



The three-body DD^*K system on the lattice EFT

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-- Experiments, Effective theories, and Lattice --

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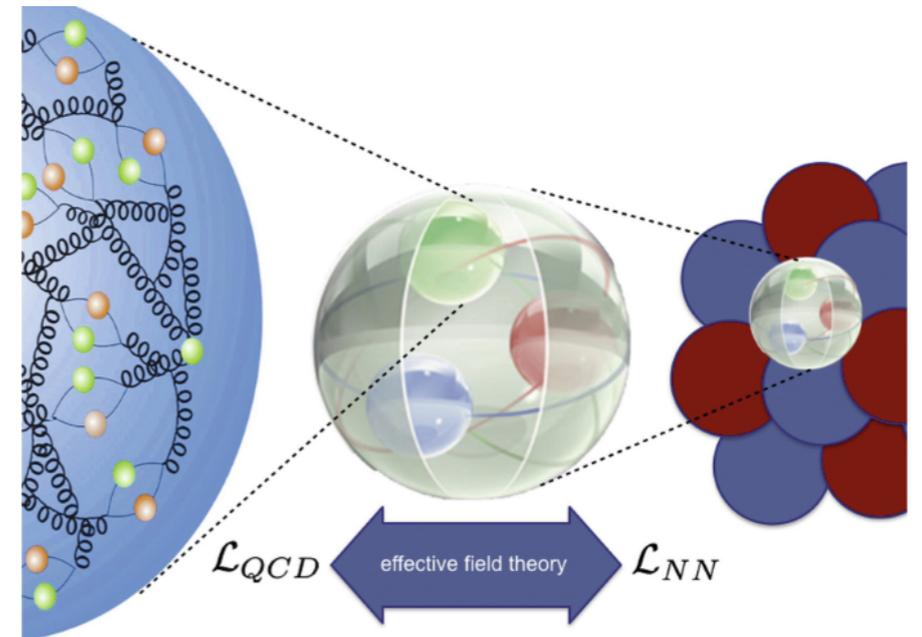
Yukawa Institute for Theoretical Physics, Kyoto University, Japan

Introduction to LEFT

Lattice EFT=Chiral EFT+Lattice+Monte Carlo

- ♦ Hadrons are effective degrees of freedom
- ♦ EFT describes hadron interactions
- ♦ Lattice spacing is set around 1 fm

	LQCD	LEFT
Degrees of freedom	quarks & gluons	hadrons
lattice spacing	~ 0.1 fm	~ 1 fm
dispersion relation	relativistic	non-relativistic
continuum limit	✓	✗
model	Lagrangian	Schrödinger
solver	path integral	matrix compute



Dean Lee, Prog.Part.Nucl.Phys., 63(2009)117,

Lähde, Meißner, “Nuclear Lattice Effective Field Theory”, Springer (2019).

Applications of LEFT

Applications in Nuclear Physics

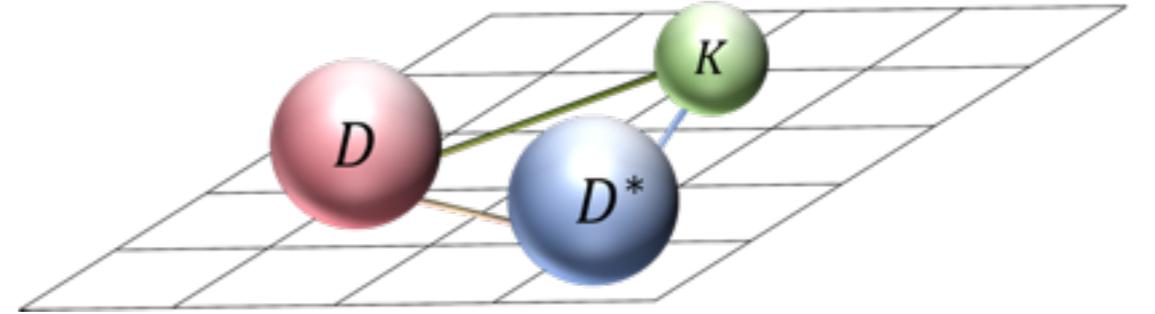
- ◆ Neutron-proton scattering [Phys.Rev.C, 98\(2018\)044002](#)
- ◆ Nuclear binding [Phys.Rev.Lett., 117\(2016\)132501](#)
- ◆ Alpha–alpha scattering [Phys.Rev.Lett., 111\(2013\)032502](#), [Nature, 528\(2015\)111](#)
- ◆ Nuclear thermodynamics [Phys.Rev.Lett., 125\(2020\)192502](#)
- ◆ Properties of nuclei [Phys.Lett.B, 797\(2019\)134863](#)
- ◆ Hoyle state [Phys.Rev.Lett., 109\(2012\)252501](#), [Nat. Commun., 14\(2023\)2777](#)

Applications of LEFT

Applications in Hadron Physics

TABLE XXXIX. Summary for heavy-flavor three-body states. Energies are in units of MeV.

