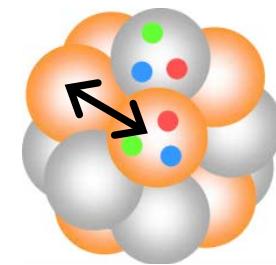
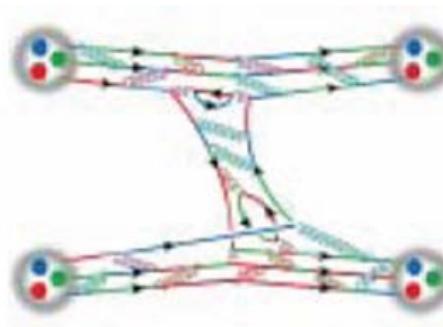
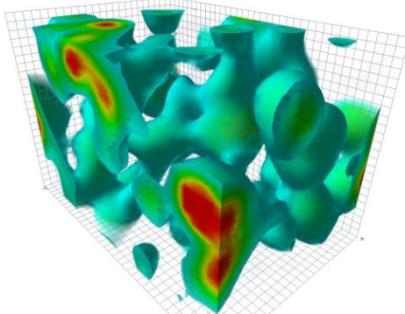
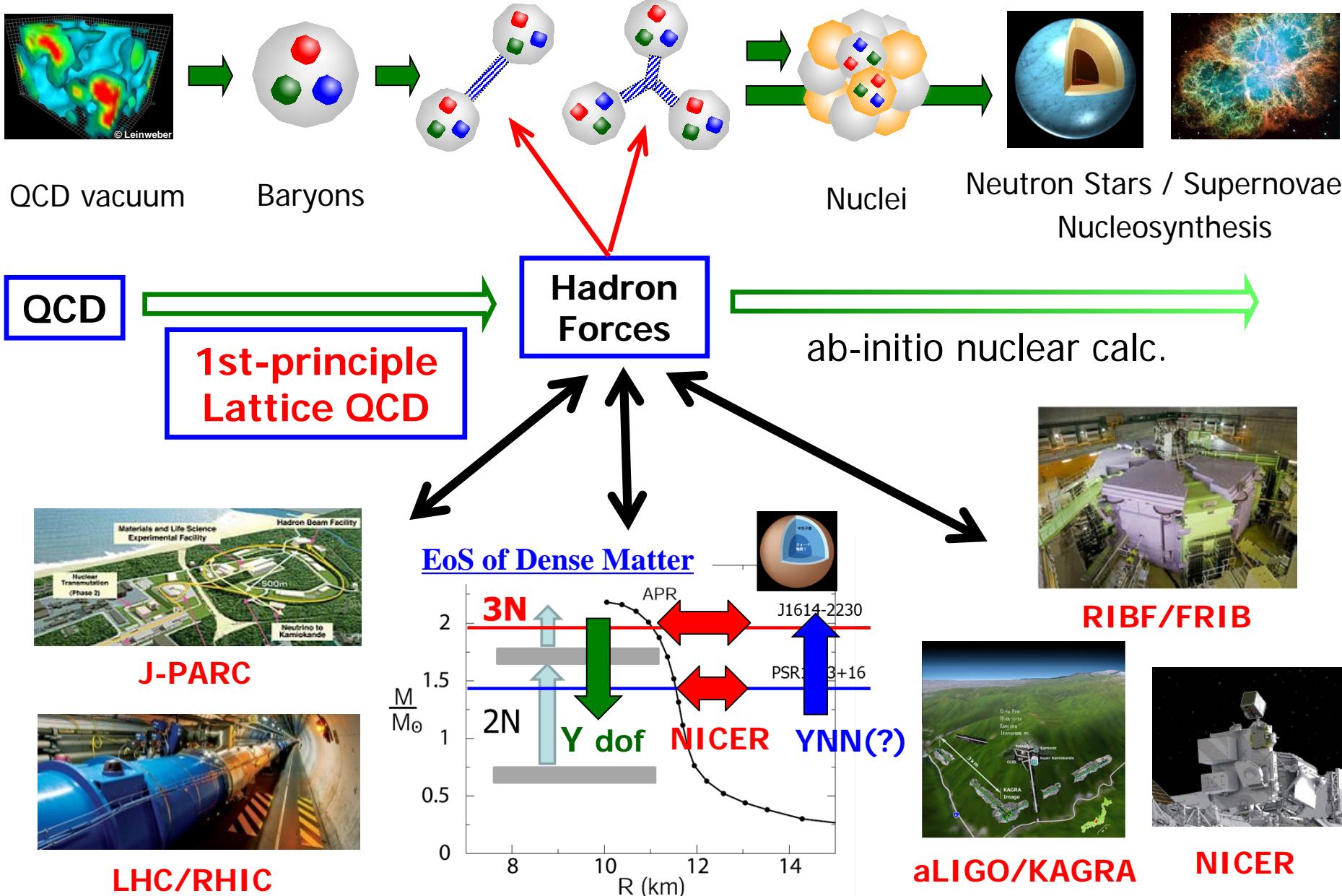


Lattice QCD studies of Hadron interactions from the HAL QCD method

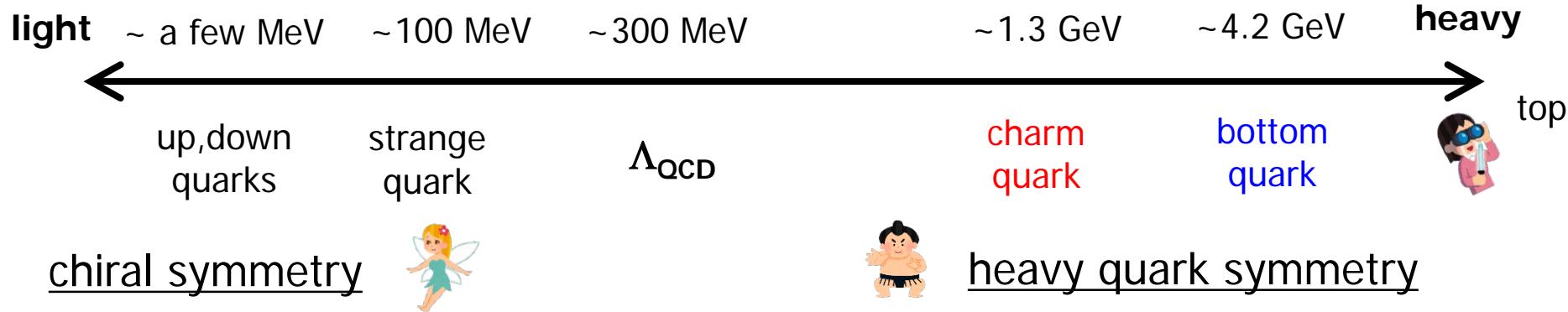
Takumi Doi
(RIKEN iTHEMS)



The Odyssey from Quarks to Universe

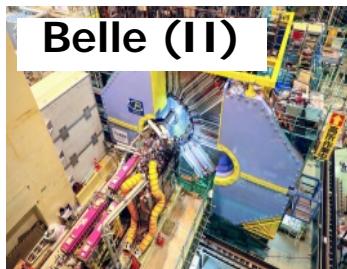
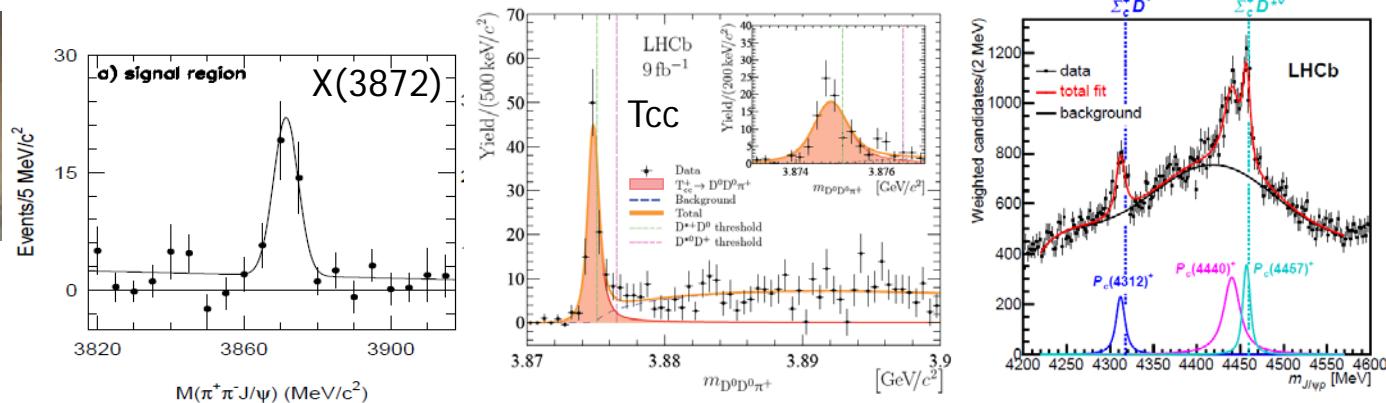
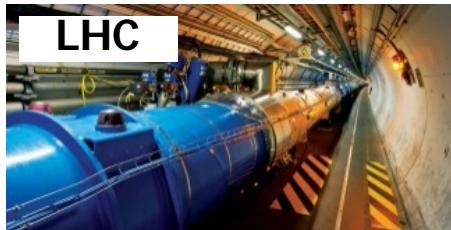


Nuclear/Hyperon Forces → Charmed/Bottomed Forces



Heavy quarks: New doorway to the mysteries of QCD

Many new exotic particles being reported!



Hadron interactions crucial to understand these "signals" !

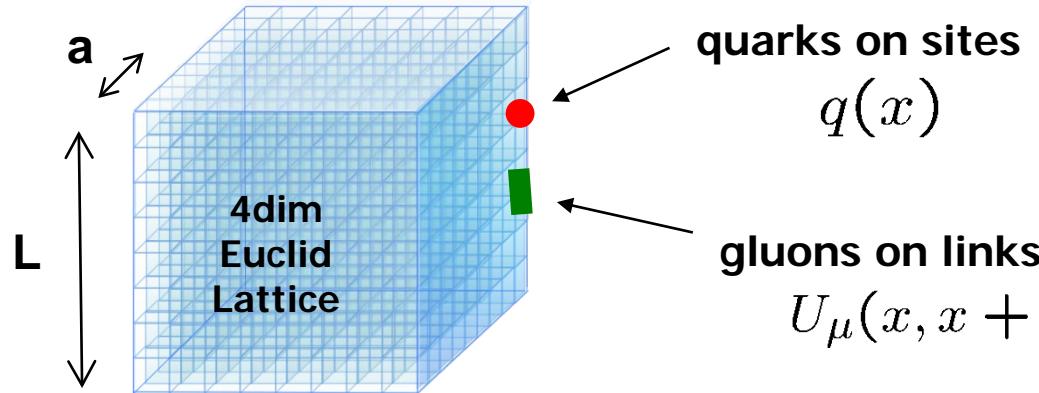
e.g., Zc(3900) from HAL LQCD → threshold cusp

Y. Ikeda et al. (HAL), PRL117(2016)242001

Lattice QCD

First-principle calculation of QCD

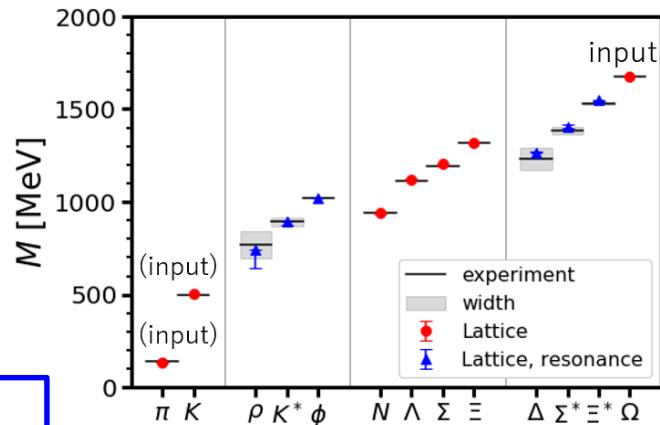
$$Z = \int dU d\bar{q} d\bar{\bar{q}} e^{-S_E}$$



K.G. Wilson
(1974)

- Regularized system (finite a and L)
- Gauge-invariance manifest
- Fully-Nonperturbative
- DoF $\sim 10^{9-10} \rightarrow$ Monte-Carlo w/ Euclid time
 - Numerical calc by supercomputers

Single hadron spectrum well reproduced
→ Next big challenge: **Interactions between hadrons**



[HAL QCD method]

- “Potential” defined through phase shifts (S-matrix)
- Nambu-Bethe-Salpeter (NBS) wave function

$$\psi(\vec{r}) = \langle 0 | H_1(\vec{x} + \vec{r}) H_2(\vec{x}) | H_1(k) H_2(-k); W \rangle$$

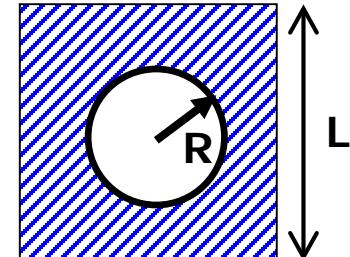
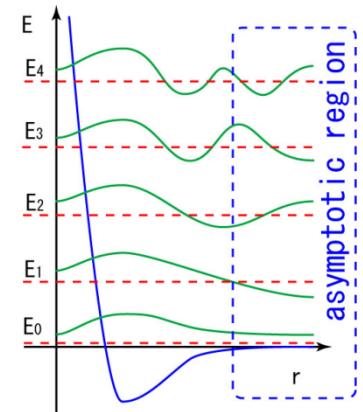
$$(\nabla^2 + k^2)\psi(\vec{r}) = 0, \quad r > R \quad W = 2\sqrt{m^2 + k^2}$$

– Wave function \leftrightarrow phase shifts

$$\psi(r) \simeq A \frac{\sin(kr - l\pi/2 + \delta(k))}{kr}$$

(below inelastic threshold)

Extended to multi-particle systems



M.Luscher, NPB354(1991)531

Ishizuka, PoS LAT2009 (2009) 119

C.-J.Lin et al., NPB619(2001)467

Aoki-Hatsuda-Ishii PTP123(2010)89

CP-PACS Coll., PRD71(2005)094504

S.Aoki et al., PRD88(2013)014036

“Potential” as a representation of S-matrix

- Consider the wave function at “interacting region”

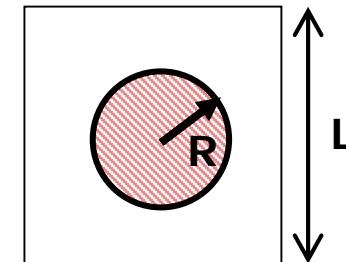
$$(\nabla^2 + k^2)\psi(\mathbf{r}) = m \int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \psi(\mathbf{r}'), \quad r < R$$

- $\mathbf{U}(\mathbf{r}, \mathbf{r}')$: faithful to the phase shift by construction
 - $\mathbf{U}(\mathbf{r}, \mathbf{r}')$: NOT an observable, but well defined
 - $\mathbf{U}(\mathbf{r}, \mathbf{r}')$: E-independent, while non-local in general
 - “Proof of Existence”: Explicit form can be given as

$$\mathbf{U}(\mathbf{r}, \mathbf{r}') = \frac{1}{m} \sum_{n,n'}^{n_{\text{th}}} (\nabla_{\mathbf{r}}^2 + k_n^2) \psi_n(\mathbf{r}) \mathcal{N}_{nn'}^{-1} \psi_{n'}^*(\mathbf{r}') \quad \mathcal{N}_{nn'} = \int d\mathbf{r} \psi_n^*(\mathbf{r}) \psi_{n'}(\mathbf{r})$$

- Non-locality → derivative expansion Okubo-Marshak(1958)

$$\mathbf{U}(\vec{r}, \vec{r}') = \begin{matrix} [V_c(r) + S_{12}V_T(r) + \vec{L} \cdot \vec{S} V_{LS}(r) + \mathcal{O}(\nabla^2)] \delta(\vec{r} - \vec{r}') \\ \text{LO} \qquad \text{LO} \qquad \text{NLO} \qquad \text{NNLO} \end{matrix}$$



Time-dependent HAL method

N.Ishii et al. (HAL QCD Coll.) PLB712(2012)437

E-indep of potential $U(r,r')$ \rightarrow (excited) scatt states share the same $U(r,r')$
They are not contaminations, but signals

Original (t-indep) HAL method

$$G_{NN}(\vec{r}, t) = \langle 0 | N(\vec{r}, t) N(\vec{0}, t) \overline{\mathcal{J}_{\text{src}}(t_0)} | 0 \rangle$$

$$R(\mathbf{r}, t) \equiv G_{NN}(\mathbf{r}, t)/G_N(t)^2 = \sum A_{W_i} \psi_{W_i}(\mathbf{r}) e^{-(W_i - 2m)t}$$

$$\int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \underline{\psi_{W_0}(\mathbf{r}')} = (\underline{E_{W_0}} - H_0) \underline{\psi_{W_0}(\mathbf{r})}$$

$$\int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \underline{\psi_{W_1}(\mathbf{r}')} = (\underline{E_{W_1}} - H_0) \underline{\psi_{W_1}(\mathbf{r})}$$

...

← Many states contribute

New t-dep HAL method

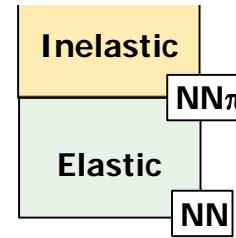
All equations can be combined as

$$\int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \underline{R(\mathbf{r}', t)} = \left(-\frac{\partial}{\partial t} + \frac{1}{4m} \frac{\partial^2}{\partial t^2} - H_0 \right) \underline{R(\mathbf{r}, t)}$$

~~G.S. saturation \rightarrow "Elastic state" saturation~~

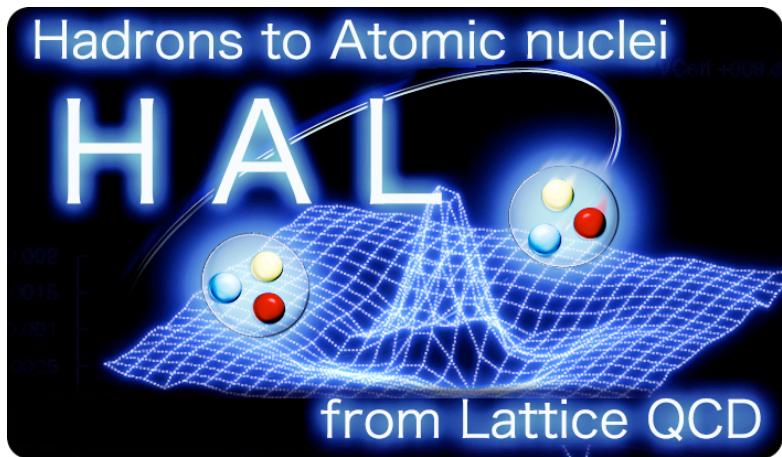
[Exponential Improvement]

Coupled Channel formalism
↔ above inelastic threshold



potential

Hadrons to **A**tomic nuclei from **L**attice QCD (**HAL** QCD Collaboration)



S. Aoki, T. M. Doi, E. Ito (Kyoto U.)
T. Aoyama (ISSP)
T. Doi, T. Hatsuda, Y. Lyu, W. A. Yamada,
L. Wang, L. Zhang (RIKEN)
F. Etminan (U. of Birjand)
Y. Ikeda, N. Ishii, P. Junnarkar, H. Nemura,
K. Sasaki (Osaka U.)
T. Inoue (Nihon U.)
K. Murakami (Science Tokyo)
K. Murase (Tokyo Metropolitan U.)
T. Sugiura (Rissho U.)
H. Tong (U. of Bonn)

「20XX年宇宙の旅」
from Quarks to Universe



+

I. Kanamori (RIKEN)
K.-I. Ishikawa (Hiroshima U.)
and many more

Challenge toward LQCD w/ physical mass

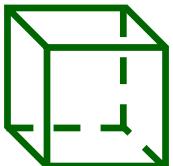
~2012



→ lighter u,d-quark masses
(=lighter pion mass M_π)

We were here

$M_\pi = 400\text{MeV}$
 $L = 3\text{fm}$



Simulation w/
Unrealistic mass

~2018



→ more challenging

Phys. point

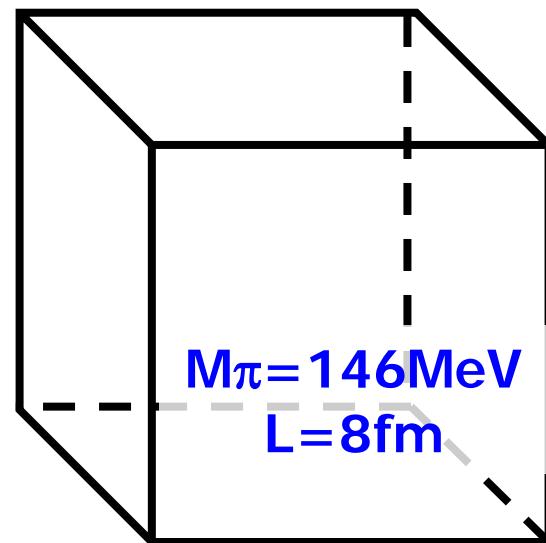
Theoretical
development

+

New supercomputer



K-computer (11PFlops)



Simulation w/
~Near physical mass

Challenge toward LQCD w/ physical mass

~2012



→ lighter u,d-quark masses
(=lighter pion mass M_π)

We were here

$M_\pi = 400\text{MeV}$
 $L = 3\text{fm}$



Simulation w/
Unrealistic mass

~2018



→ more challenging

Phys. point

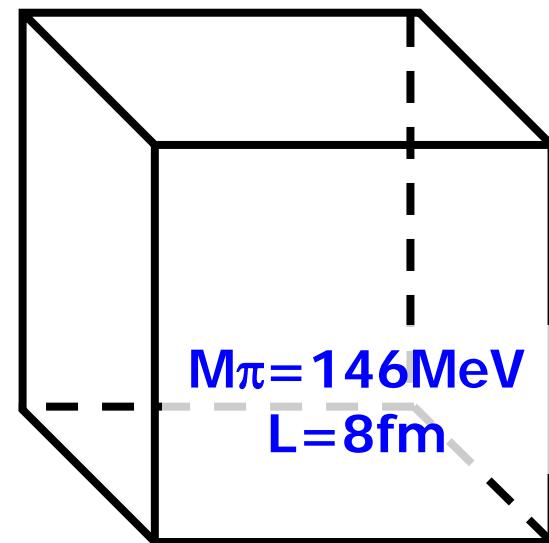
Theoretical
development

+

New supercomputer



Fugaku (440PFlops)



