama, M. Isaka, T. Doi, and T. Hatsuda, PRC 106, 064318 (2022).

K. Sasaki et al., NPA, 121737 (2019).

potential	$E_B (\Xi^0\alpha)$ [MeV]	$E_B (\Xi^-\alpha)$ [MeV]
$V_{\text{HAL}}$	(Unbound)	0.47
$V_{\text{strong}} = 2 * V_{\alpha\Xi}$	1.15	2.08
$V_{\text{weak}} = V_{\alpha\Xi} / 2$	(Unbound)	0.18



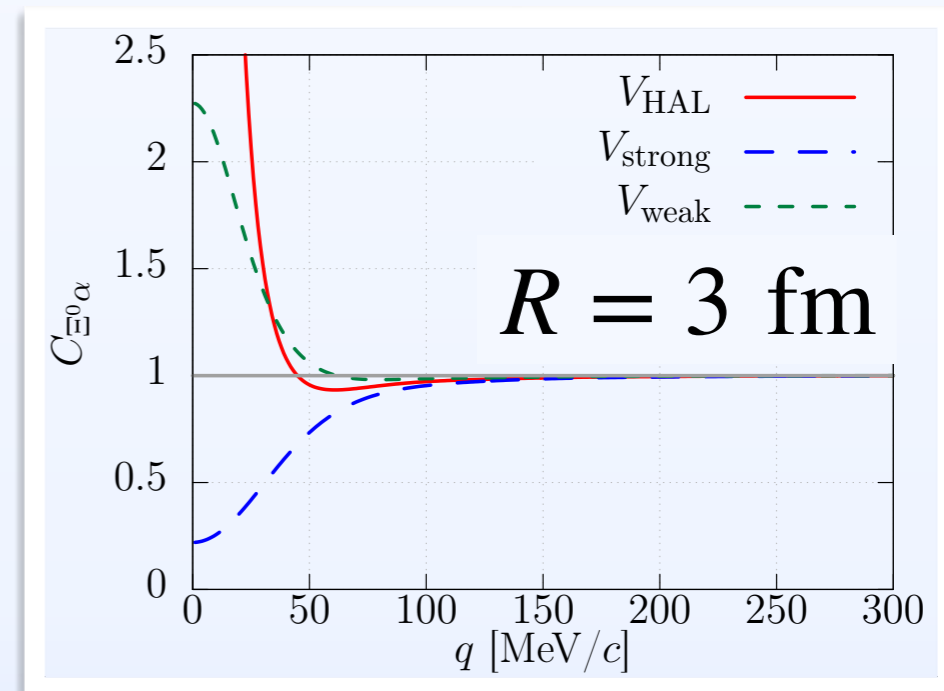
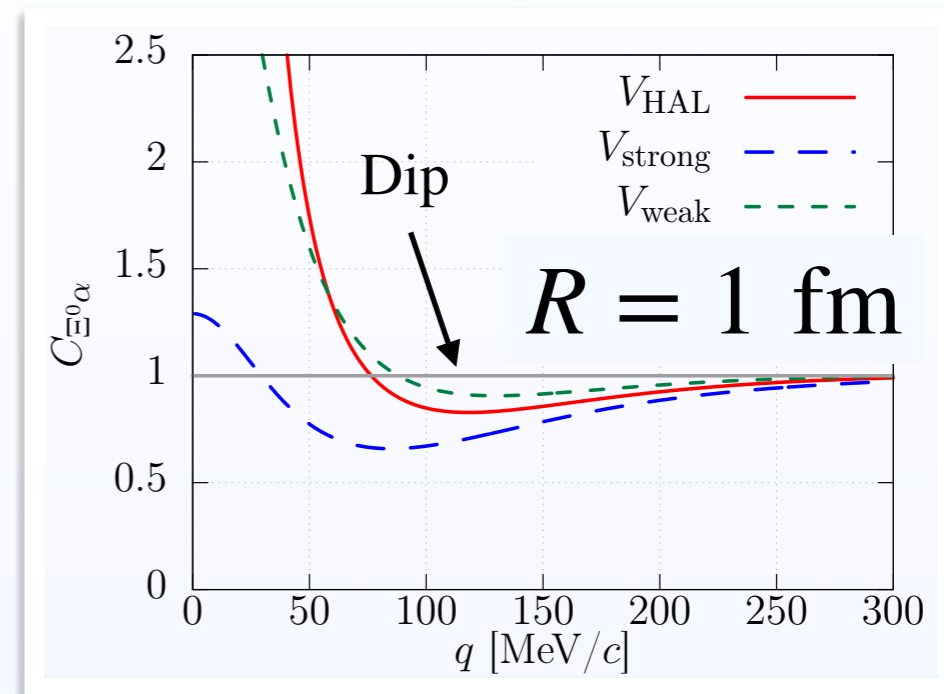
- Check  $V/E_B$  dependence of  $C_{\Xi\alpha}$

# $E\alpha$ correlation

## $E^0\alpha$ correlation

potential	$EB$ [MeV]
$V_{\text{HAL}}$	(Unbound)
$V_{\text{strong}}$	1.15
$V_{\text{weak}}$	(Unbound)

- $V_{\text{strong}}$ : Typical source size dependence with bound state
    - Suppression for large  $R$
    - Enhancement and dip for for small  $R$
  - $V_{\text{HAL}}, V_{\text{weak}}$ : strong enhancement
    - consistent with No  ${}^5_{\text{E}}\text{H}$
  - Dip in  $q \sim 100$  MeV/ $c$  for  $V_{\text{HAL}}$  and  $V_{\text{weak}}$ 
    - Suppression by repulsive core?
- 
- Source size dependence can
  - Effect of detailed potential shape?