Components	$I(J^P)$	Results (Method)	Decay modes
DNN	$\frac{1}{2}(0^-)$	BS $\sim 3500 - 15i$ (FCA, V) [836]	$\Lambda_c \pi^- p, \Lambda_c p$ [836]
$NDK, ND\bar{K},$ NDD	$\frac{1}{2}(\frac{1}{2}^+)$	BS $\sim 3050, 3150, 4400$ (FCA) [837]	†
DD^*N	$\frac{1}{2}(\frac{1}{2}^+, \frac{3}{2}^+)$	BS $\sim 4773.2, 4790.7$ (GEM) [838]	$T_{cc}p, DDp + \pi(\gamma), \Xi_{cc} + \pi(\gamma),$ charmed baryon + charmed meson [838]
DD^*N	$\frac{3}{2}(-)$	difficult to form bound states (GEM) [838]	†
$DK\bar{K}$	$\frac{1}{2}(0^-)$	D -like state ~ 2845.5 (FCA) [821], D -like state ~ 2900 (QCDSR, χ F) [839]	$\pi\pi D$ [821]
DKK	$\frac{1}{2}(0^-)$	no bound state (FCA) [821]	†
$\bar{D}\bar{K}\Sigma_c$	$1(\frac{1}{2}^+)$	BS ~ 4738.6 (GEM) [840]	$D\bar{K}, D_s \Sigma_c$ [840]
$D^{(*)}$ multi ρ	...	several $D_J^{(*)}$ states (FCA) [841, 842]	†
$\rho D\bar{D}$	$0(?), 1(?)$	BS $\sim 4241 - 10i, [4320 - 13i, 4256 - 14i]$ (FCA) [843]	†
DDK	$\frac{1}{2}(0^-)$	BS ~ 4162 (GEM) [273], 4140 (χ F) [819], 4160 (FV) [820]	$DD_s^*, D^* D_s$ [826]
$D\bar{D}K$	$\frac{1}{2}(0^-)$	BS ~ 4181.2 (GEM) [822], 4191 (FCA) [825]	$D_s \bar{D}^*, J/\psi K$ [822]
DD^*K	$\frac{1}{2}(1^-)$	BS ~ 4317.9 (BO) [823]	†
$D\bar{D}^*K$	$\frac{1}{2}(1^-)$	BS ~ 4294.1 (GEM) [822], 4317.9 (BO) [823], 4307 (FCA) [824]	$D_s^{(*)} \bar{D}^{(*)}, J/\psi K^*$ [823, 844]
$D^* D^* \bar{K}^*$	$\frac{1}{2}(0^-, 1^-, 2^-)$	BS $\sim [4850 - 46i, 4754 - 50i],$ (FCA) [845] [4840 - 43i, 4755 - 50i]	$D^* D^* \bar{K}^*,$ $D^* D^{(*)} \bar{K}^*,$ [845] $[D^* D^* \bar{K}^*, D^* D^{(*)} \bar{K}^*]$
$\bar{D}\bar{D}^* \Sigma_c$	$1(\frac{1}{2}^+, \frac{3}{2}^+)$	BS $\sim 6292.3, 6301.5$ (GEM) [829]	$J/\psi p \bar{D}^{(*)}, T_{cc} \Lambda_c \pi$ [829]
$J/\psi K\bar{K}$	$0(1^-)$	$Y(4260) \sim 4150 - 45i$ (χ F) [481]	†
DDD^*	$\frac{1}{2}(1^-)$	BS ~ 5742.2 (GEM) [833]	$DDD\pi(\gamma)$ [833]
$DD^* D^*$	$\frac{1}{2}(0^-, 1^-, 2^-)$	several loosely bound states (GEM) [834]	charmed mesons + ... [834]
$D^* D^* D^*$	$\frac{1}{2}(0^-, 1^-, 2^-, 3^-)$	several loosely bound states (GEM) [834]	charmed mesons + ... [834]
$D^* D^* D^{(*)}$	$\frac{1}{2}(0^-, 1^-, 2^-)$	BS $\sim 5790.9 - 49.8i, 5990.2, 5989.4$ (FCA) [835]	
$D^* D^* D^{(*)}$	$\frac{3}{2}(-)$	difficult to form bound states (GEM) [834]	†
$D^* D^* \bar{D}$	$\frac{1}{2}(2^-)$	BS ~ 5879 (F) [846]	†
$D^* D^* \bar{D}^*$	$\frac{1}{2}(3^-)$	BS ~ 6019 (F) [846]	†
$\Omega_{ccc} \Omega_{ccc} \Omega_{ccc}$	$?(\frac{3}{2}^+)$	no bound state (GEM) [847]	†
$\Xi_{cc} \Xi_{cc} \bar{K}$	$\frac{1}{2}(0^-)$	BS ~ 7641.8 (GEM) [848]	†



