

Phase diagram of QCD matter with magnetic field: domain-wall Skyrmion chain in chiral soliton lattice

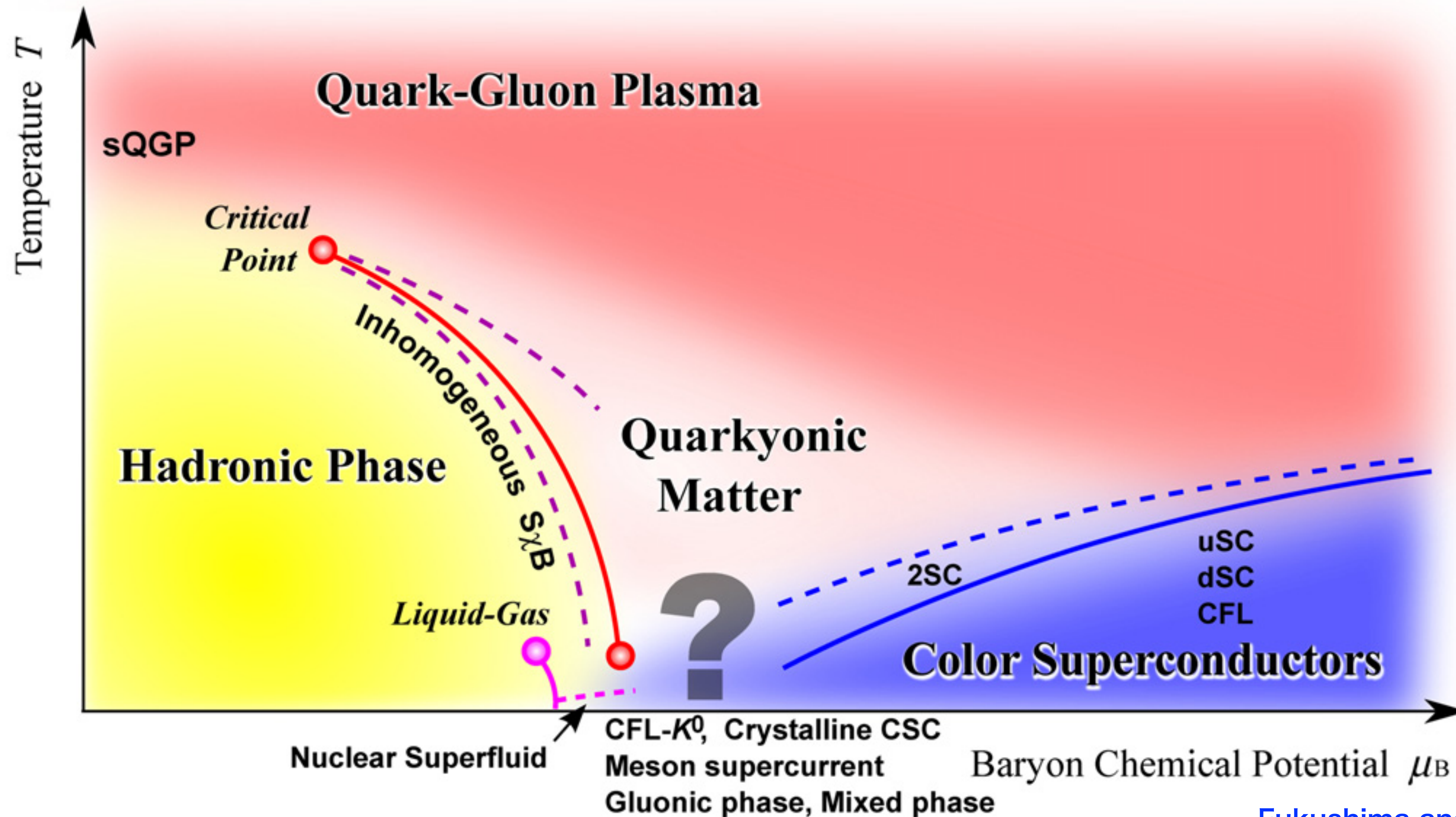