# $E\alpha$ correlation

- Detailed potential dependence

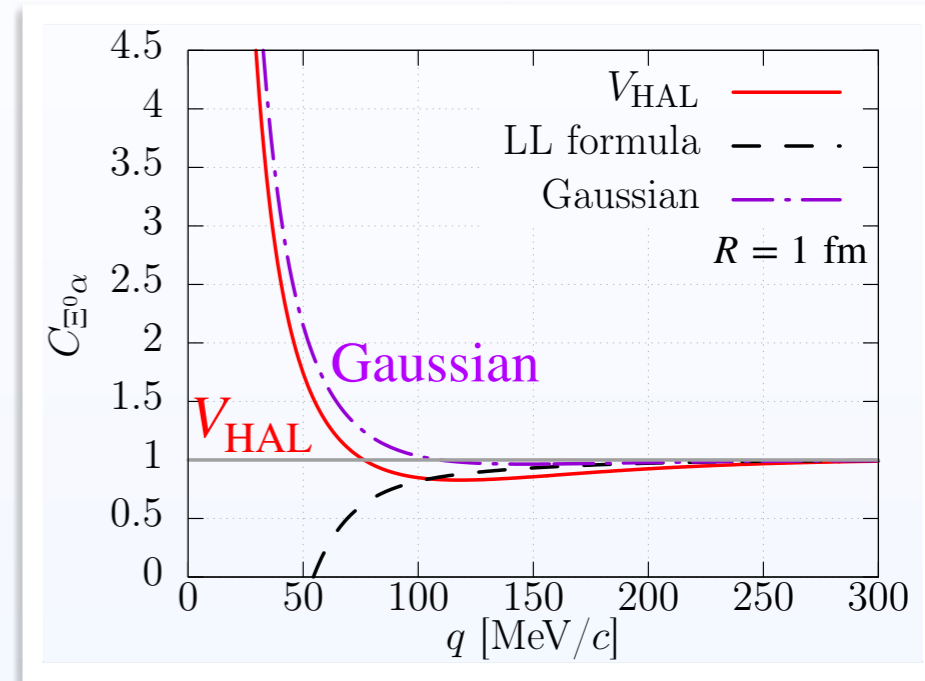
- Compare the **folding potential results** with simpler models
- **Purely attractive Gaussian potential**

$$V_{\text{Gaussian}}(r) = V_0 \exp(-r^2/b^2),$$

- Larger  $C(q)$  than the folding potentials
- No dip structure at  $q \sim 100$  MeV/c



Repulsive core causes dip in  $C_{E\alpha}$ !

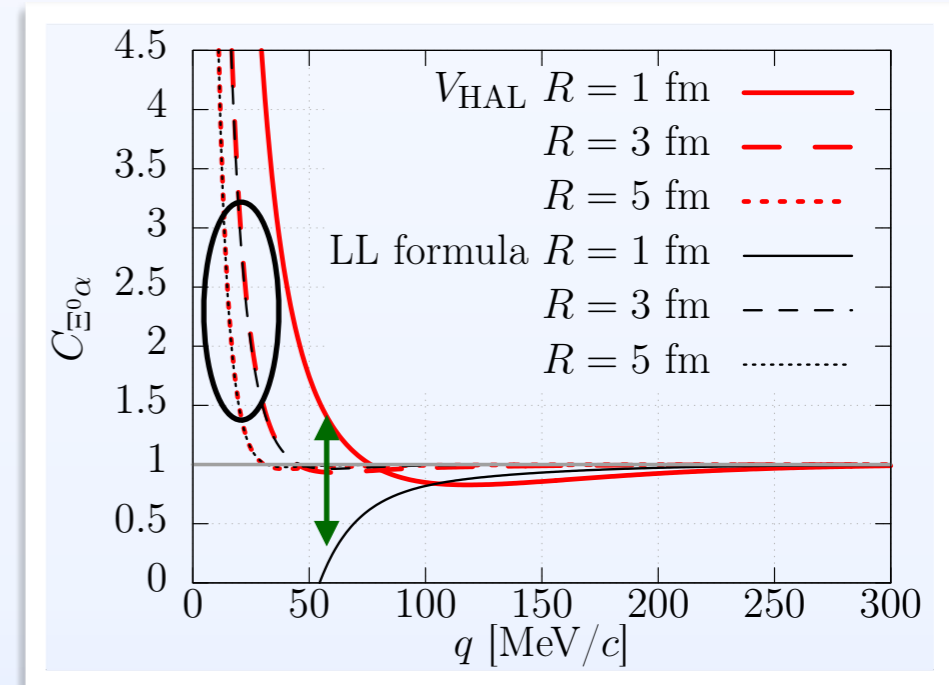


- Lednicky-Lyuboshitz (LL) formula

R. Lednicky, et al. Sov. J. Nucl. Phys. 35(1982).

- approximation by asymptotic wave function  
—> Good description for short range potential
- **Large deviation due to the large effective range for small source**