- ◆ Gaussian expansion method (GEM)
- ◆ QCD sum rule (QCDSR)
- ◆ Born-Oppenheimer approximation
- ◆ Fixed center approximation (FCA)
- ◆ Faddeev equation (F)

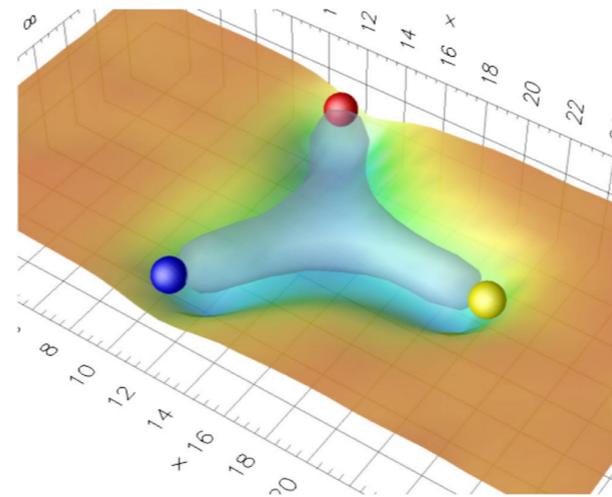
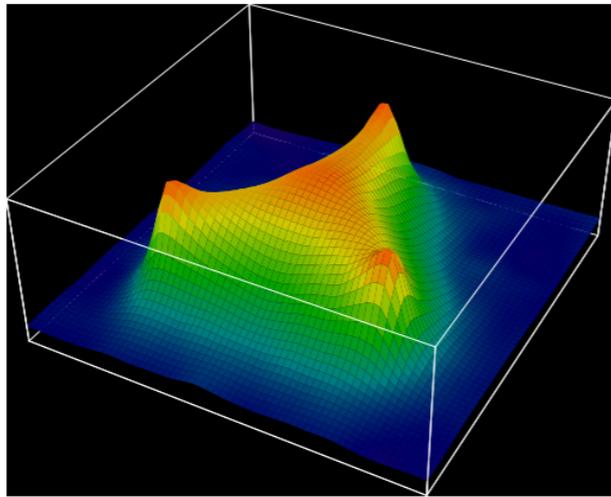
Without 3-body force!