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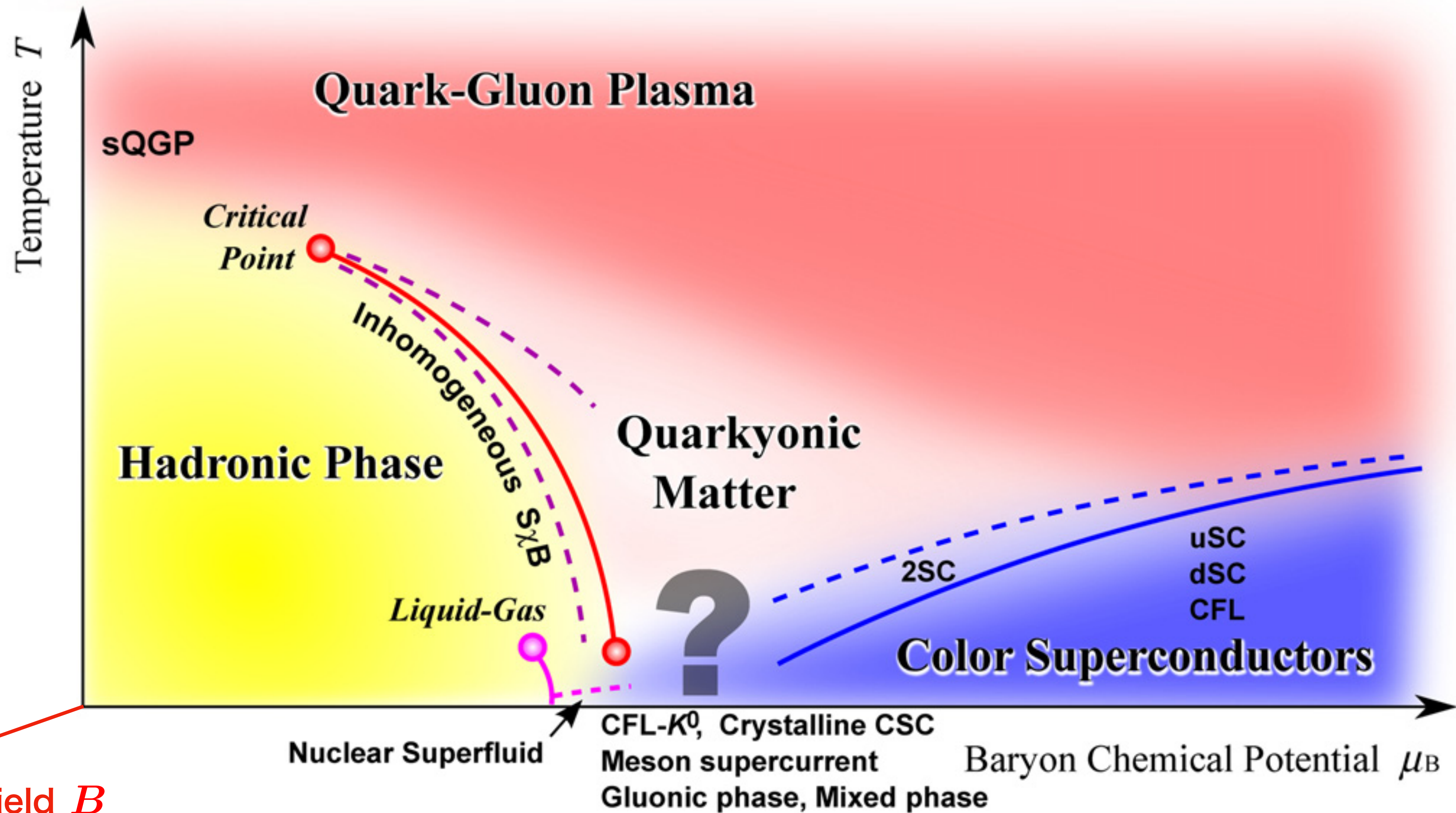
Compact stars in QCD phase diagram 2024/10/11