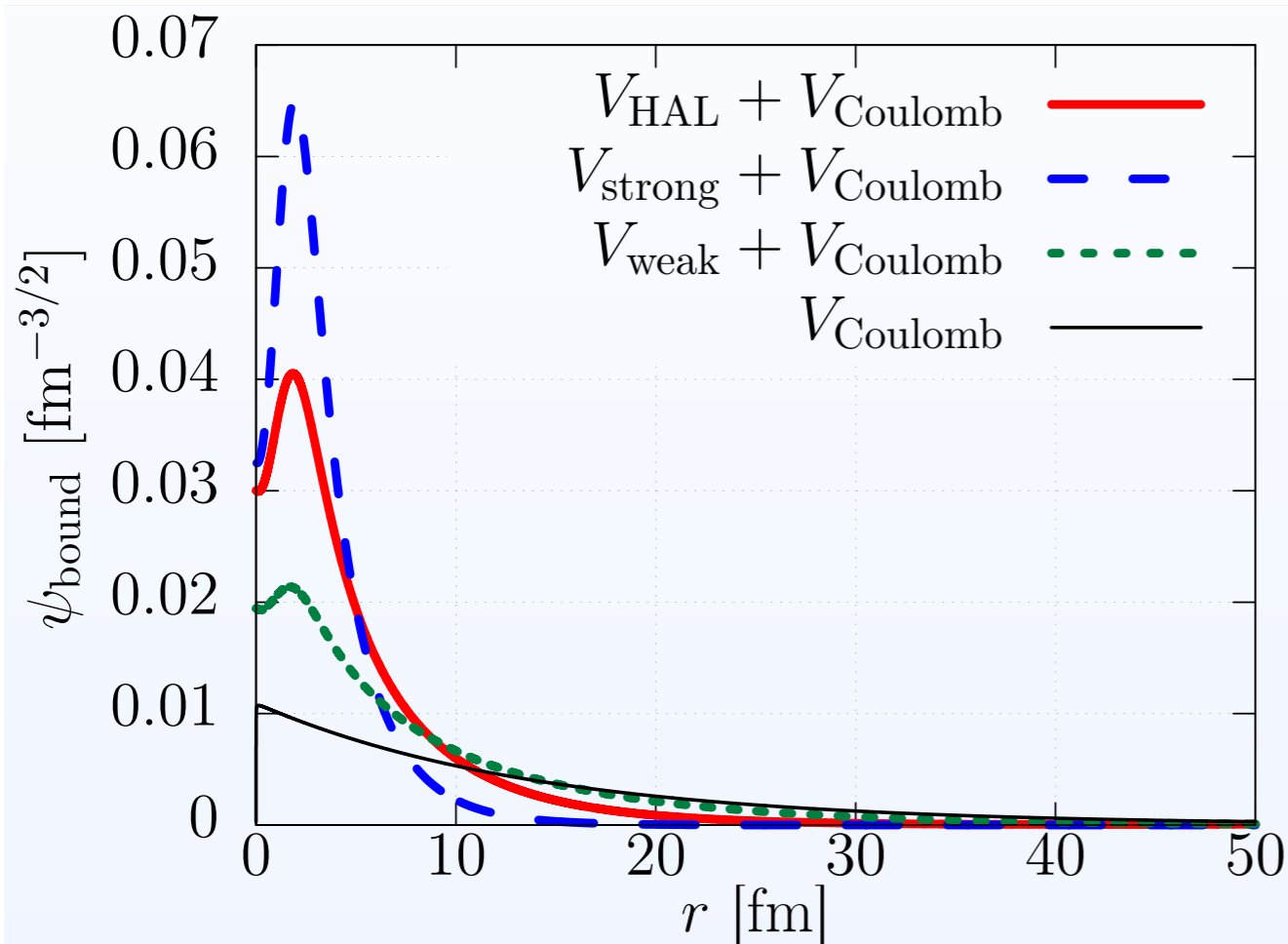
$$r_e = 4.5 \text{ fm} \quad (V_{\text{HAL}})$$



LL formula does not work for  $C(q)$  from small source.

# $\Xi\alpha$ correlation

- $\Xi^- \alpha$  bound state and Coulomb effect



potential	$\langle V_{\text{short}} \rangle$	$\langle V_{\text{Coulomb}} \rangle$	$B$
<b>V<sub>HAL</sub></b>	<b>-0.93</b>	-0.63	0.47
<b>V<sub>strong</sub></b>	<b>-4.36</b>	-0.94	2.08
<b>V<sub>weak</sub></b>	-0.14	<b>-0.36</b>	0.18

[MeV]

- $V_{\text{HAL}}$  and  $V_{\text{strong}}$  : W.f. strongly localized in strong int. range.  
→ Short range int. is dominant.
- $V_{\text{weak}}$  : long range tail similar to pure Coulomb case  
→ Coulomb int. is dominant.