Talk by Lisheng Geng this morning

Liu et al., arXiv:2404.06399

Advantages of LEFT

- Include three-body force directly
- ◆ Three-body force in baryon from LQCD



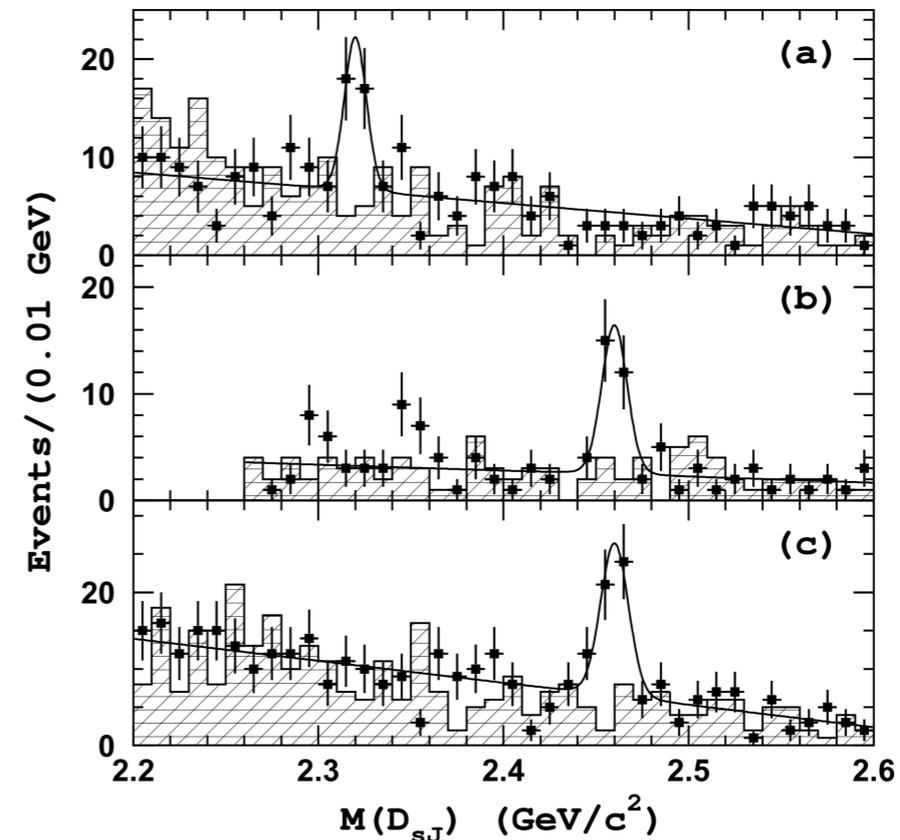
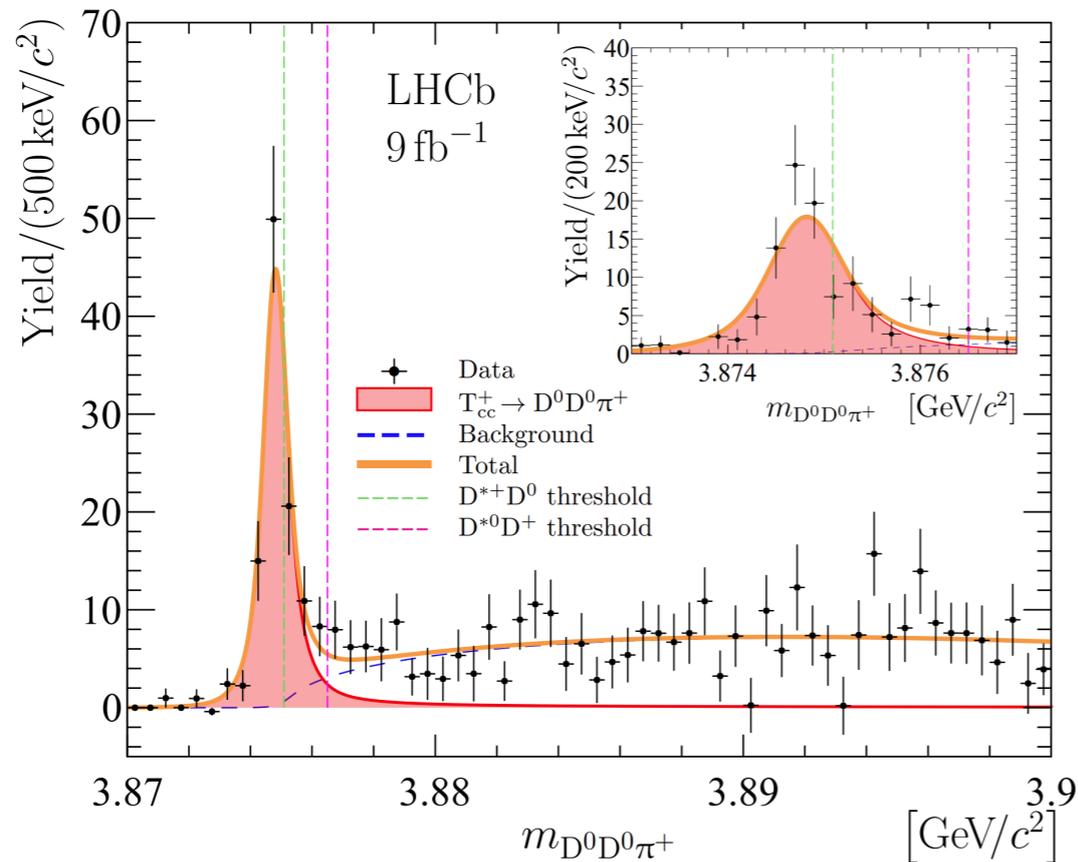
H. Ichie et al., Nucl. Phys.A 721(2003)C899-C902, F. Bissey et al., Phys.Rev.D 76(2007)114512

- ◆ Three-body force makes nucleus system ren. group invariant

Λ (MeV)	250	275	300	325	350	375	400	Exp.
c_E	5.170	2.763	1.538	0.890	0.561	0.412	0.380	
$E_{2NF}({}^3\text{H})$	-6.17(4)	-6.63(4)	-7.05(2)	-7.39(2)	-7.64(1)	-7.77(1)	-7.78(1)	-8.482
$E_{2NF+3NF}({}^3\text{H})$	-8.482	-8.482	-8.489	-8.485	-8.483	-8.483	-8.483	-8.482
$E_{2NF}({}^4\text{He})$	-30.6(7)	-30.3(6)	-30.7(4)	-30.0(4)	-29.8(4)	-29.4(4)	-29.2(4)	-28.34
$E_{2NF+3NF}({}^4\text{He})$	-29.8(7)	-29.5(6)	-29.9(4)	-29.2(4)	-29.0(4)	-28.6(4)	-28.4(4)	-28.34
$E_{2NF}({}^{16}\text{O})$	-144.0(21)	-135.1(14)	-136.3(11)	-139.1(9)	-140.6(8)	-141.7(8)	-141.8(9)	-127.6
$E_{2NF+3NF}({}^{16}\text{O})$	-135.8(20)	-124.8(14)	-124.5(11)	-126.3(9)	-127.3(8)	-128.1(8)	-128.1(8)	-127.6