JHEP 12 (2023) 032

QCD phase diagram

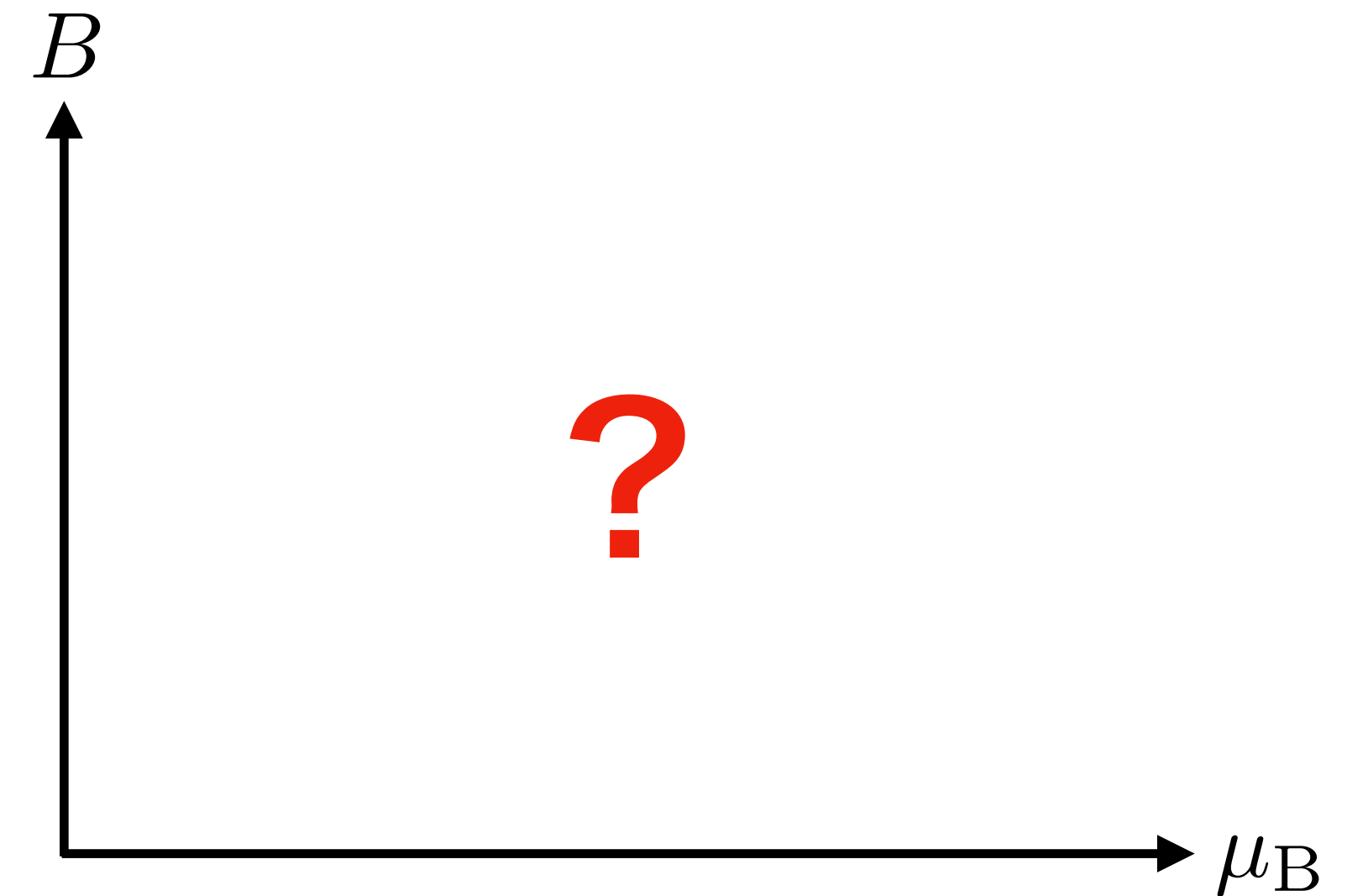


QCD phase diagram with B



Questions

- How is the phase diagram modified by B and μ_B ?
- Consider the region where B and μ_B are small, and chiral perturbation theory is valid.
- Temperature and isospin chemical potential are not considered.



- **Skyrmion** plays an important role to determine the phase structure.
- Since pions do not carry baryon number, nothing seems to happen even if μ_B is considered.