# $E\alpha$ correlation

- $E^- \alpha$  correlation

potential	$EB$ [MeV]
$V_{\text{HAL}}$	0.47
$V_{\text{strong}}$	2.16
$V_{\text{weak}}$	0.18

- Coulomb int. added:  
 —> Strong int. effect appear as deviation from pure Coulomb

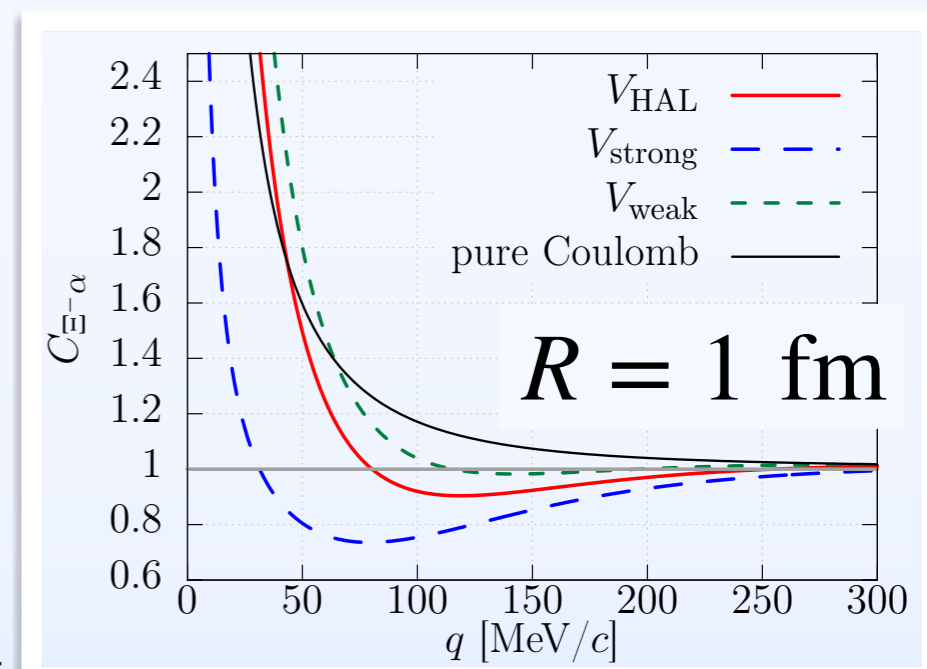
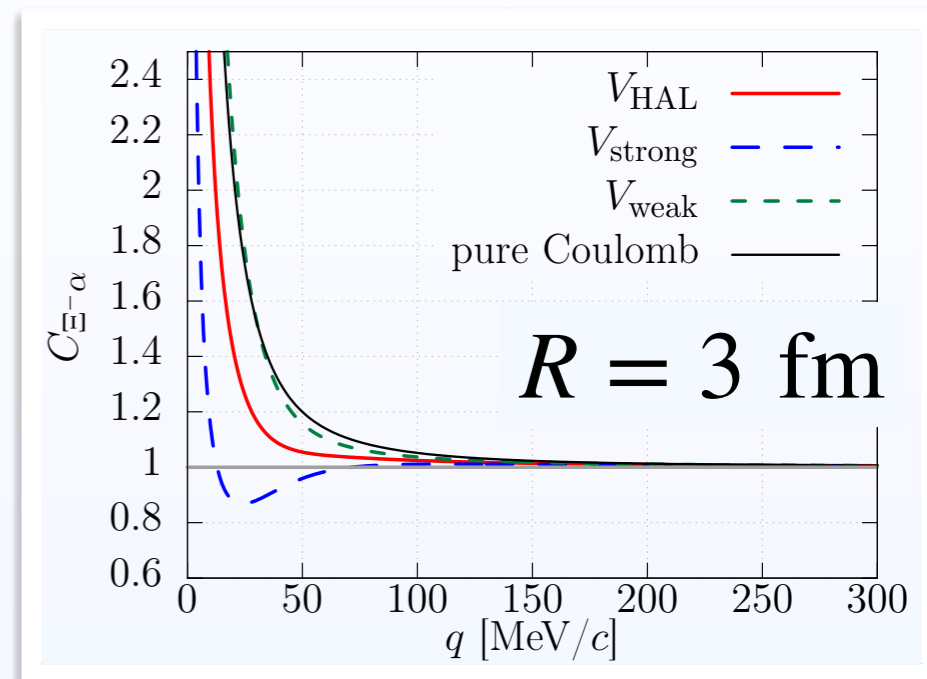
- $V_{\text{strong}}$  and  $V_{\text{weak}}$ : Coulomb enhancement added to  $C_{E^0\alpha}$