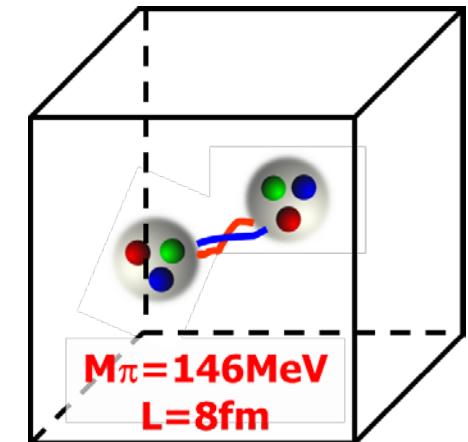
Simulation w/
~Near physical mass

LQCD calc near the physical point

- **Nf = 2 + 1 gauge configs**

- clover fermion + Iwasaki gauge w/ stout smearing
- $V=(8.1\text{fm})^4$, $a=0.085\text{fm}$ ($1/a = 2.3 \text{ GeV}$)
- **$m_\pi = 146 \text{ MeV}, m_K = 525 \text{ MeV}$**
- #traj $\sim= 2000$ generated
- (Quenched) Charm quark w/ RHQ action

PACS Coll., PoS LAT2015, 075



Y. Namekawa (PACS), PoS LAT2016, 125

- Nuclear/**Hyperon forces** + Charmed forces to be predicted

- Central/tensor forces in S, D-waves

S=0	S=-1	S=-2	S=-3	S=-4	S=-5	S=-6
NN	NΛ, NΣ	ΛΛ, ΛΣ, ΣΣ, NΞ	ΛΞ, ΣΞ, NΩ	ΞΞ	ΞΩ	ΩΩ



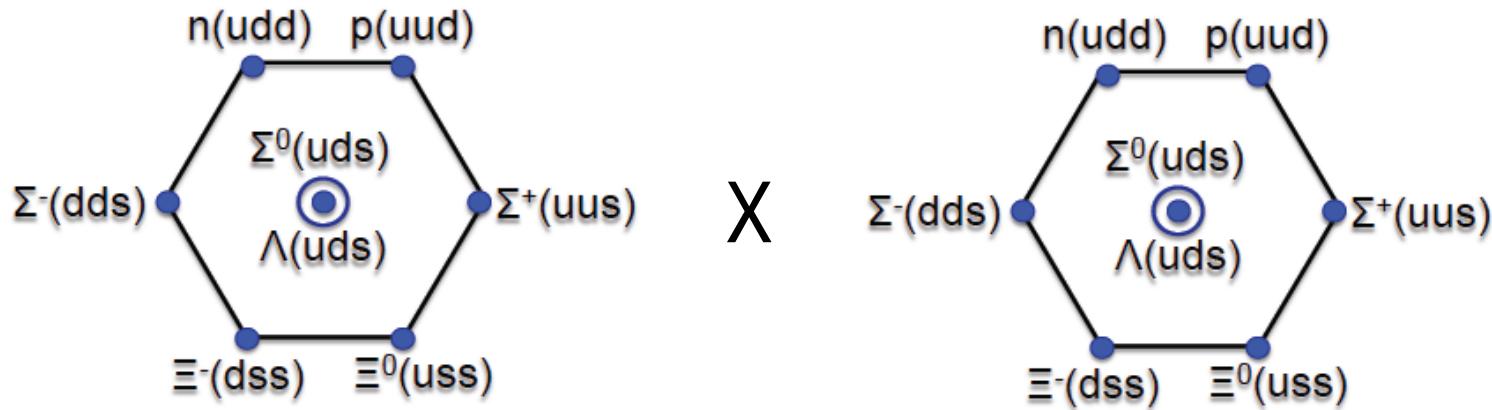
EXP
rich data

LQCD
better S/N

Baryon-Baryon Interactions

Birds-eye View

classification w/ flavor SU(3)-irrep base



$$8 \times 8 = \underbrace{27}_{\text{symmetric}} + \underbrace{8s + 1}_{\text{anti-symmetric}} + \underbrace{10^* + 10 + 8a}_{\text{NN channel}}$$

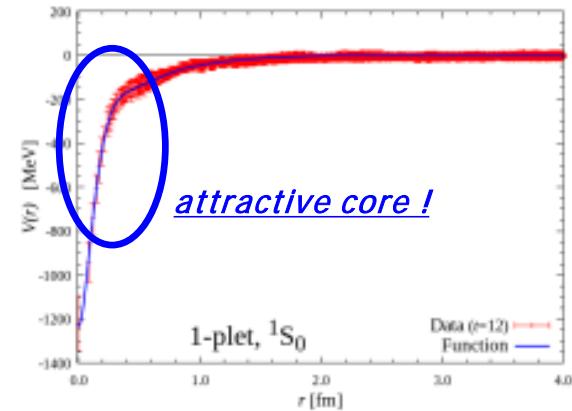
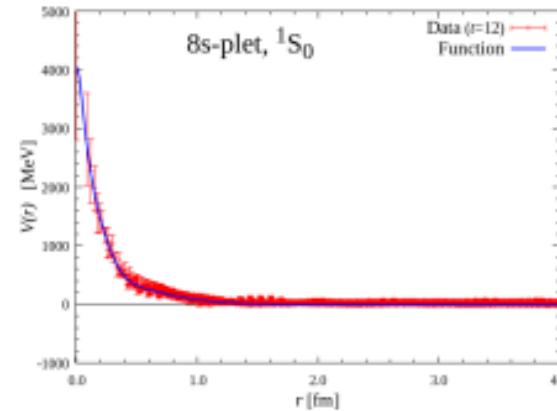
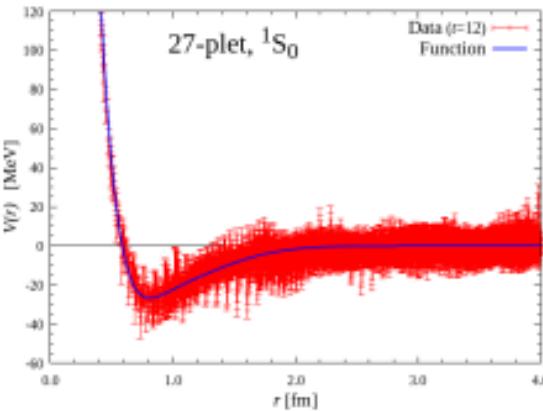
c.f. Exact SU(3) limit LQCD calc @ heavy masses

Diagonal Potentials in SU(3)f-irrep base in S=-2

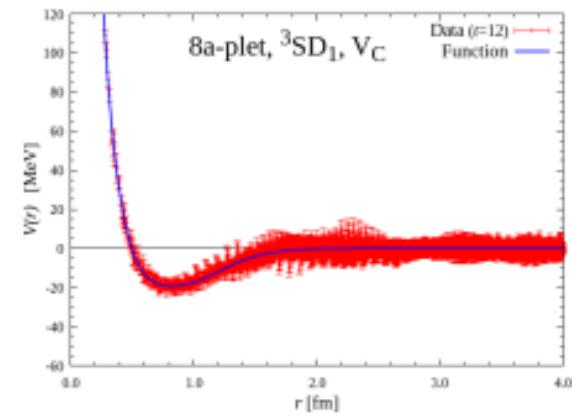
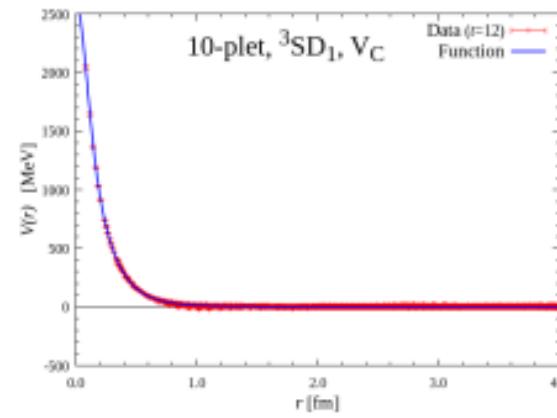
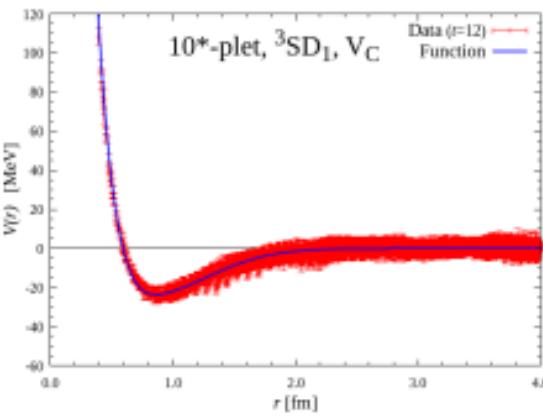
(only) S= -2 can access all irreps
off-diag pot relatively small

T.Inoue (HAL), AIP Conf. Proc. 2130 (2019) 020002

$1S_0$



$3S_1 - 3D_1$



27,10*:
NN-type

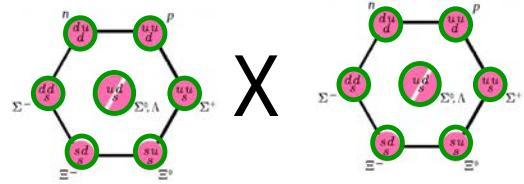
8s,10:
strong repulsive core

1s: deep attractive pocket
8a: weak repulsive core

Quark Pauli repulsion + OGE for short range

M.Oka et al., NPA464(1987)700

Candidates of di-baryons



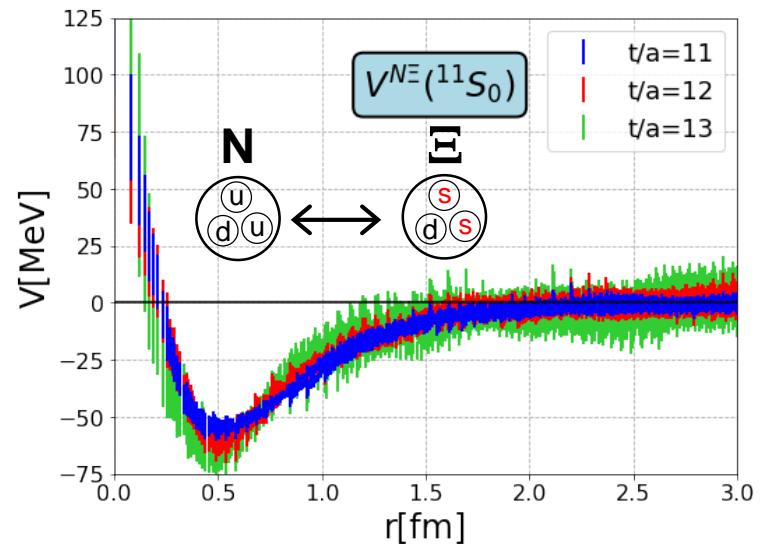
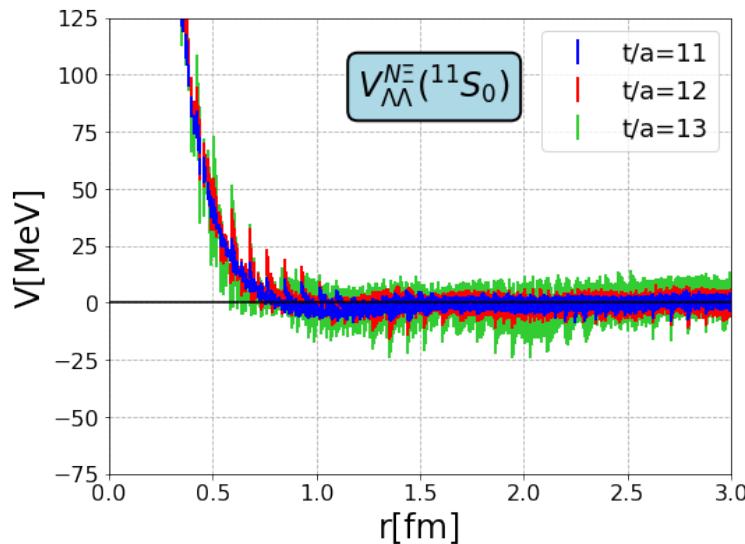
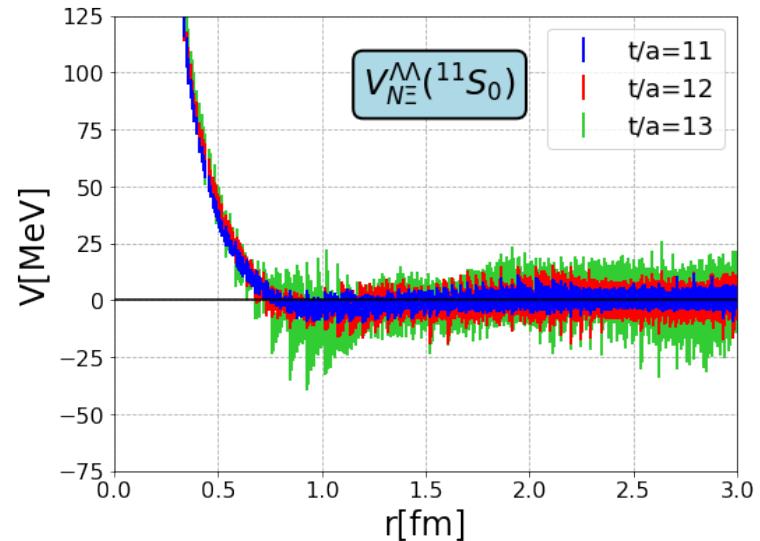
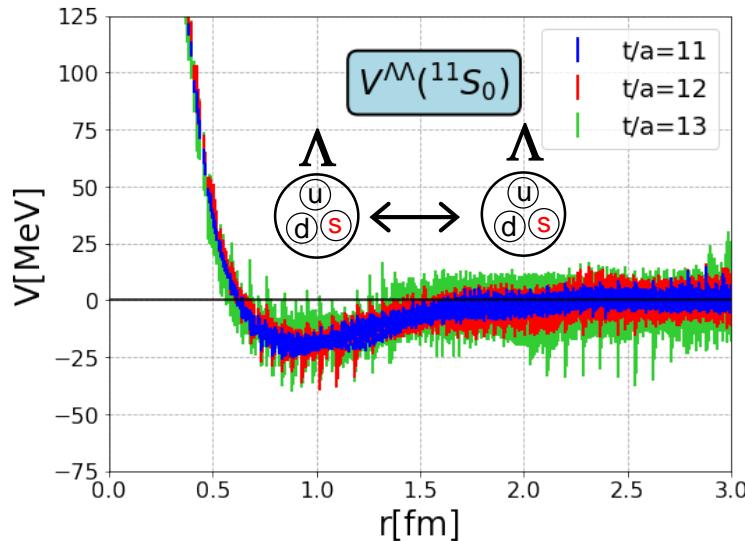
$$8 \times 8 = 27 + 8s + 1 + 10^* + 10 + 8a$$

dineutron, $\Xi\Xi$ etc. H-dibaryon Deuteron
 $(J=0)$ $(J=0)$ $(J=1)$

There may also exist $S = -2$ hypernuclei relevant to these strong attractions

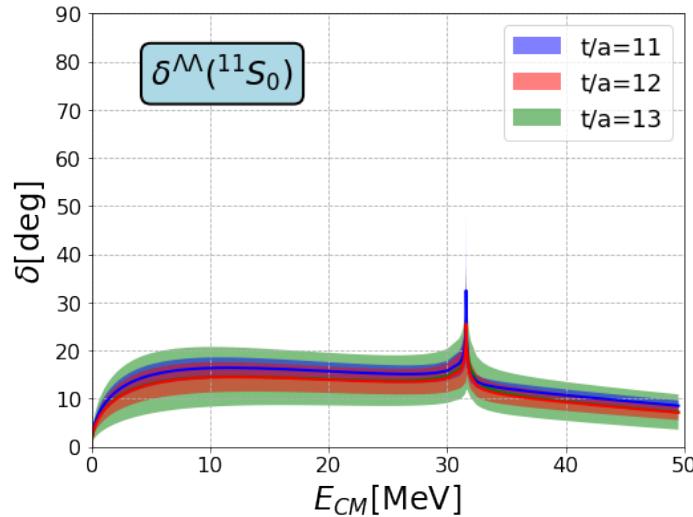
→ Detailed study w/ SU(3) breaking effects
(particle-base)

$\Lambda\Lambda$, $N\Xi$ (effective) 2x2 coupled channel analysis



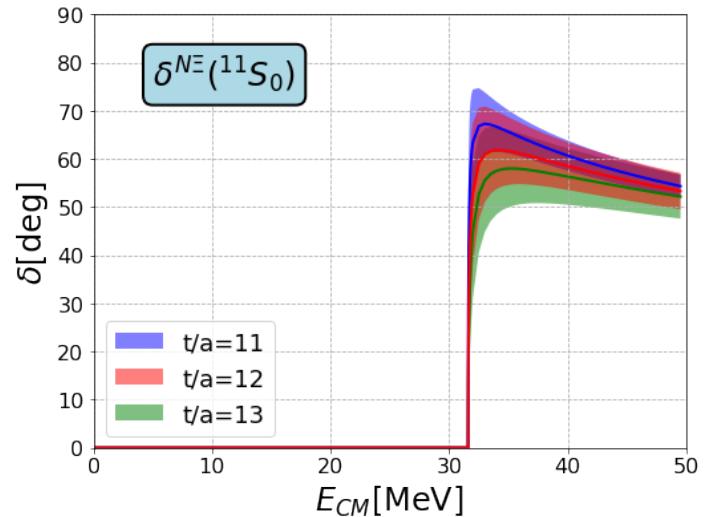
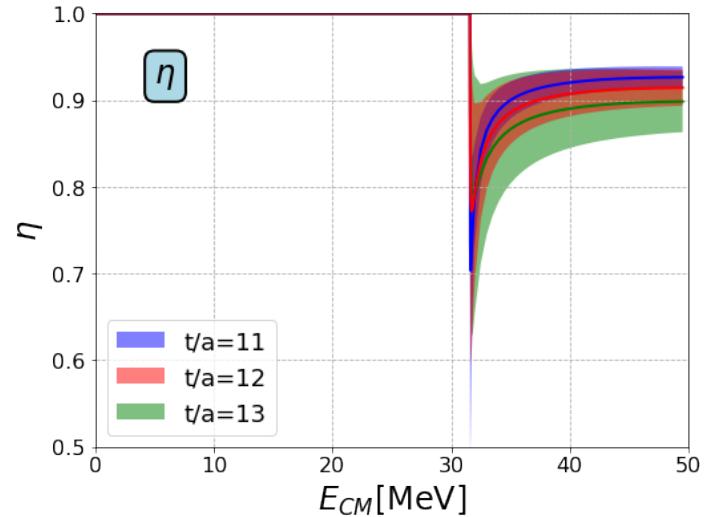
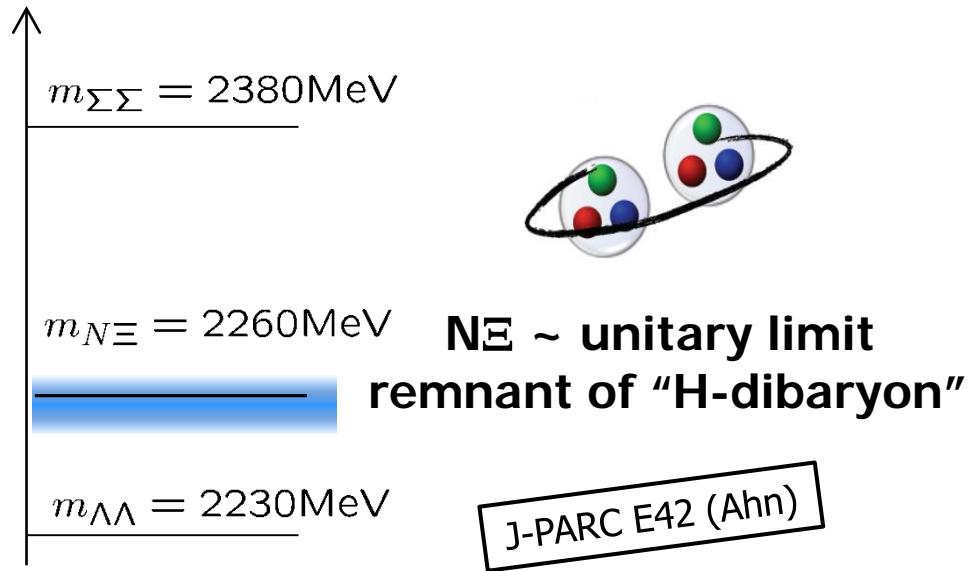
$N\Xi$ (1S_0) channel is attractive
 $N\Xi-\Lambda\Lambda$ coupling is small

$\Lambda\Lambda$, $N\Xi$ (effective) 2x2 coupled channel analysis



$$a_0 = -0.81(23)(+0.00/-0.13) \text{ [fm]}$$

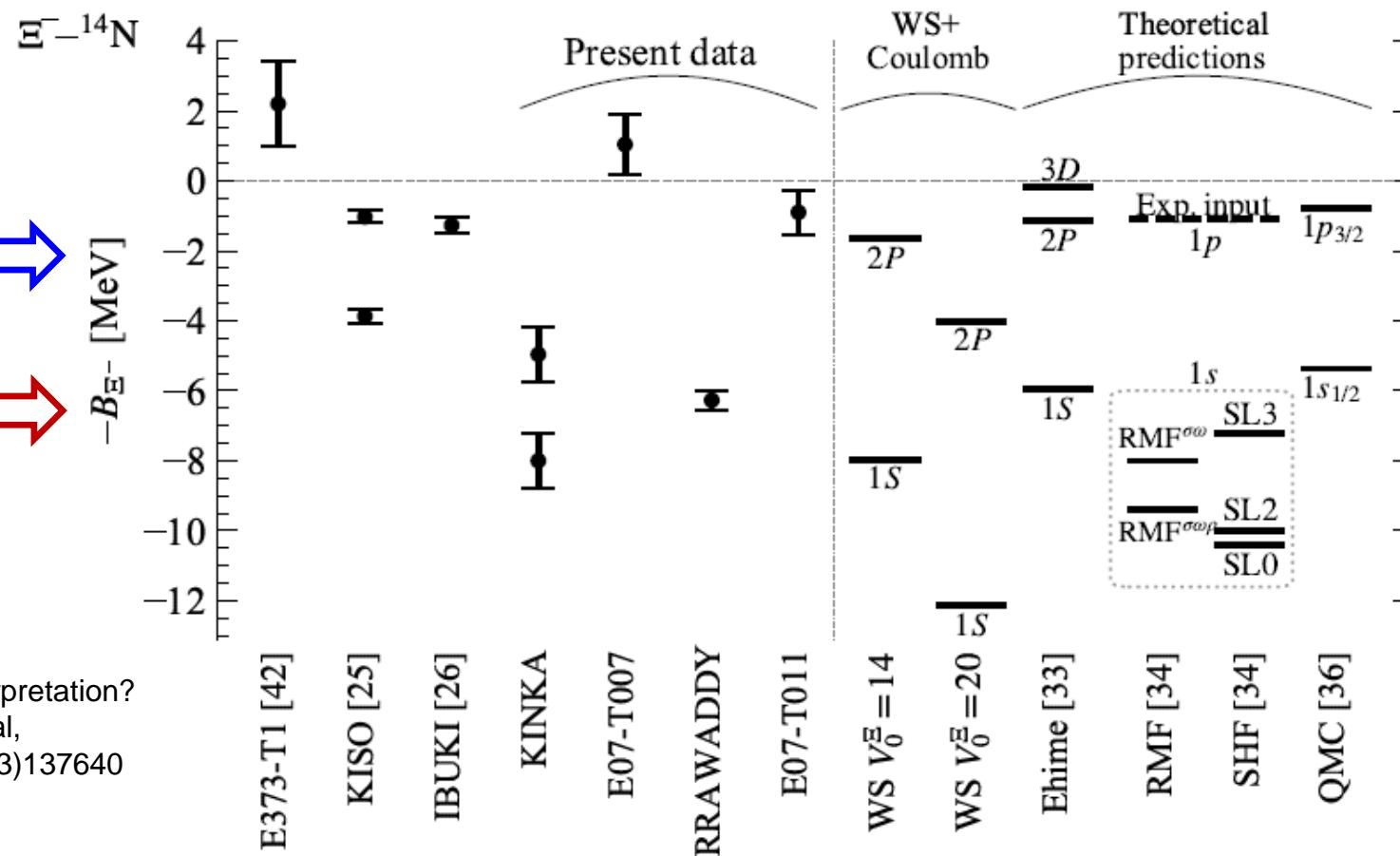
$$r_{\text{eff}} = 5.47(78)(+0.09/-0.55) \text{ [fm]}$$



(N.B. $N\Xi = 1\text{rep } 50\%, 27\text{rep } 30\% \text{ in SU}(3)$)

Recent experimental progress on Ξ -Hypernuclei

Excited state? 
Ground state? 