The motivation of the DD^*K system

- The observation of T_{cc}^+ , $D_{s0}^*(2317)$, $D_{s1}(2460)$ in experiment



R. Aaij et al. (LHCb), Nature Phys. 18 (2022)751 P. Krokovny et al. (Belle), Phys. Rev. Lett. 91(2003) 262002

- ♦ Close to the DD^* , DK , D^*K thresholds
- ♦ Could be hadronic molecular candidates
- ♦ Be used to constraint the two-body force

The motivation of the DD^*K system

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Meng et al., PRD104(2021)051502,

Agaev et al., NPB975(2022)115650,

Feijoo et al., PRD104(2021)114015,

Yan et al., PRD105(2022)014007,

Albaladejo et al., PLB829(2022)137052,

Du et al., PRD105(2022)014024,

Padmanath et al., PRL129(2022)032002,

Chen et al., PLB833(2022)137391,

Lyu et al., PRL131(2023)161901,.....

Guo et al., EPJA40(2009)171,

Liu et al., PRD87(2013)014508,

Guo et al., PRD98(2018)014510,

Liu et al., PRD109(2024)5,

Kim et al., PTEP2024(2024)073D01,

Gil-Dominguez et al., PLB843(2023)137997,

Asokan et al., EPJC83(2023)850,

Fu et al., EPJA58(2022)70,

Kong et al., PRD104(2021)094012,.....

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Two-body interactions

- The observation of T_{cc}^+ , $D_{s0}^*(2317)$, $D_{s1}(2460)$ in experiment

- ♦ DD^* interaction: LO+OPE

Du et al, PRD105(2022)014024



- ♦ DK interaction: LO+NLO

Guo et al, EPJA40(2009)171



- ♦ D^*K interaction: LO+NLO



- ♦ Single particle regulator is used to obtain a better ren. Group invariant

Lu et al, arXiv:2308.14559

Two-body interactions

- The observation of T_{cc}^+ , $D_{s0}^*(2317)$, $D_{s1}(2460)$ in experiment

- DD^* interaction: LO+OPE

Du et al, PRD105(2022)014024

$$V_{DD^*}^{\text{Con}} = v_0 \epsilon \cdot \epsilon^* \quad V_{DD^*}^{\text{OPE}}(\mathbf{q}) = -\frac{3g^2}{4f_\pi^2} \frac{\epsilon \cdot \mathbf{q} \epsilon^* \cdot \mathbf{q}}{q^2 + \mu^2}$$

- DK interaction: LO+NLO

Guo et al, EPJA40(2009)171

$$V_{\text{LO}}^{DK}(p_i) = \frac{-1}{2f_\pi^2} (p_1 \cdot p_2 + p'_1 \cdot p'_2 + p_1 \cdot p'_2 + p_2 \cdot p'_1)$$

$$V_{\text{NLO}}^{DK}(p_i) = -\frac{8M_K^2}{3f_\pi^2} h_1 + \frac{4}{f_\pi^2} \left(h_3 p_2 \cdot p'_2 + h_5 (p_1 \cdot p_2 p'_1 \cdot p'_2 + p_1 \cdot p'_2 p_2 \cdot p'_1) \right)$$

- D^*K interaction: LO+NLO

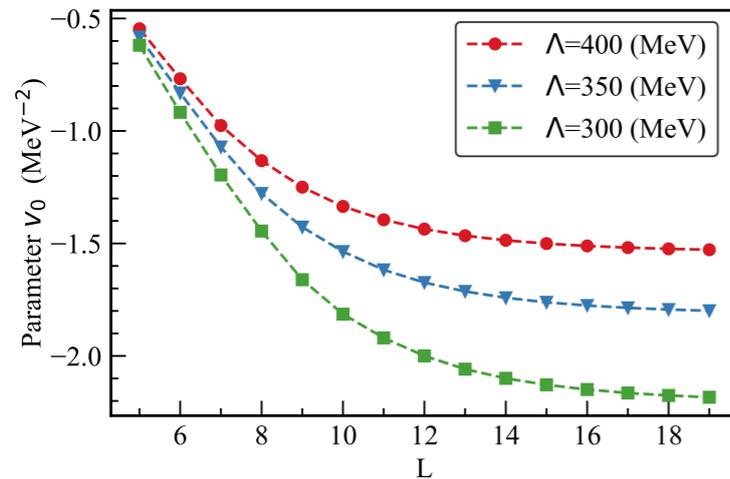
$$V_{\text{NLO}}^{D^*K}(p_i) = \left(-\frac{8M_K^2}{3f_\pi^2} h_1^* + \frac{4}{f_\pi^2} \left(h'_3 p_2 \cdot p'_2 + h_5^* (p_1 \cdot p_2 p'_1 \cdot p'_2 + p_1 \cdot p'_2 p_2 \cdot p'_1) \right) \right) \epsilon \cdot \epsilon^*$$