- $V_{\text{HAL}}$ :  $C(q)$  with  $R = 3$  fm turns to be suppressed  
 —> Typical source size dependence with bound state

➔  ${}^5_{E^0}\text{H}$  can be distinguished by the source size dependence

- Dip structure at  $q \sim 100$  MeV/c for  $R = 1$  fm

➔ Repulsion core effect can be investigated with small source



# Summary

- Femtoscopic study on the hadron interaction
  - Direct approach to the low-energy interaction
  - Sensitive to the near-threshold resonance
- $K^-p$  correlation
  - Chiral SU(3) model give the good agreement with the various  $K^-p$  data
  - Finite deviation in small source indicates the stronger coupling
- $\Xi\alpha$  correlation function
  - Existence of  ${}^5_{\Xi}H$  can be tested with the source size dependence
  - Dip structure at intermediate momentum by the repulsive core

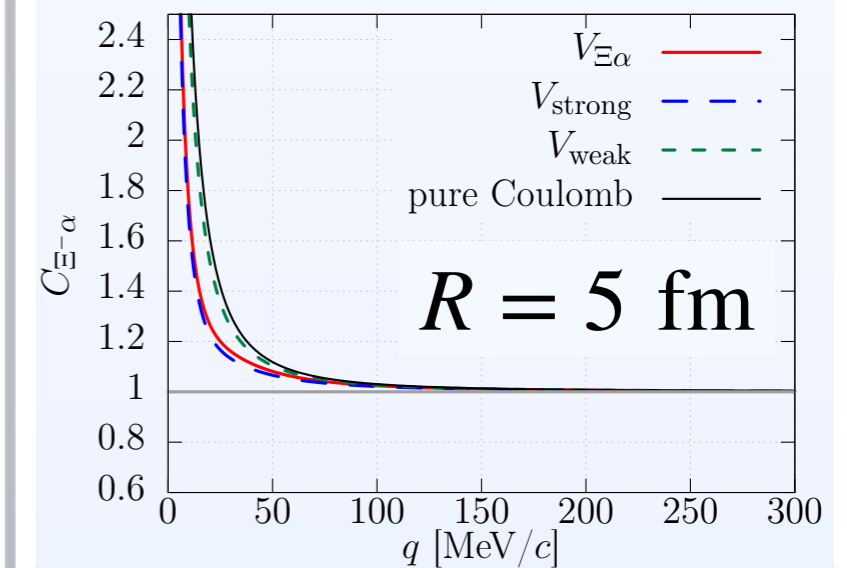
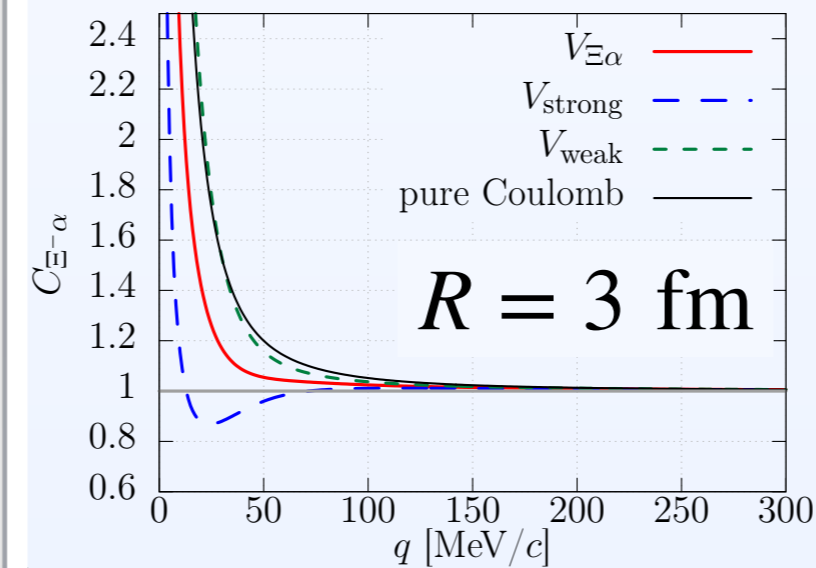
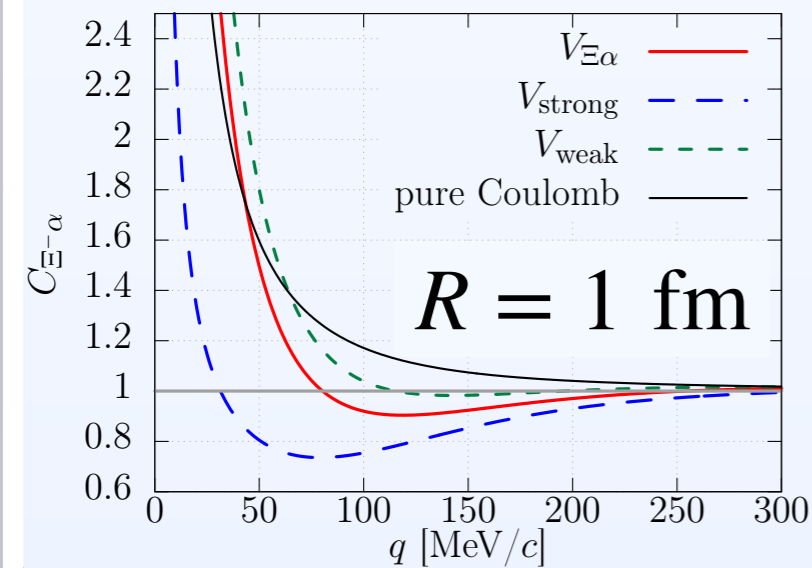
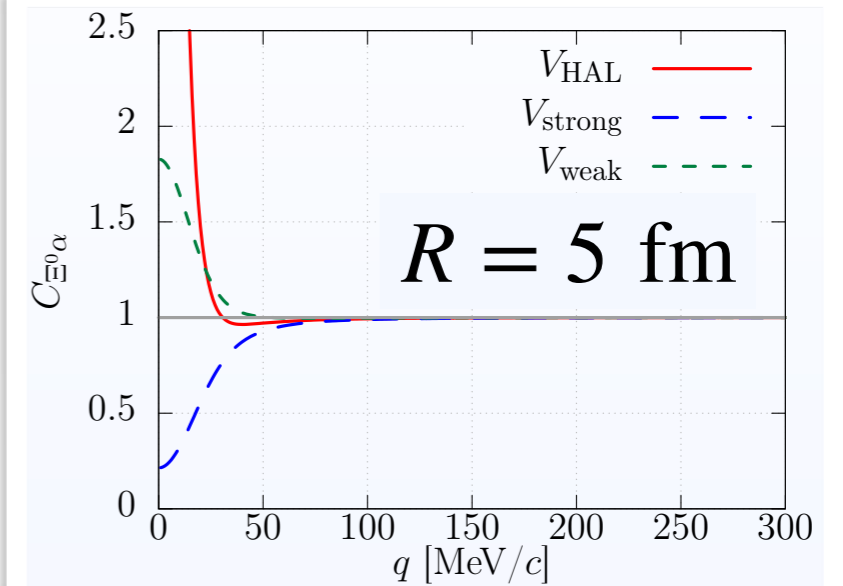
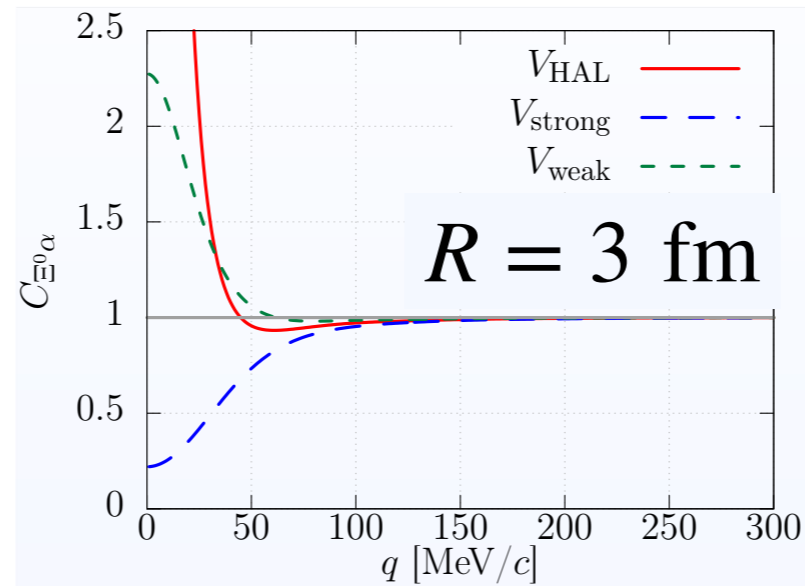
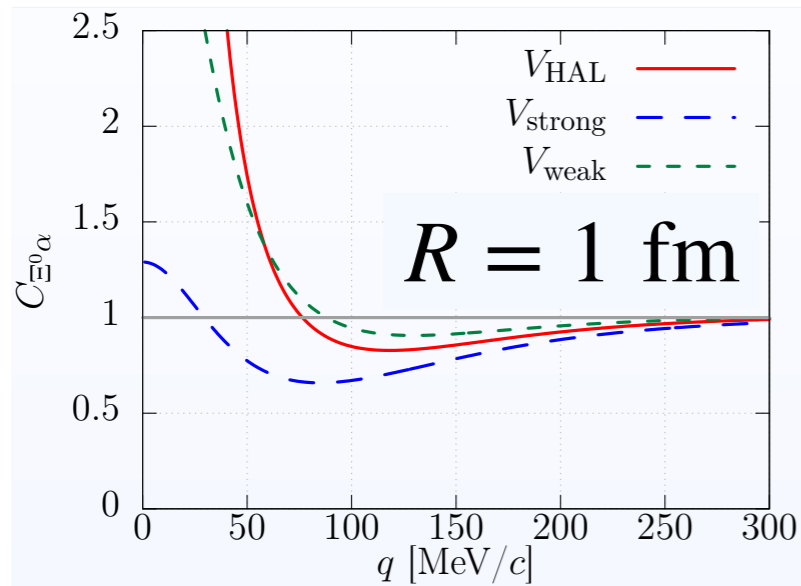
*Thank you for your attention!*