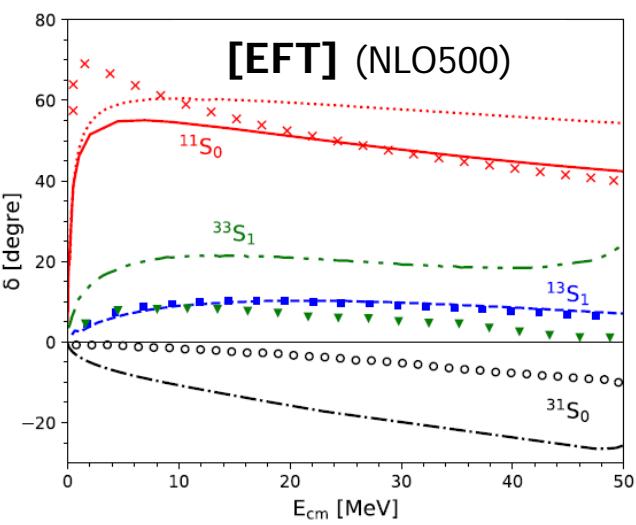
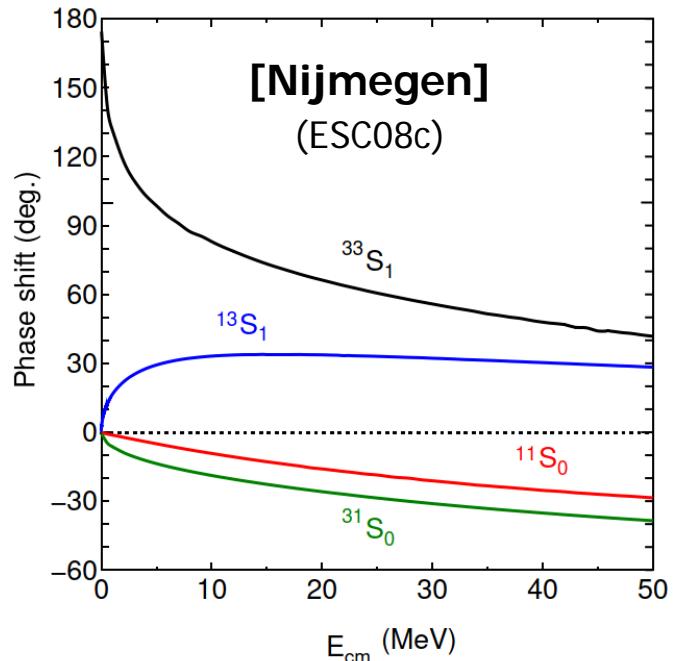
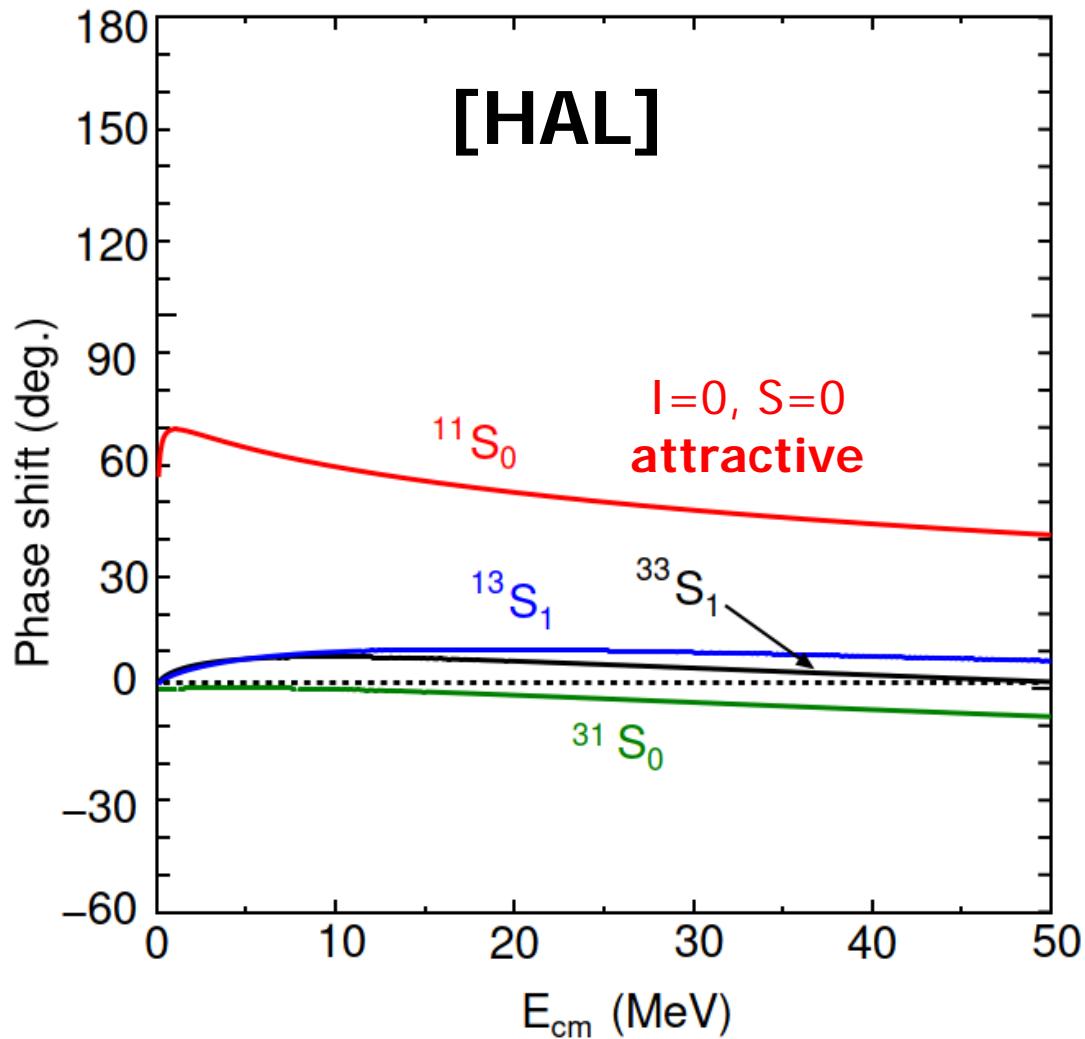
Another interpretation?
Friedman-Gal,
PLB837(2023)137640

J-PARC E07, E70,
more in HEF-EX

Attractive $N\Xi$ -int well established
Small $N\Xi-\Lambda\Lambda$ coupling indicated

M. Yoshimoto et al.,
PTEP2021, 073D02

NE scattering phase shifts and spin-isospin dep



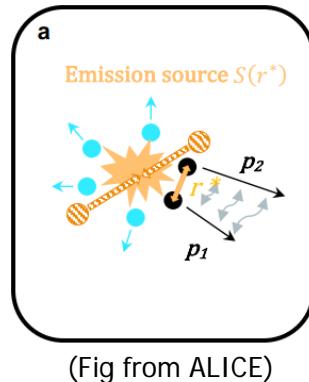
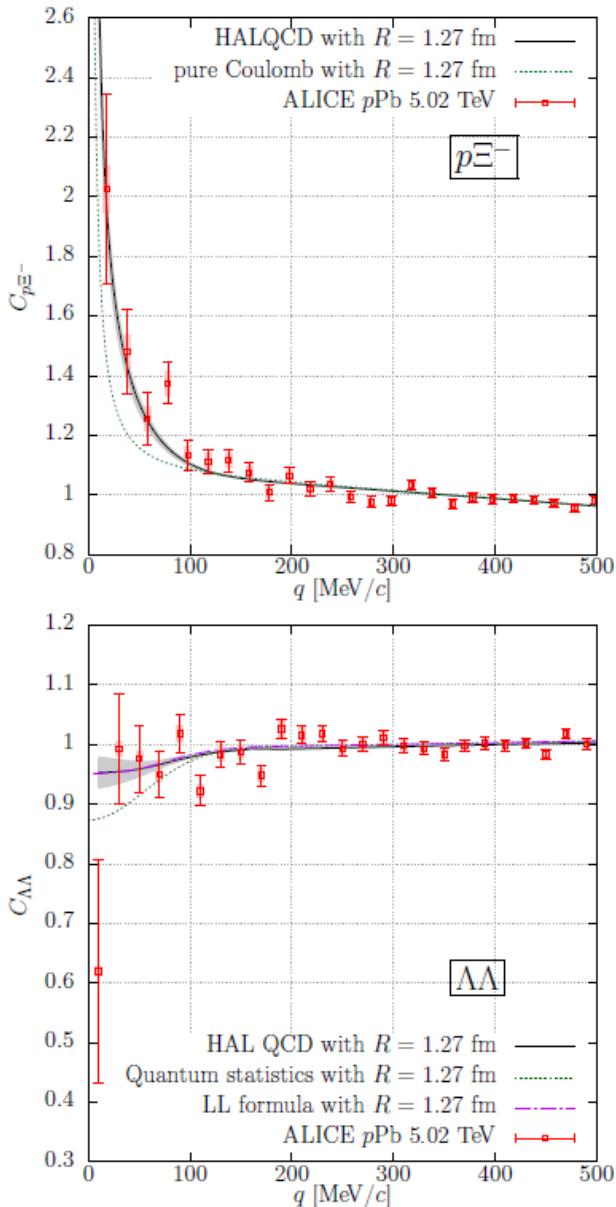
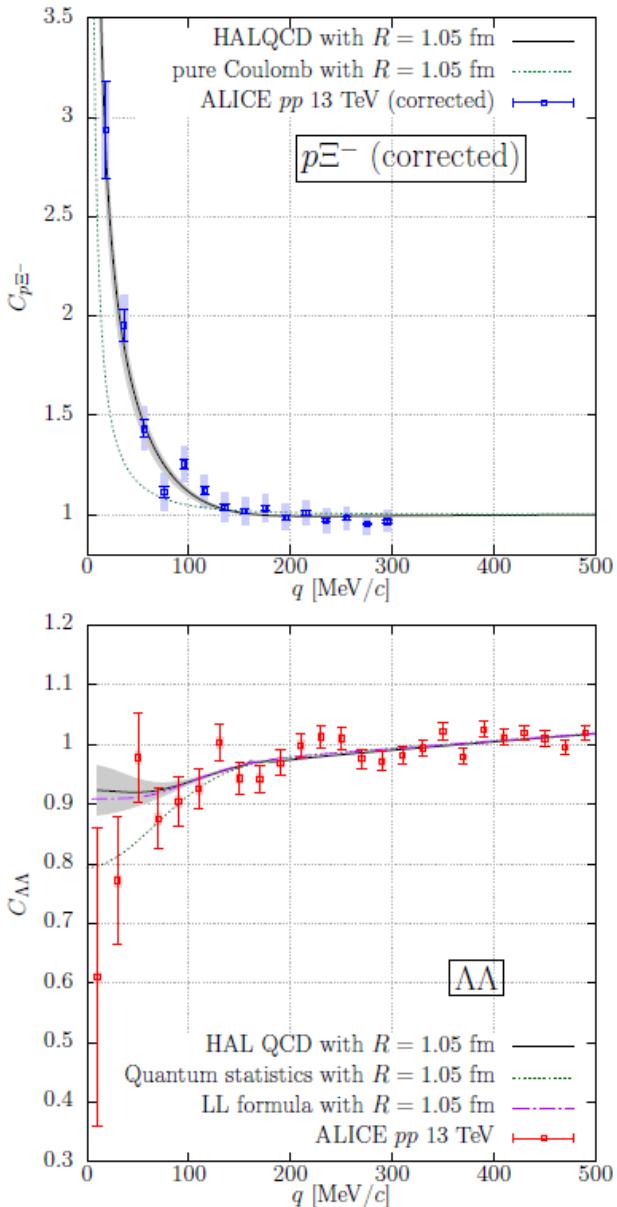
K. Sasaki et al., (HAL Coll.), NPA998(2020)12137

E. Hiyama et al., PRL124(2020)092501

H. Le et al., EPJA57(2021)12

EFT (lines), HAL (points)

Femtoscopy from nucleus collisions



$\overleftarrow{p\Xi^-}$

LQCD prediction confirmed
by experiment!

$\overleftarrow{\Lambda\Lambda}$

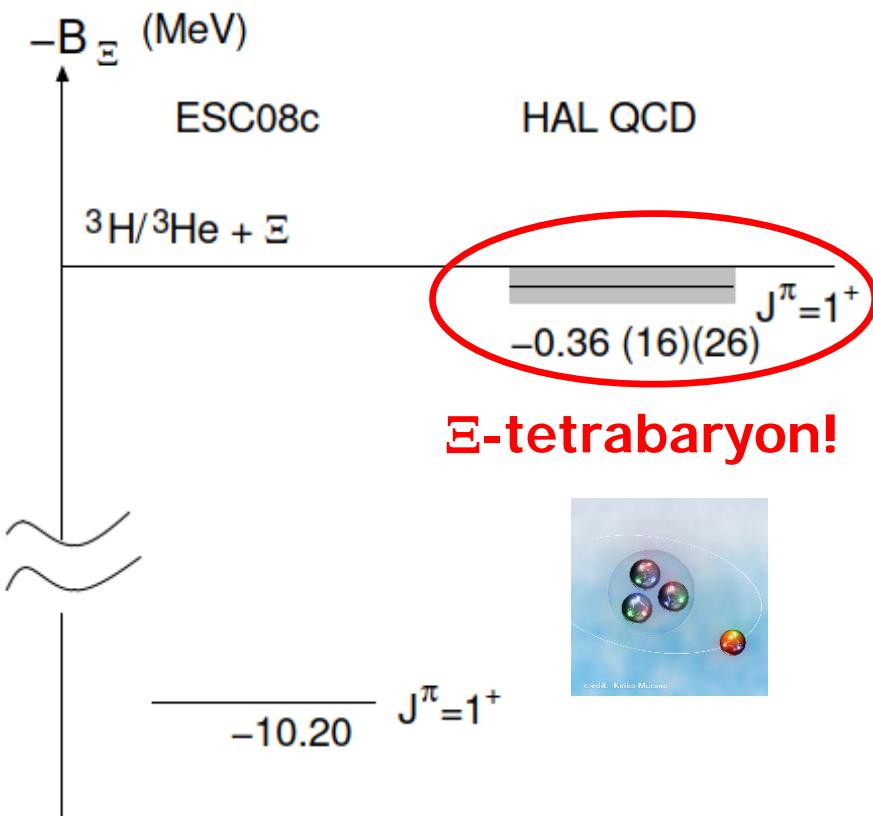
Y. Kamiya et al.,
PRC105(2022)014915

See also ALICE Coll., PLB797(2019)134822,
PRL123(2019)112002, Nature 588(2020)232

S= -2 Light Hypernuclei from LQCD

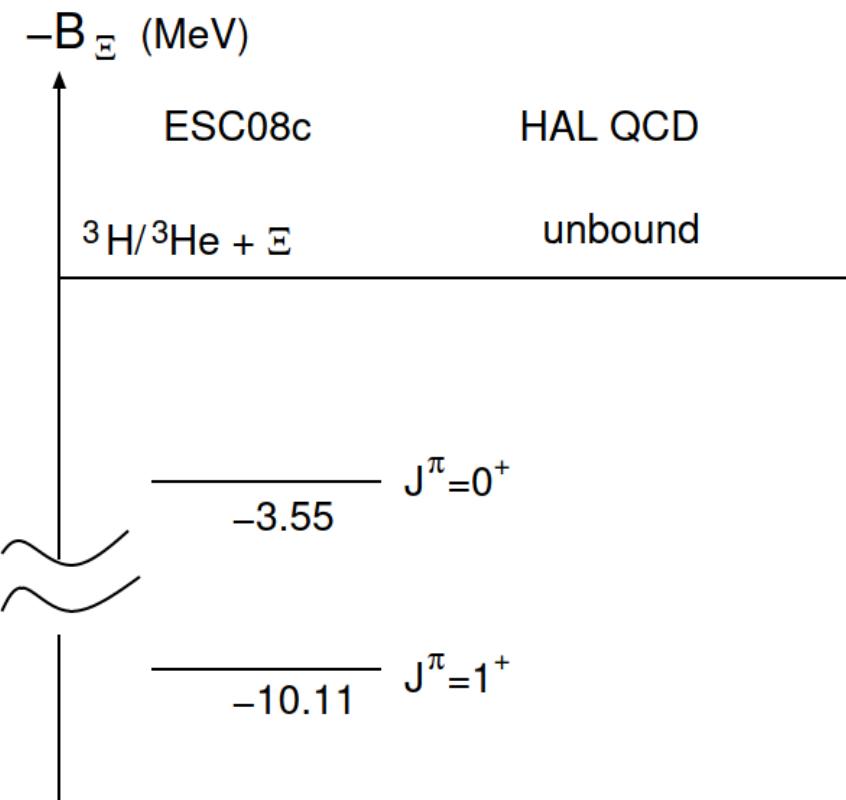
E. Hiyama et al., PRL124(2020)092501

a) NNN Ξ (T=0)



LQCD pot after
chiral extrapolation used
(via physical values of m_π , m_K)

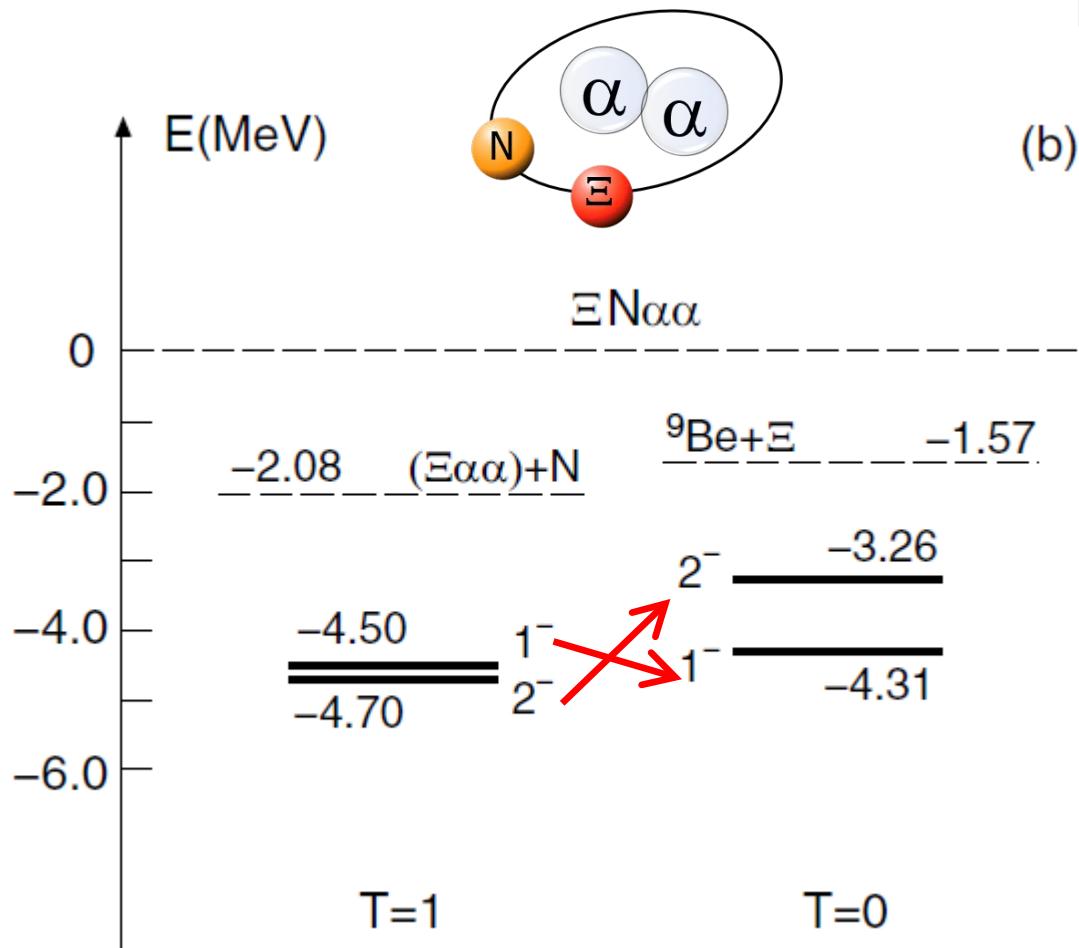
b) NNN Ξ (T=1)



c.f. EFT \rightarrow bound states
(B.E. = 0.6-0.7 MeV)
in T=1, $J^\pi = 1^+, 0^+$
due to stronger ${}^{33}\text{S}_1$

