

Moats, moats, everywhere

Z, Nussinov, M. Ogilvie, L. Pannullo, RDP,
F. Rennecke, S. Schindler & M. Winstel
= NOPPRS W arXiv: 2410.xxxxx

Washington
University in St. Louis



Zohar Nussinov



Mike Ogilvie



JUSTUS-LIEBIG-
UNIVERSITÄT
GIESSEN

Fabian Rennecke

UNIVERSITÄT
BIELEFELD

Laurin Pannullo



GOETHE
UNIVERSITÄT
FRANKFURT AM MAIN

Marc Winstel



BROOKHAVEN
NATIONAL LABORATORY



Rob Pisarski

MIT →
Los Alamos
NATIONAL LABORATORY
EST. 1943



Stella Schindler & Frodo

Critical End Point (CEP):

Asakawa & Yazaki 199

Rajagopal, Shuryak & Stephanov 198, 199
+ ...

$$\vec{\Phi} = (\sigma, \vec{\pi}) \quad m_{\pi} = 0$$

$$\mathcal{L}_{\text{eff}} = \frac{1}{2} (\partial_{\mu} \vec{\Phi})^2 + m^2 \vec{\Phi}^2 + \lambda (\vec{\Phi}^2)^2 + \kappa (\vec{\Phi}^2)^3$$

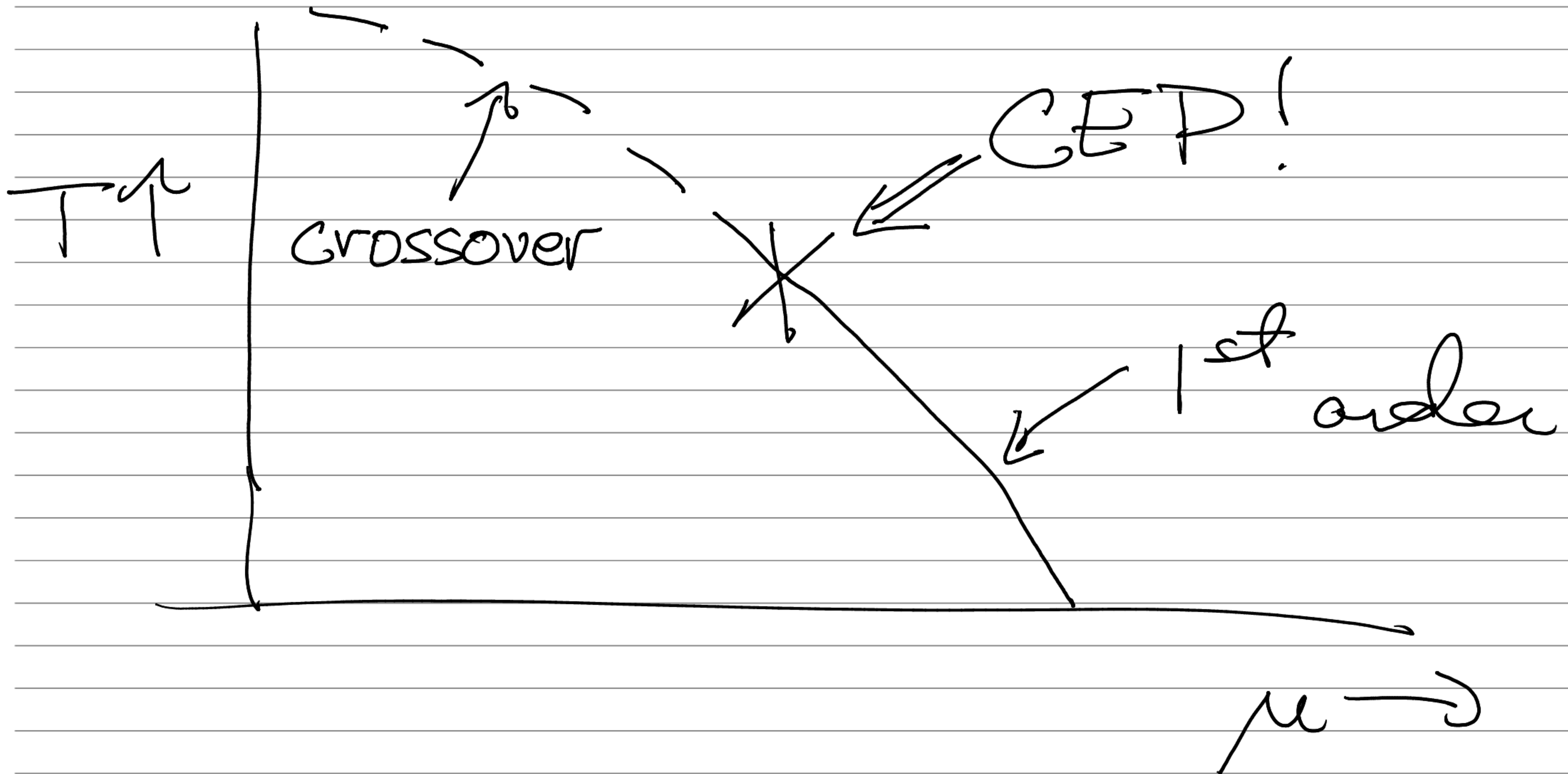
Usually: $\lambda > 0$, $m^2 < 0$ - broken
> - symmetric
= - 2nd order

BUT: $\lambda < 0$ ($\kappa > 0$) ok \Rightarrow 1st order

$\lambda = 0$ = tri-critical point

But $m_H \neq 0$! Still, can have $\lambda < 0$

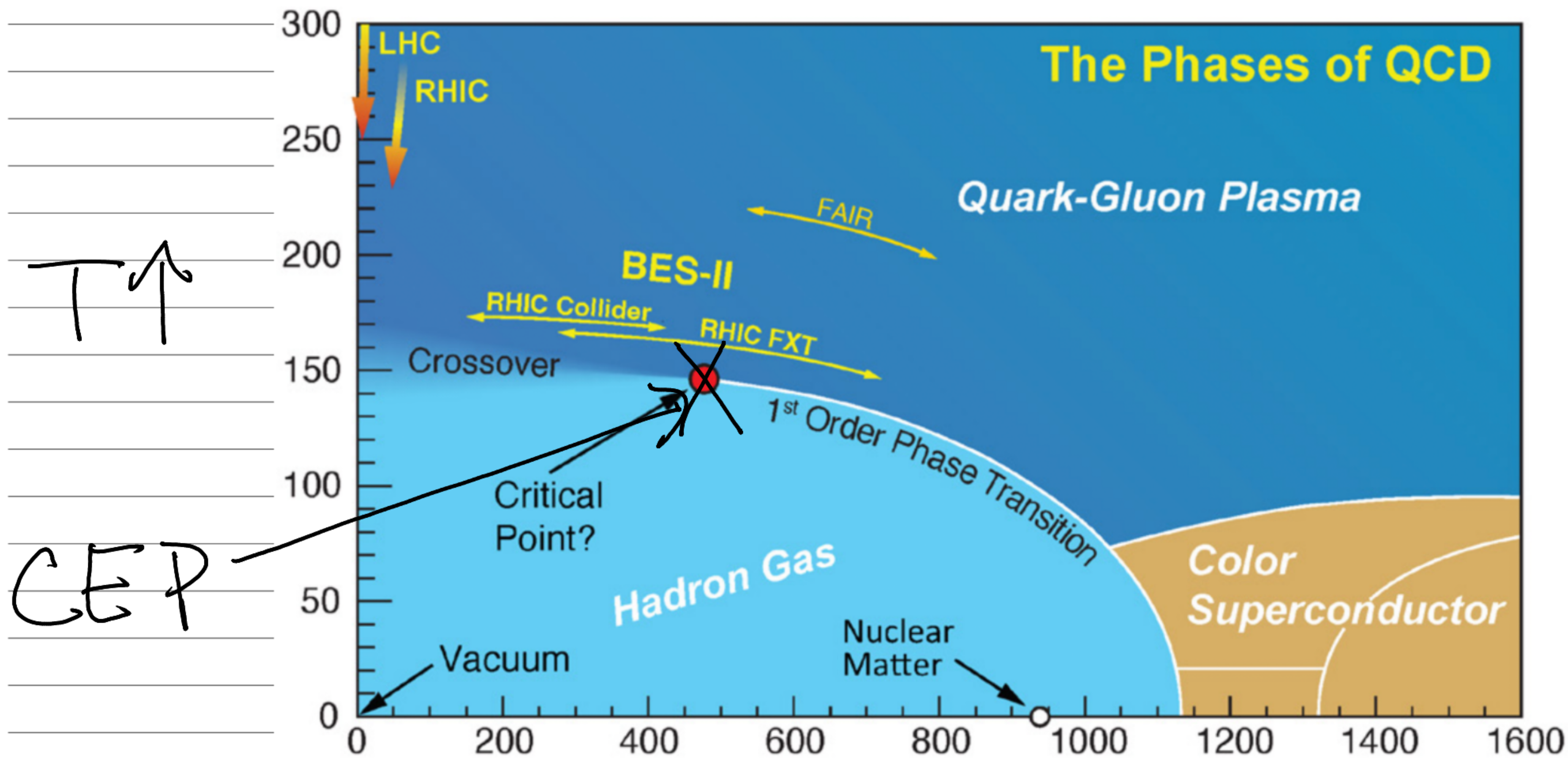
$\lambda = 0 = \text{CEP}$, True 2nd order PT



CEP & QCD phase diagram

Can diagnose CEP from fluctuations

Stephanov '08 + ...



CEP

μ →

CEP: Functional Renormalization Group = FRG

Wei-jie Fu, Pawłowski, Pennocke 1909.02991

$$T_{\text{CEP}} \sim 110 \quad \mu_{\text{CEP}}^{\text{Baryon}} \sim 630 \quad \pm 5\%$$

All later values consistent

eg. Shah + ... 2410, 16206

Basar 2312.06952:

Lee-Yang edge sing. + Padé

No direct evidence from lattice:

compute moments $\frac{\partial^{2n} p(T, \mu)}{\partial \mu^{2n}} \Big|_{\mu=0}$ & extrapolate

But: in vacuum, CEP - width?
 σ meson is heavy & broad

CEP: $m_\sigma = 0$, $m_\pi \neq 0$

Going from vacuum \rightarrow CEP \Rightarrow

$m_\sigma = m_\pi$ only very close to $(T_{\text{CEP}}, \mu_{\text{CEP}})$

Suggests narrow $\Delta T, \Delta \mu$. FRG: $\Delta T, \Delta \mu \approx \pm 2$ MeV!

Affects Fluc's in broader range 1906.00936

Probed by BES @ RHIC, CBM @ FAIR

CEP: "Is that all there is, let's keep ..."

Old story: chiral spirals (spatially inhomogeneous crystal)

Here, new story: moats!

Moats!



Bodiam Castle in East Sussex

By WyrldLight.com, CC BY-SA 3.0, <https://commons.wikimedia.org/w/index.php?curid=7910287>

Moats in physics

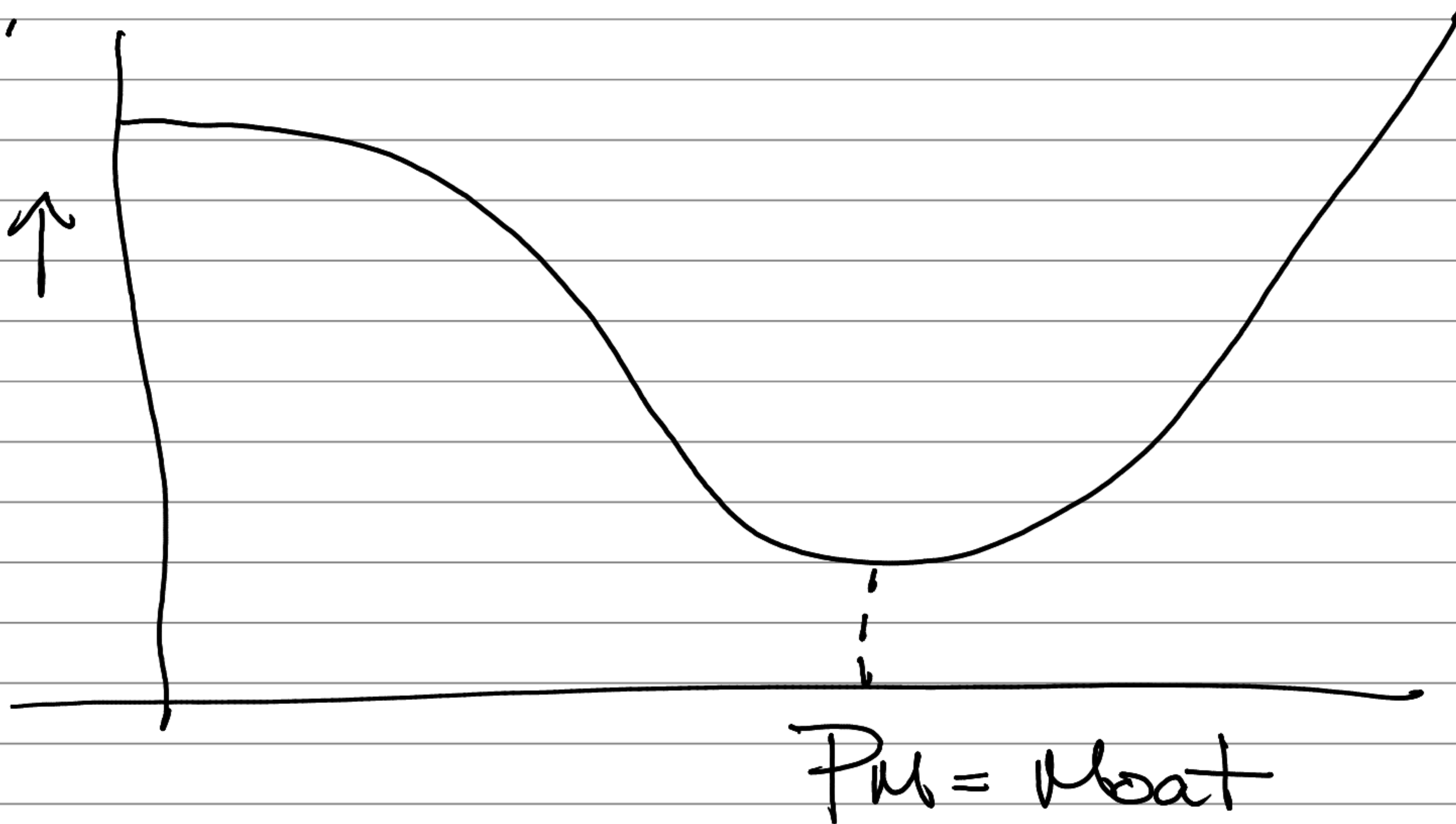
Term from GM: MANY systems.