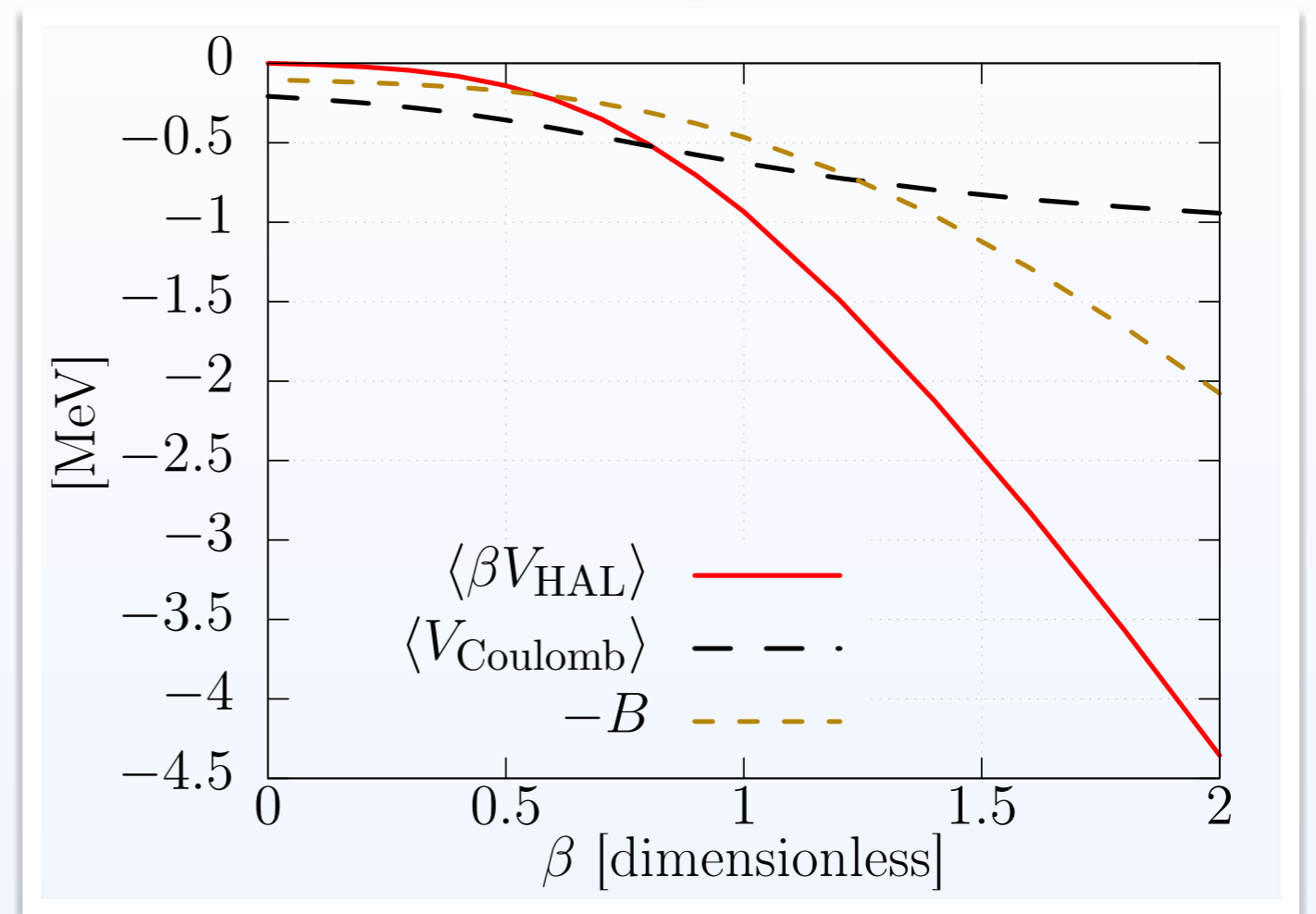
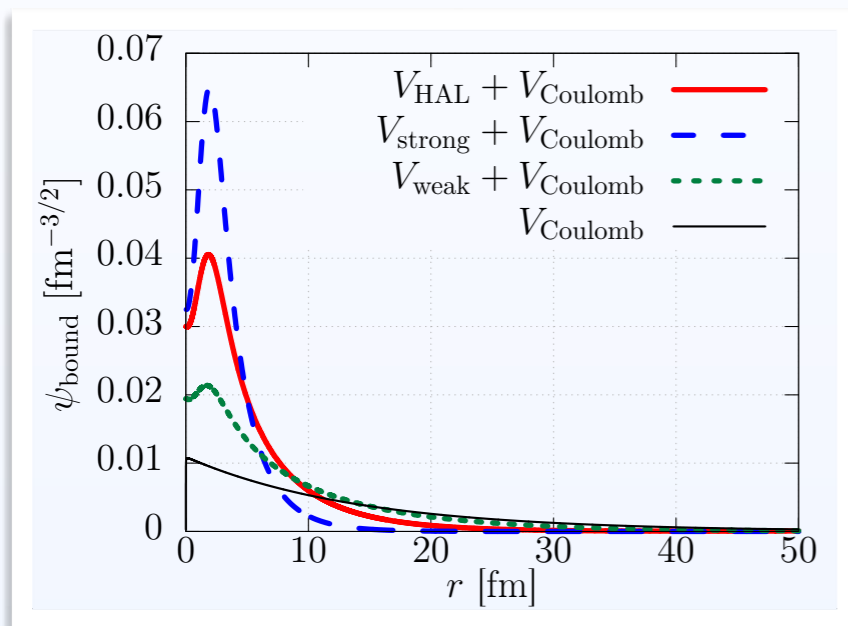