- Single particle regulator is used to obtain a better ren. Group invariant

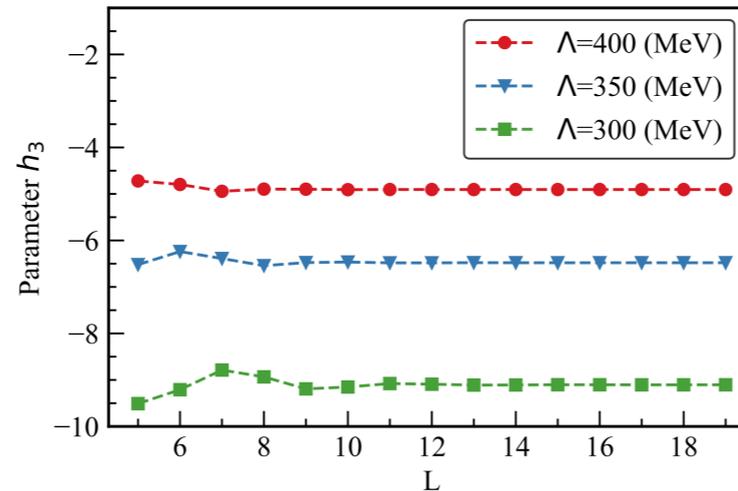
Lu et al, arXiv:2308.14559

The two-body parameters

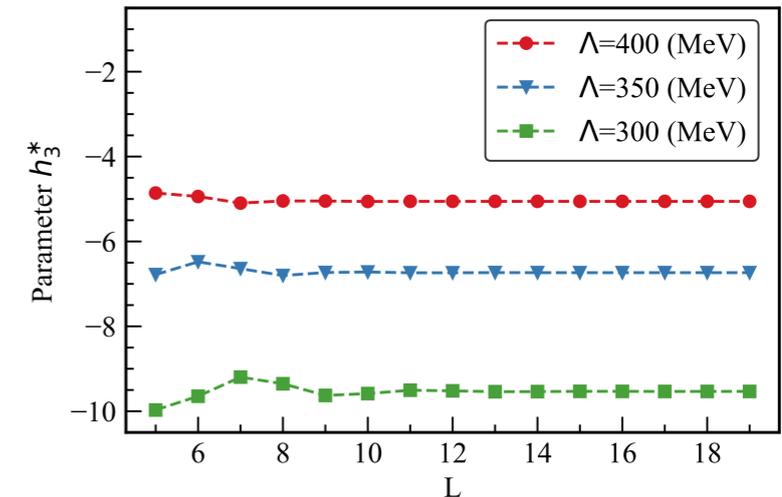
$$T_{cc}^+ \Rightarrow v_0$$



$$D_{s_0}^*(2317) \Rightarrow h_3$$



$$D_{s_1}(2460) \Rightarrow h_3^*$$



- ◆ Cubic lattice $L^3 = 5^3 \dots 19^3$
- ◆ Cutoff $\Lambda = 300, 350, 400$ MeV in the regulator
- ◆ Lattice spacing $a = 1/200$ MeV ~ 0.99 fm
- ◆ v_0 converges slow \Leftarrow long-ranged force
- ◆ $\Lambda = 400$ MeV converges quickly