NEED $T \neq 0$ or $\mu \neq 0$ - lose Lorentz inv.

$$L_{\text{eff}} = (\partial_0 \Phi)^2 + \sum (\partial_i \Phi)^2 + \frac{1}{M^2} (\partial_i^2 \Phi)^2 + V(\Phi)$$

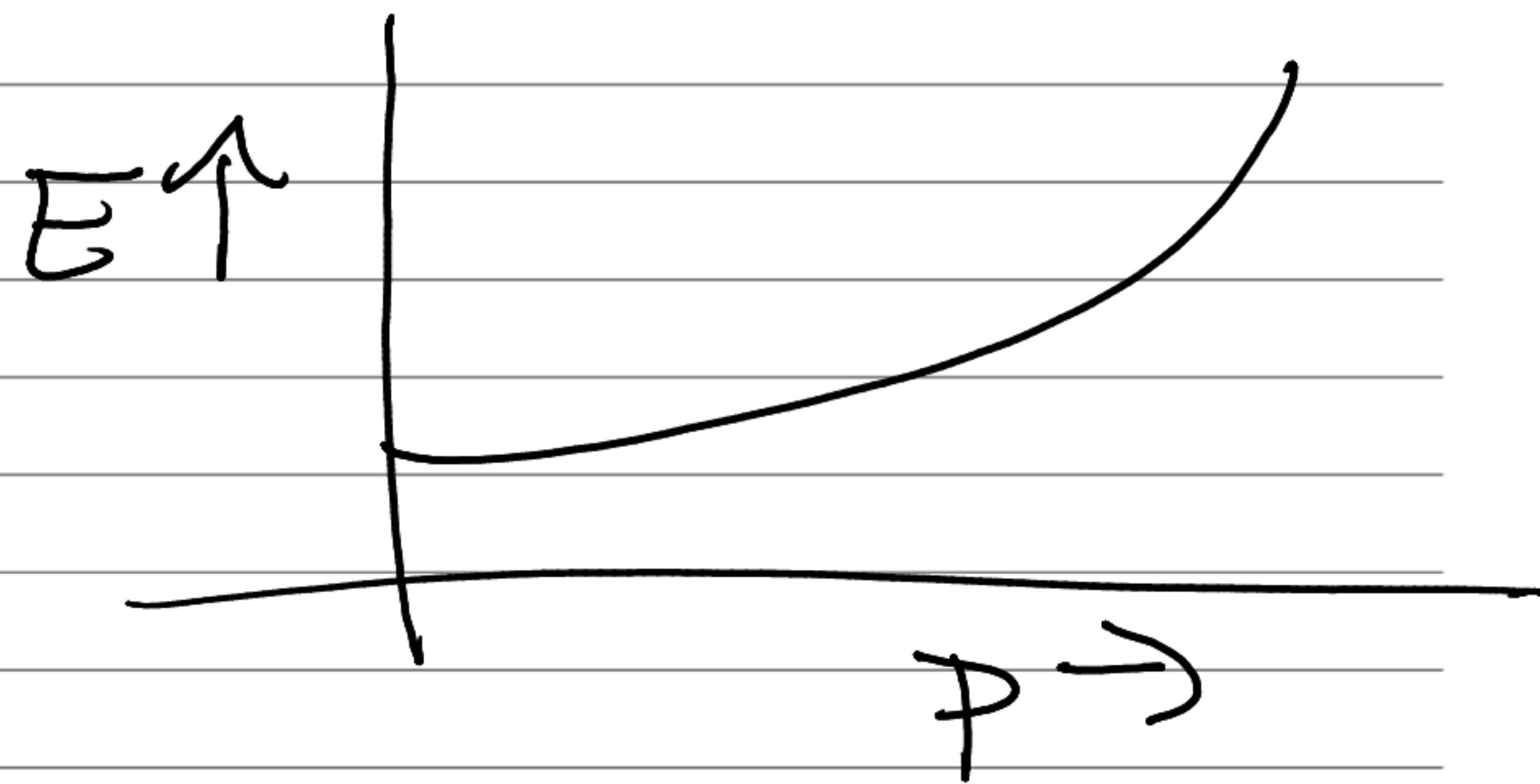
If $\sum > 0$:

$E(\Phi) \uparrow$

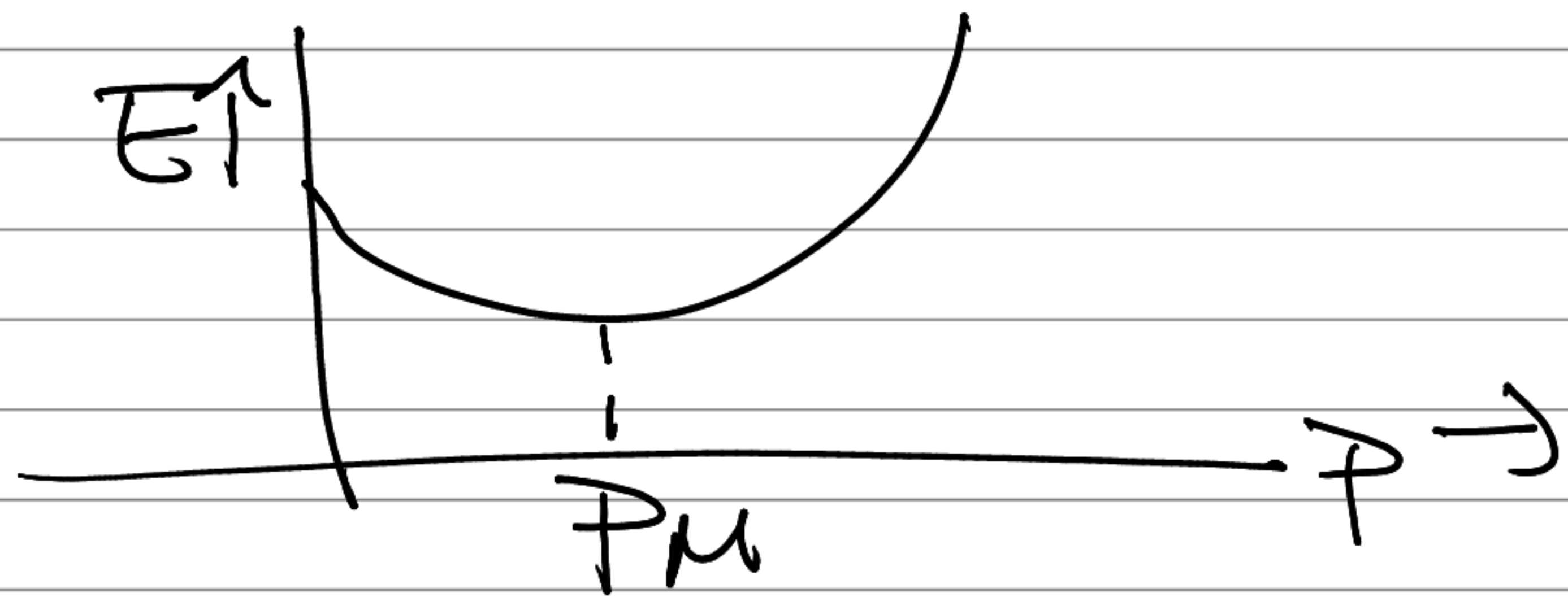


Basic physics

"Normal" phase:
Lightest particle
@ $p=0$

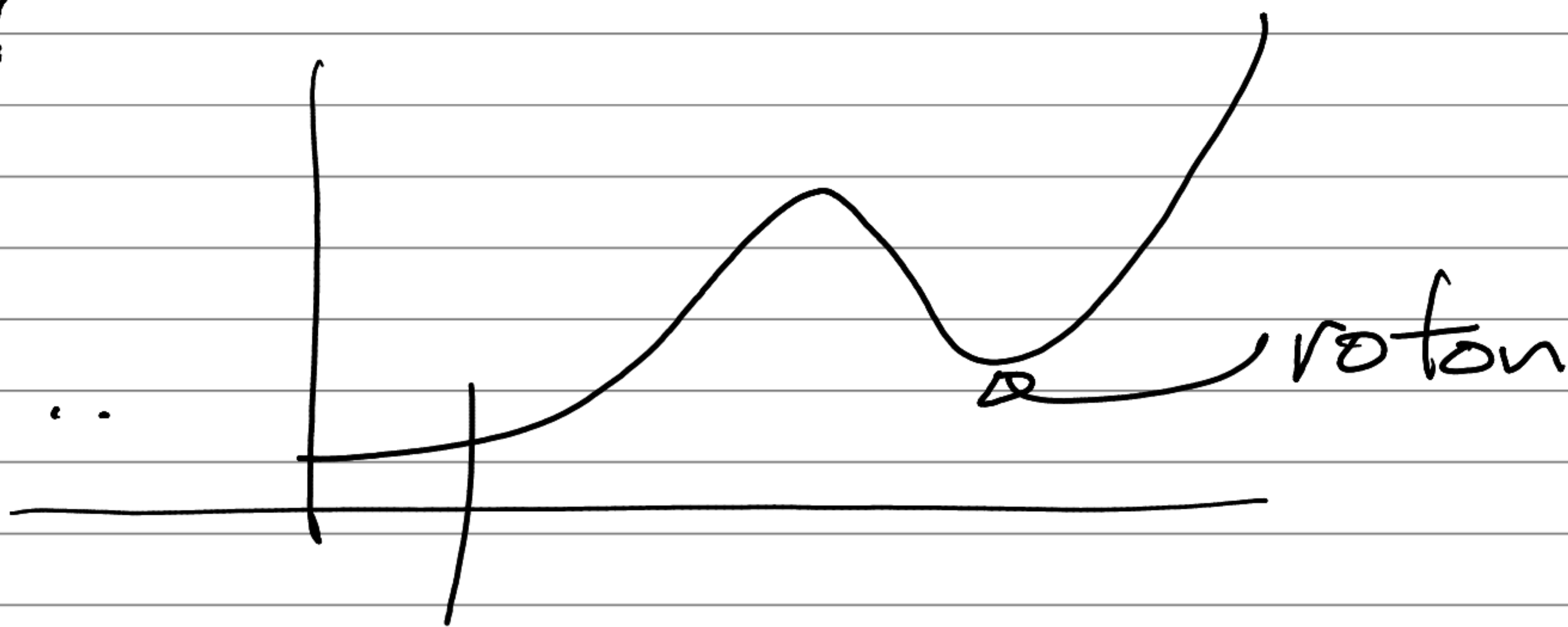


"Moat": Lightest
@ $p_M \neq 0$



Not roton:

Many
possibilities...

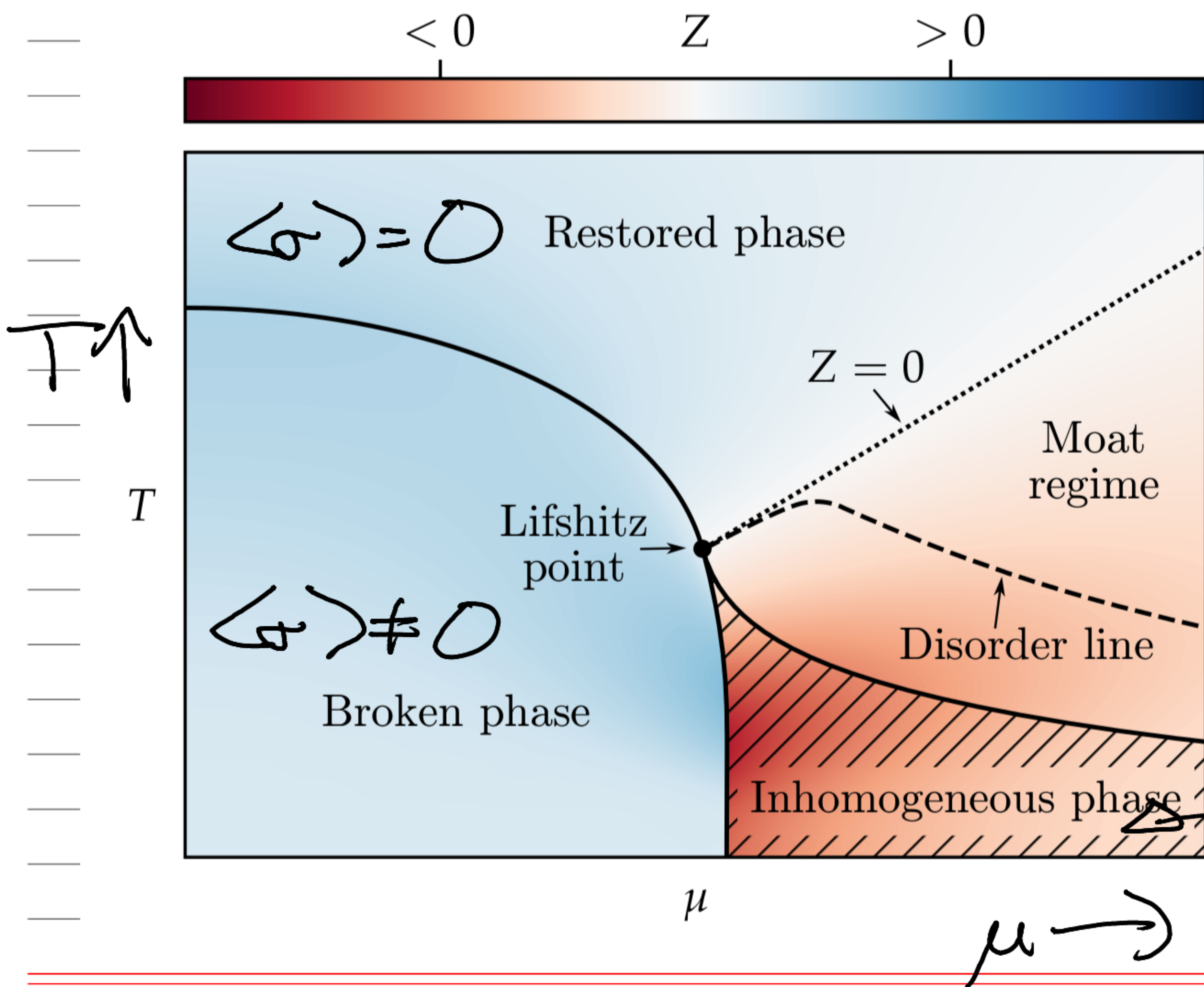


Moats in Gross-Neveu

GN: 1+1 dim, $\mathcal{L}_{GN} = \bar{\psi} i \not{\partial} \psi + g^2 (\bar{\psi} \psi)^2$

Asymp. free, soluble as $N = \# \text{ fermions} \rightarrow \infty$
 Dynamical mass gen. \Rightarrow χ sym. broken

KPRSW 2112.07024



Moat regime ~~*~~
~~*~~ HUGE

1+1 dim's \Rightarrow no CEP,
 Lifschitz point

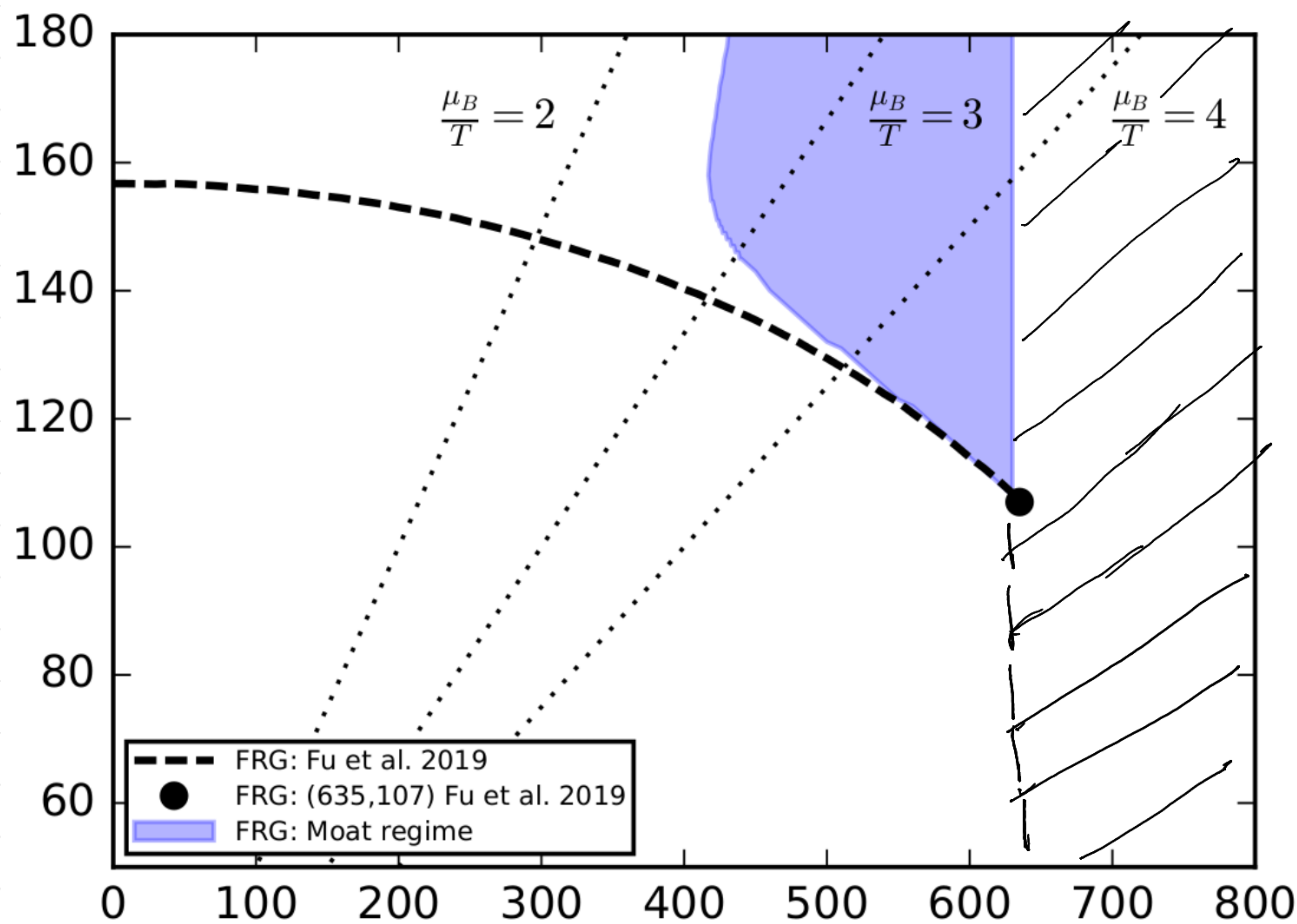
think crystal

Moats in QCD

Wei-jie Fu, Pawlowski, RDP, Pennecke, Rui Wen, Shi Yin

2411.xxxx?