Chiral perturbation theory

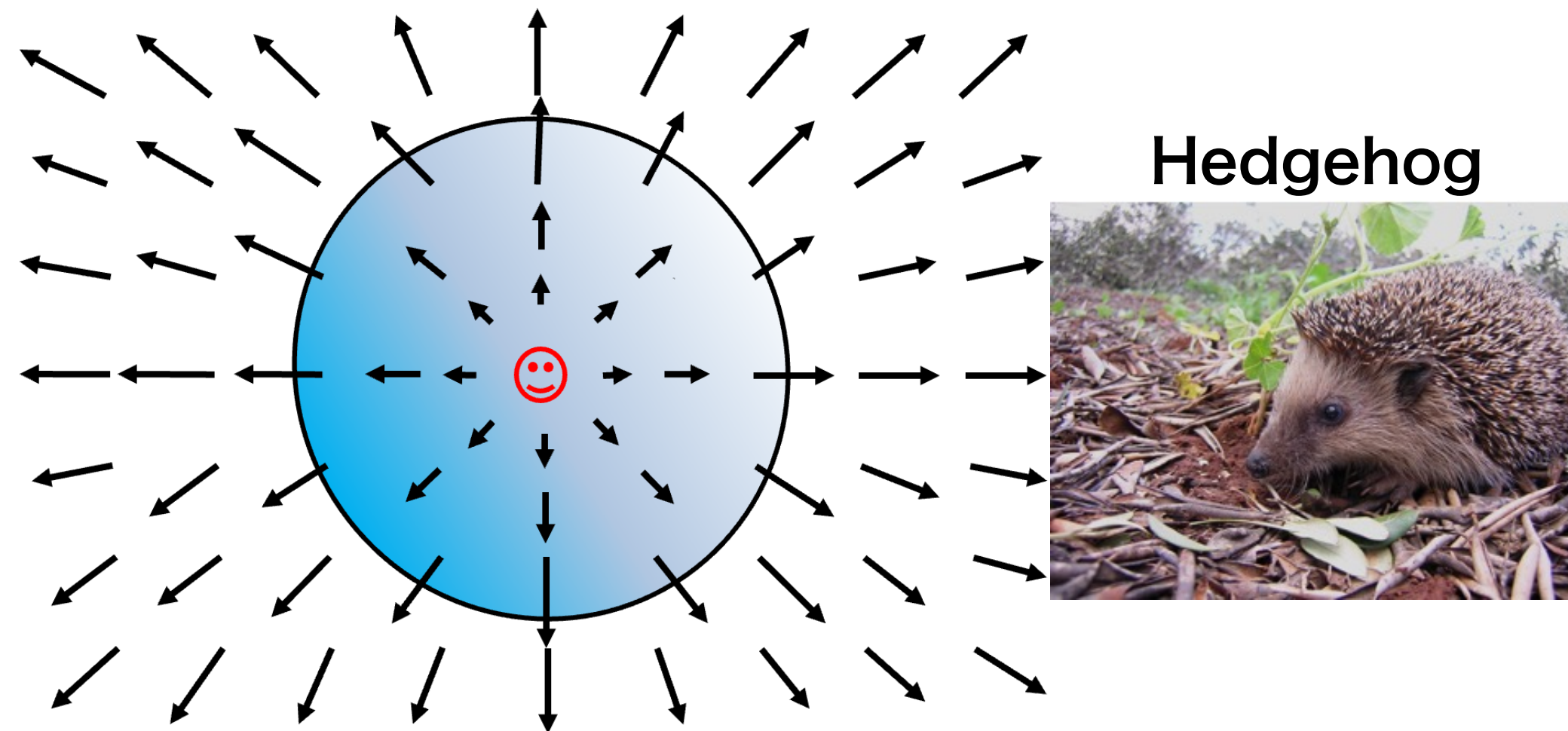
- **Order parameter is the chiral condensate:** $\langle \bar{q}q \rangle = |\langle \bar{q}q \rangle| \Sigma$
- **Nambu-Goldstone boson:** $\Sigma = \exp(i\sigma_a \phi_a)$, $\phi_a \equiv \pi_a / f_\pi$
- **Effective Lagrangian:**
$$\mathcal{L}_{\text{ChPT}} = \frac{f_\pi^2}{4} \text{tr} (D_\mu \Sigma D^\mu \Sigma^\dagger) - \frac{f_\pi^2 m_\pi^2}{4} (2 - \Sigma - \Sigma^\dagger)$$
$$D_\mu \Sigma = \partial_\mu \Sigma + iA_\mu [Q, \Sigma], \quad Q = \text{diag}(2/3, -1/3)$$

Skyrmion

- Can the baryons be made by pions (rather than quarks)?

- **Baryon as soliton = Skyrmion**

Skyrme (1961)



- Topological charge :

$$j_B^\mu = \frac{\epsilon^{\mu\nu\alpha\beta}}{24\pi^2} \text{tr}(\Sigma\partial_\nu\Sigma^\dagger\Sigma\partial_\alpha\Sigma^\dagger\Sigma\partial_\beta\Sigma^\dagger)$$

- \mathbb{R}^3 surrounds the configuration space of the pions $S^3 : \pi_3(S^3)$

ChPT w/ topological terms

- Baryon current couples to $U(1)_B$ gauge field: $\mathcal{L}_B = -A_B^\mu j_{B\mu}$, $A_B^\mu = (\mu_B, \mathbf{0})$
- “trial and error” $U(1)_{em}$ gauging while preserving baryon number conservation.

$$j_B^\mu = -\frac{\epsilon^{\mu\nu\alpha\beta}}{24\pi^2} \text{tr}\{L_\nu L_\alpha L_\beta - 3ie\partial_\nu [A_\alpha Q(L_\beta + R_\beta)]\} \quad L_\mu \equiv \Sigma\partial_\mu\Sigma^\dagger, R_\mu \equiv \partial_\mu\Sigma^\dagger\Sigma$$

Skyrimon charge
 $U(1)_{em}$ gauged part

$Q = \text{diag}(2/3, -1/3)$

Son and Stephanov (2008); Goldstone and Wilczek (1981); Witten (1983)

Chiral Soliton Lattice

- SG theory w/ total derivative

$$\mathcal{H} = \frac{f_\pi^2}{2} (\partial_z \phi_3)^2 + f_\pi^2 m_\pi^2 (1 - \cos \phi_3) - \frac{e\mu_B}{4\pi^2} B \partial_z \phi_3$$

- Consider only $\pi_0 : \Sigma = e^{i\phi_3 \tau_3}$
- z-dependence of ϕ_3 is nontrivial.