S= -2 ($\Xi N\alpha\alpha$) Hypernuclei from LQCD

E. Hiyama et al., PRC106(2022)064318



$A=10 (\Xi N\alpha\alpha)$ is found to be bound

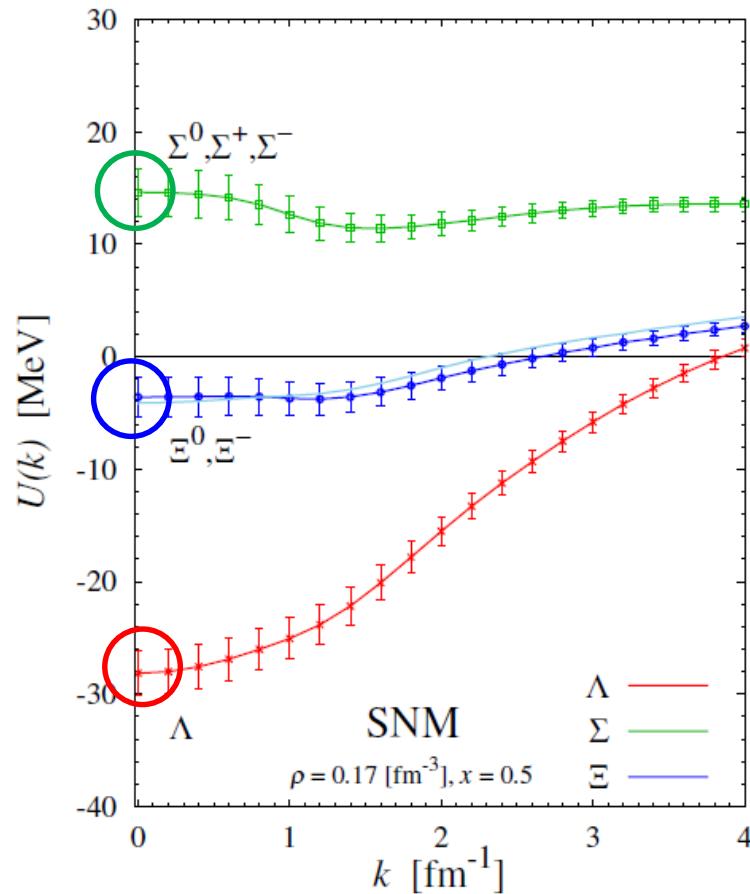
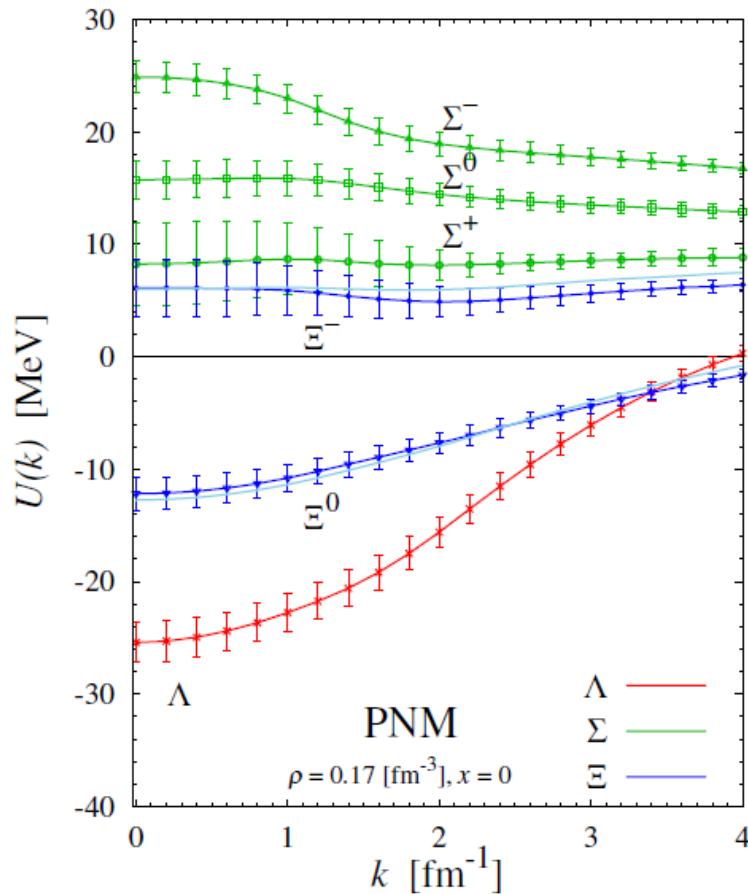
Inversion of spin-doublet between $T=1, 0$ due to spin-isospin dep of ΞN potential

Can be studied by
 $(K^-, K^+), (K^-, K^0)$ reaction on ${}^{10}\text{B}$

$A=6 (\Xi N\alpha)$ is unlikely to be bound

"Super-super heavy nuclei": Dense matter from LQCD

Hyperon single-particle potential



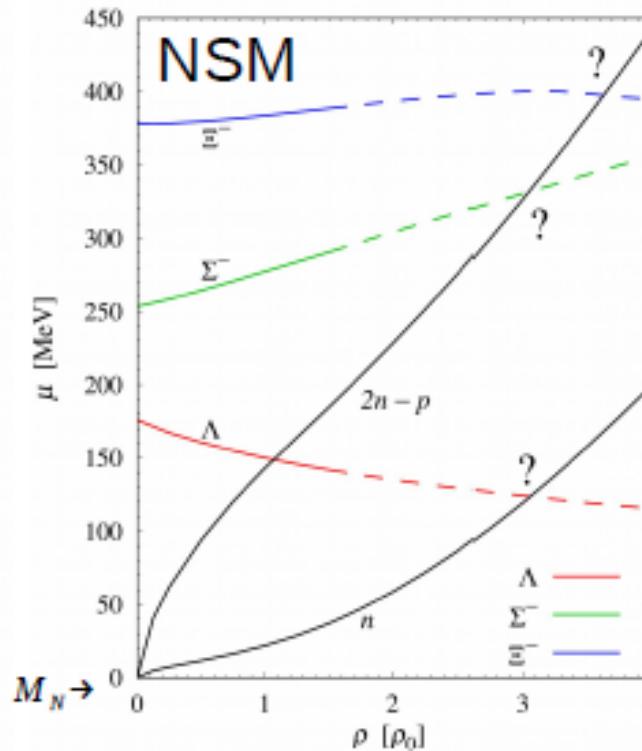
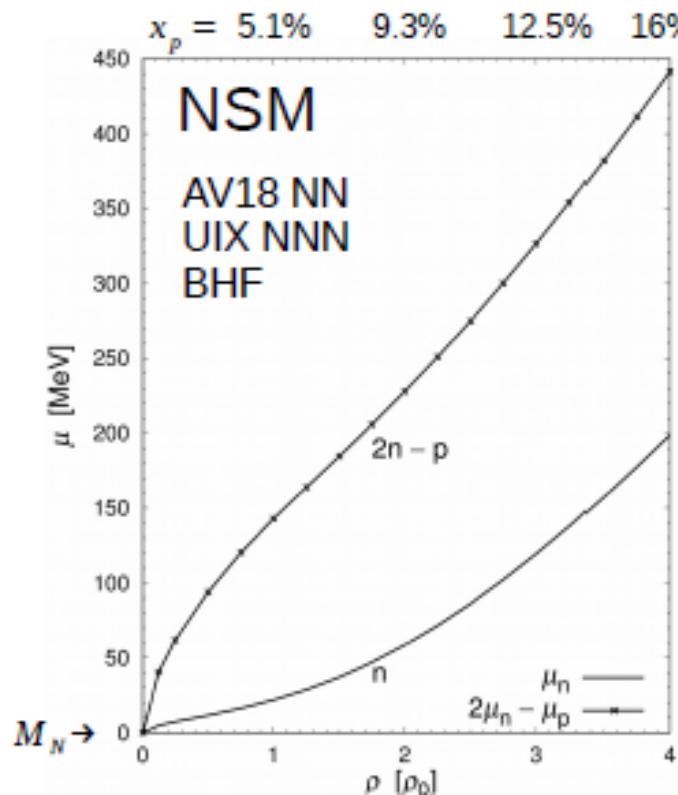
- Results are compatible with experimental suggestion.

$$U_{\Lambda}^{\text{Exp}}(0) \simeq -30, \quad U_{\Xi}(0)^{\text{Exp}} \simeq -10, \quad U_{\Sigma}^{\text{Exp}}(0) \geq +20 \quad [\text{MeV}]$$

attraction attraction small repulsion

(YN/YY pot from SU(3)f-irrep diag used)

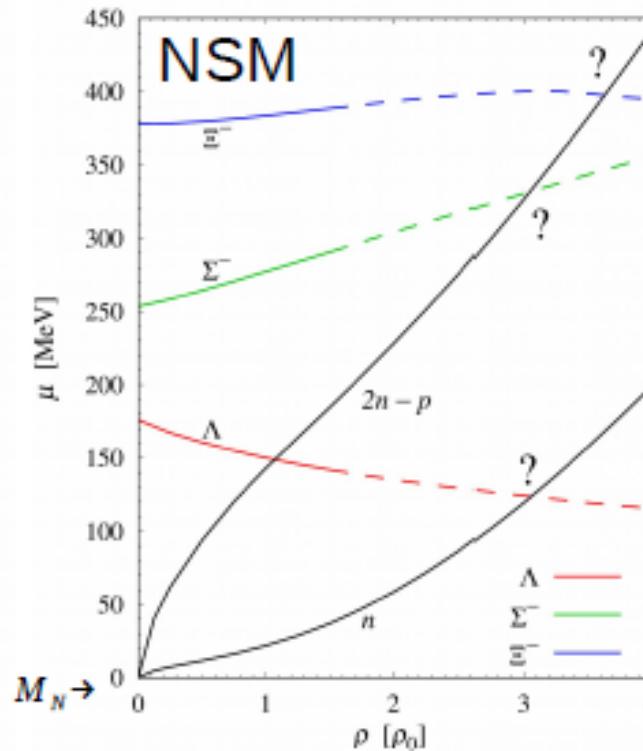
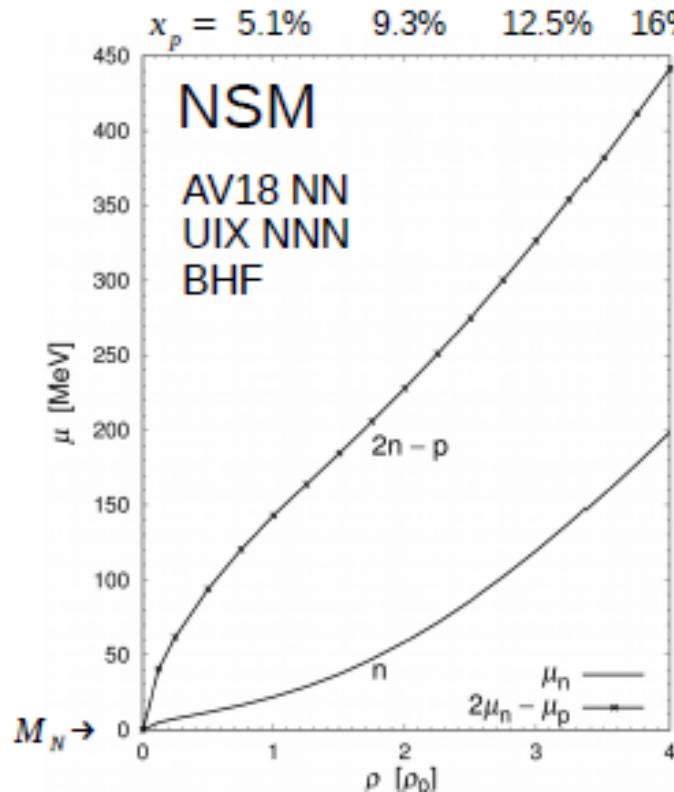
Hyperon onset in NSM (just for fun)



Preliminary

- Result indicate Λ , Σ^- , Ξ^- appear around $\rho = 3.0 - 4.0 \rho_0$
- However,
 - $YN^{L=1,2\dots}$ and YNN force could be important at high density.
 - We may need to compare with more sophisticated μ_n , μ_p than BHF.⁵⁴

Hyperon onset in NSM (just for fun)



- Result indicate Λ , Σ^- , Ξ^- appear around $\rho = 3.0 - 4.0 \rho_0$
- However,
 - $YN^{L=1,2,\dots}$ and YNN force could be important
 - We may need to compare with more S-wave/LS/BHF

[T. Inoue]

[Challenges]
Precision for $|S| \leq 1$
3-baryon forces
P-wave/LS forces

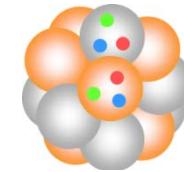
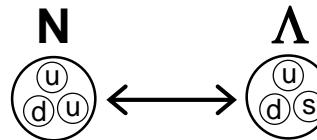
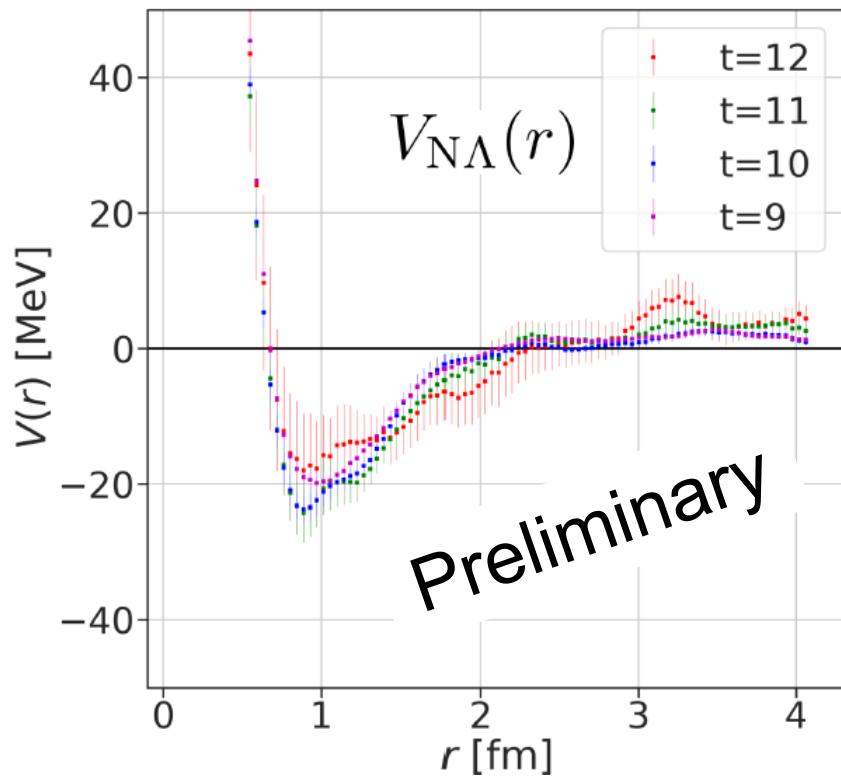
BHF₅₄

NΛ potentials @ physical point

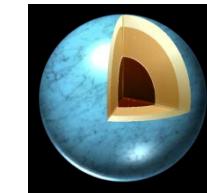
$$\begin{aligned} m_\pi &\neq 146 \text{ MeV} \\ &= 137 \text{ MeV} \end{aligned}$$

Effective central force in 1S_0 channel

[T. M. Doi]



Properties of Λ hyper-nuclei



Structure of neutron stars

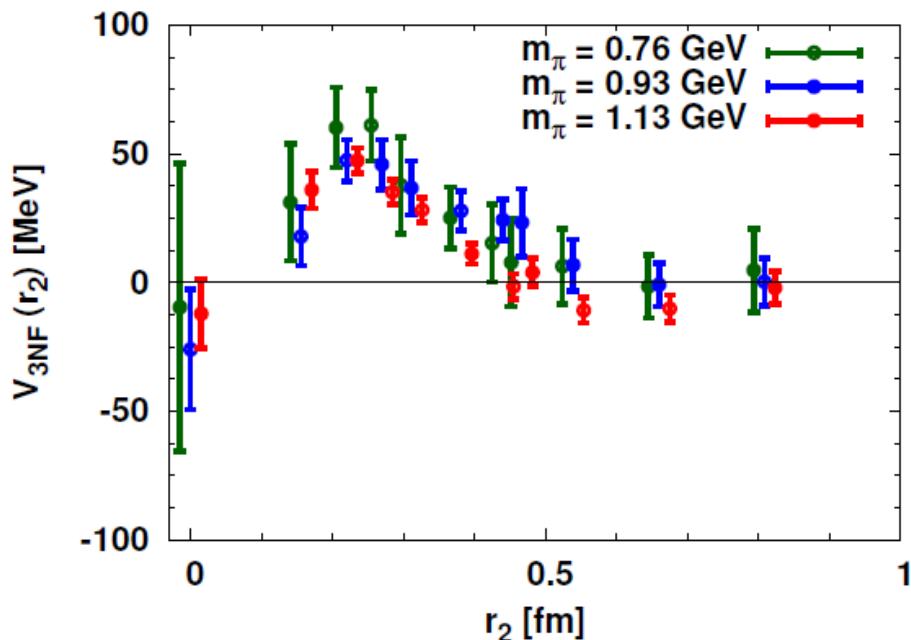
Challenge: Large statistical noises &
contaminations from unwanted (inelastic) states

Possibly Deep Learning is useful to overcome this issue?

→ related talk by L. Wang (Oct 30)

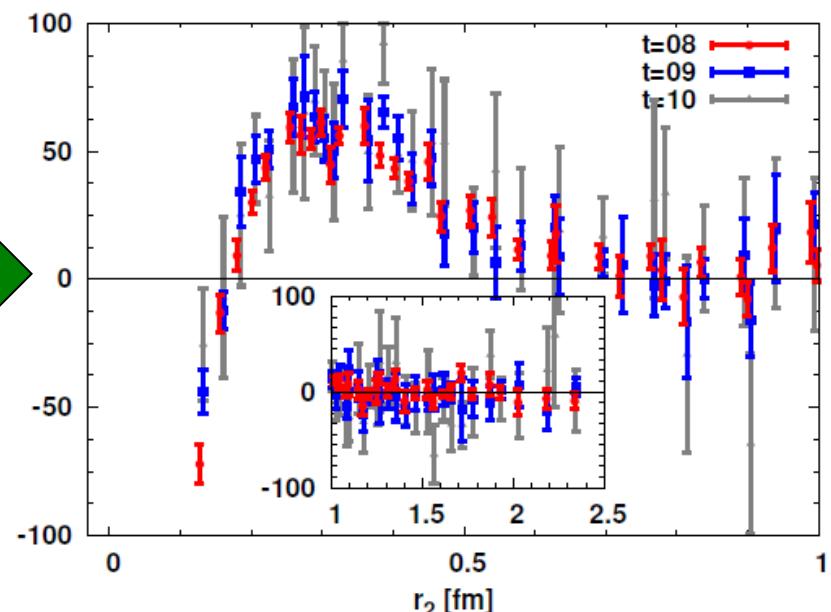
3N-forces (3NF)

Nf=2, $m\pi=0.76-1.1$ GeV



Triton channel

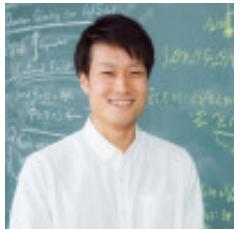
Nf=2+1, $m\pi=0.51$ GeV



Magnitude of 3NF is similar for all masses

Range of 3NF tend to be enlarged for $m(\pi)=0.5\text{GeV}$

Next challenge: **Calc of P-wave 2BF** : better subtraction of 2BF in 3-body systems
YNN (w/o or w/ P-wave 2BF) for structure of neutron stars



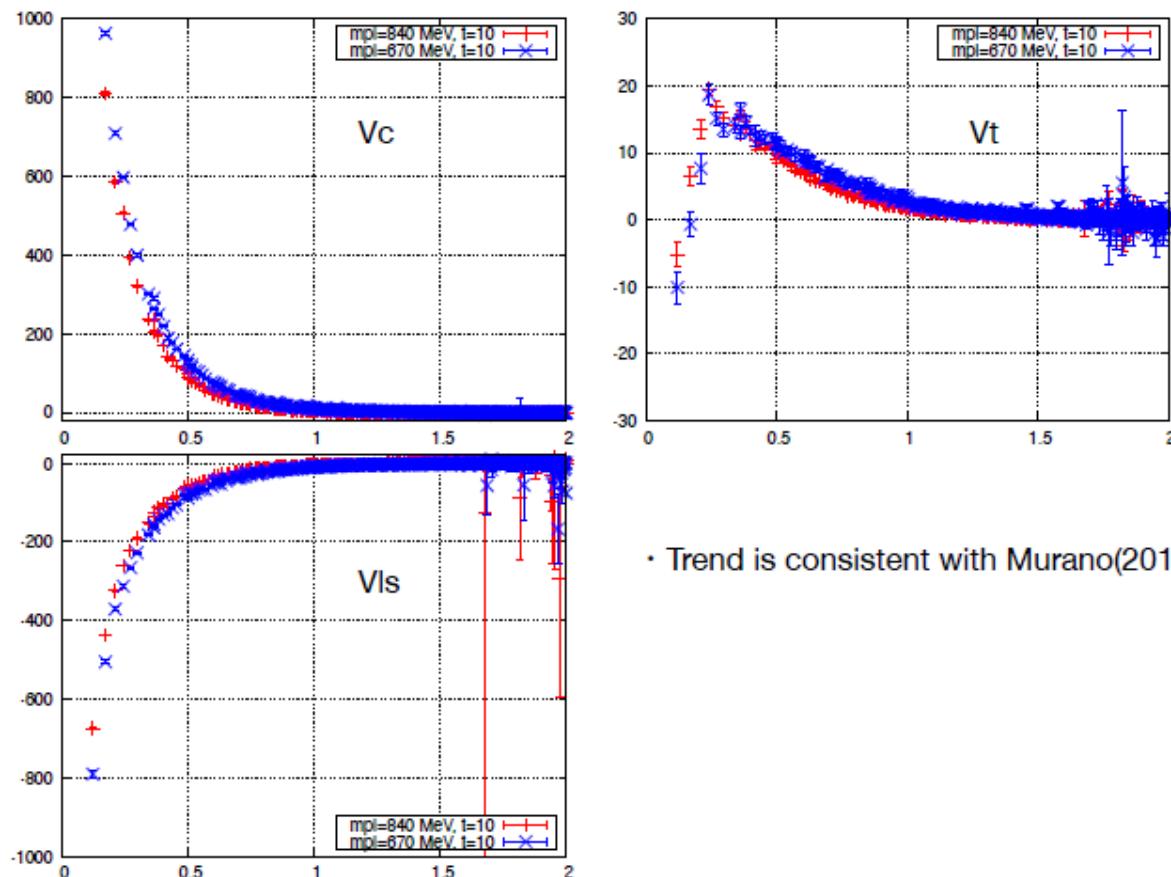
Interactions in P-wave channel

- NN interaction in P-wave @ SU(3), $m(\text{PS})=670, 840 \text{ MeV}$
- fLapH method for efficient calc