Using FRG: Cannot calculate (yet) for $\mu > \mu_{\text{CEP}}$



But find
~~* HUGE *~~
moat regime
for $\mu < \mu_{\text{CEP}}$
 $T > T_{\text{CEP}}$!
Surely extends
to $\mu > \mu_{\text{CEP}}$
 $T \gtrsim T_{\text{CEP}}$

Dileptons in a moat

NOPPRSW: original $\mathcal{L} =$

$$\mathcal{L}_{\text{eff}} = (\partial_0 \phi)^2 + \sum (\partial_i \phi)^2 + \frac{1}{M^2} (\partial_i^2 \phi)^2 + V(\phi)$$

Gauge invariance $\Rightarrow \partial_i \rightarrow D_i$

$$\mathcal{L}_{\text{eff}}^0 = (D_0 \phi)^2 + \sum (D_i \phi)^2 + \frac{1}{M^2} (D_i^2 \phi)^2 + V$$

Added one new operator = dimension 6

Eight other operators = dim. 6

8 new operators

No Lorentz inv. ($T, \mu \neq 0$). $D_0 \neq D_i$

Six boring operators dim. 6:

$$(D_0 \bar{\phi})^2 |\bar{\phi}|^2, (D_i \bar{\phi})^2 |\bar{\phi}|^2$$

$$(\bar{\phi}, D_0 \bar{\phi})^2, (\bar{\phi}, D_i \bar{\phi})^2,$$

$$F_{0i}^2 |\phi|^2, F_{ij}^2 |\phi|^2$$

Two novel:

$$\text{ie } F_{0i} (D_0 \phi)^* D_i \phi - \text{c.c.}, \text{ ie } F_{ij} (D_i \phi)^* D_j \phi$$

$C =$ charge conjugation: $F_{\mu\nu} \rightarrow -F_{\mu\nu}, \phi \leftrightarrow \phi^*$

\Rightarrow novel terms C-even.

Dilepton rate

Rate $\sim \left| \begin{array}{c} m \\ \bar{m} \end{array} \right\rangle \left[\text{diagram: } \begin{array}{c} \text{---} \\ \diagup \quad \diagdown \\ \text{---} \end{array} \right] \left| \begin{array}{c} l^+ \\ l^- \end{array} \right\rangle^2$

$\sim \propto \text{Im} \Pi^{\mu\mu}(\omega, \vec{P}) \quad \xrightarrow{\text{PM of } l^+ l^-}$

Take $\vec{P} = 0 \Rightarrow$ track to back $l^+ l^-$ in rest frame

$\text{Im} \Pi \sim \text{Im} \sim \left[\text{diagram: } \begin{array}{c} \text{---} \\ \text{---} \end{array} \right] \left[\text{diagram: } \begin{array}{c} \text{---} \\ \text{---} \end{array} \right] \quad \xrightarrow{\mathcal{L}_{\text{eff}}^0}$

+ $\text{Im} \left[\text{diagram: } \begin{array}{c} \text{---} \\ \text{---} \end{array} \right]$

6 boring op's

= tadpoles = 0

+ $\text{Im} \left[\text{diagram: } \begin{array}{c} \text{---} \\ \text{---} \end{array} \right]$

2 novel op's = 0

@ $\vec{P} = 0$ (not obvious)

Back to back $Q^+ Q^-$

$$L_{\text{eff}}: \Delta^{-1} = p_0^2 + \frac{(\vec{p}^2)^2}{\mu^2} + 2\vec{p}^2 + m^2$$

$= \Gamma^j(P, k)$

$$= ie (2k+p)^j \left(2 + \frac{1}{\mu^2} \underbrace{(E^2 + (E + \vec{p})^2)}_{|D \not{p}|^2} \right)$$

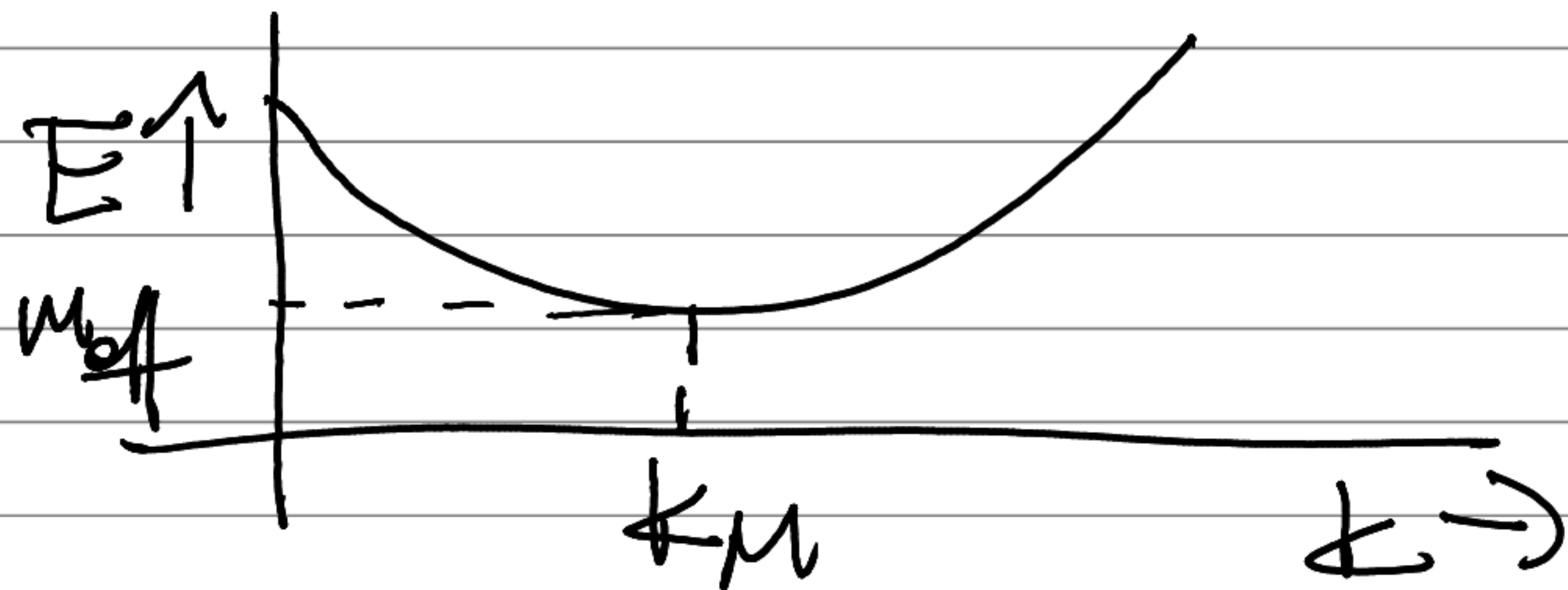
Satisfies Ward id.:

$$P^\mu \Gamma^\mu = ie (\Delta^{-1}(k+p) - \Delta^{-1}(k))$$

$$\Rightarrow P^\mu \Pi^{\mu\nu} = 0 \quad \text{which is good.}$$

Bottom of the Moat

If $z < 0$



$$E^2 = \frac{(k^2)^2}{M^2} + z k^2 + m^2 = \frac{1}{M^2} (k^2 - k_\mu^2)^2 + m_{\text{eff}}^2$$

About the bottom of the moat,

$$k \approx k_\mu + \delta k$$

$$E \approx m_{\text{eff}} + \frac{2k_\mu^2}{m_{\text{eff}} M} (\hat{k} \cdot \delta k)^2 + \dots$$

= quadratic in $(\delta k)^2$. Of course, , ,

van Hove singularity?

Near threshold, $\omega = 2m_{\text{eff}} + \delta\omega$

$$\delta(\omega - 2m_{\text{eff}}) \sim \delta(\delta\omega + \dots (\delta k)^2)$$

$$\frac{1}{\omega - 2m_{\text{eff}}} \delta(\delta k)$$

= (van Hove) sing.! BUT:

$$\Gamma_{ij}(P, k) \sim 2ie k^i \left(\frac{4k_M}{M^2} \mathbf{k} \cdot \delta \mathbf{k} \right)$$

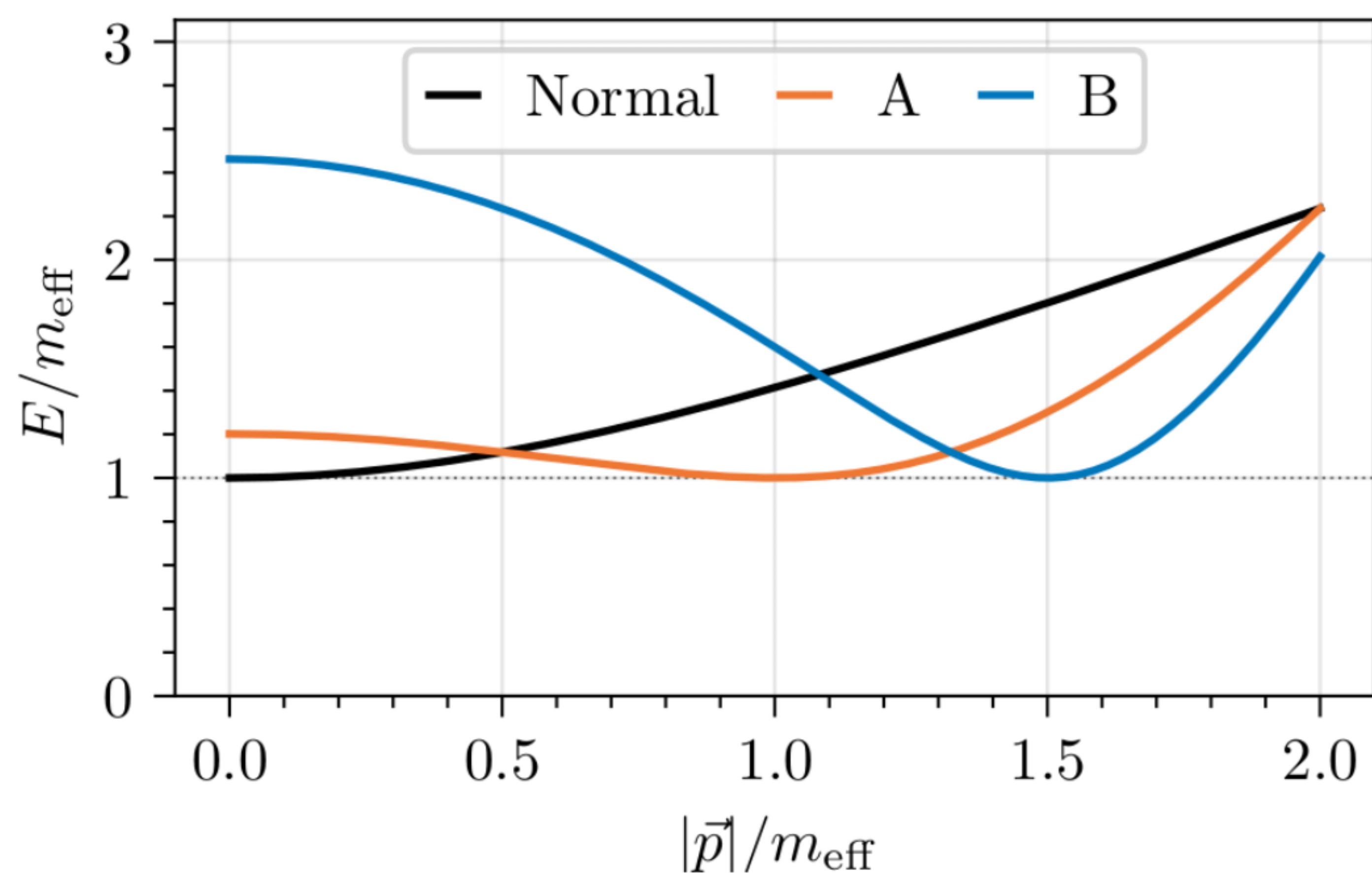
2 vertices \Rightarrow NO sing.

Moaton-Moaton production is enhanced
relative to normal π 's!