- EOM : $\partial_z^2 \phi_3 = m_\pi^2 \sin \phi_3$

- Energy : $E = \int_{-\infty}^{\infty} dz \mathcal{H} = 8m_\pi^2 f_\pi - \frac{e\mu_B B}{2\pi}$

- Critical B : $B_{\text{CSL}} = \frac{16\pi m_\pi f_\pi^2}{e\mu_B}$

$B\hat{z}$

$\pi_0 \text{ DW} : \phi_3 = 4 \tan^{-1} e^{m_\pi z}$

Chiral Soliton Lattice

- SG theory w/ total derivative

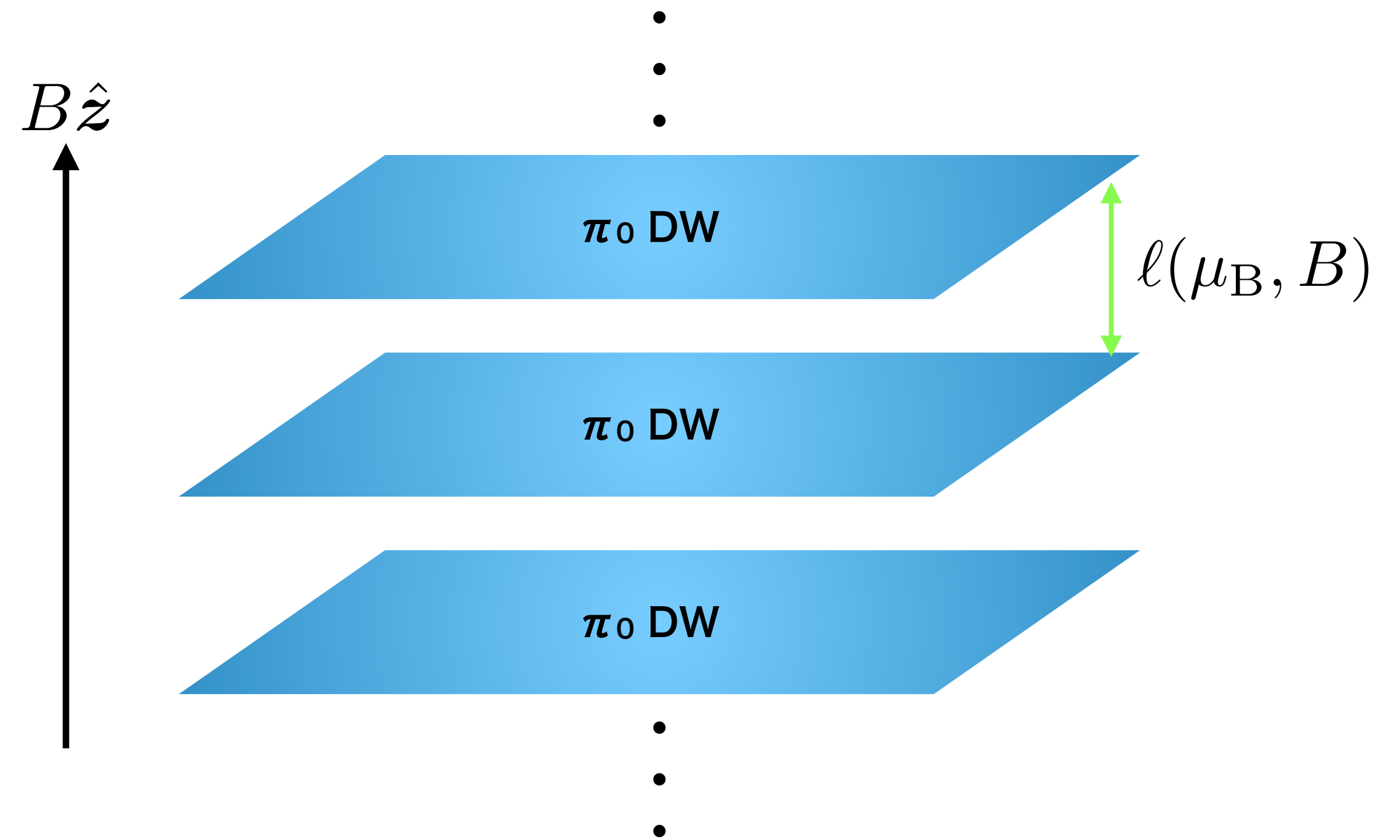
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- Pack many DWs in ground state!
- Impossible to pack due to the repulsive force.