T. Sugiura

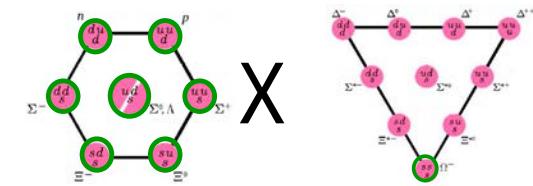
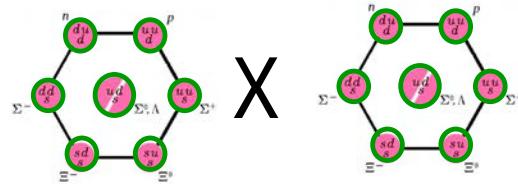
Potential: mpi dependence

21/30



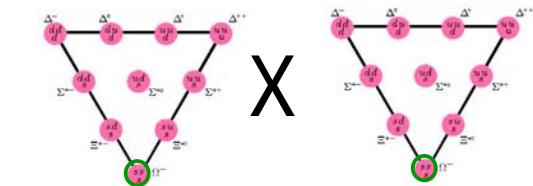
- Trend is consistent with Murano(2013)

Candidates of di-baryons



$$8 \times 10 = 35 + 8 + 10 + 27$$

NΩ (J=2) Goldman et al. ('87)
 Oka ('88)



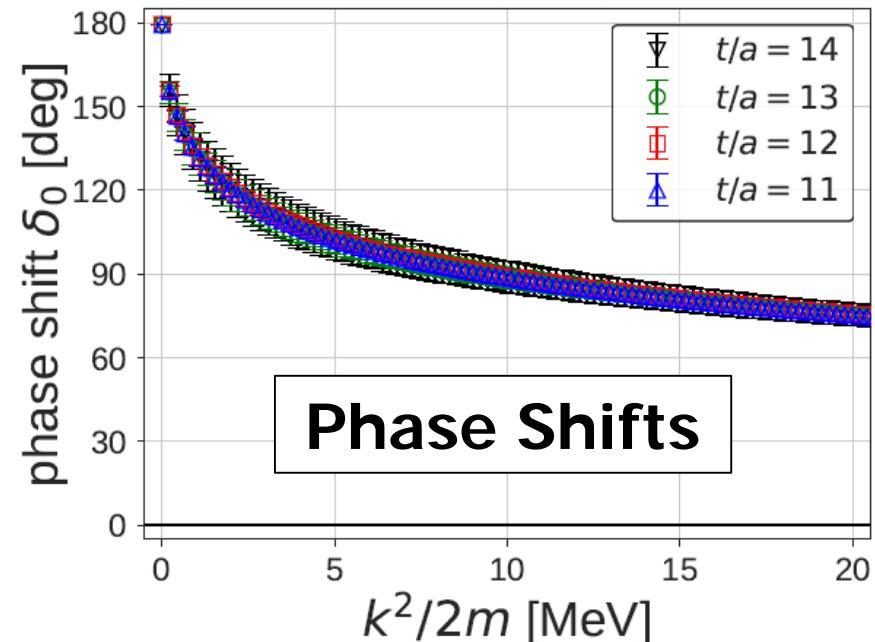
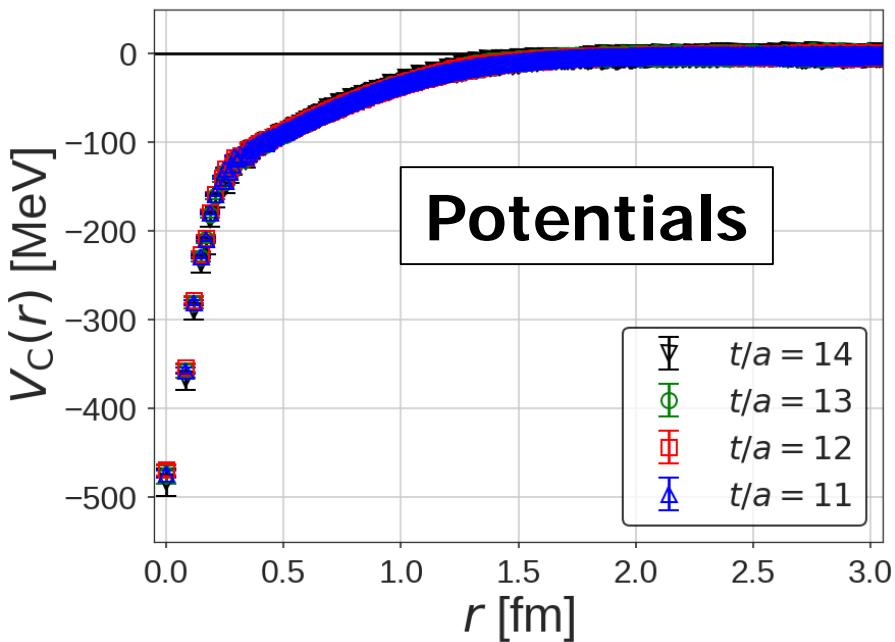
$$10 \times 10 = 28 + 27 + 10^* + 35$$

$\Omega\Omega$ ($J=0$) $\Delta\Delta$ ($J=3$)

Zhang et al. ('97)

Dyson-Xuong ('64)
Kamae-Fujita ('77)
Oka-Yazaki ('80)

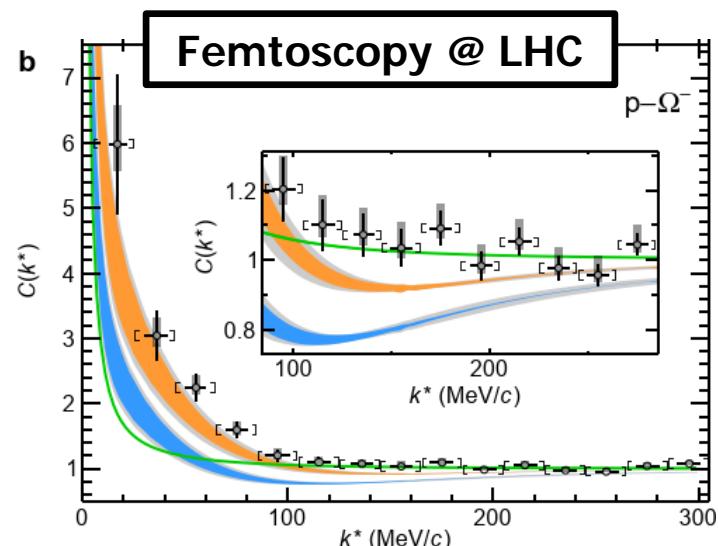
N Ω system (5S_2)



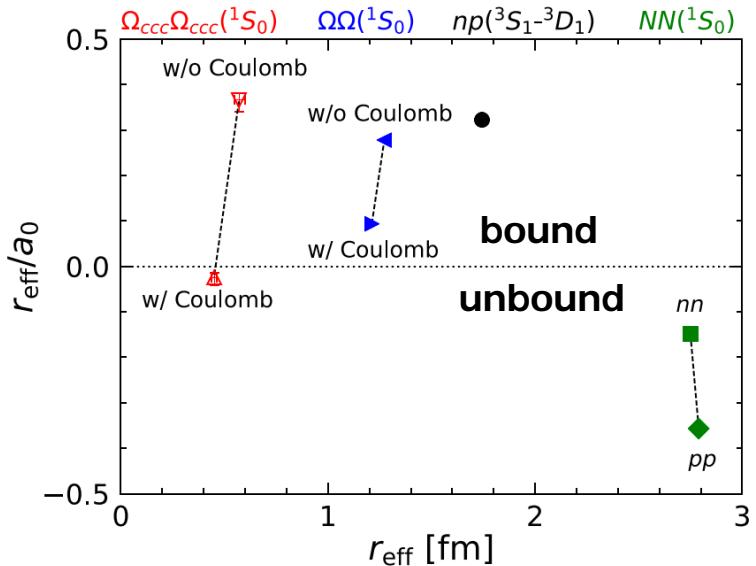
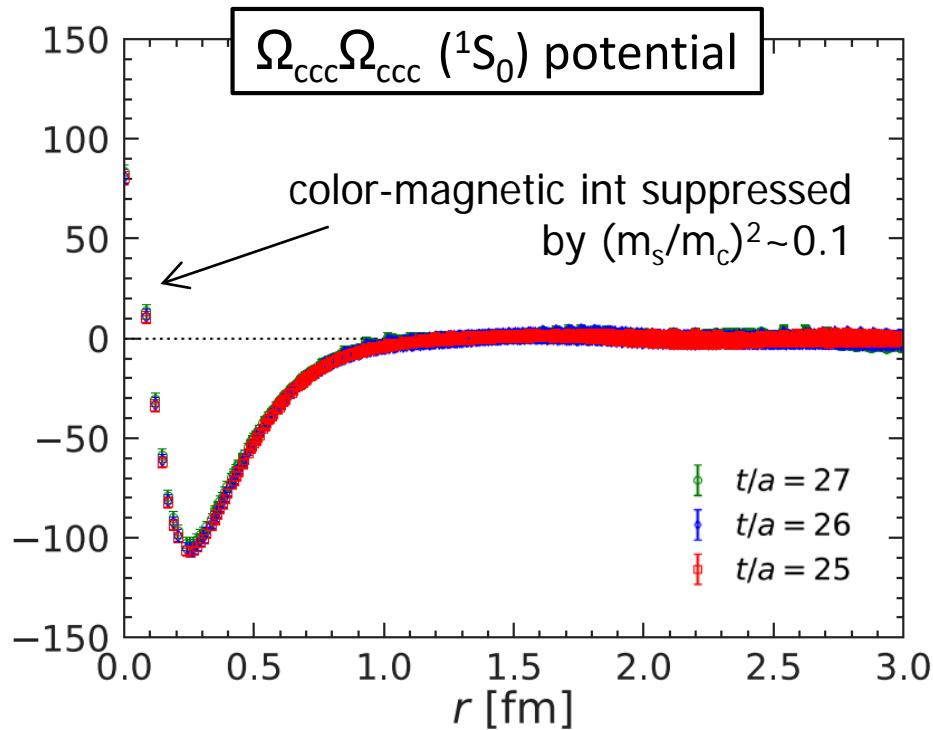
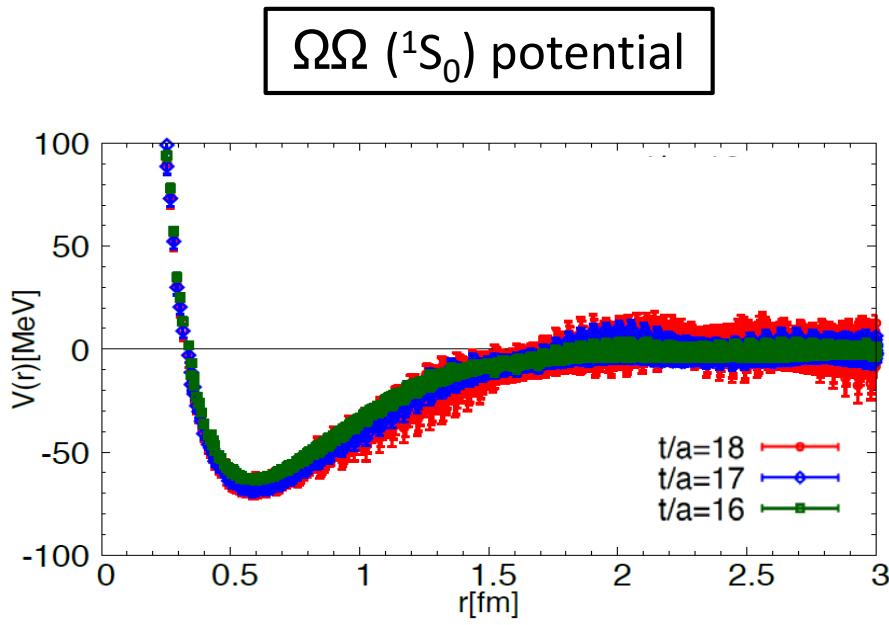
(Quasi) Bound state
[~ Unitary limit]

$$B_{N\Omega} = 1.54(0.30)(^{+0.04}_{-0.10}) \text{ MeV}$$

$$B_{p\Omega^-} = 2.46(0.34)(^{+0.04}_{-0.11}) \text{ MeV}$$



Most Strange/charming Dibaryons : $\Omega\Omega$, $\Omega_{ccc}\Omega_{ccc}$



di-Omegas near unitarity

$\Omega\Omega$ could be searched in LHC RUN3

S. Gongyo et al. (HAL Coll.), PRL120(2018)212001
 Y. Lyu, H. Tong et al., PRL127(2021)072003

Meson-Meson and Meson-Baryon Interactions

Related talks by

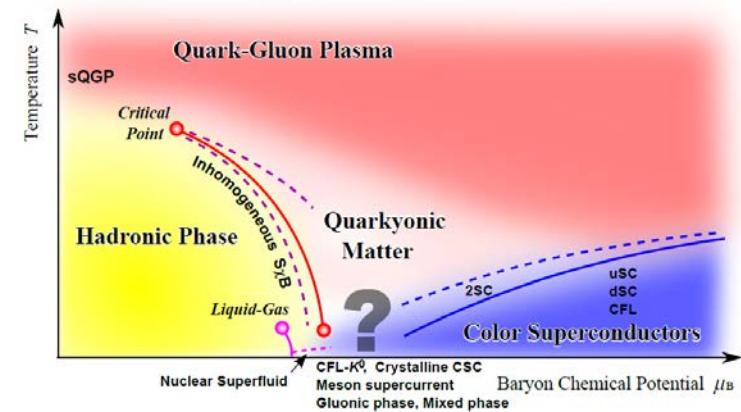
- Y. Lyu (Oct 18) : D-Dbar* for T_{cc}
- K. Murakami (Oct 28) : Kbar-N for $\Lambda(1405)$
- W. A. Yamada (Oct 29) : Dbar-N for Penta(?)

N ϕ system

Medium effects on hadron properties

↔ partial restoration
of chiral symmetry

Hayano-Hatsuda, ReV. Mod. Phys. 82 (2010) 2949

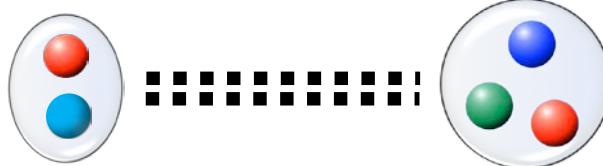


Fukushima-Hatsuda, RPP74(2011)014001

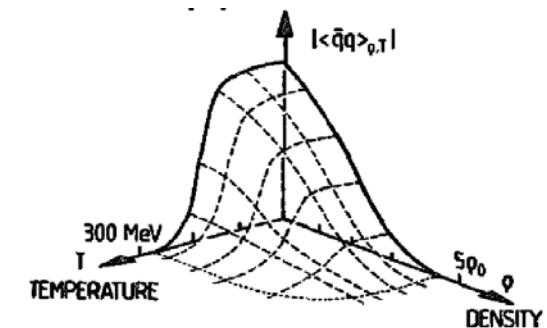
ϕ -meson is one of the ideal probes

Brown-Rho, RRL66(1991)2720
Hatsuda-Lee, PRC46(1992)R34
Gubler-Weise, NPA954(2016)125

How color-dipole interacts w/ other hadrons?



2-pion-exchange?



Weise, NPA553(1993)59

Dipole-Dipole int: H. Fujii and D. Kharzeev, PRD60(1999)114039
Dipole-Nucleon int: J. Castella and G. Krein, PRD98(2018)014029

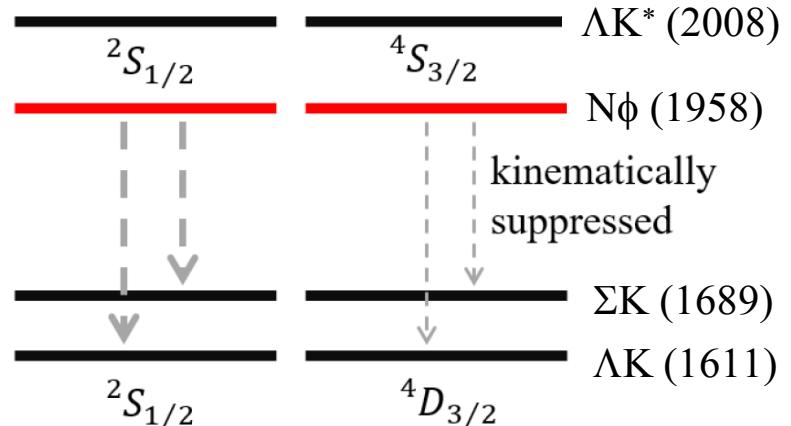
N ϕ system

[PDG]

N ϕ interaction from LQCD

ΛK , ΣK thresholds are open

Coupled channel calc expensive



$J=3/2$: ΛK , ΣK suppressed by D-wave

$J=1/2$: ΛK , ΣK S-wave mixing, no suppression

(3-body decay modes,
 ϕ -decay modes neglected)

(ωN , ρN also neglected by OZI)

→ We calculate N ϕ ($J=3/2$) w/ single-channel approximation

Y. Lyu et al., PRD106(2022)074507

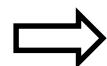
Hadron	Lattice [MeV]	Expt. [MeV]
π	146.4(4)	138.0
K	524.7(2)	495.6
ϕ	1048.0(4)	1019.5
N	954.0(2.9)	938.9

s-sbar annihilation neglected
 $\phi \rightarrow K\bar{K}$ forbidden at this mass
 $\phi \rightarrow 3\pi$ found to be negligible



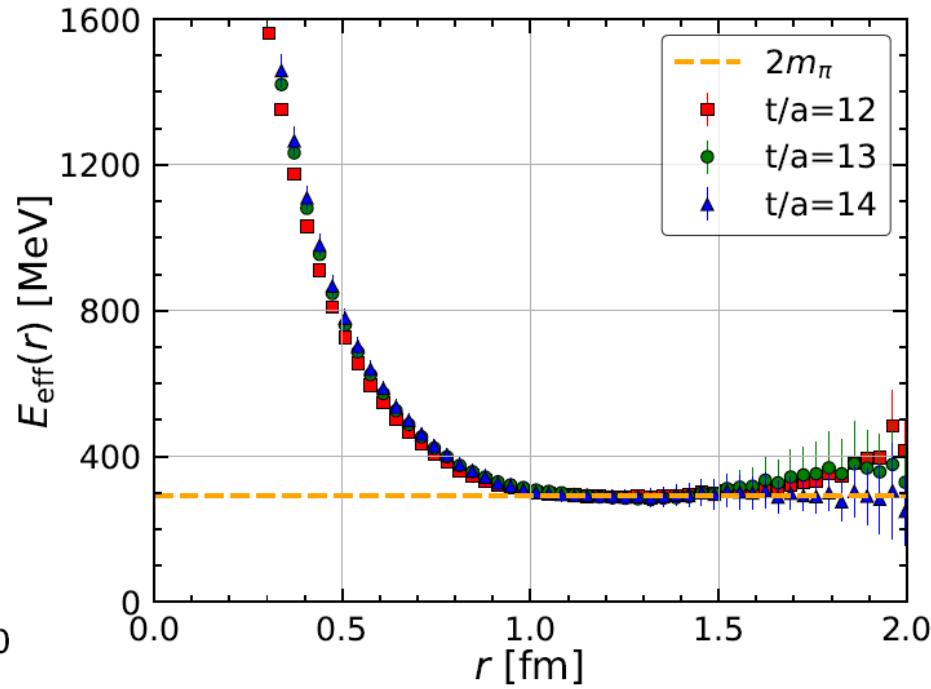
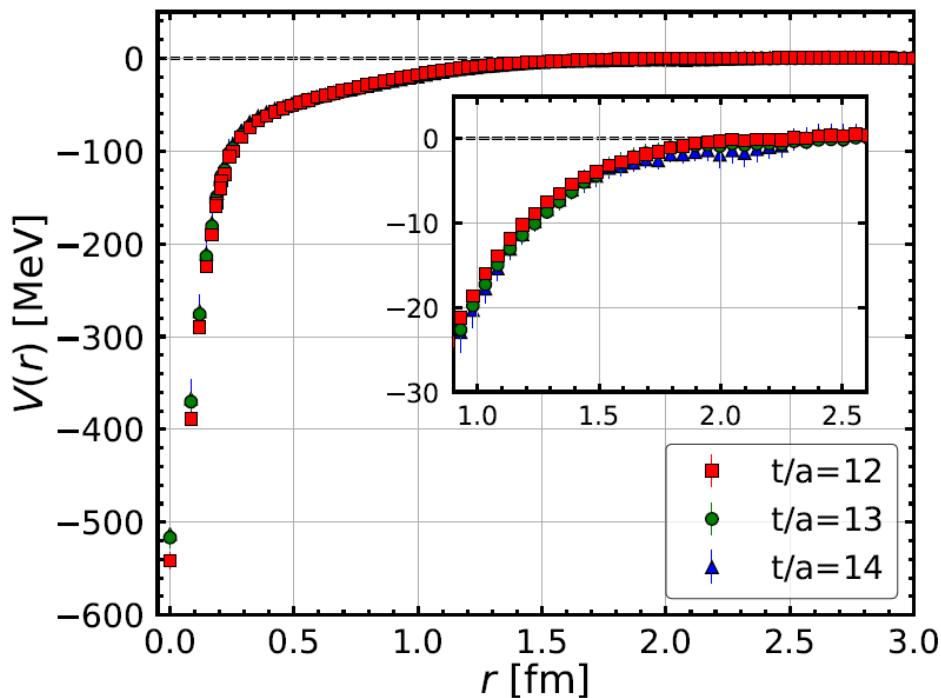
N ϕ system ($^4S_{3/2}$)

Potential



Tail structure of potential

Y. Lyu

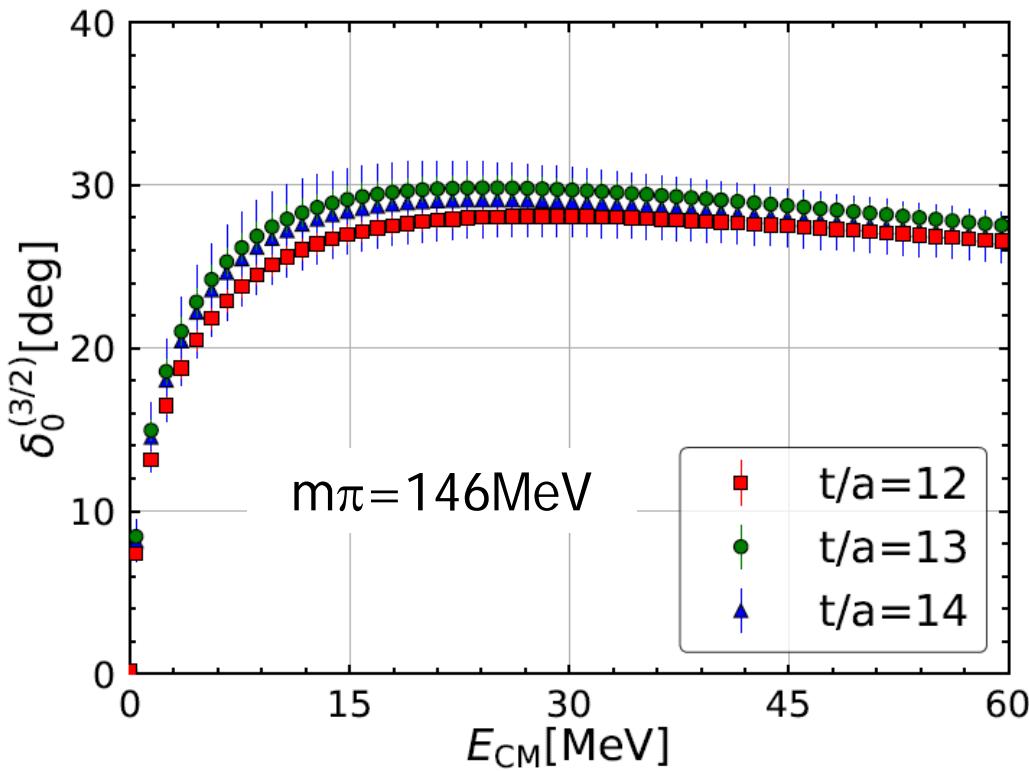


$$V(r) \xrightarrow{r \rightarrow \infty} -\alpha \frac{e^{-2m_\pi r}}{r^2} \implies E_{\text{eff}}(r) = -\frac{\ln[-V(r)r^2/\alpha]}{r} \xrightarrow{r \rightarrow \infty} 2m_\pi$$

Potential is attractive at all distances

Tail is consistent w/ two-pion exchange (TPE) !