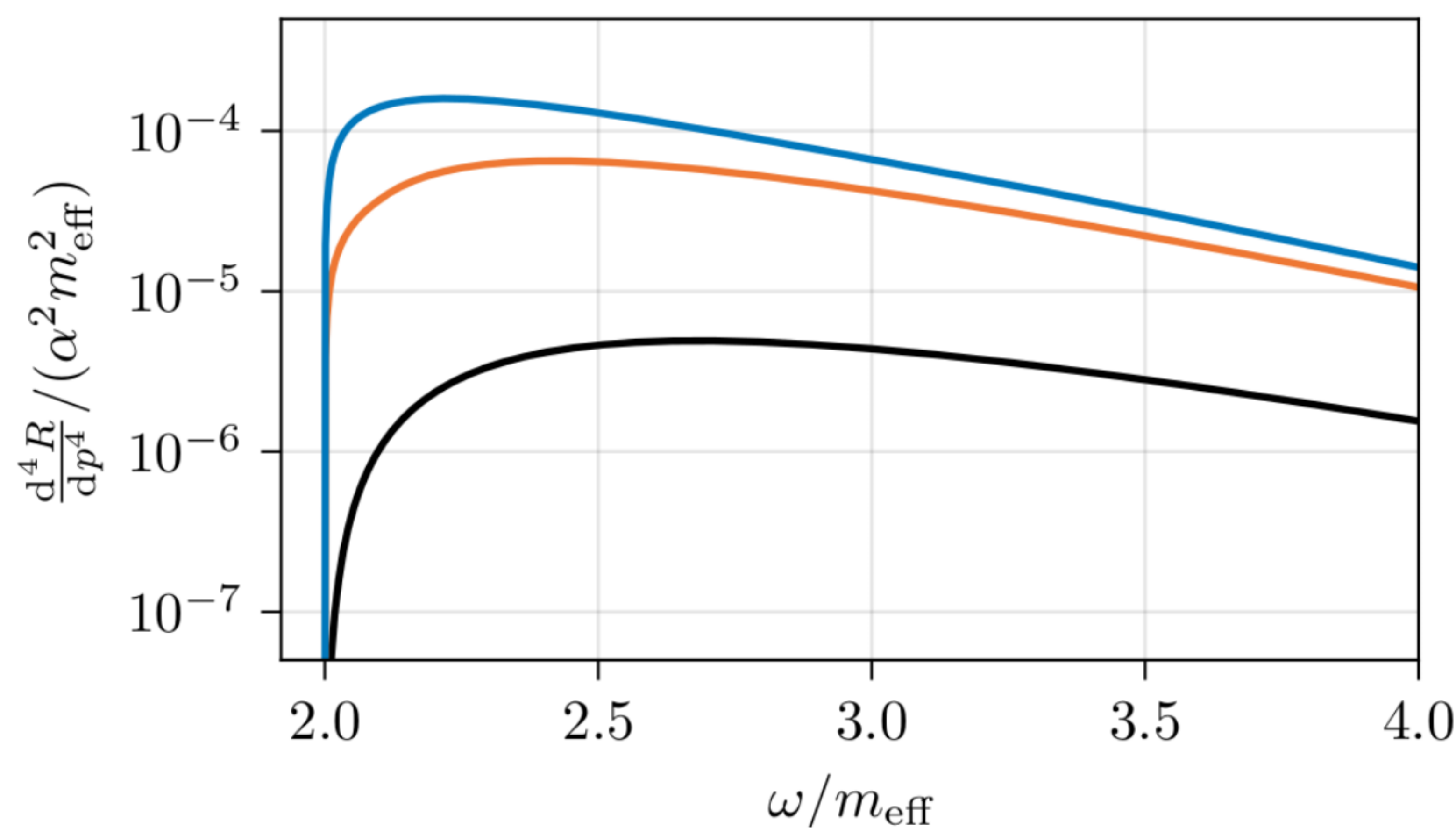
But no van Hove sing.

Moatons enhance Q^+Q^-

Disp. relation



Q^+Q^- rate

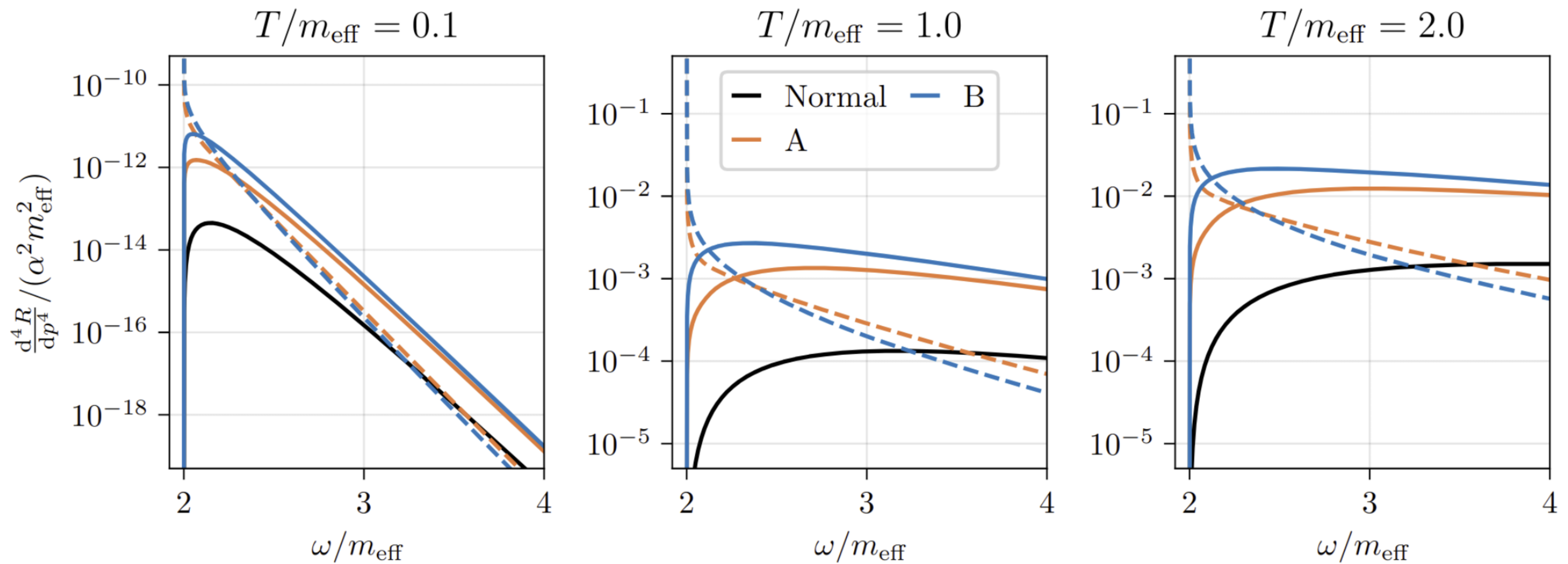


W.B. m_{eff} is NOT constant,
function of T & μ !

Without cancel Γ^μ , van Hove sing's

Hayashi & Tsue 2407,08523

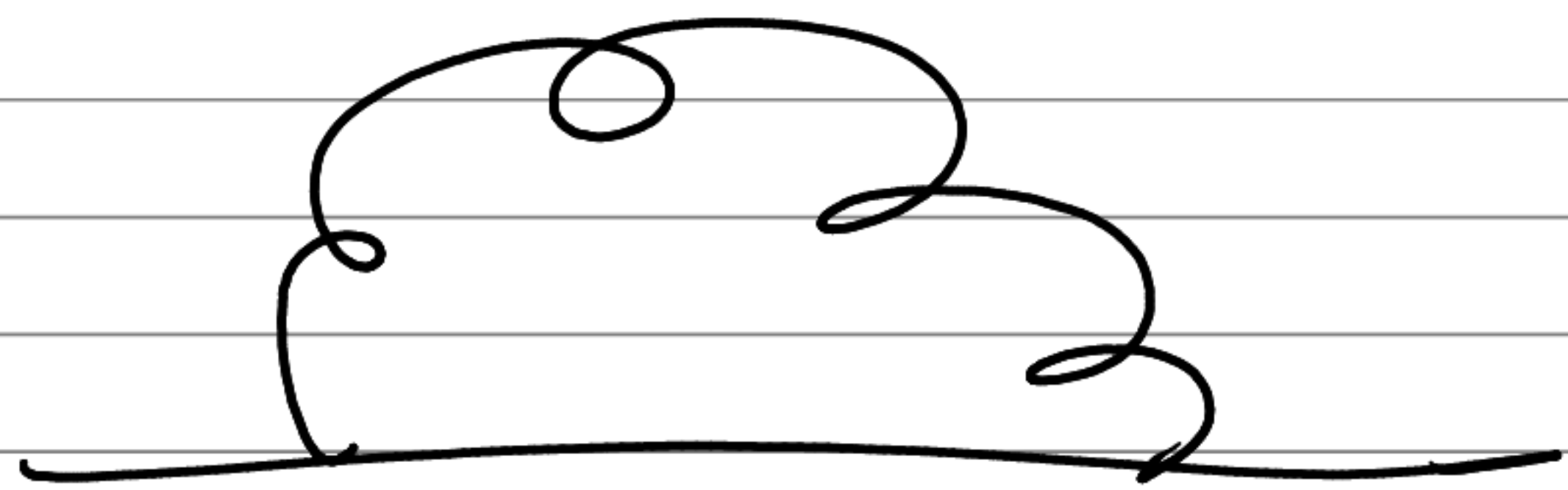
l^+l^- from chiral spiral; do not gauge $(\partial^2\phi)^2$
van Hove sing's appear. But $\int d^4x \Pi^{\mu\nu} \neq 0$.



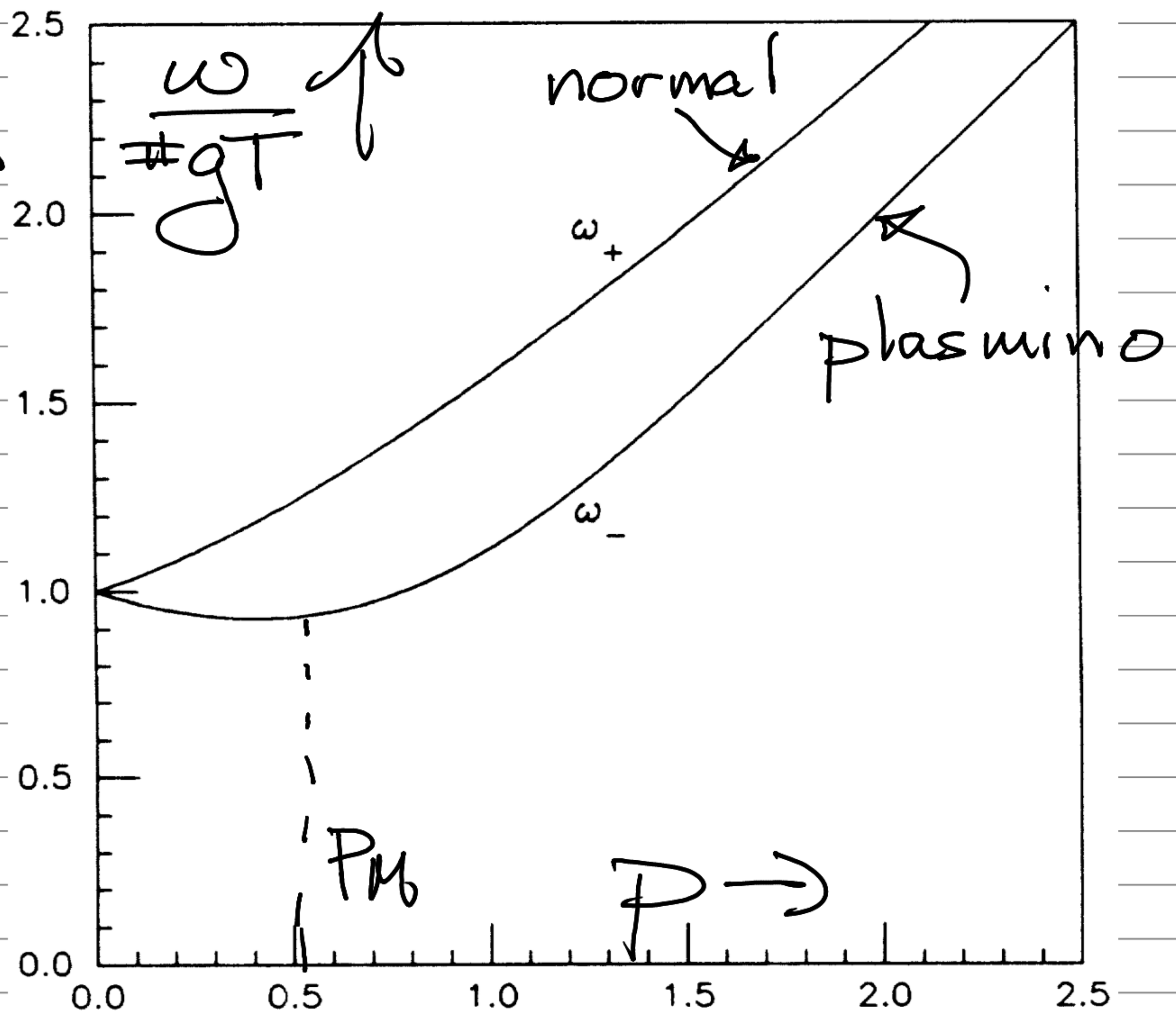
van Hove sing.'s?

Do vanH sing.'s always cancel? No!
"Plasmino" in Hard Thermal Loops

To $\gg T$,
massless quarks
get thermal
"mass" $\sim gT$



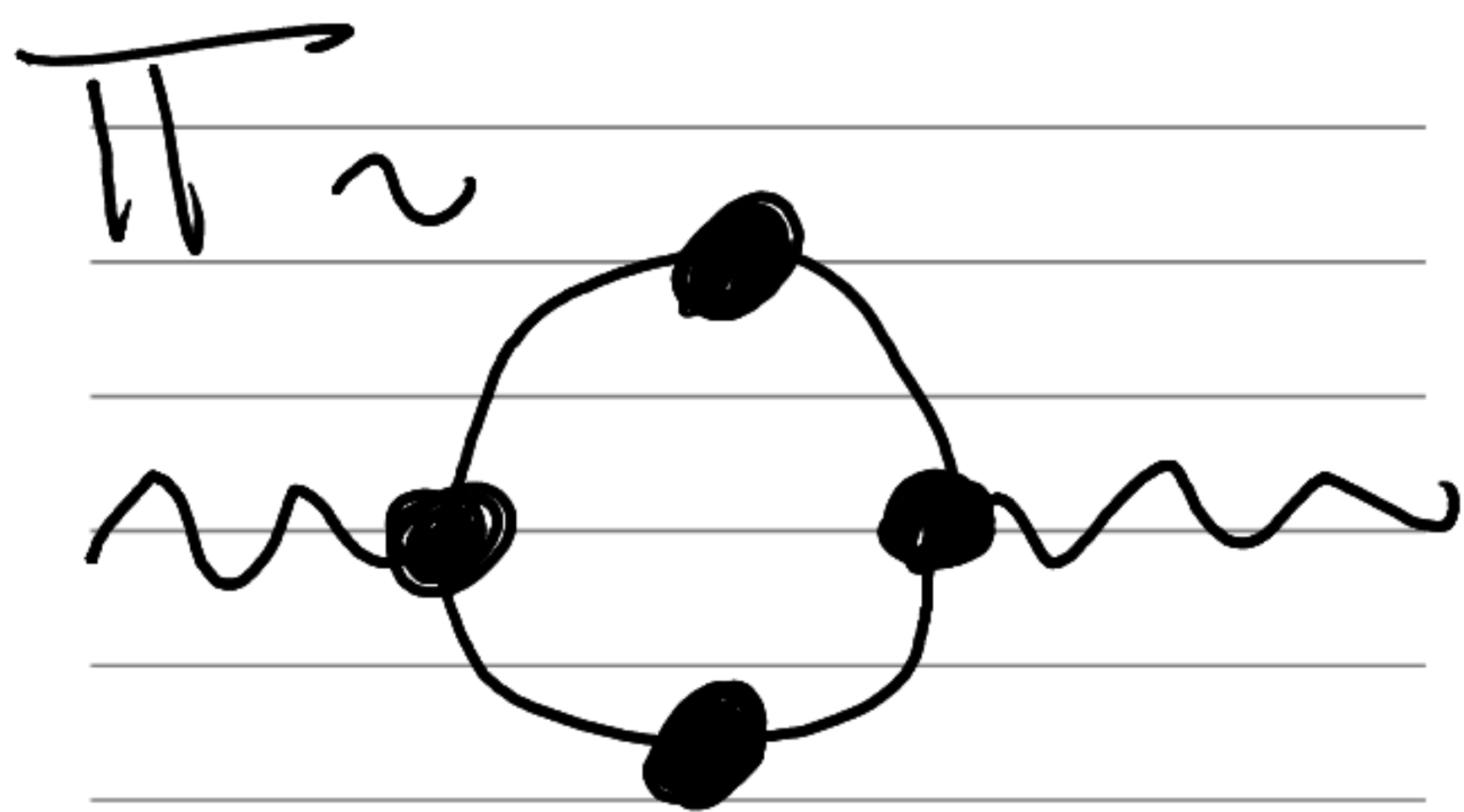
Kilinov & Webber
'81 '82



Plasminos do give van H sing.s!