Dautry and Nyman (1979); Hatsuda (1986); Son and Stephanov (2008); Nishiyama, Kawasaki and Tatsumi (2015); Brauner and Yamamoto (2017)

Non-Abelian soliton

- The single soliton:

$$\Sigma_0 = e^{i\sigma_3\theta}, \quad \theta = 4\tan^{-1}e^{m_\pi z}$$

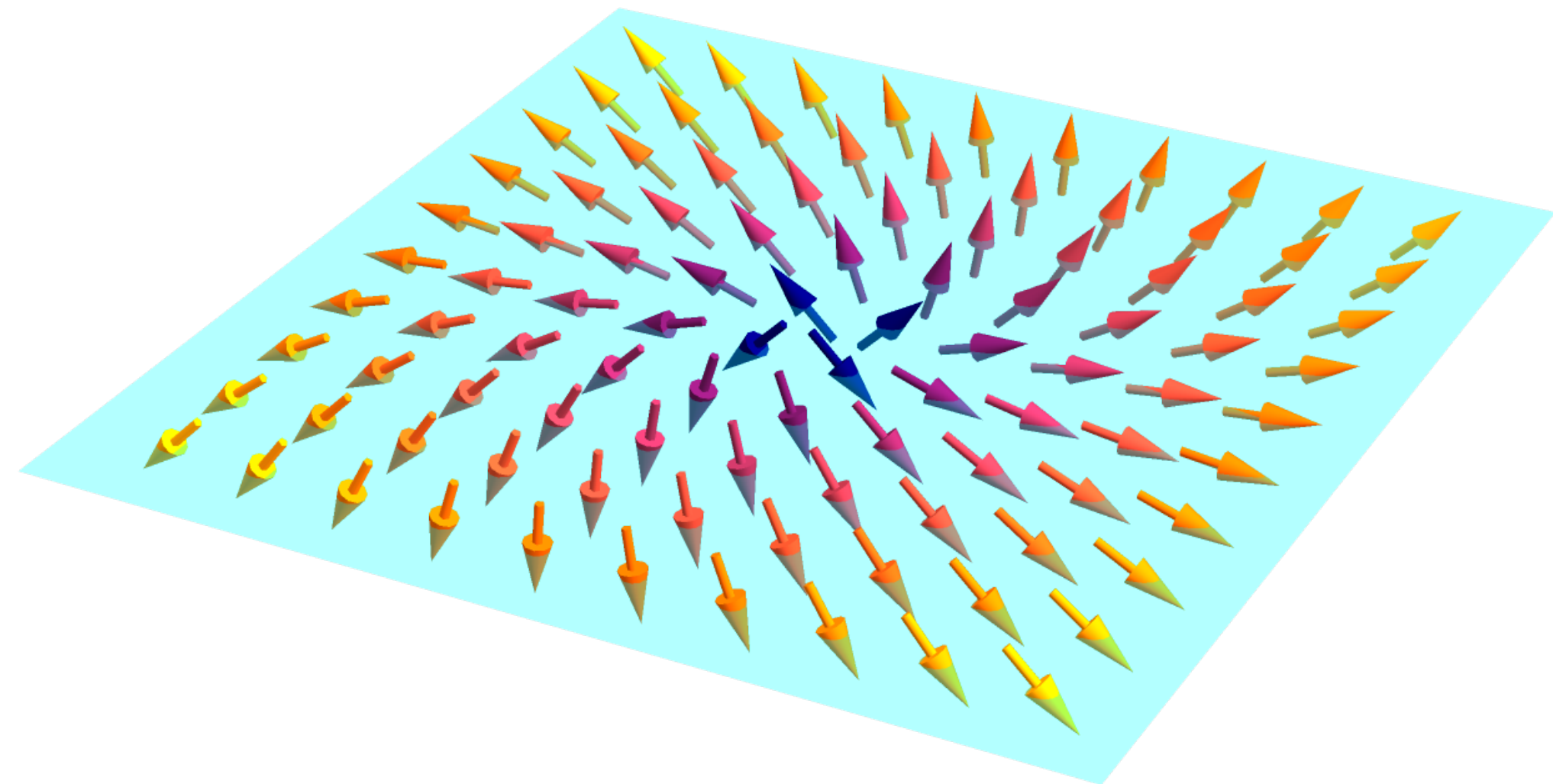
- More general solution :

$$\Sigma = g\Sigma_0g^\dagger = \exp(i\theta g\tau_3g^\dagger)$$

- SSB of $SU(2)_V \rightarrow U(1)$

- Σ_0 is invariant under $g = e^{i\tau_3\theta}$

S² moduli on the domain wall



- The collective coordinate : $\phi \in \mathbb{C}^2, \quad \phi^\dagger\phi = 1$