The background features a decorative pattern of swirling lines in shades of purple, blue, and orange, set against a dark purple background. A horizontal band of a lighter purple color runs across the middle of the page, containing the text.

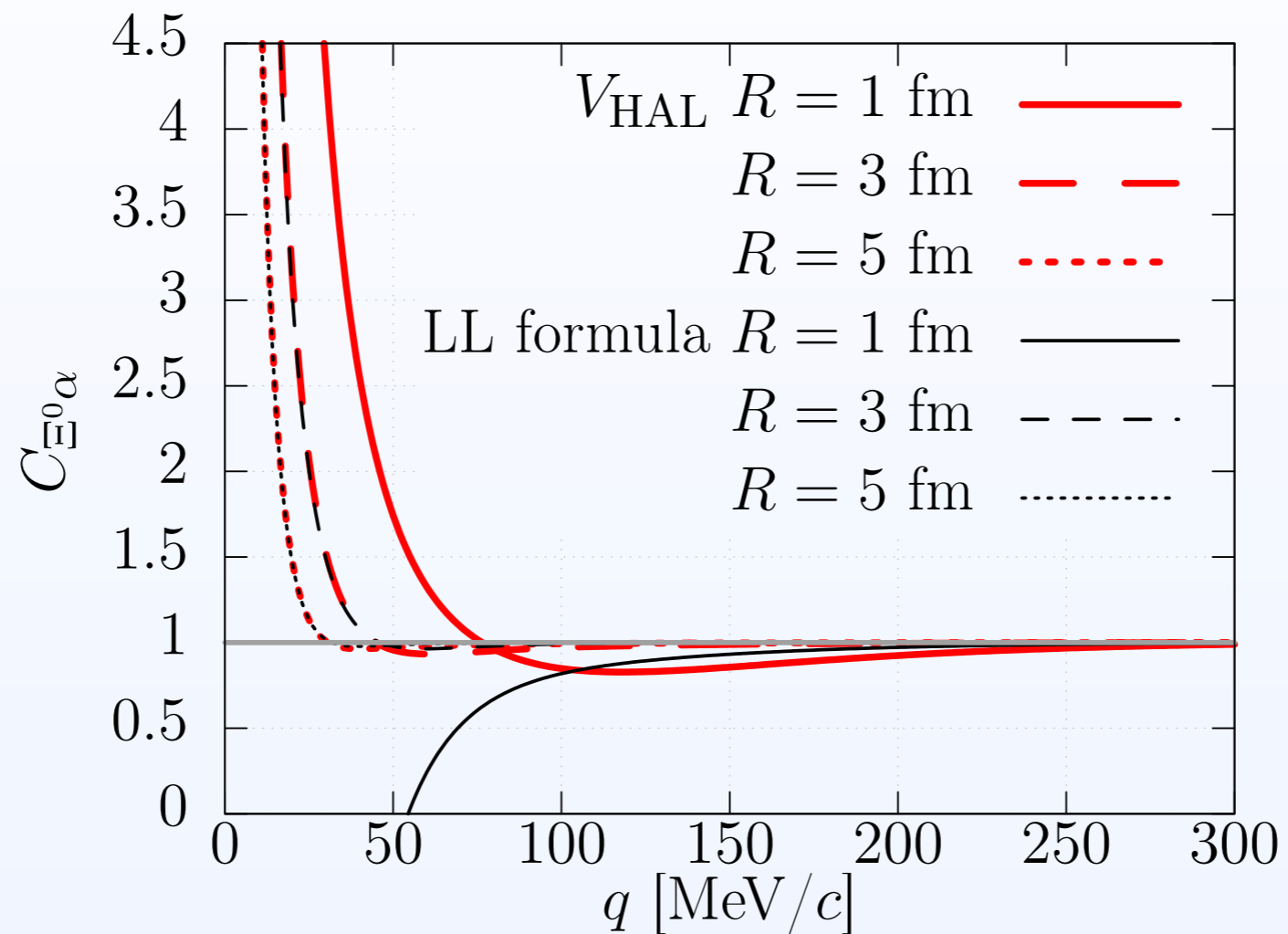
*Thank you!*

# $\Xi\alpha$ correlations





# LL formula for $R = 1, 3, 5$ fm



# Gaussian potential for $V_{\text{HAL}}$