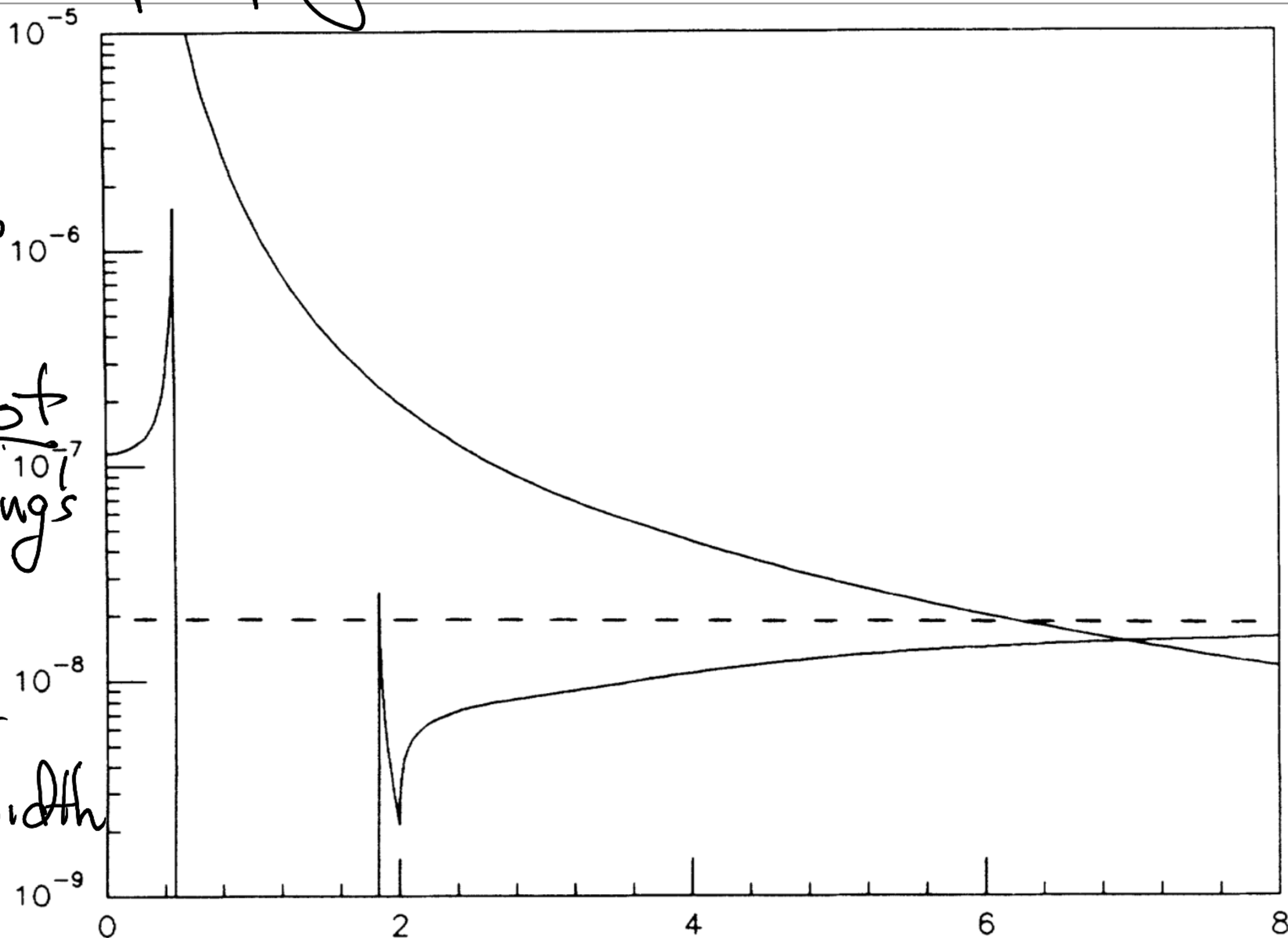
Braaten, RDP & Yuan 1990

HTL resummed propagators and vertices



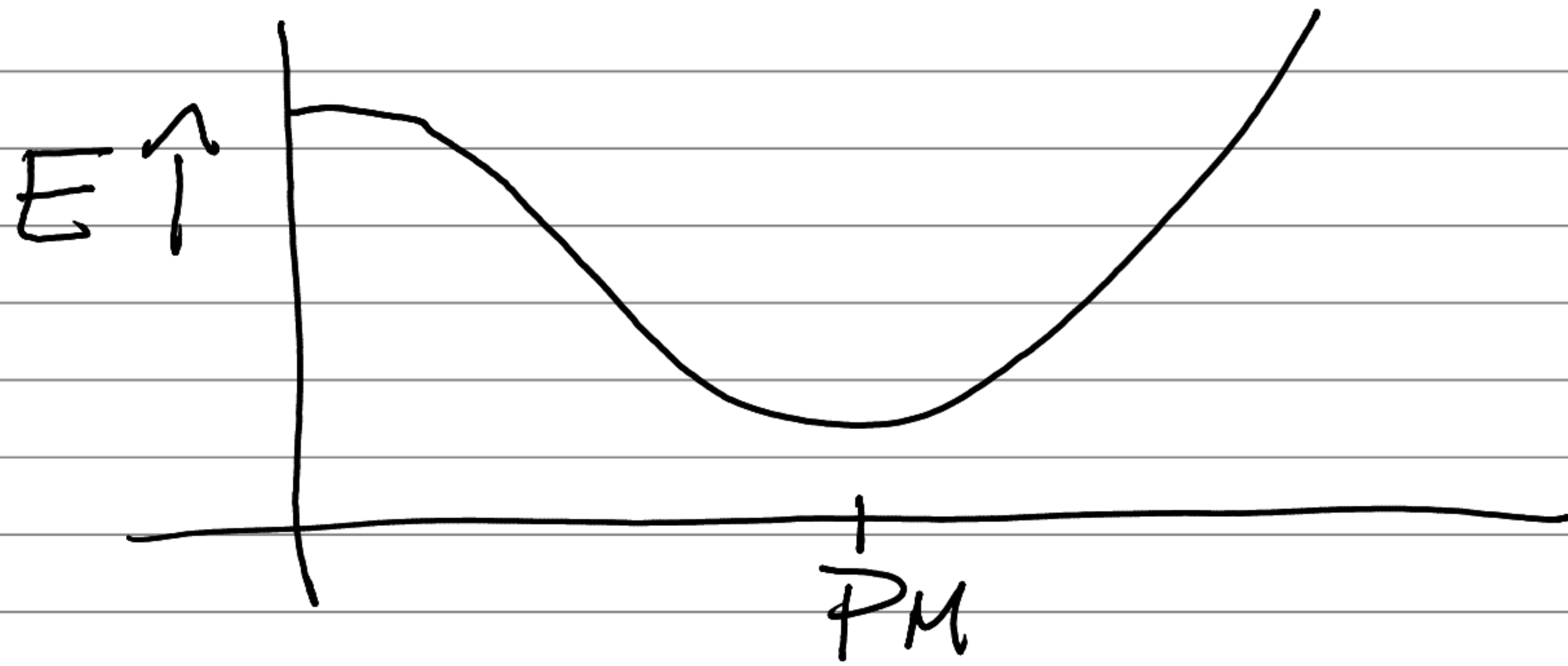
Vertices do not cancel van H sing^s

van H sing,
smoothed out
by thermal width
anyway

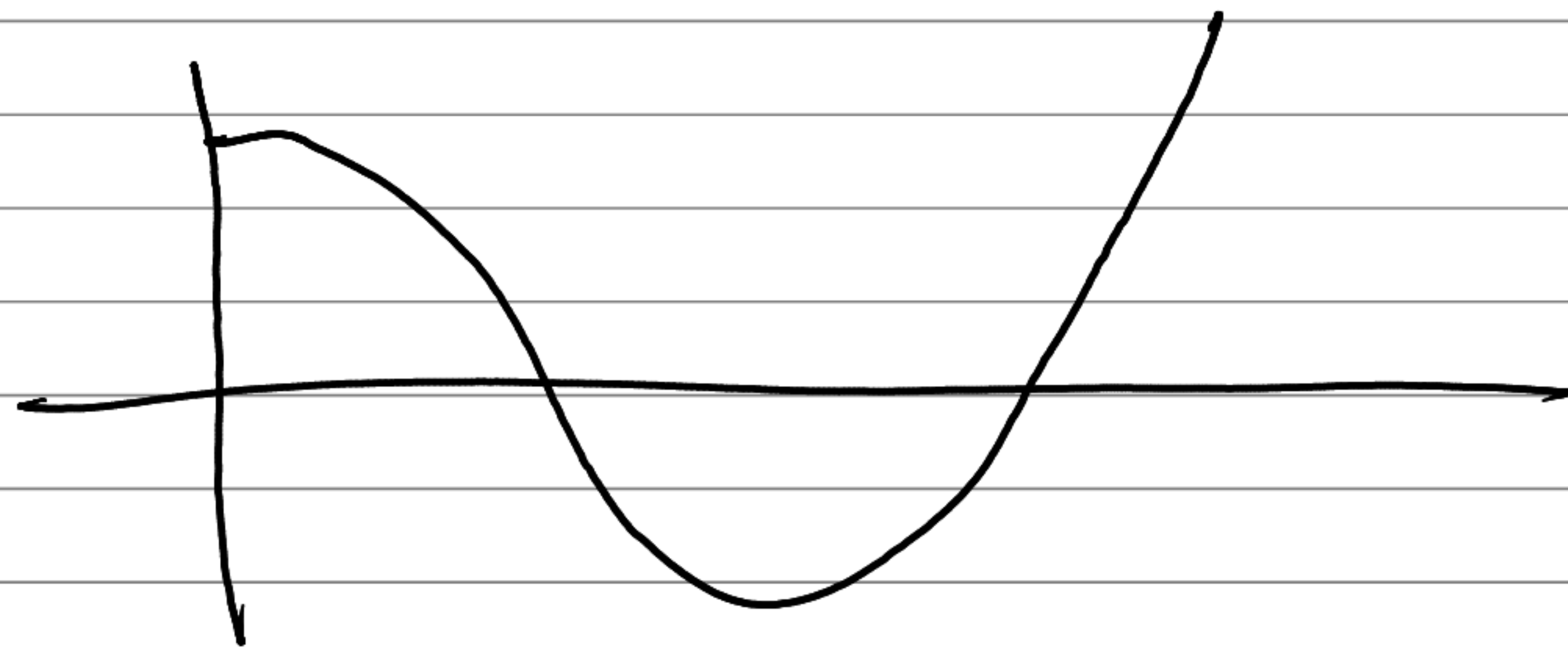


Moats & χ spirals

Moat:



What if
 $E < 0$?



\Rightarrow spatially inhomogeneous condensate
usually 1-dim. (spont. breaks rot. sym.)
 $O(1)$: kink crystal
 $O(N \geq 2)$: χ spiral
seen in
Gross-Neveu

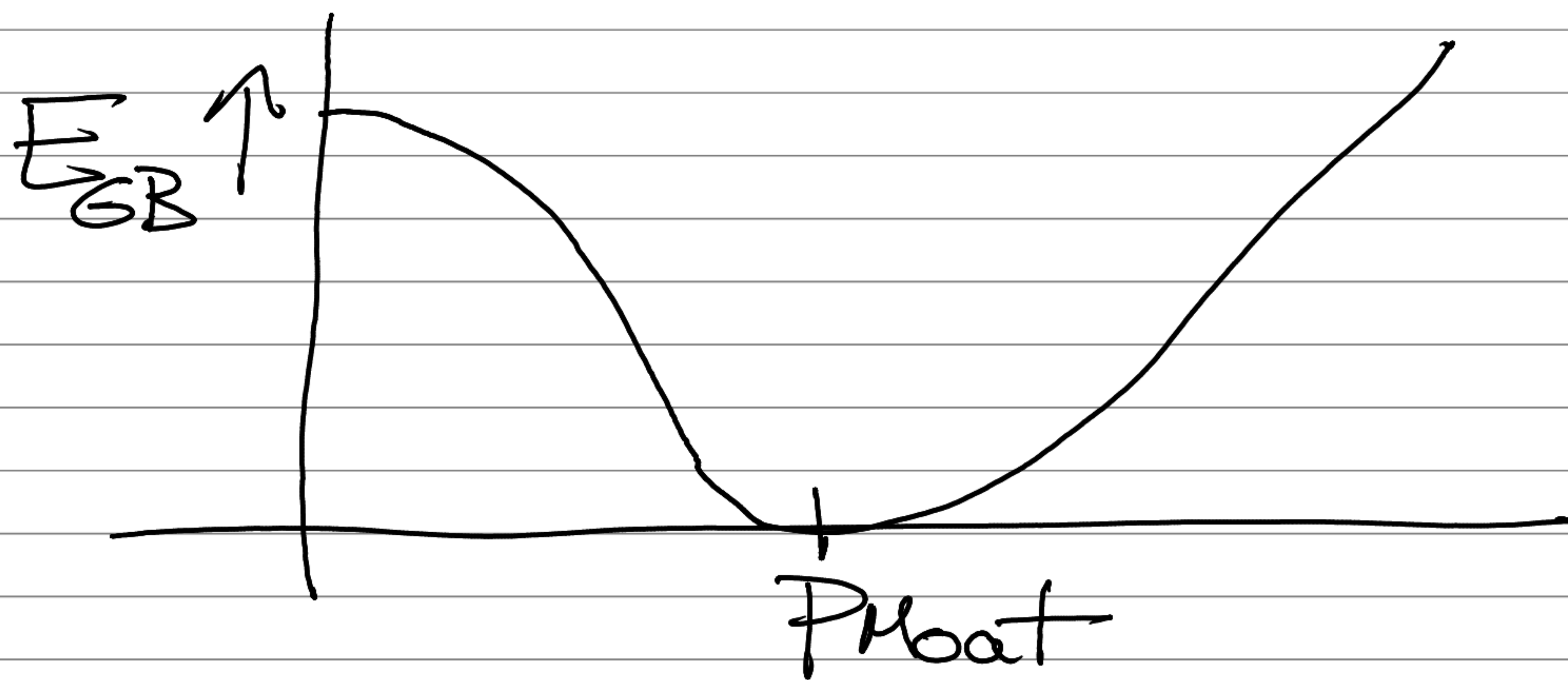
GB's in a Moat

For $O(N > 2)$, χ spiral has Goldstone Bosons (GB)

At what P_{GB} is $E_{GB} = 0$?

RDP, S. Valgushev & A. Tsvetlik 2005, 10259

With specific ansatz for χ spiral,



GB's are massless @ P_{μ} !