The three-body interactions

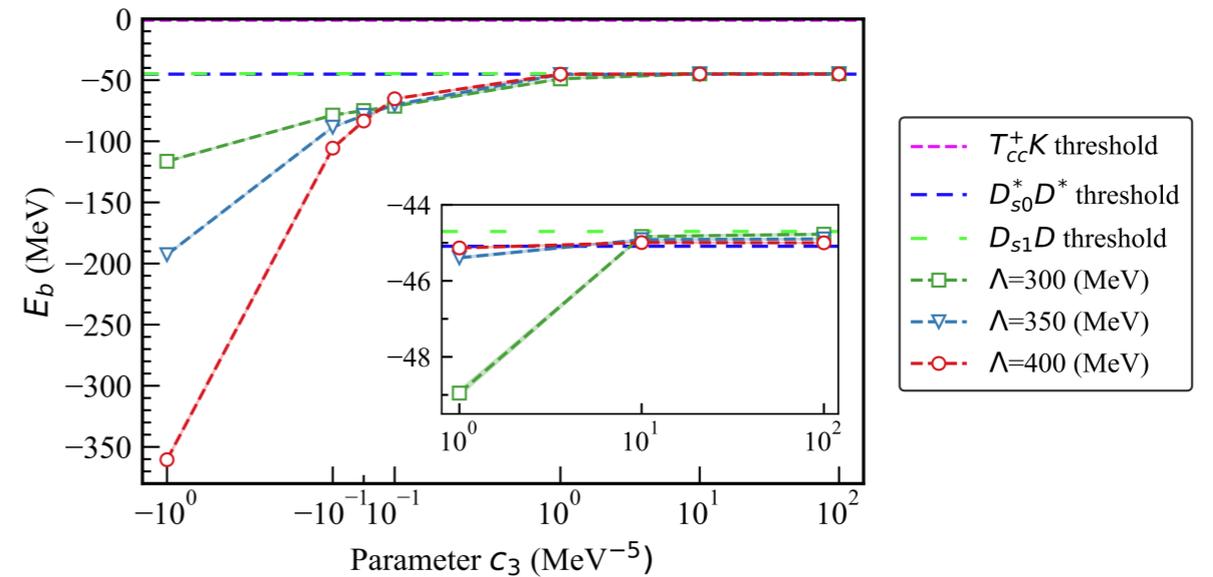
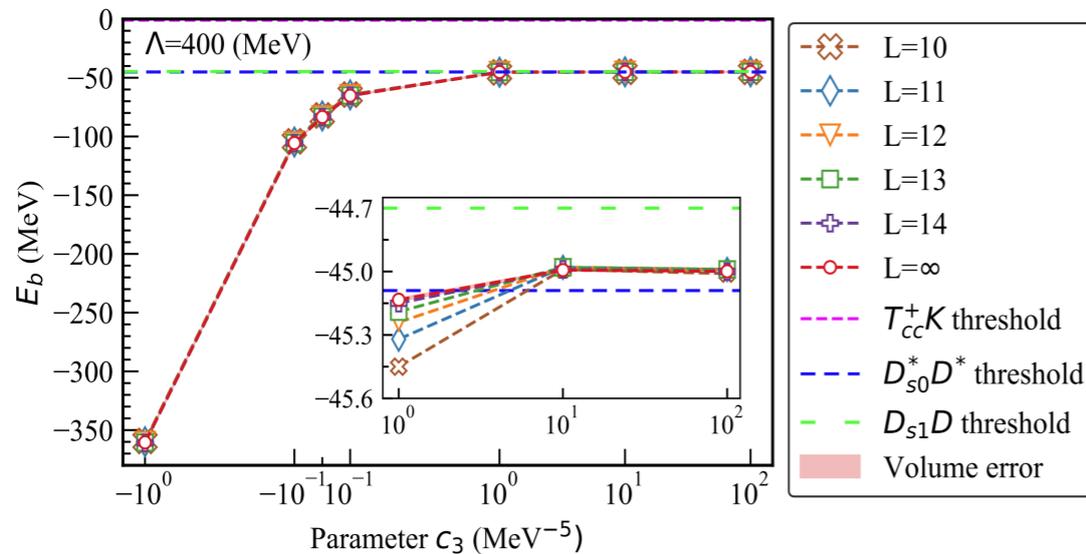
- DD^*K Lag. $\mathcal{L} = c_3 \left\langle H \mathcal{D}_\mu H^\dagger H \mathcal{D}^\mu H^\dagger \right\rangle + c'_3 \left\langle H \mathcal{A}_\mu H^\dagger H \mathcal{A}^\mu H^\dagger \right\rangle$

- DD^*K three-body force

$$V_{DD^*K}(p_i) = \frac{c_3}{4f_\pi^2} (p_1 \cdot p_3 + p_1 \cdot p'_3 + p_2 \cdot p_3 + p_2 \cdot p'_3 + p'_1 \cdot p_3 + p'_1 \cdot p'_3 + p'_2 \cdot p_3 + p'_2 \cdot p'_3) \epsilon \cdot \epsilon^*$$

- DD^*K binding energy

Zhang et al., arXiv:2409.01325



- Extrapolate to infinite volume

Meng et al., PRD98(2018)014508

$$\frac{\Delta E}{E_T} = -(\kappa L)^{-3/2} \sum_{i=1}^3 C_i \exp(-\mu_i \kappa L)$$