Nitta (2015); Eto and Nitta (2015)

$$g\sigma_3g^\dagger = 2\phi\phi^\dagger - 1$$

EFT for S^2 moduli

• **Effective Lagrangian :** $\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{const}} + \mathcal{L}_{\text{kin}} + \mathcal{L}_{\text{topo}}$ [Eto, KN and Nitta, JHEP 12 \(2023\) 032](#)

• **Kinetic term :** $\mathcal{L}_{\text{kin}} = \mathcal{C}(\ell)[(\phi^\dagger D_\alpha \phi)^2 + D^\alpha \phi^\dagger D_\alpha \phi]$

• **Topological terms :** $\mathcal{L}_{\text{topo}} = -2\mu_B q + \frac{e\mu_B}{2\pi} \epsilon^{03jk} \partial_j [A_k (1 - n_3)]$ O(3) nonlinear sigma model
 $n_a \equiv \phi^\dagger \sigma_a \phi \quad |\mathbf{n}| = 1$

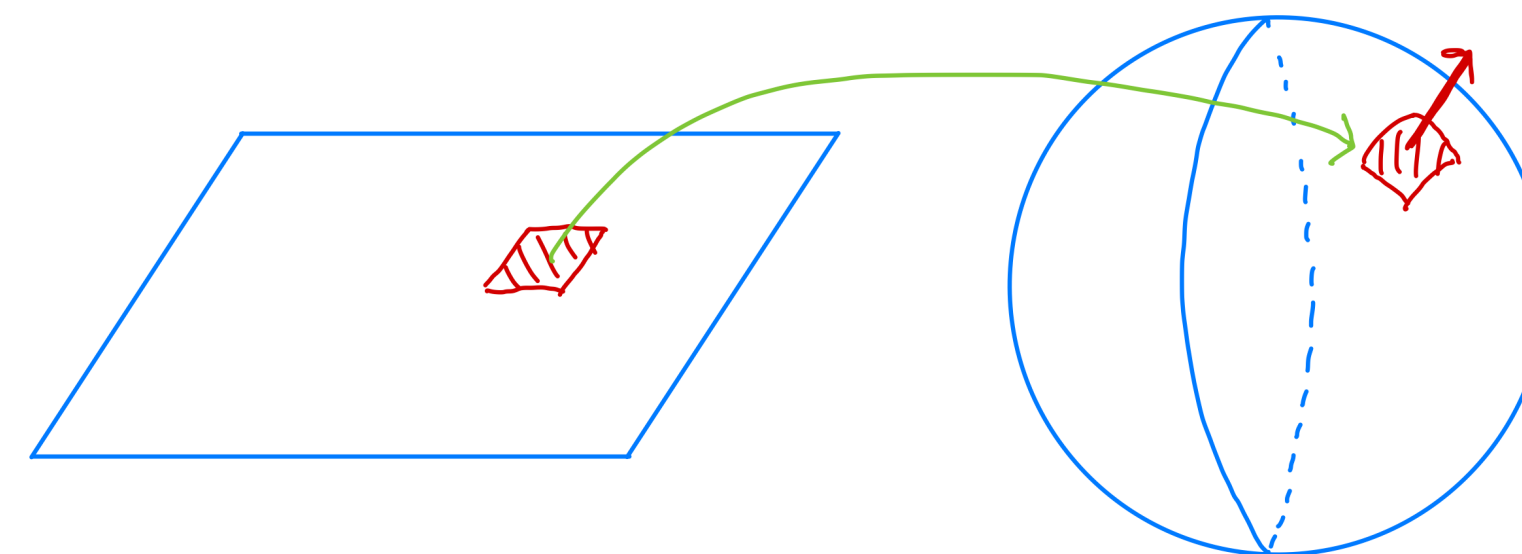
- The red term stabilizes the configuration with finite k!

$\pi_2(S^2)$ topological charge (counting how many times xy plane covers S^2 moduli)

$$k = \int d^2x q$$

$$= \frac{1}{4\pi} \int \mathbf{n} \cdot \left(\frac{\partial \mathbf{n}}{\partial x} dx \times \frac{\partial \mathbf{n}}{\partial y} dy \right)$$

$\in \mathbb{Z}$



Bogomol'nyi bound

- Baby Skyrmion naturally appears when minimizing the Hamiltonian.

$$\mathcal{H}_{\text{DW}} = \frac{\mathcal{C}(\kappa)}{4} \partial_i \mathbf{n} \cdot \partial_i \mathbf{n} + 2\mu_B q - \frac{e\mu_B}{2\pi} \epsilon^{03jk} \partial_j [A_k (1 - n_3)]$$