GB's \Rightarrow Quantum π Liquid

E_{GB} quadratic about \bar{p} ; $\bar{p} \approx \bar{p}_M + \delta p$

$$\sim T \int \frac{d^3 p}{(\delta p)^2} \sim T \bar{p}_M^2 \int \frac{d\delta p}{|\delta p|^2}$$

GB's disorder the system; linear IR divergence
Like Quantum Spin Liquid \Rightarrow
"Quantum π Liquid"

$$\langle \phi(x) \phi(0) \rangle \underset{x \rightarrow \infty}{\sim} e^{-m_{\text{real}} x} \cos(m_{\text{imag}} x)$$

$m_{\text{imag}} \neq 0$: defines "disorder line"

Distinct from disorder of X -spiral by trans. fluctuations
Tong-Gyu Lee + ... 1504, 03185

HBT in a Moat

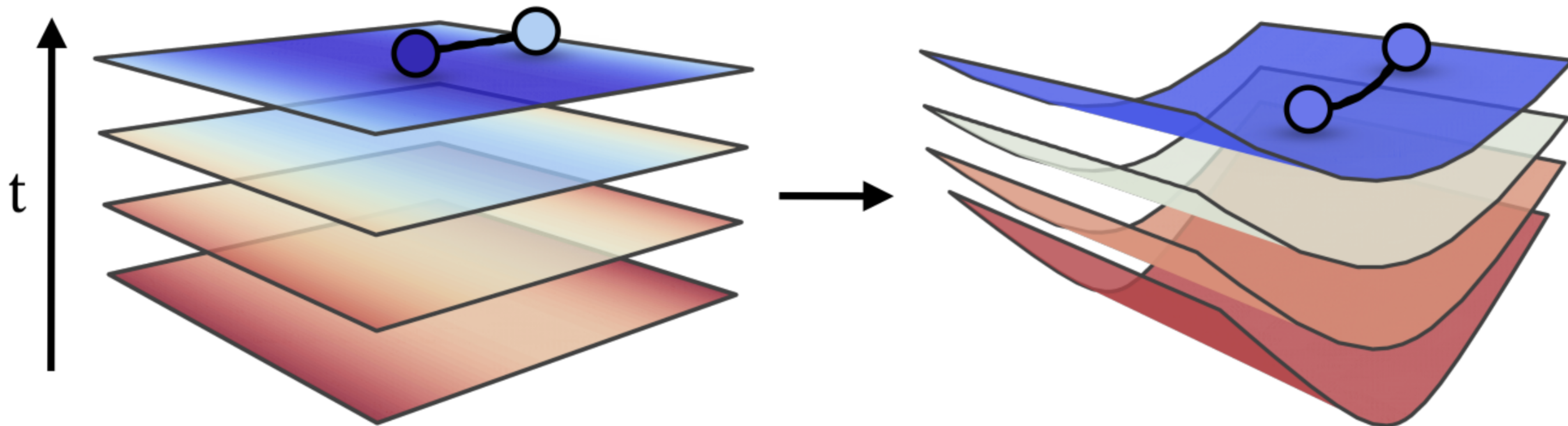
F. Pennecke, RDP, D. Rischke 2301.11484

Hanbury-Brown-Twiss (HBT) interferometry

Consider emission not surface @ constant time

But on curved hypersurface
defined by same T & μ :

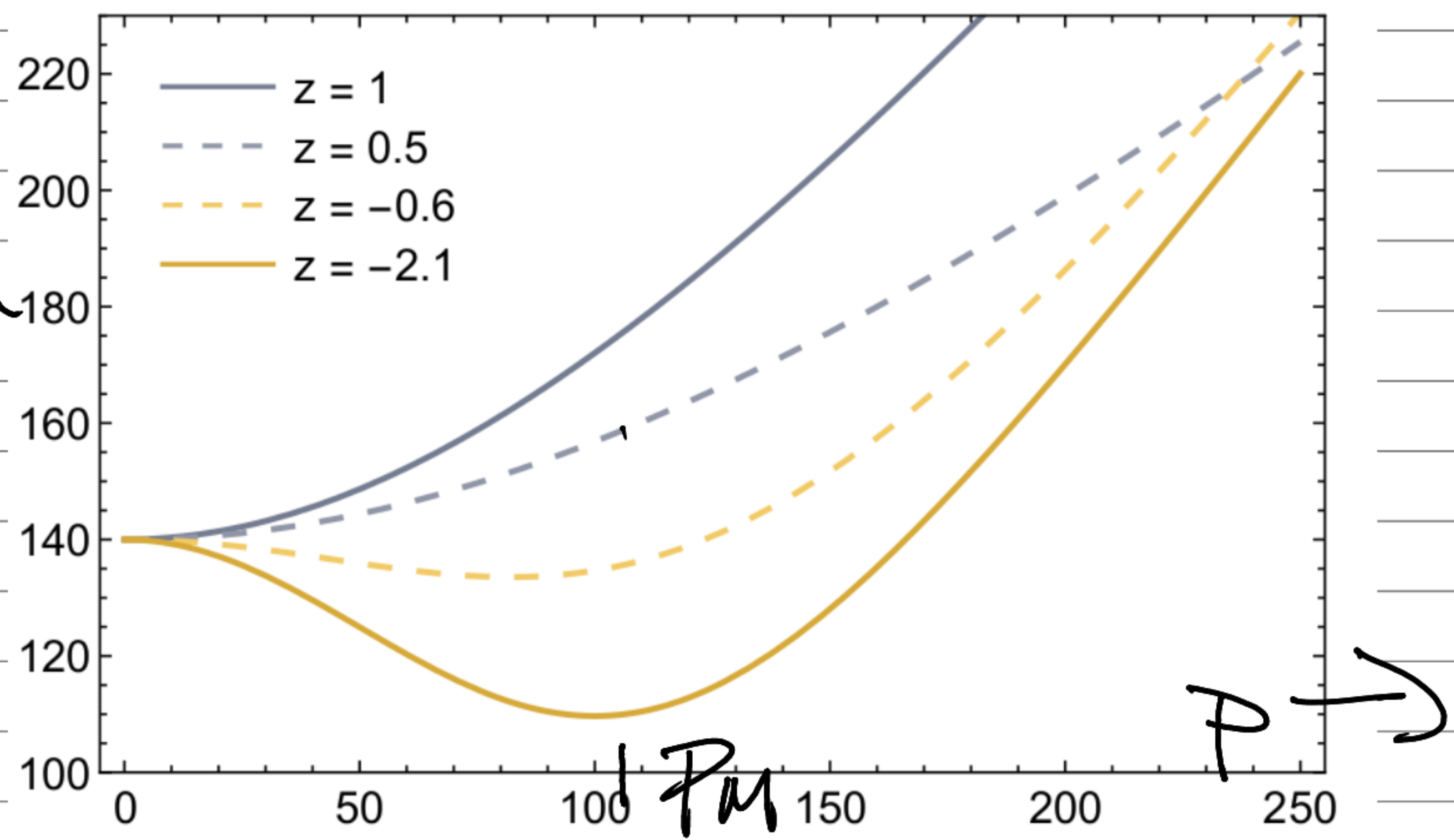
Cooper-Frye



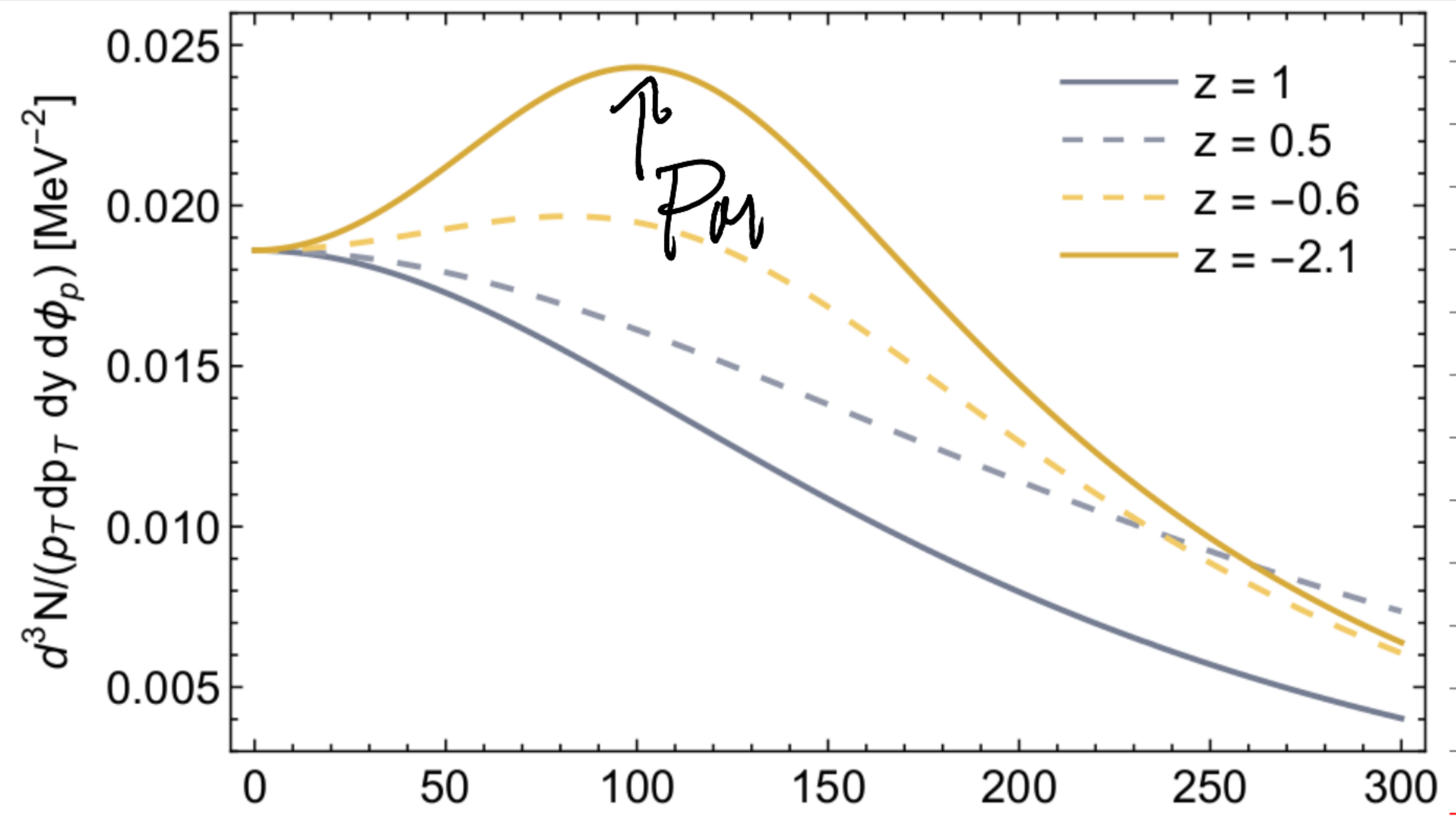
Pions in a Moat

Assume π 's emitted
in a moat,

$E(p) \uparrow$



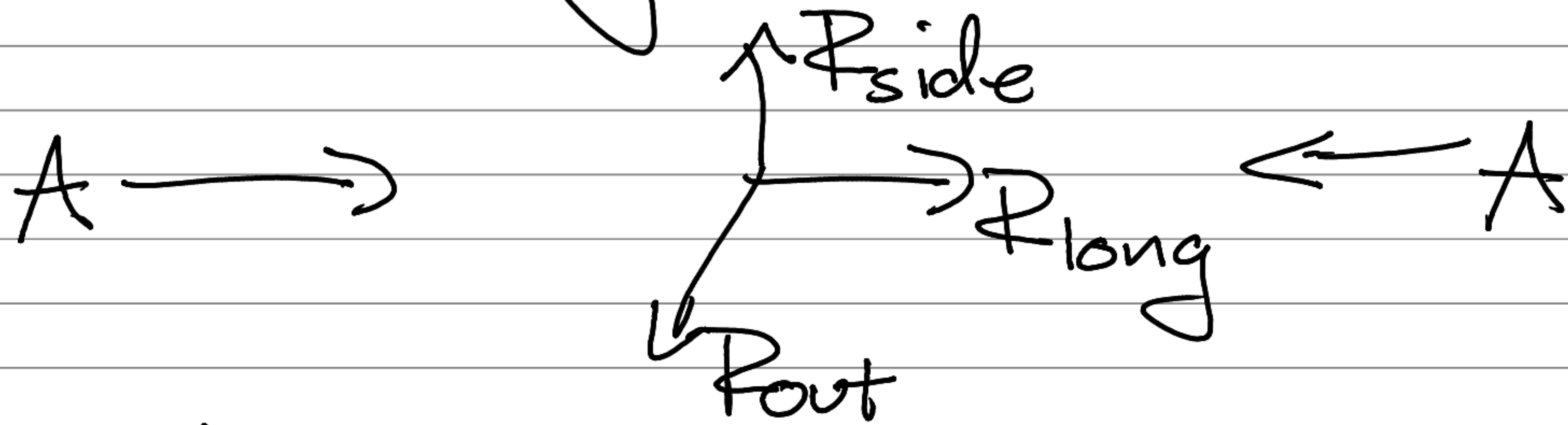
Then single
particle fnc,
has a maximum
@ p_m