Phase shifts of $N\phi$ ($^4S_{3/2}$)



Fit the potential and solve Schrodinger eq. in infinite V

$$V_{\text{fit}}(r) = \sum_{i=1,2} a_i e^{-(r/b_i)^2} + a_3 m_\pi^4 f(r; b_3) \frac{e^{-2m_\pi r}}{r^2}$$

↑
EFT by Castella-Krein

Fit w/ $m_\pi = m_\pi(\text{LQCD})$

Semi-quantitative extrapolation
by $m_\pi \rightarrow m_\pi(\text{phys})$

m_π [MeV]	$a_0^{(3/2)}$ [fm]	$r_{\text{eff}}^{(3/2)}$ [fm]
146.4	$-1.43(23)_{\text{stat}}(^{+36}_{-06})_{\text{syst}}$	$2.36(10)_{\text{stat}}(^{+02}_{-48})_{\text{syst}}$
138.0	$\simeq -1.25$	$\simeq 2.49$

c.f. ALICE Coll.

$$a_0^{(\text{spin-ave})} = -0.85(34)(14) \text{ fm}$$

$$r_{\text{eff}}^{(\text{spin-ave})} = 7.85(1.54)(0.26) \text{ fm}$$

New avenue for Hadron interactions by Combined analysis of LQCD & femtoscopy

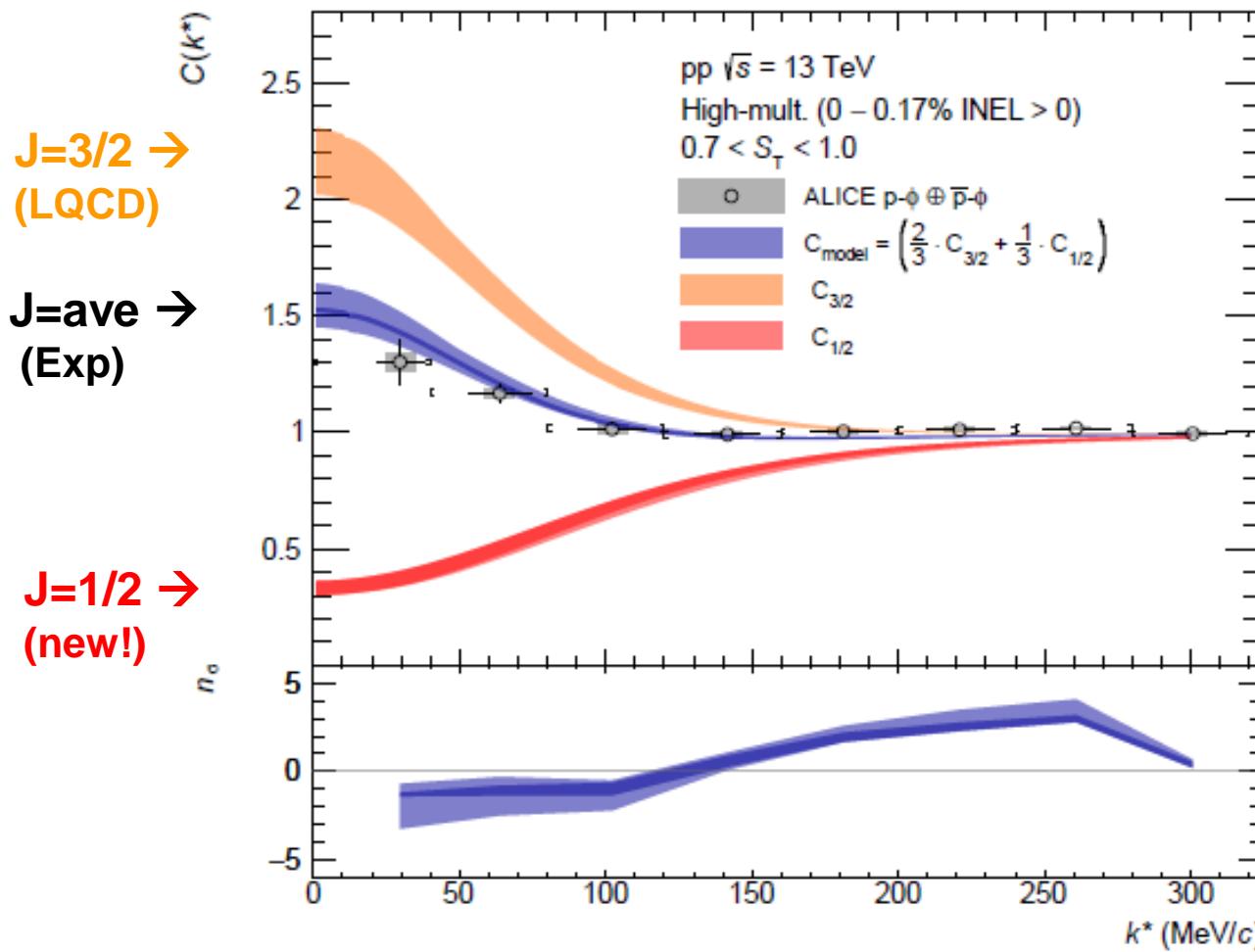
$N\phi$ $J=3/2 \leftarrow$ LQCD calc possible (ΛK , ΣK mixings suppressed by D-wave)

$N\phi$ $J=1/2 \leftarrow$ LQCD challenging (ΛK , ΣK mixings are S-wave)

$N\phi$ spin average \leftarrow ALICE femtoscopy possible
(but spin projection difficult)

→ By combining LQCD and ALICE femtoscopy,
we can extract $N\phi$ $J=1/2$

N ϕ system ($^2S_{1/2}$) from LQCD + Exp

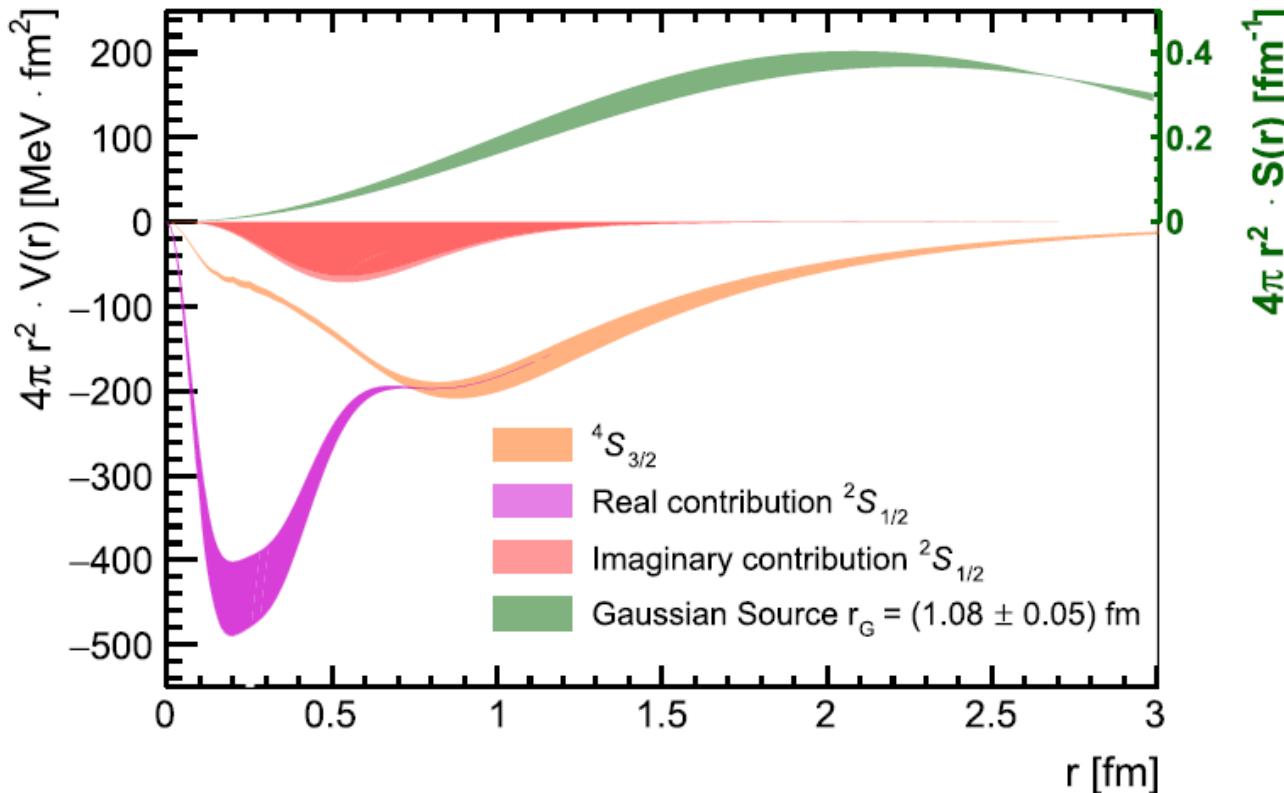


Correlation function measured by ALICE

Fit the femtoscopy data w/ parameters in potential in J=1/2

N ϕ system ($^2S_{1/2}$) from LQCD + Exp

Potential (w/ phase space factor of $4\pi r^2$)



Indication of N ϕ bound state in J=1/2

B.E. = 12.8-56.1 MeV

$$\text{Re } f_0^{(1/2)} = -1.54^{+0.53}_{-0.53}(\text{stat.})^{+0.16}_{-0.09}(\text{syst.}) \text{ fm},$$

$$\text{Im } f_0^{(1/2)} = 0.00^{+0.35}_{-0.00}(\text{stat.})^{+0.16}_{-0.00}(\text{syst.}) \text{ fm},$$

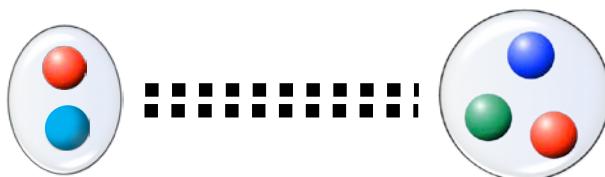
$$\text{Re } d_0^{(1/2)} = +0.39^{+0.09}_{-0.09}(\text{stat.})^{+0.02}_{-0.03}(\text{syst.}) \text{ fm},$$

$$\text{Im } d_0^{(1/2)} = 0.00^{+0.00}_{-0.04}(\text{stat.})^{+0.00}_{-0.02}(\text{syst.}) \text{ fm}.$$

Exploring hadron interactions in broader scope w/ charm degrees of freedom

- $N\text{-J}/\psi, N\text{-}\eta_c (N\text{-}c\bar{c})$ interactions $\longleftrightarrow N\phi (N\text{-}s\bar{s})$

How color-dipole interacts w/ other hadrons?



Dipole-Dipole int: H. Fujii and D. Kharzeev, PRD60(1999)114039
Dipole-Nucleon int: J. Castella and G. Krein, PRD98(2018)014029

- $N\Omega_{ccc}$ interactions $\longleftrightarrow N\phi, N\Omega (=N\Omega_{sss})$

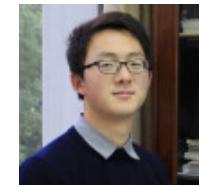
$N\Omega$ is quasi bound \leftarrow issue of open channels, $\Lambda\Xi, \Sigma\Xi$

T. Iritani et al. (HAL Coll.), PLB792(2019)284

$N\Omega_{ccc}$ is ideal system $\leftarrow N\Omega_{ccc}$ is below $\Lambda_c\Xi_{cc}, \Sigma_c\Xi_{cc}$ thresholds

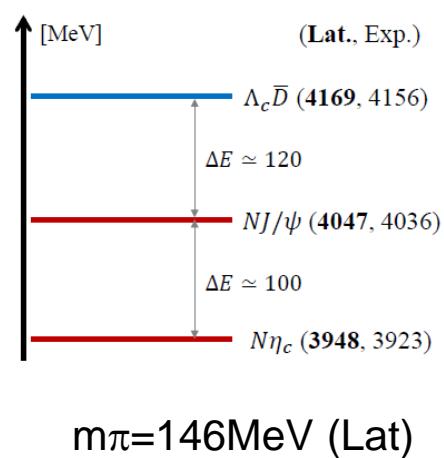
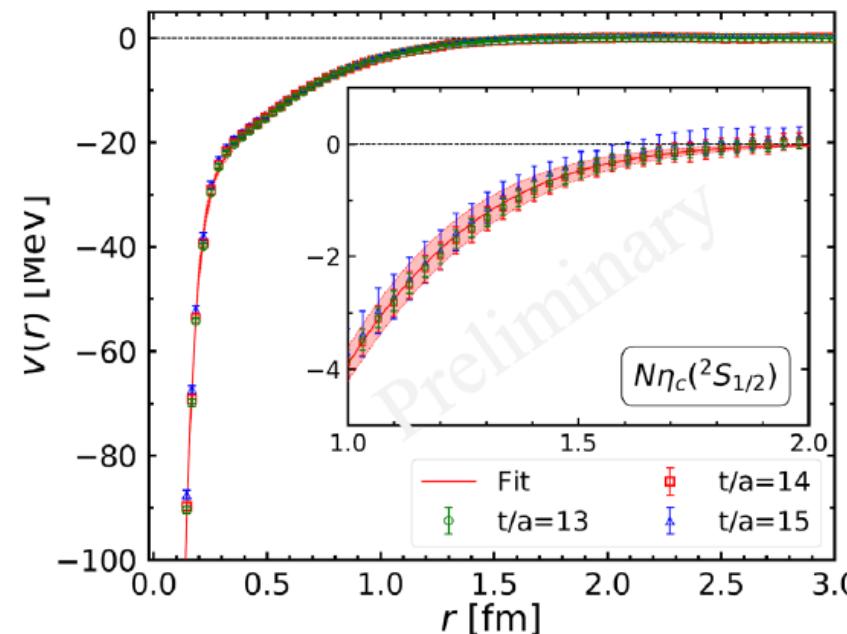
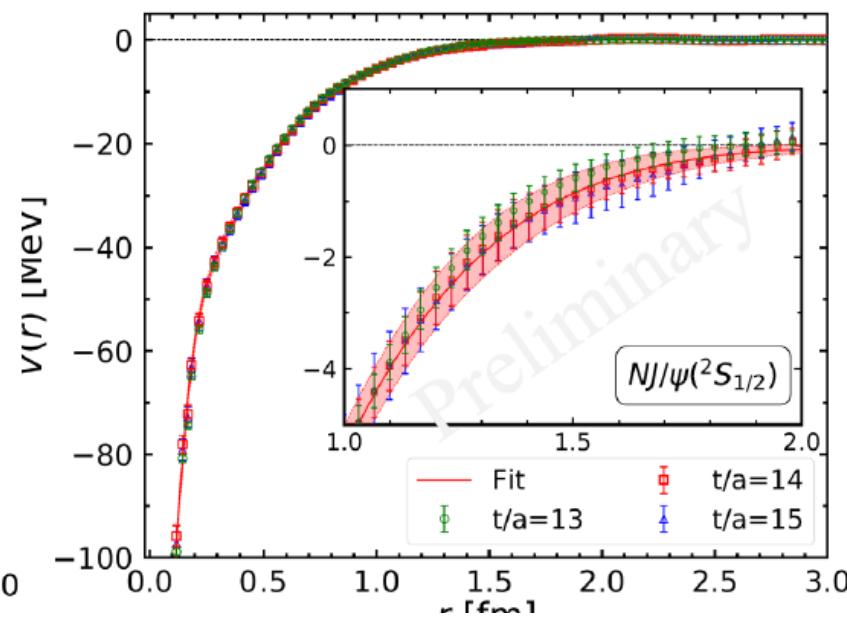
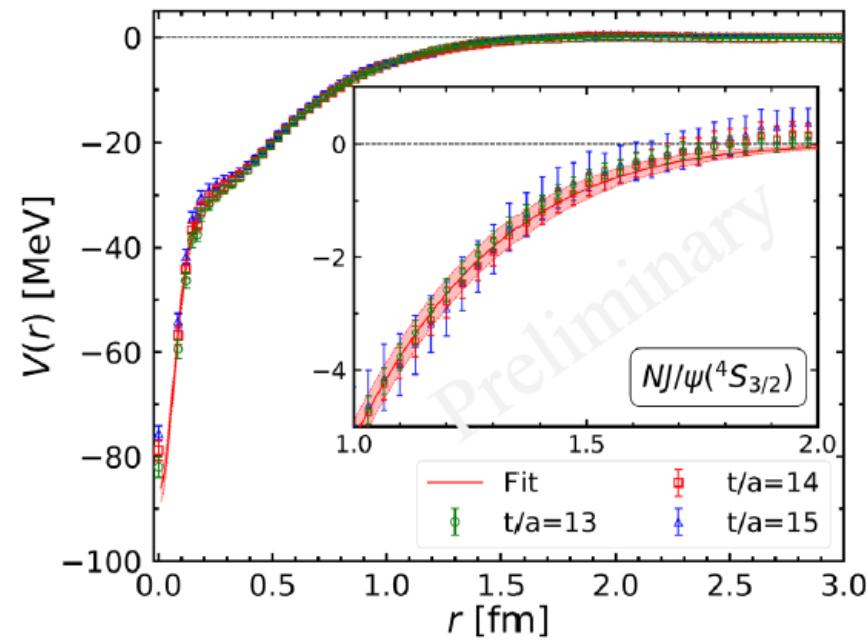
- $\Xi_{cc}-\Xi_{cc}$ interactions \longleftrightarrow superflavor partner of T_{cc}