- Switch off three-body force, the result is consistent with that in

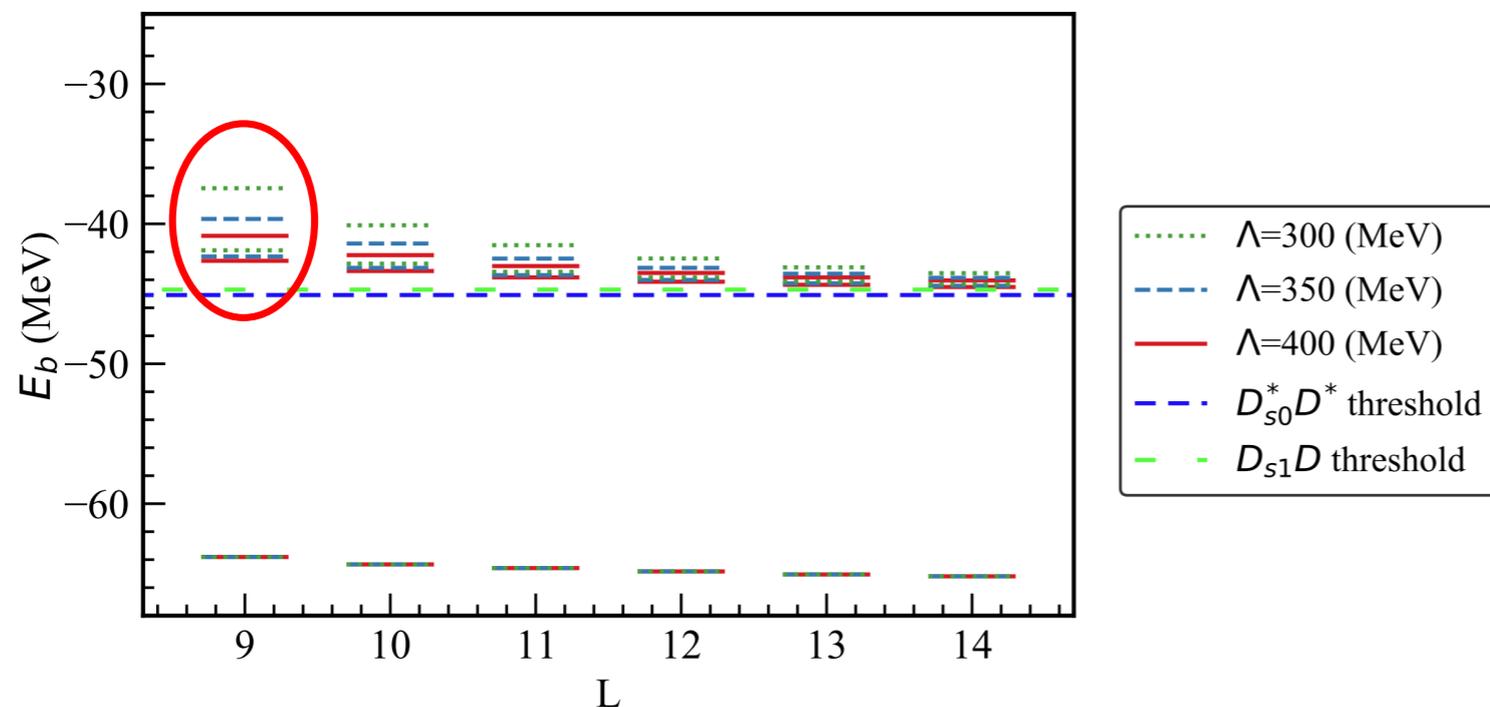
Ma et al., CPC43(2019)014102

The first excited state

- No experimental data \Rightarrow binding energy with $\Lambda = 400$ MeV as input
- The parameter c_3 at various cubic

Λ (MeV)	Parameter	L						State
		9	10	11	12	13	14	
400	c_3 (MeV ⁻⁵)	0.100	0.100	0.100	0.100	0.100	0.100	Input
350		0.170	0.162	0.164	0.164	0.163	0.163	Fitted
300		0.328	0.305	0.281	0.278	0.281	0.280	Fitted

Zhang et al., arXiv:2409.01325



- $E_{\Lambda_1}^{\text{excited}} - E_{\Lambda_2}^{\text{excited}}$ decreases
- ρ -type and λ -type excitation
- The standard angular momentum and parity projection technique is used

• $J^P = 1^- \quad \Leftarrow$

$$|\Psi_A\rangle = \frac{d_n}{24} \sum_{i=1}^{24} \chi_n(\Omega_i) R(\Omega_i) |\Psi_0\rangle$$

Lu et al., PRD90(2014)034507

Summary and outlook

- The number of particles is conserved
- Find a $J^P = 1^- DD^*K$ bound state with binding energy in the $(-84, -44)$ MeV region
- We checked the ren. group invariance
- Two energy levels of excited state are ρ -type and λ -type
- Lattice EFT has more advantage for many hadrons system

Thank you very much for your attention!