Exp'y, could be
obscured by
emission of π 's
not in a moat

Moat HBT

HBT: interferometry btwn identical particles



$R = \pi\pi$ pair,

if single particle dist. func. peaked

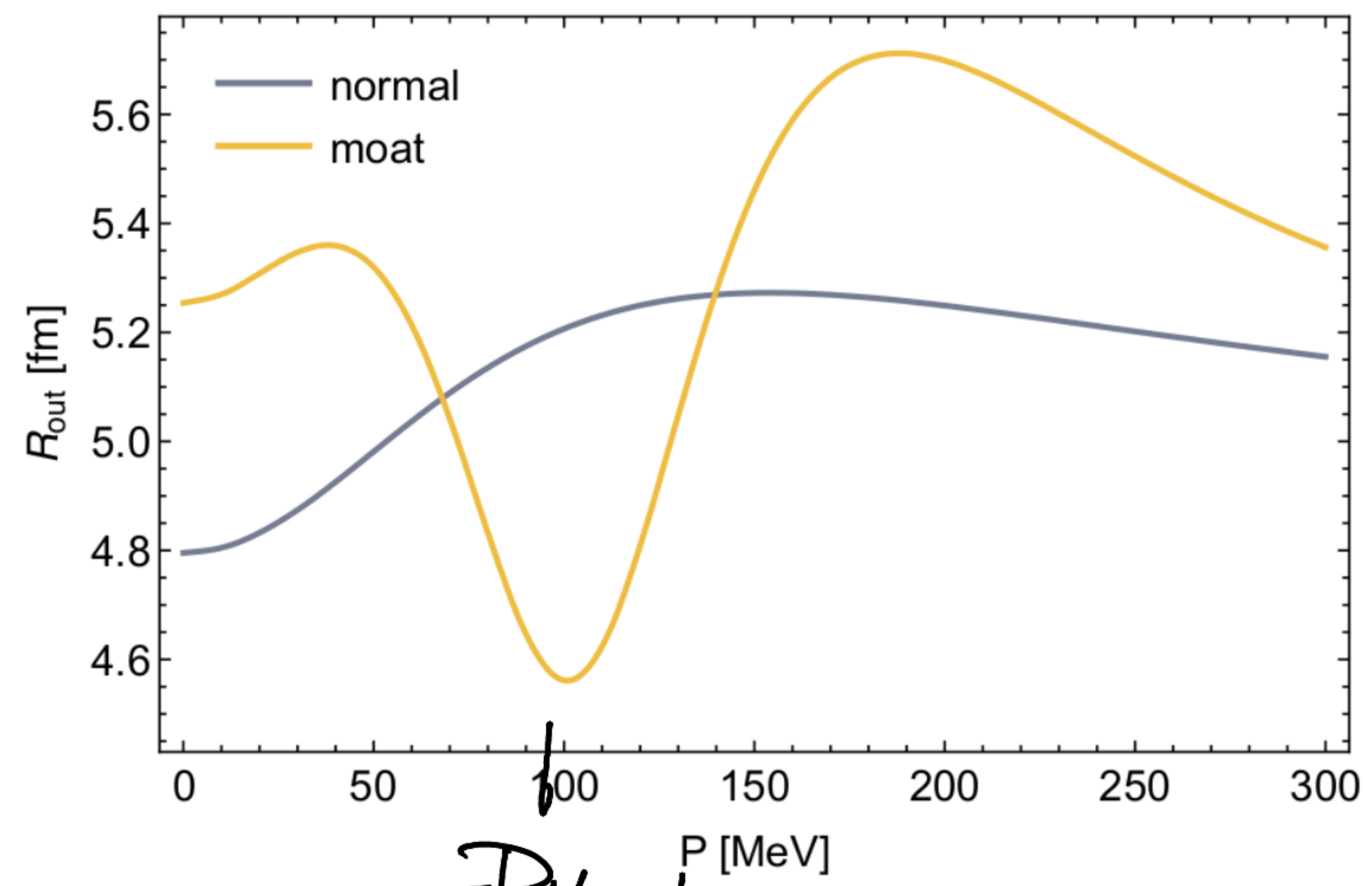
@ P_{Moat} , most show up in

two particle correlation func's

Root in a Moat

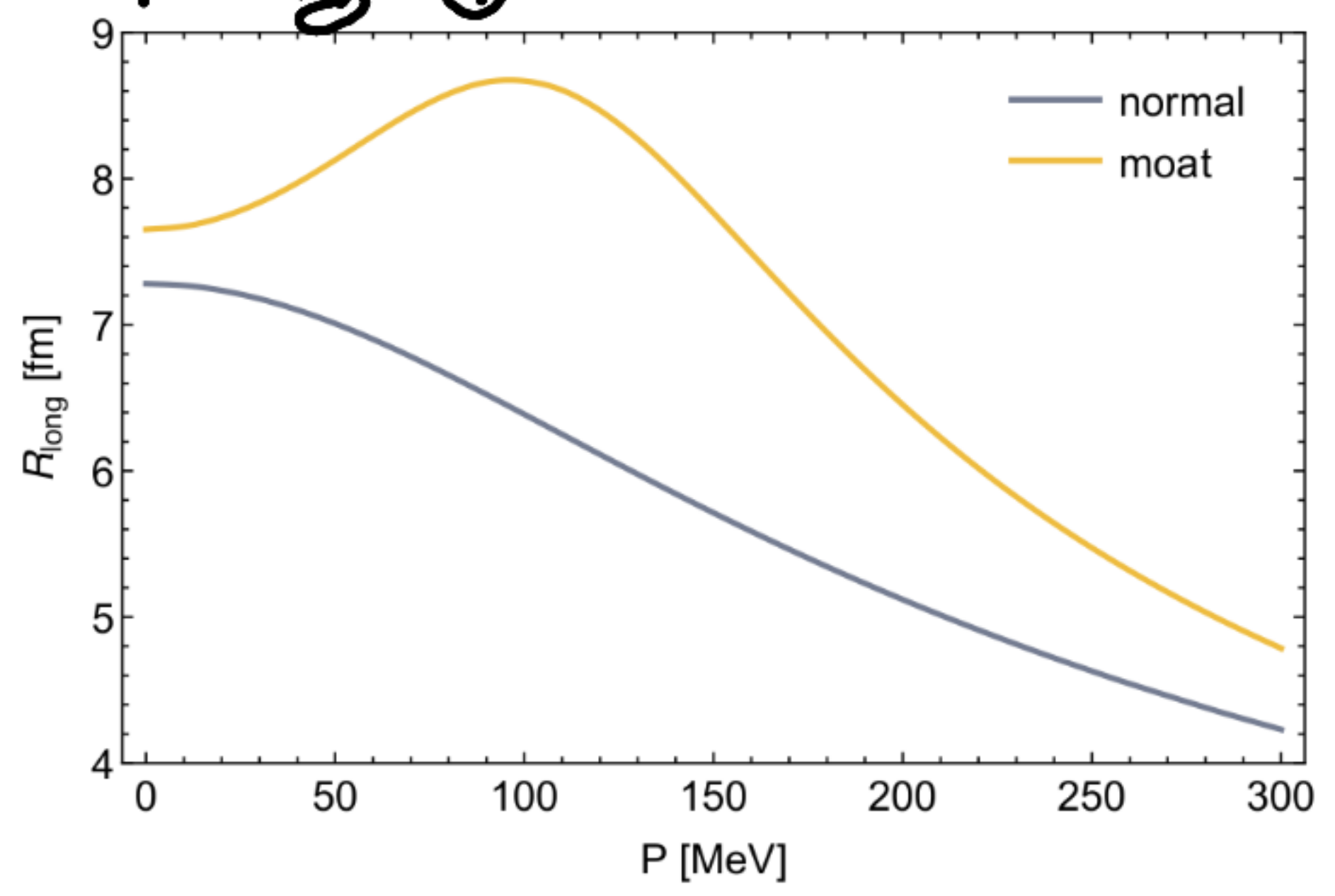
HBT radius where 2-particle corr. func. falls to $\frac{1}{2}$ maximal size

$R_{out} \uparrow$



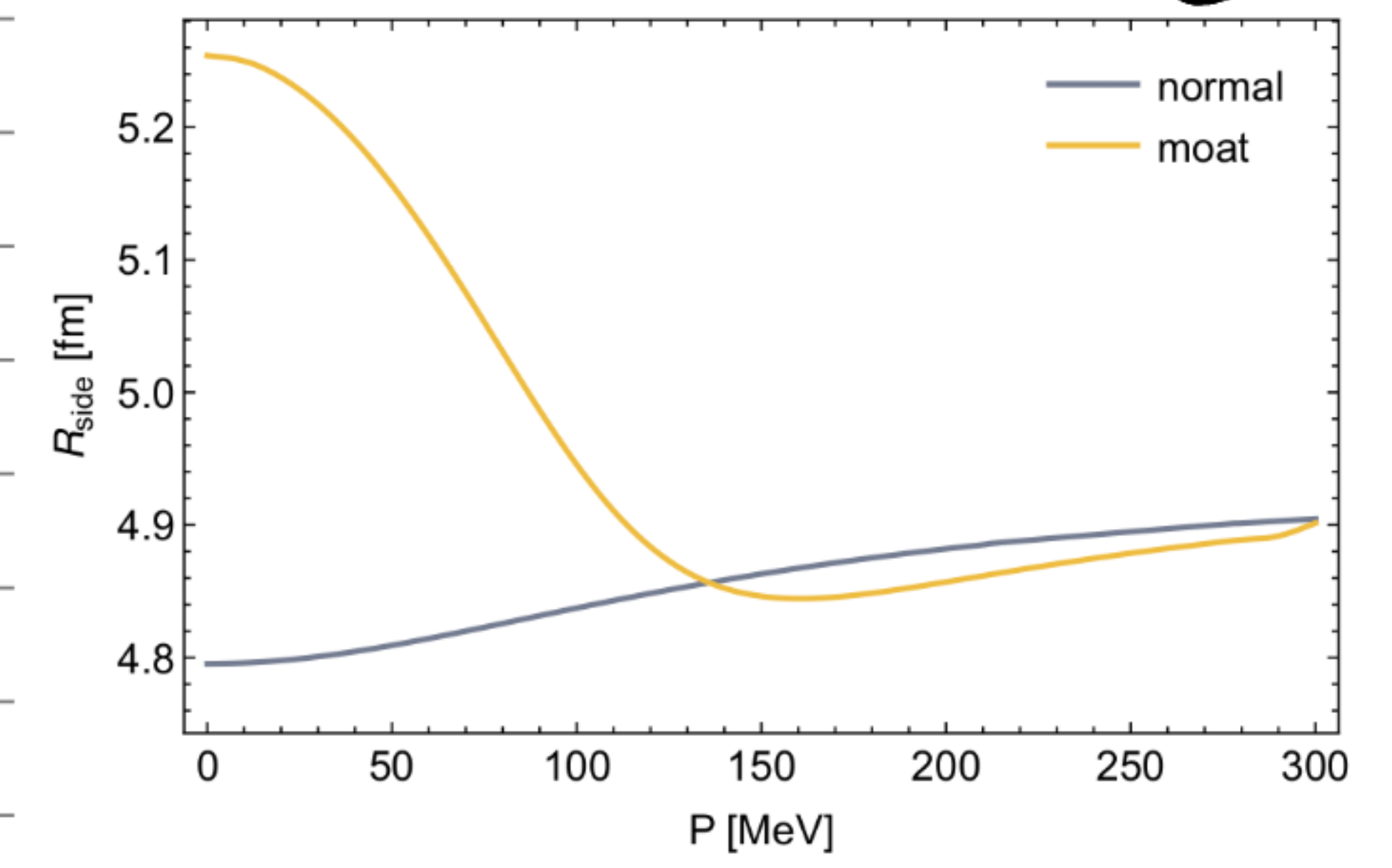
R_{out} signal moat dramatic

$R_{long} \downarrow$

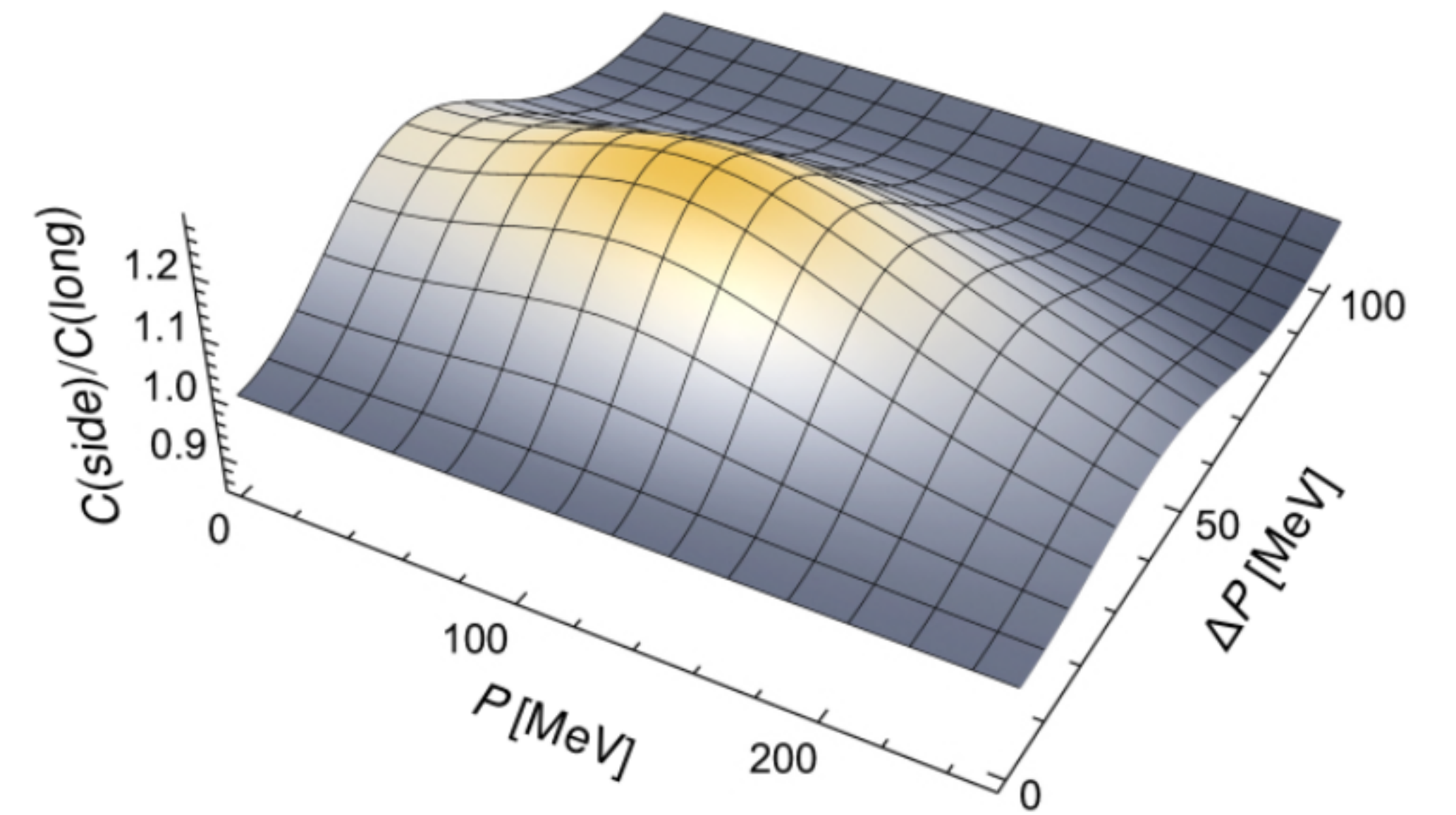
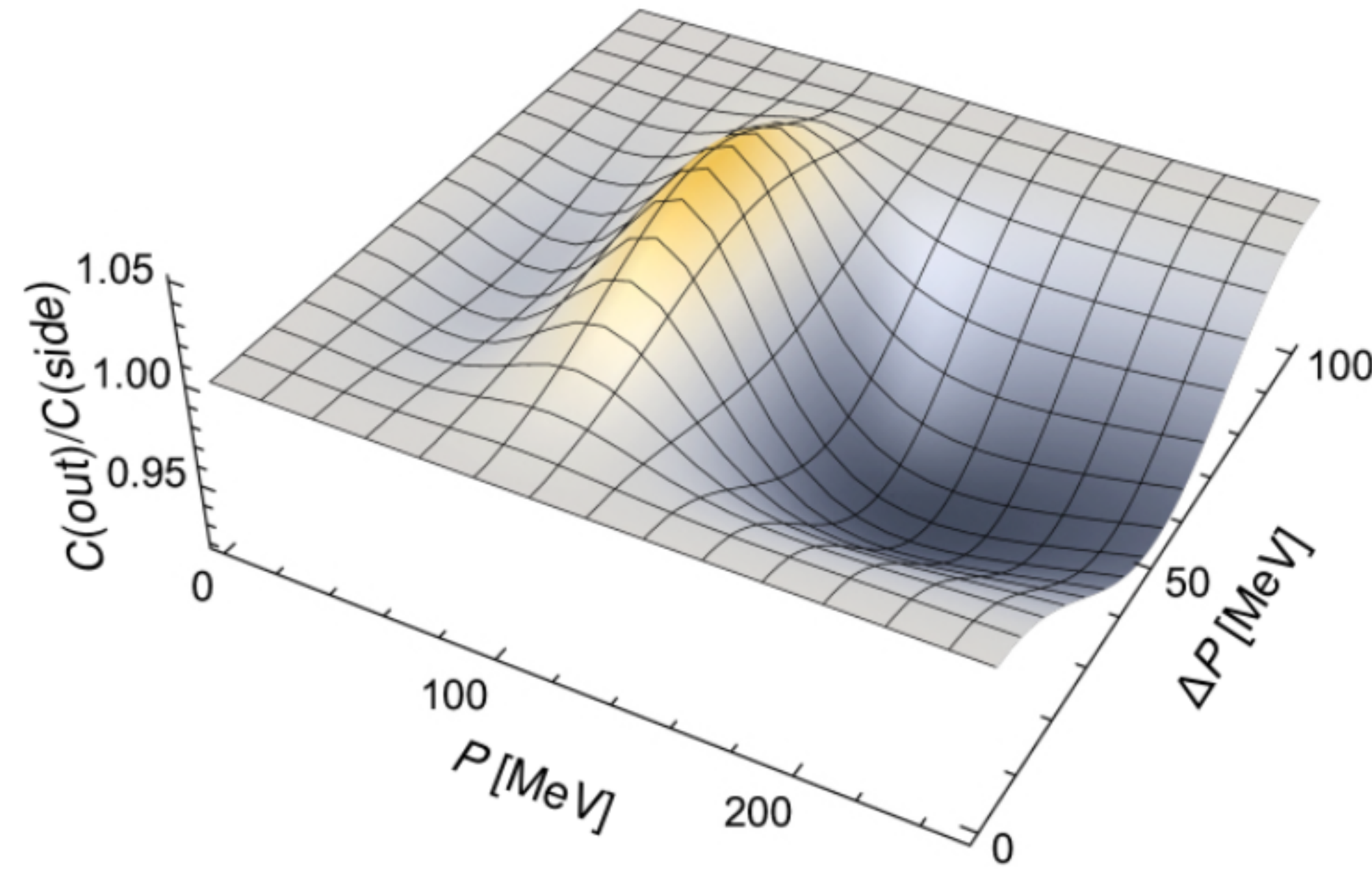
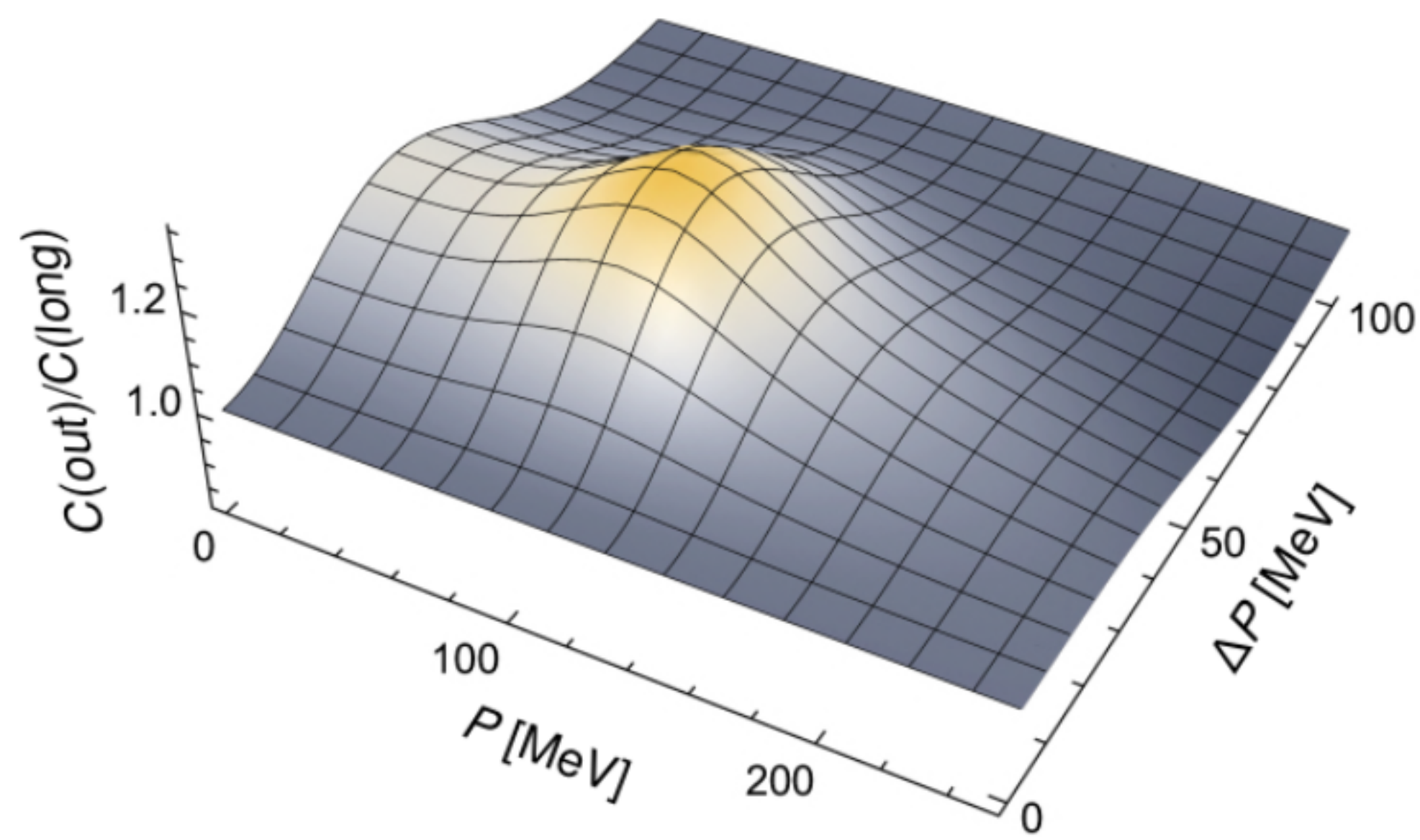