$N-c\bar{c}$ potentials near phys point

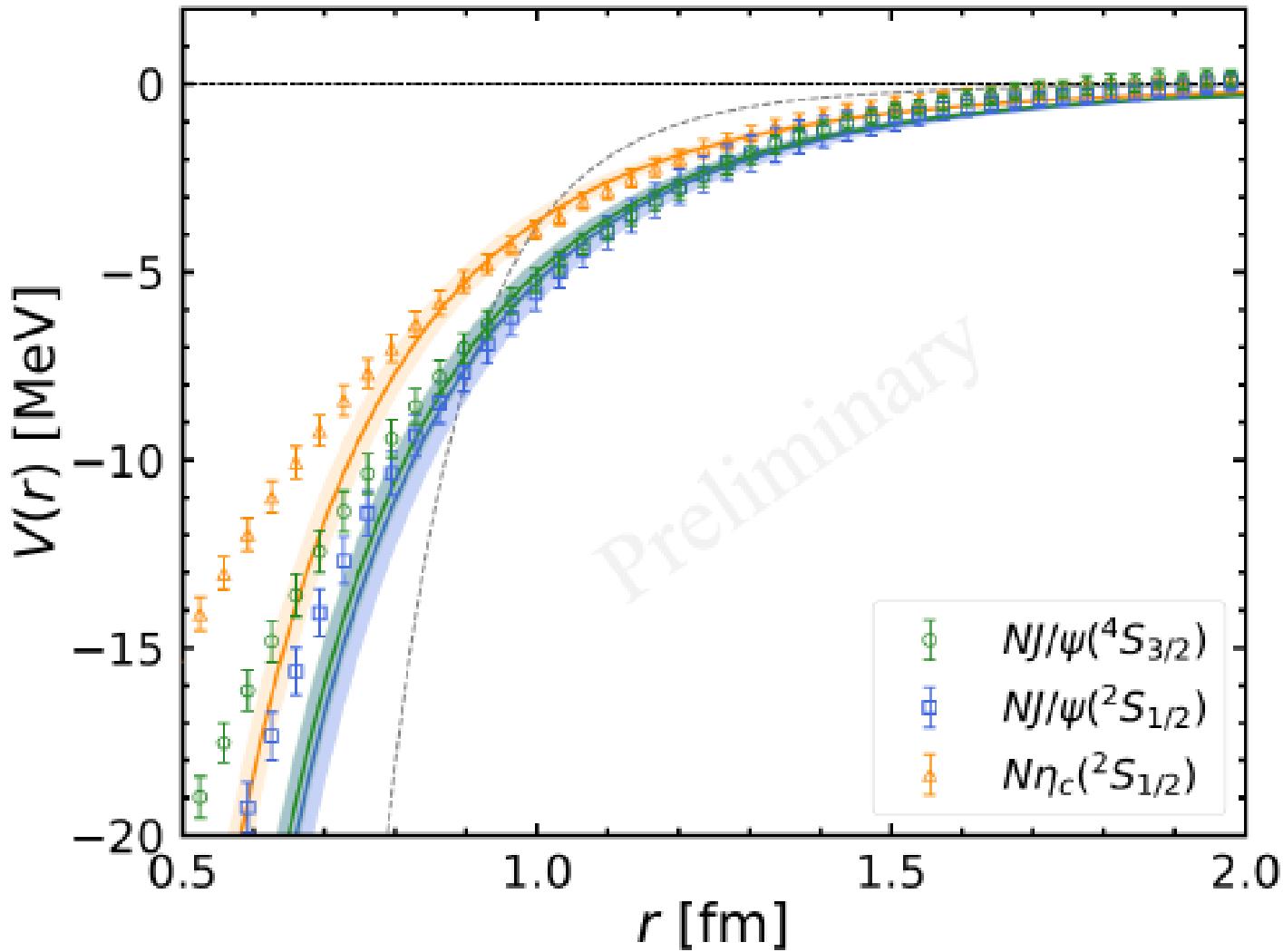


Y. Lyu

Slides from
Lattice 2024



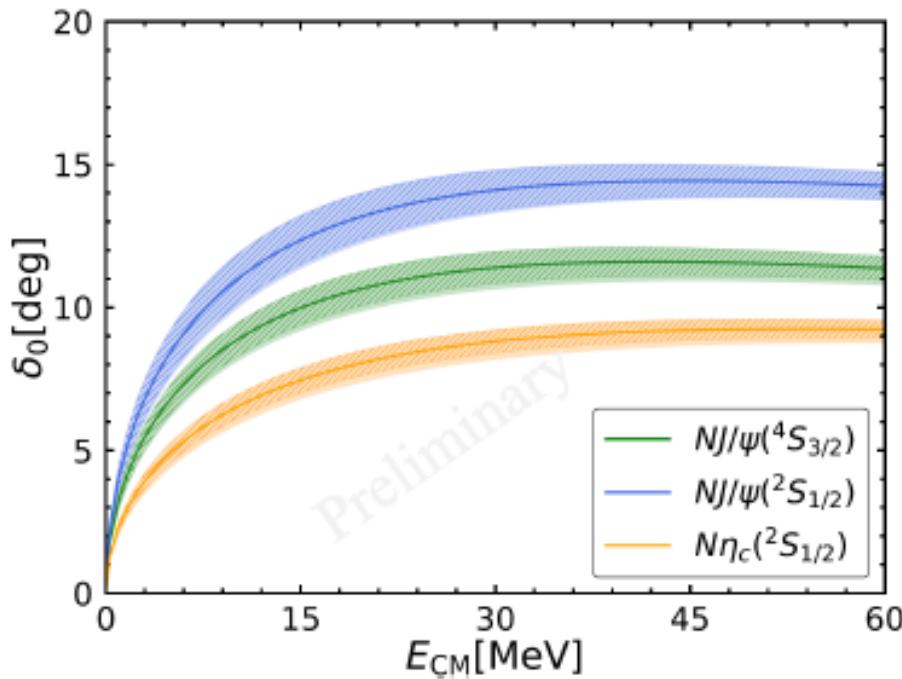
Long-range potentials of $N-c\bar{c}$



The long-range potentials are consistent with the two-pion-exchange (TPE)

Physical observables

➤ Scattering phase shifts



$$k \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_{\text{eff}} k^2$$

channel	a_0 [fm]	r_{eff} [fm]
$NJ/\psi(^4S_{3/2})$	$0.30(2) \begin{pmatrix} +0 \\ -2 \end{pmatrix}$	$3.25(12) \begin{pmatrix} +6 \\ -9 \end{pmatrix}$
$NJ/\psi(^2S_{1/2})$	$0.38(4) \begin{pmatrix} +0 \\ -3 \end{pmatrix}$	$2.66(21) \begin{pmatrix} +0 \\ -10 \end{pmatrix}$
$N\eta_c(^2S_{1/2})$	$0.21(2) \begin{pmatrix} +0 \\ -1 \end{pmatrix}$	$3.65(20) \begin{pmatrix} +0 \\ -6 \end{pmatrix}$

➤ A direct phenomenological application

- The J/ψ mass modification in nuclear medium is related to the spin-averaged scattering length of N - J/ψ scattering

A. Hayashigaki, Prog. Theor. Phys. 101 (1999)

$$\delta m_{J/\psi} \simeq -\frac{2\pi(m_N + m_{J/\psi})}{m_N m_{J/\psi}} a_{J/\psi}^{\text{spin-av}} \rho_{\text{nm}} = -19(3) \text{ MeV}$$

N- Ω_{ccc} interactions on the phys point

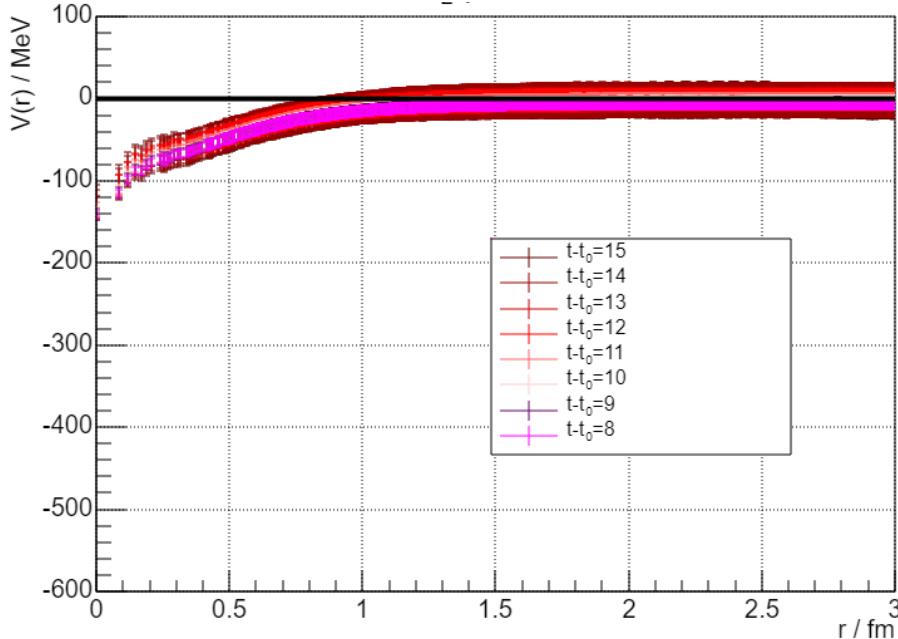


Results for the potentials

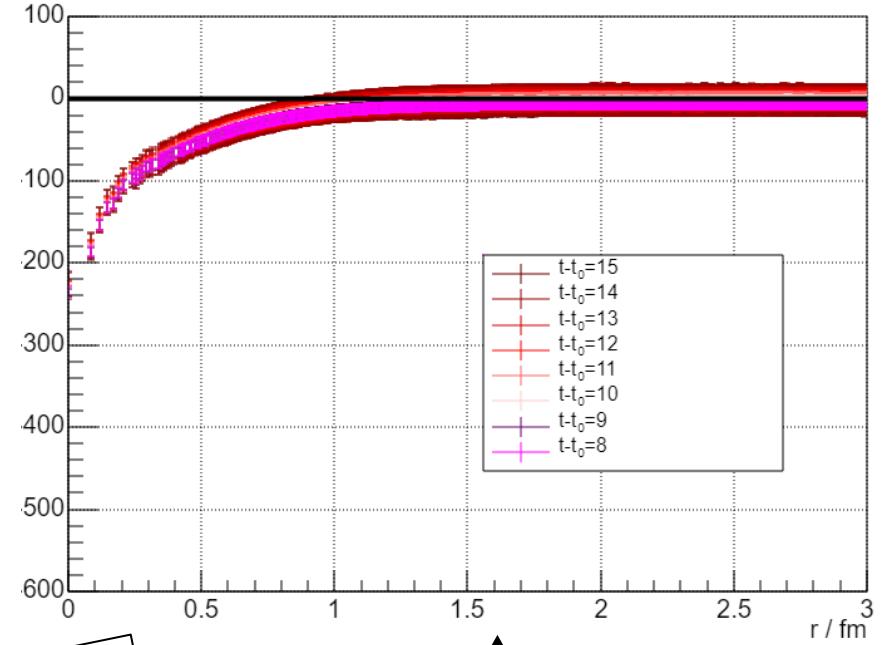
$m\pi=137\text{MeV}$ (Lat)

L. Zhang

5S_2



3S_1



Preliminary

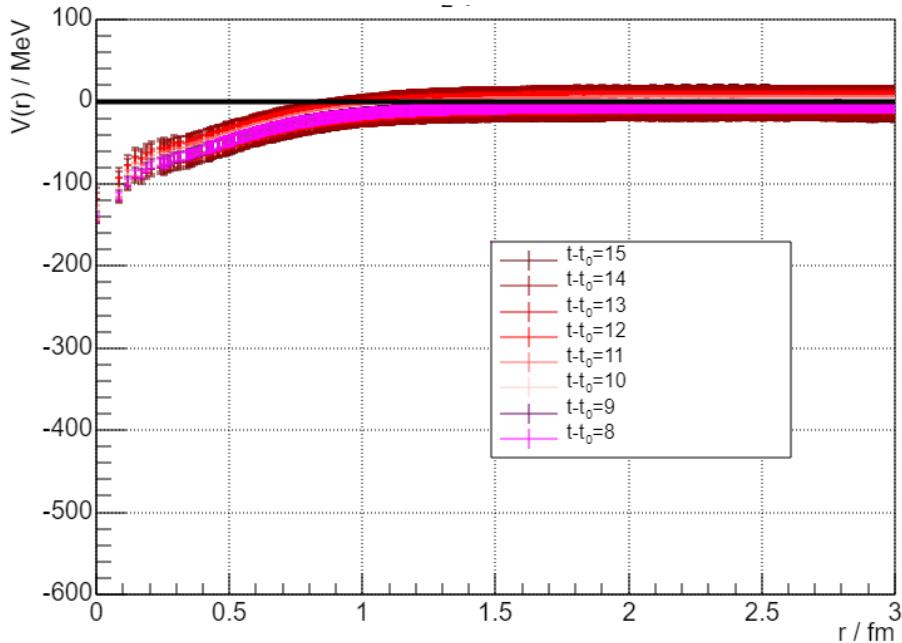


channel difficult to
be accessed by $N\Omega$

TPE or not? To be studied

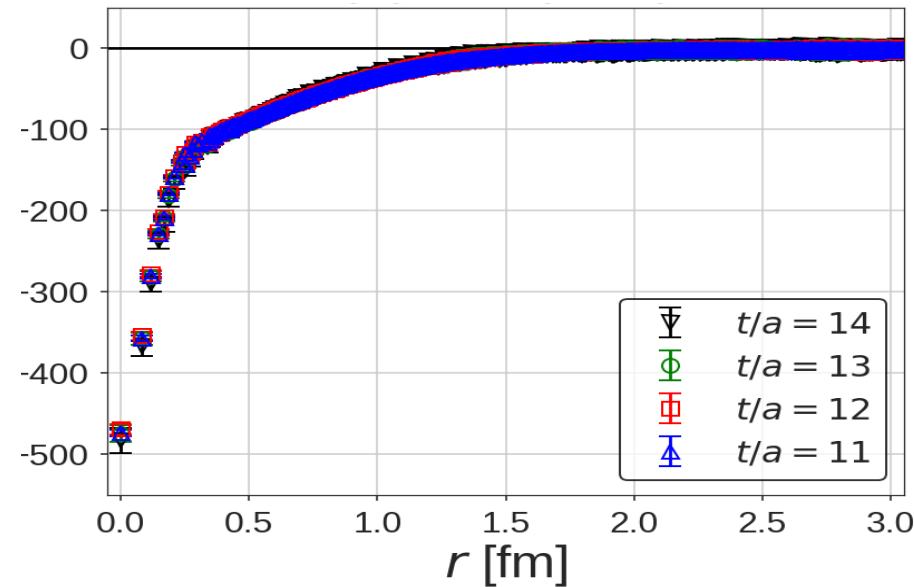
N- Ω_{ccc} vs. N- Ω interactions

$N\Omega_{ccc} ({}^5S_2)$



$m\pi=137\text{MeV}$ (Lat)

$N\Omega ({}^5S_2)$



$m\pi=146\text{MeV}$ (Lat)

T. Iritani et al. (HAL Coll.), PLB792(2019)284

Λ_c -N interactions near the phys point



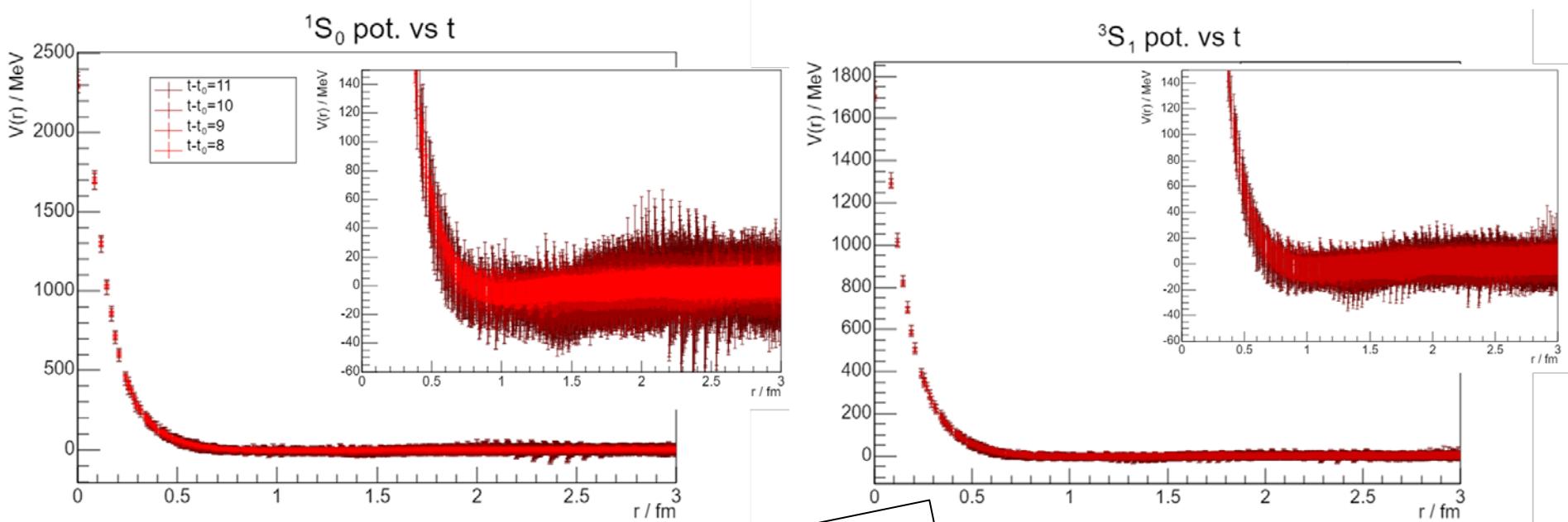
Possibility of Λ_c hypernuclei?

L. Zhang

c.f. T. Miyamoto et al., (HAL Coll.) NPA971(2018)113
 $m(\pi)=0.41\text{--}0.70 \text{ GeV}$

Impact on EFT

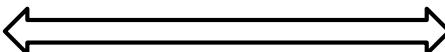
Haidenbauer-Krein (2018, 2021), J. Song et al. (2020)



Preliminary

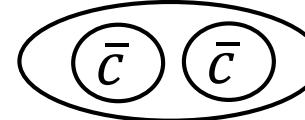
Superflavor partner of Tcc ?

Heavy quark



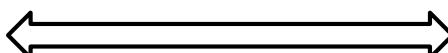
Same color charge

Heavy anti-diquark



Georgi-Wise, PLB243(1990)279
Savage-Wise, PLB248(1990)177

Tcc (D-D^{*})



Pentaquark state?

(Dbar^(*)-Ξ_{CC}^(*))

(cbar u) - (cc d)

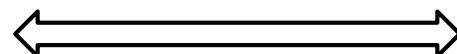
Tcc bar: (cbar u) - (cbar d)

$I(J^P) = 0(1^+)$

Asanuma-Yamaguchi-Harada, arXiv:2311.04695

$I(J^P) = 0(1/2^-)$

OBEP + superflavor sym → Bound state could exist



Hexaquark state?

(Ξ_{CC}-Ξ_{CC})

(cc u) - (cc d)

$I(J^P) = 0(1^+)$

Physical point simulation of Ξ_{cc} - Ξ_{cc} interactions

- F-conf is used for phys point calc on Fugaku
- Ξ_{cc} - Ξ_{cc} interactions

$$\begin{array}{lll} I(J^P) = 0(1^+) & {}^3S_1\text{-}{}^3D_1 \text{ channel} & \xleftarrow{\hspace{1cm}} \text{Candidate of} \\ & & \text{T}_{cc} \text{ partner} \\ I(J^P) = 1(0^+) & {}^1S_0 \text{ channel} & \end{array}$$

- Ξ_{ss} - Ξ_{ss} (= usual Ξ - Ξ) interactions
 - We can use this system as a “reference” to probe the role of heavy quark

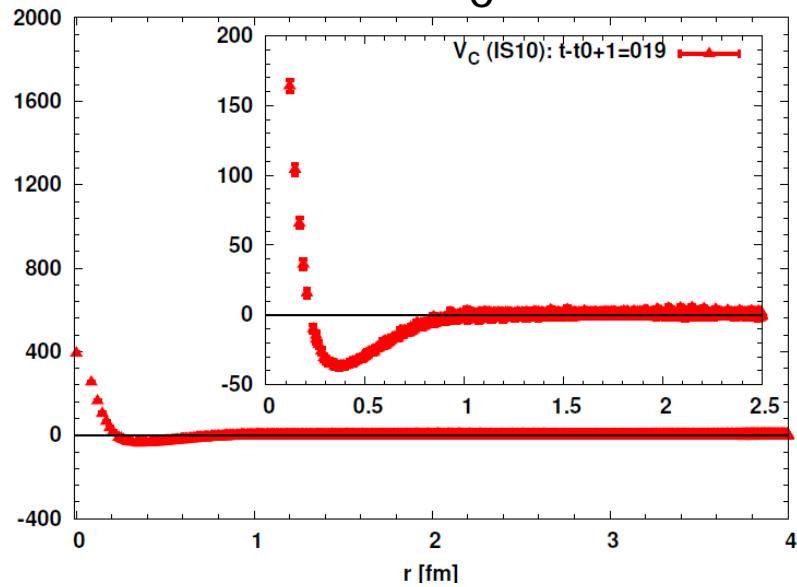
$$\begin{array}{lll} I(J^P) = 0(1^+) & {}^3S_1\text{-}{}^3D_1 \text{ channel} & \xleftarrow{\hspace{1cm}} \text{SU(3)f 10-plet} \\ I(J^P) = 1(0^+) & {}^1S_0 \text{ channel} & \xleftarrow{\hspace{1cm}} \text{SU(3)f 27-plet} \end{array}$$

$\Xi_{\text{CC}} - \Xi_{\text{CC}}$ interactions

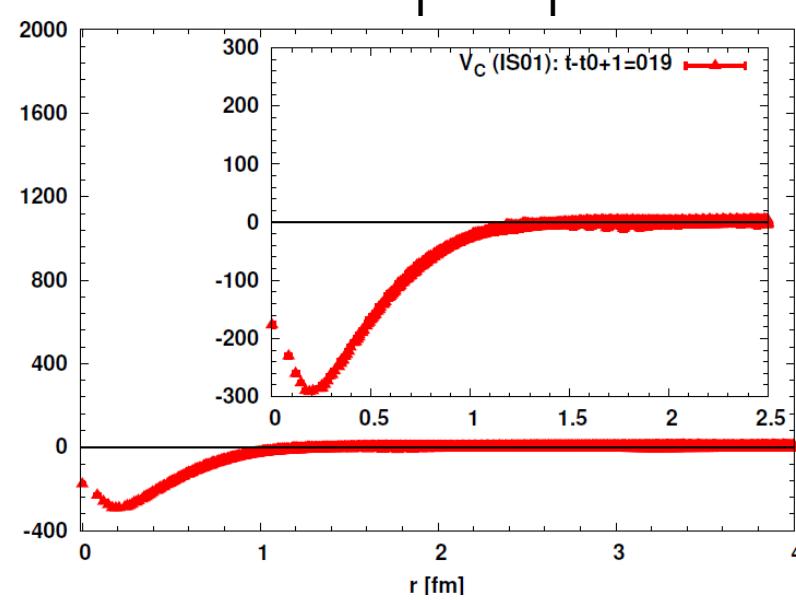
1S_0

$^3S_1 - ^3D_1$

V(r) [MeV]



V(r) [MeV]



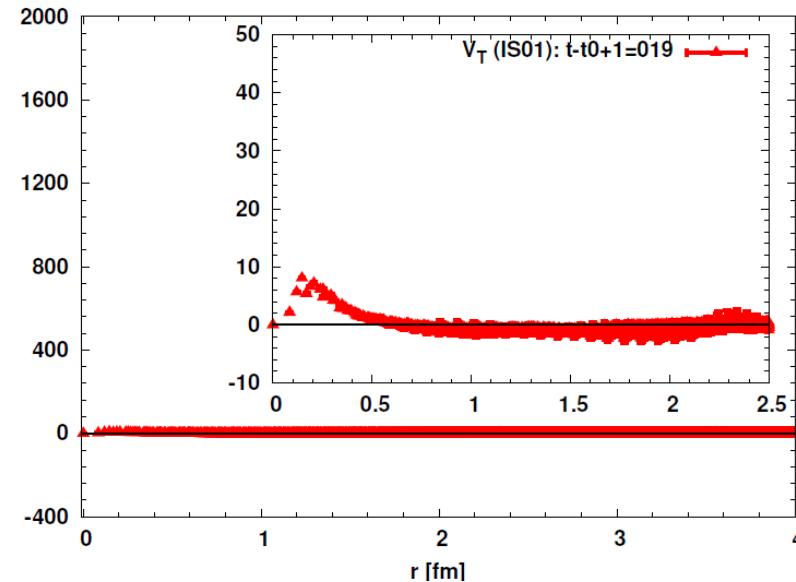
Central

Tensor

1S_0 : Repulsive core
+ attractive pocket
 3SD_1 : Strong attraction
& weak tensor force

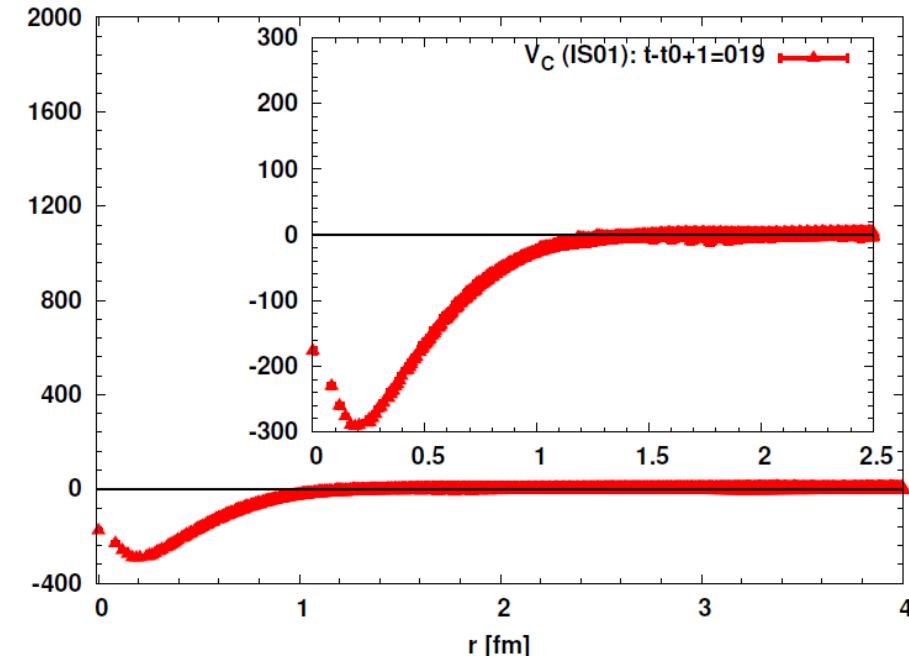
$$M(\Xi_{\text{CC}}) = 3642.9(1.4) \text{ MeV}$$

V(r) [MeV]

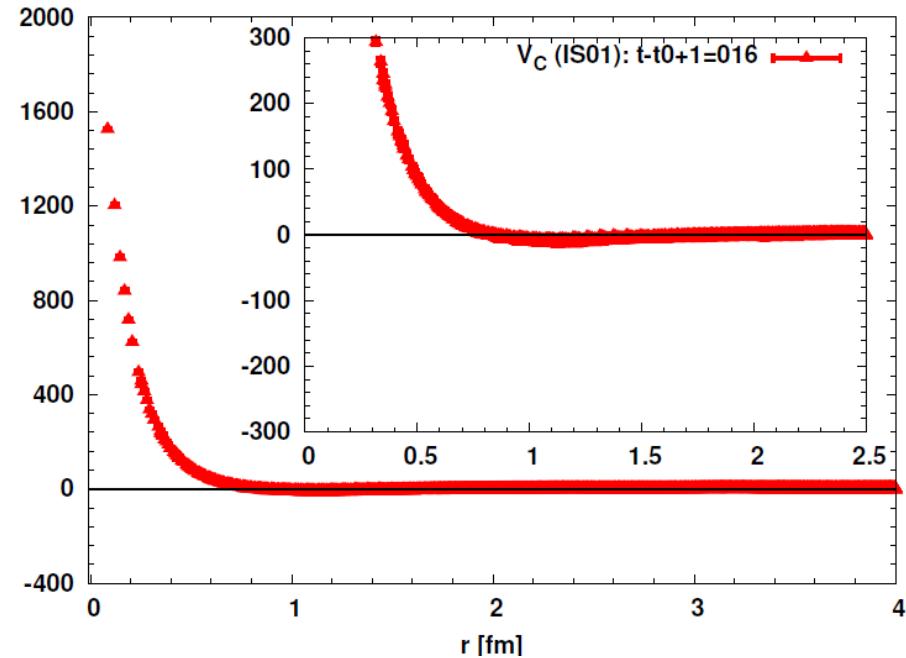


Comparison of central potential in 3S_1 channel

$\Xi_{cc} - \Xi_{cc}$



$\Xi_{ss} - \Xi_{ss}$



Completely different!

(N.B. There exists open channel, $\Lambda_c \Omega_{ccc}$, below $\Xi_{cc}\Xi_{cc}$ channel)

Challenge toward LQCD w/ physical mass

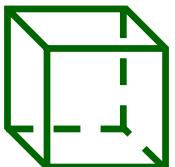
~2012



→ lighter u,d-quark masses
(=lighter pion mass M_π)

We were here

$M_\pi = 400\text{MeV}$
 $L = 3\text{fm}$



Simulation w/
Unrealistic mass

Theoretical
development
+
New supercomputer



Fugaku (440PFlops)

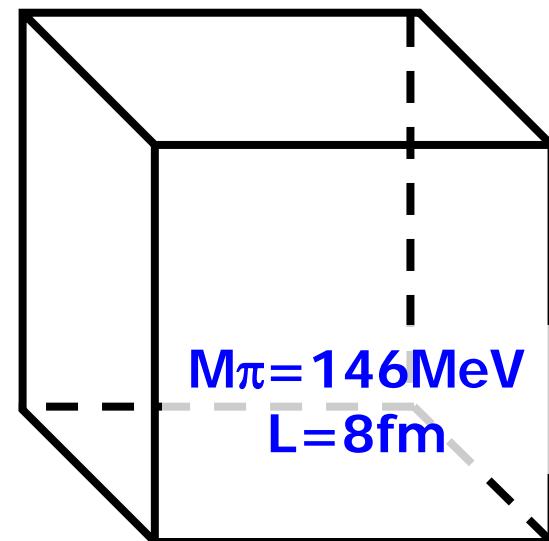
~2018



2023 !



Phys. point



Simulation w/
~Near physical mass

Challenge toward LQCD w/ physical mass

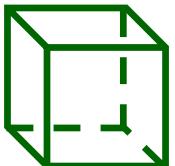
~2012



→ lighter u,d-quark masses
(=lighter pion mass M_π)

We were here

$M_\pi = 400\text{MeV}$
 $L = 3\text{fm}$



Simulation w/
Unrealistic mass

Theoretical
development

+

New supercomputer



Fugaku (440PFlops)

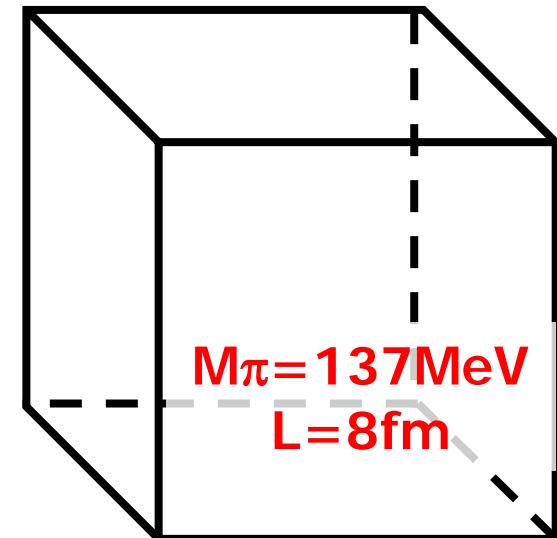
~2018



2023 !

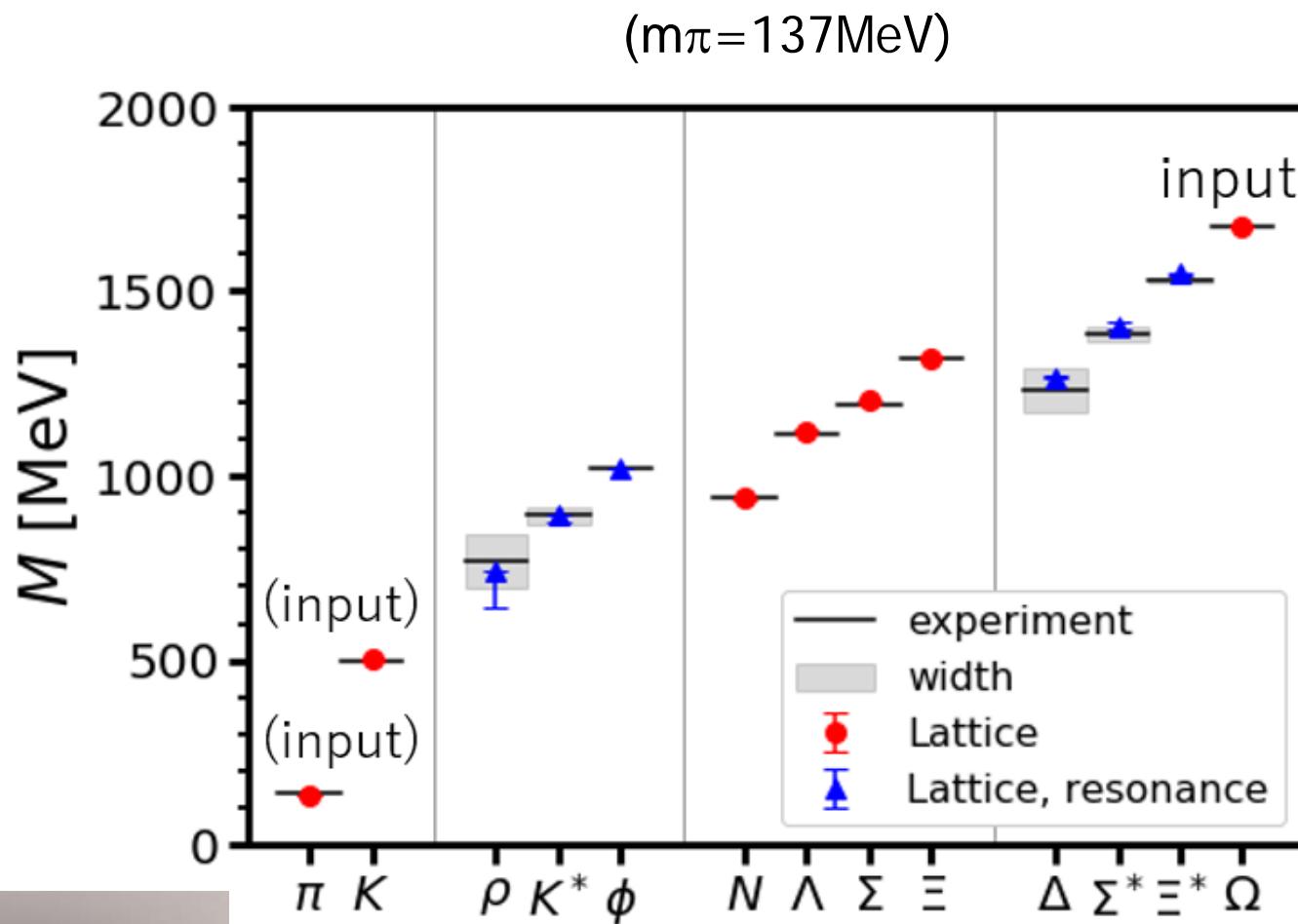


Phys. point



Simulation w/
the physical mass

LQCD simulations at the physical quark masses



E. Itou



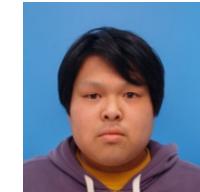
T. Aoyama



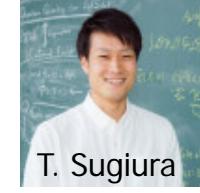
T. M. Doi



Y. Lyu



K. Murakami



T. Sugiura



Fugaku (440PFlops)

HAL QCD Coll., PRD in press, arXiv:2406.16665

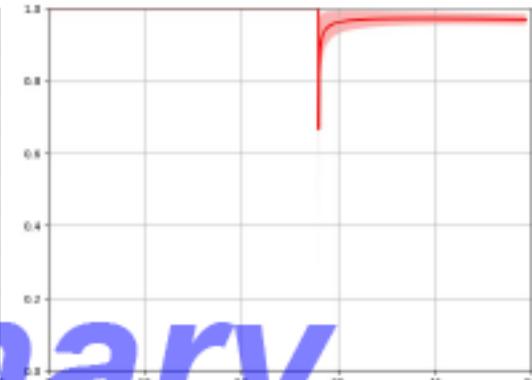
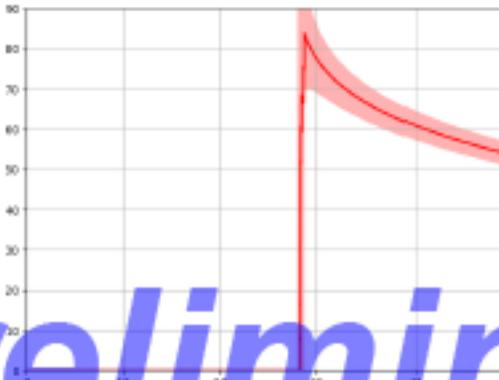
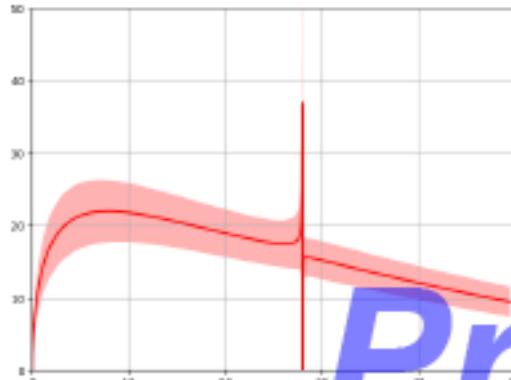
Calc of Hadron Interactions in progress!

Fate of H-dibaryon on the physical point

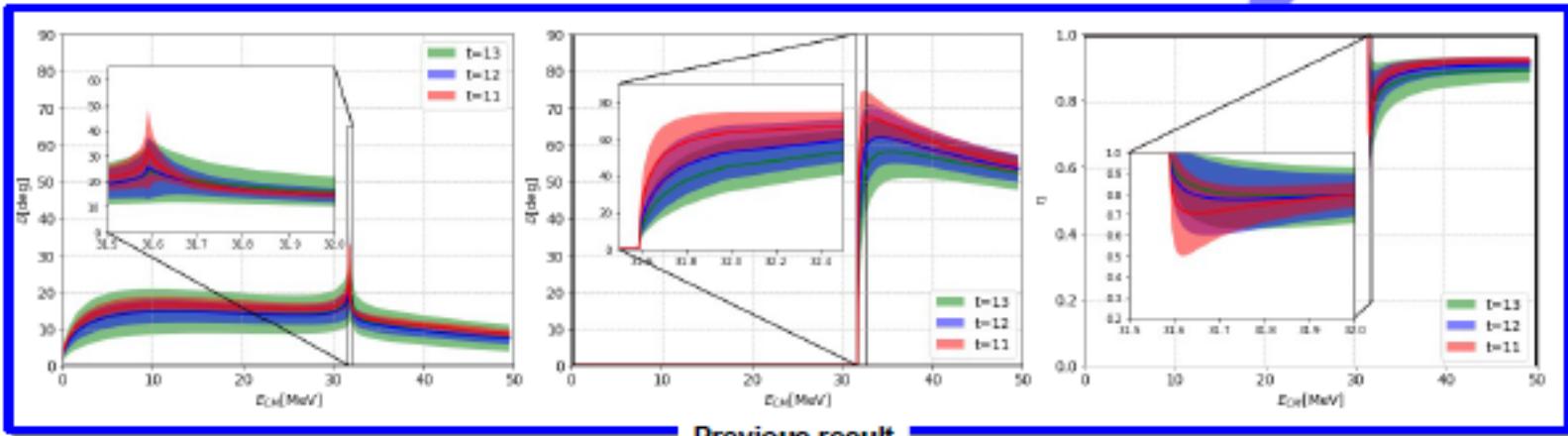
($M\pi=137\text{MeV}$)

$\Lambda\Lambda$ and $N\Xi$ phase shift and inelasticity

t=12

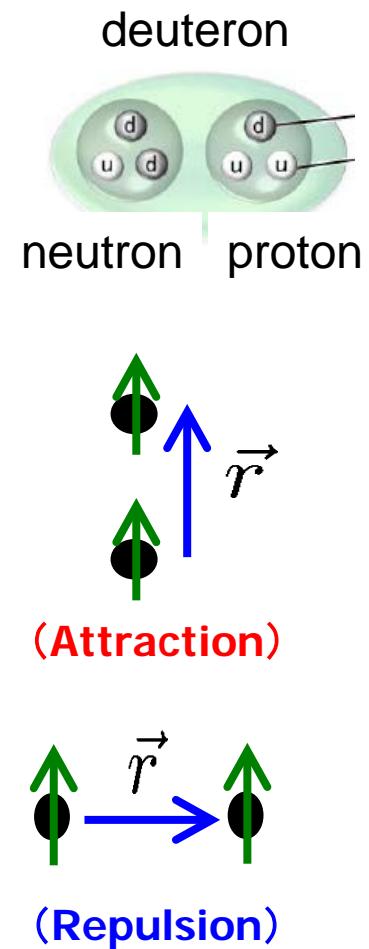
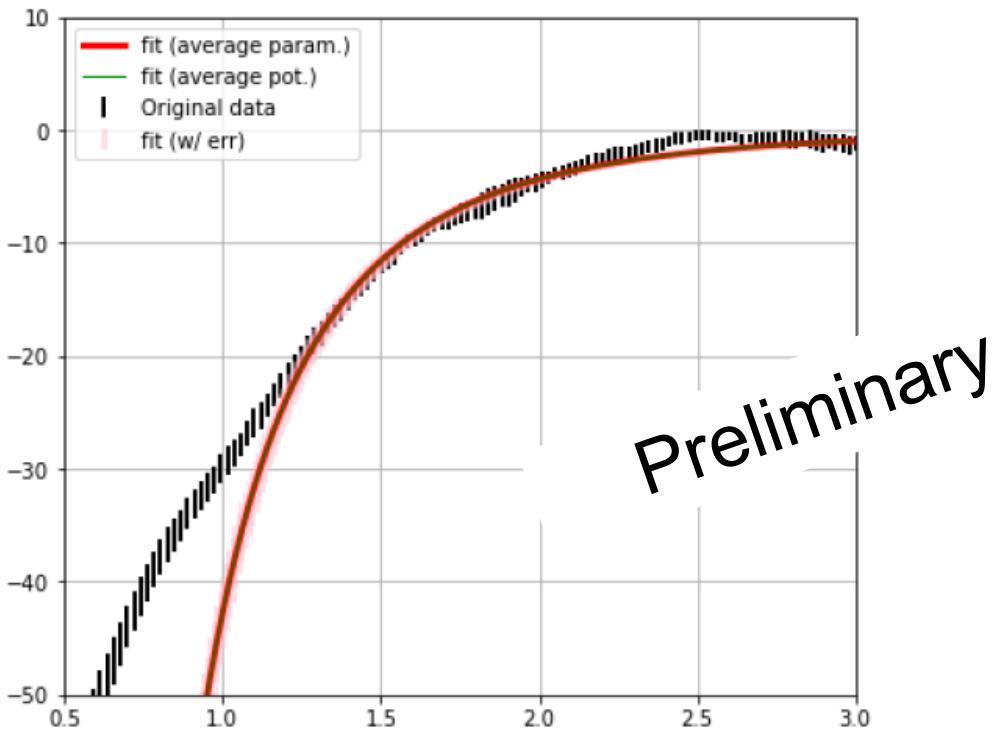


Preliminary



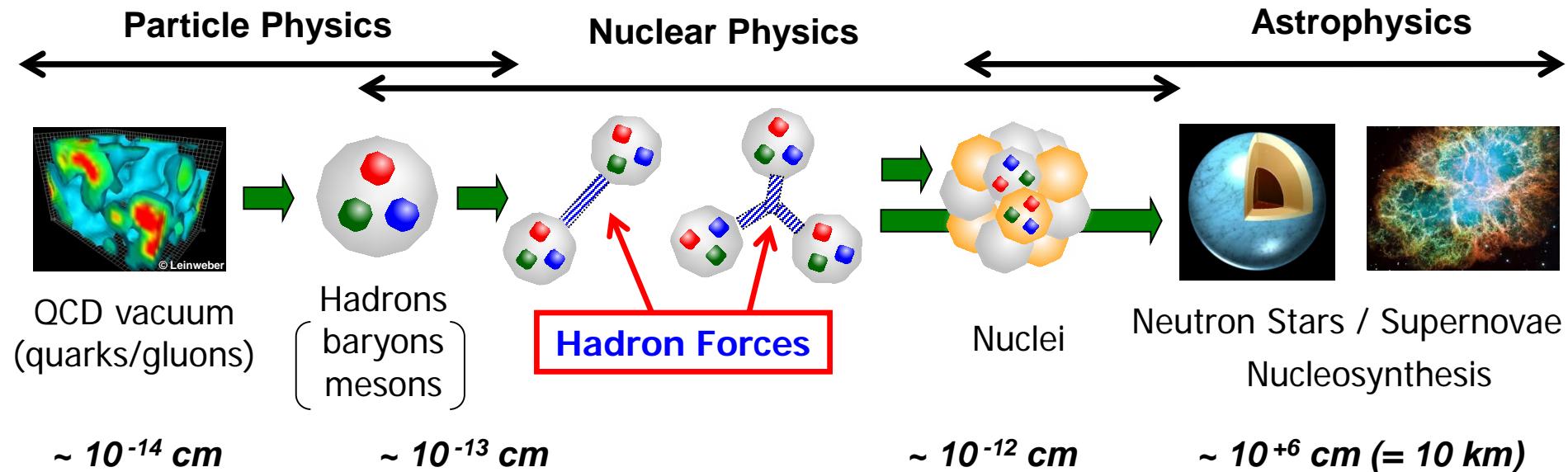
NN potentials @ physical point

Tensor force: one of the most important potentials
characteristics of Yukawa's pion theory



Consistent w/ tail structure of one-pion exchange potential

20XX : A Space Odyssey from Quarks to Nuclei & Universe



Nuclear Physics from QCD
New Era is dawning !

J-PARC (HEF-ex)
LHC (ALICE, LHCb)
Belle II
BES III
HIAF
RIBF/FRIB
LIGO/Virgo/KAGRA