R_{moat}

$R_{side} \downarrow$



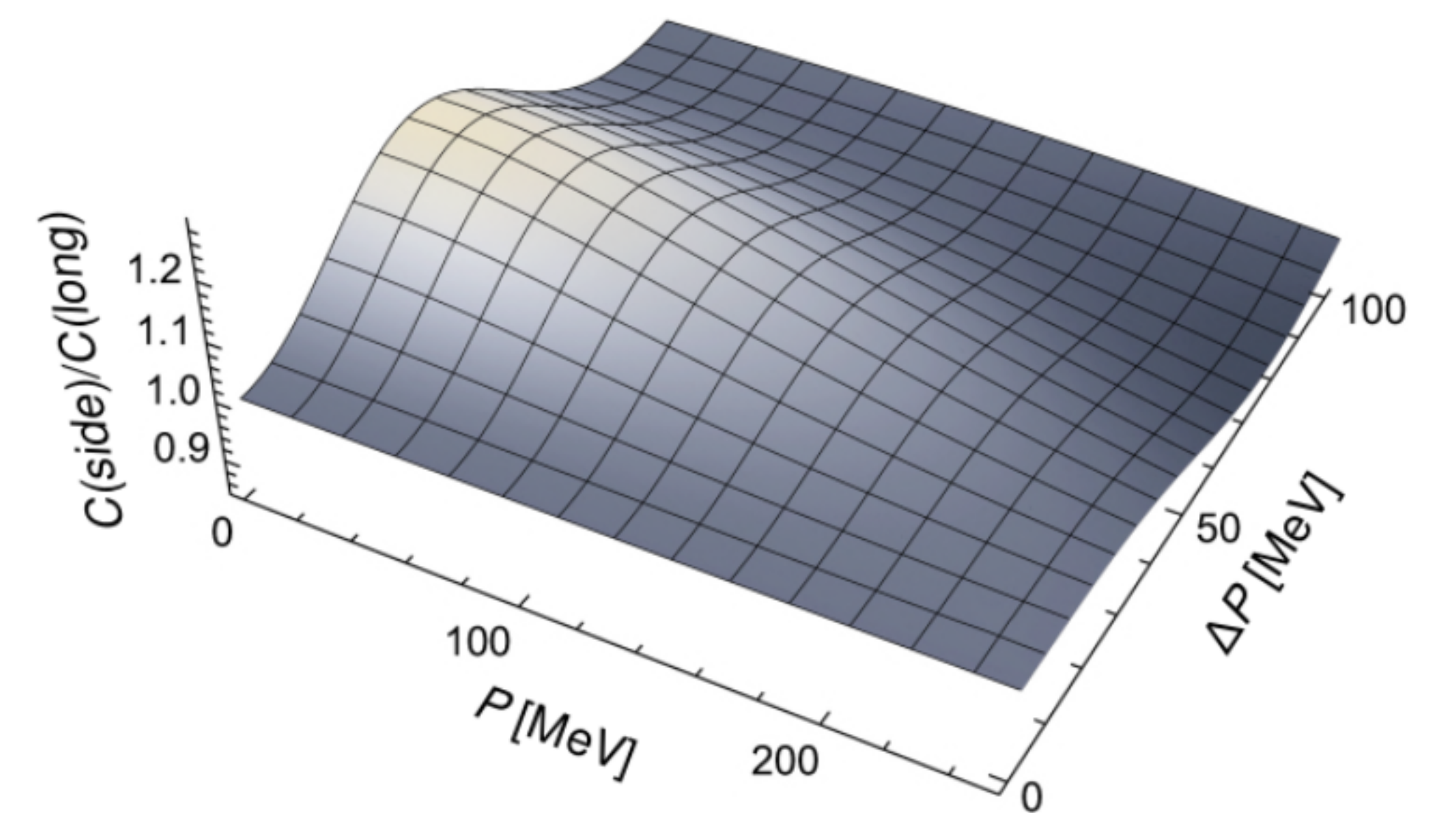
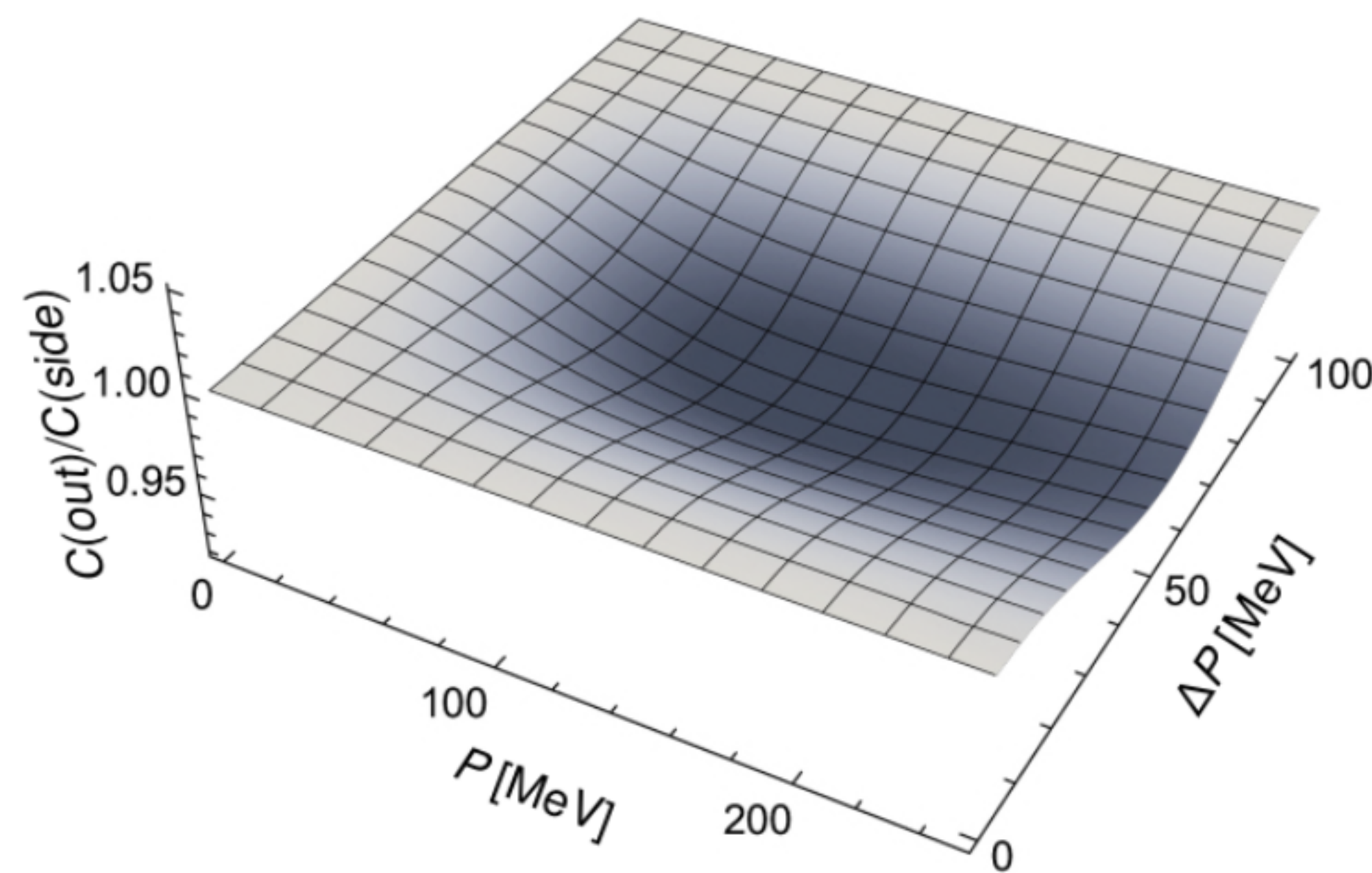
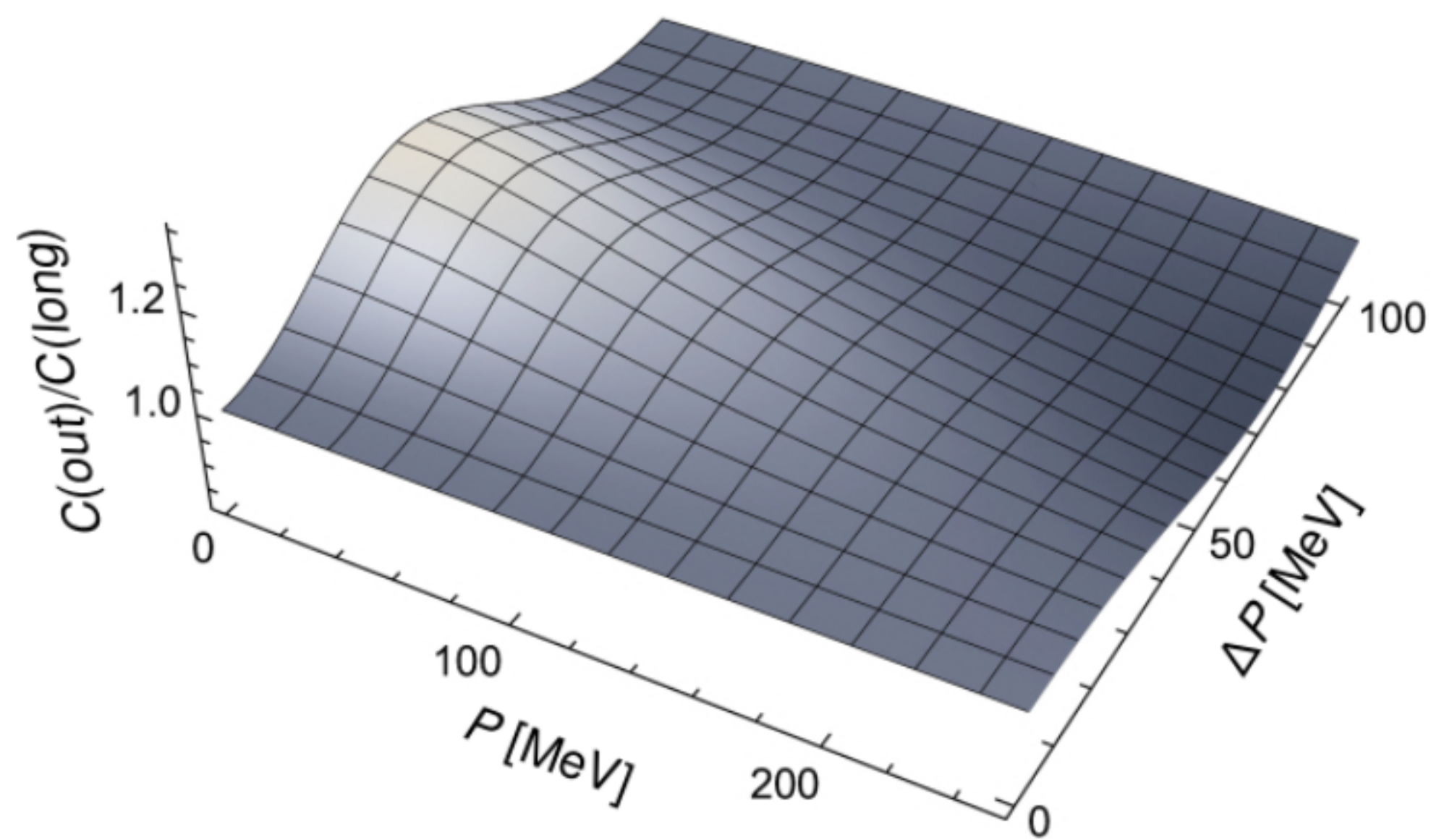
Ratio of C's in a Moat Moaty π 's



$C_{\text{out}} / C_{\text{long}}$

$C_{\text{out}} / C_{\text{side}}$

$C_{\text{side}} / C_{\text{long}}$



Normal π 's