- Completing the square of the kinetic term is useful!

$$\begin{aligned} (\partial_i \mathbf{n})^2 &= \frac{1}{2} \underbrace{(\partial_i \mathbf{n} \pm \epsilon_{ij} \mathbf{n} \times \partial_j \mathbf{n})^2}_{= 0} \pm 8\pi q \geq \pm 8\pi q \\ &\rightarrow \text{BPS equation} \rightarrow \text{Baby Skyrmion!} \end{aligned}$$

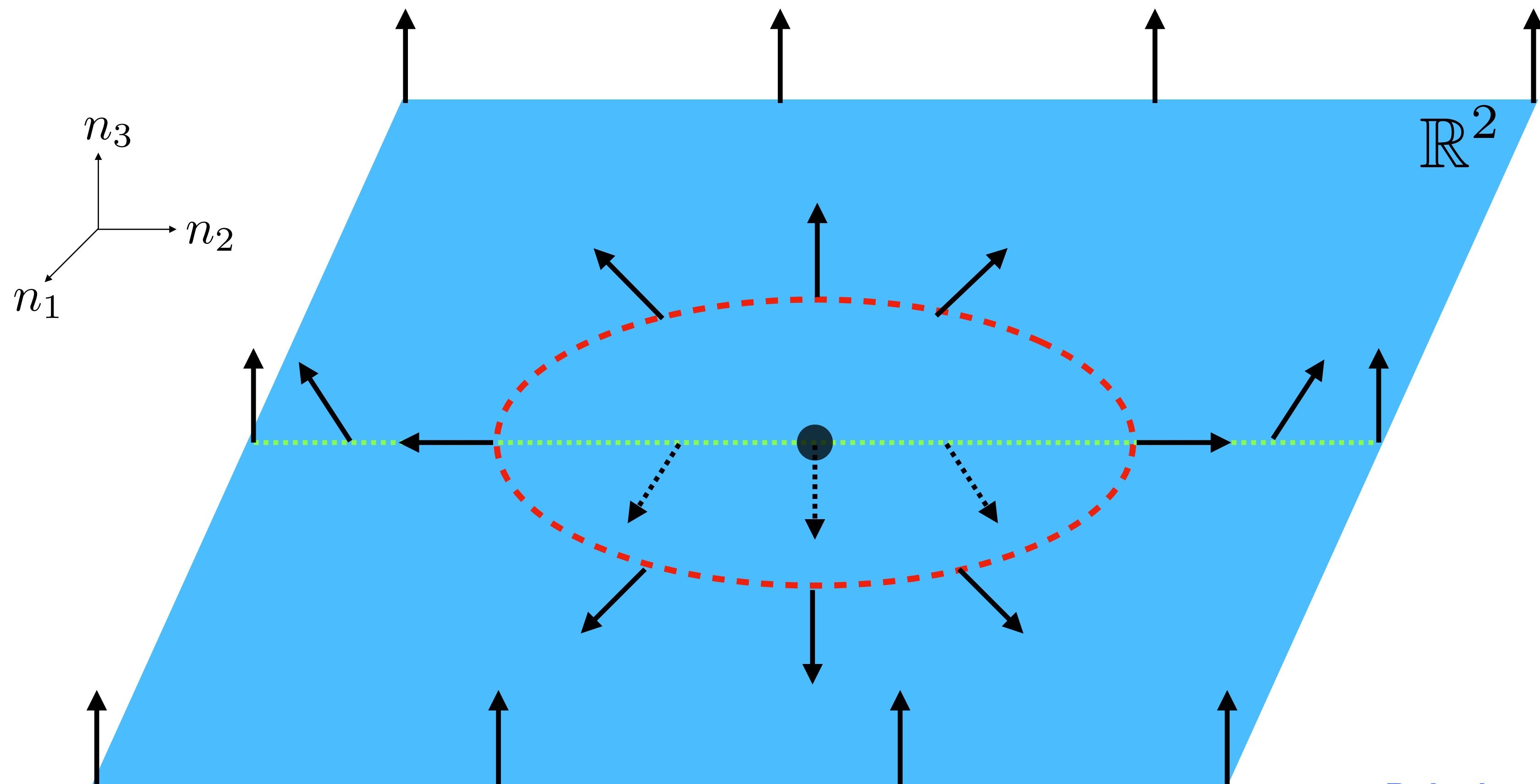
- **Critical μ_B :** $E_{\text{DW}} \geq \underbrace{2\pi\mathcal{C}(\kappa)|k| + 2\mu_B k}_{\text{negative}} - \underbrace{\frac{e\mu_B}{2\pi} \int d^2x \epsilon^{03jk} \partial_j [A_k (1 - n_3)]}_{\text{positive}}$

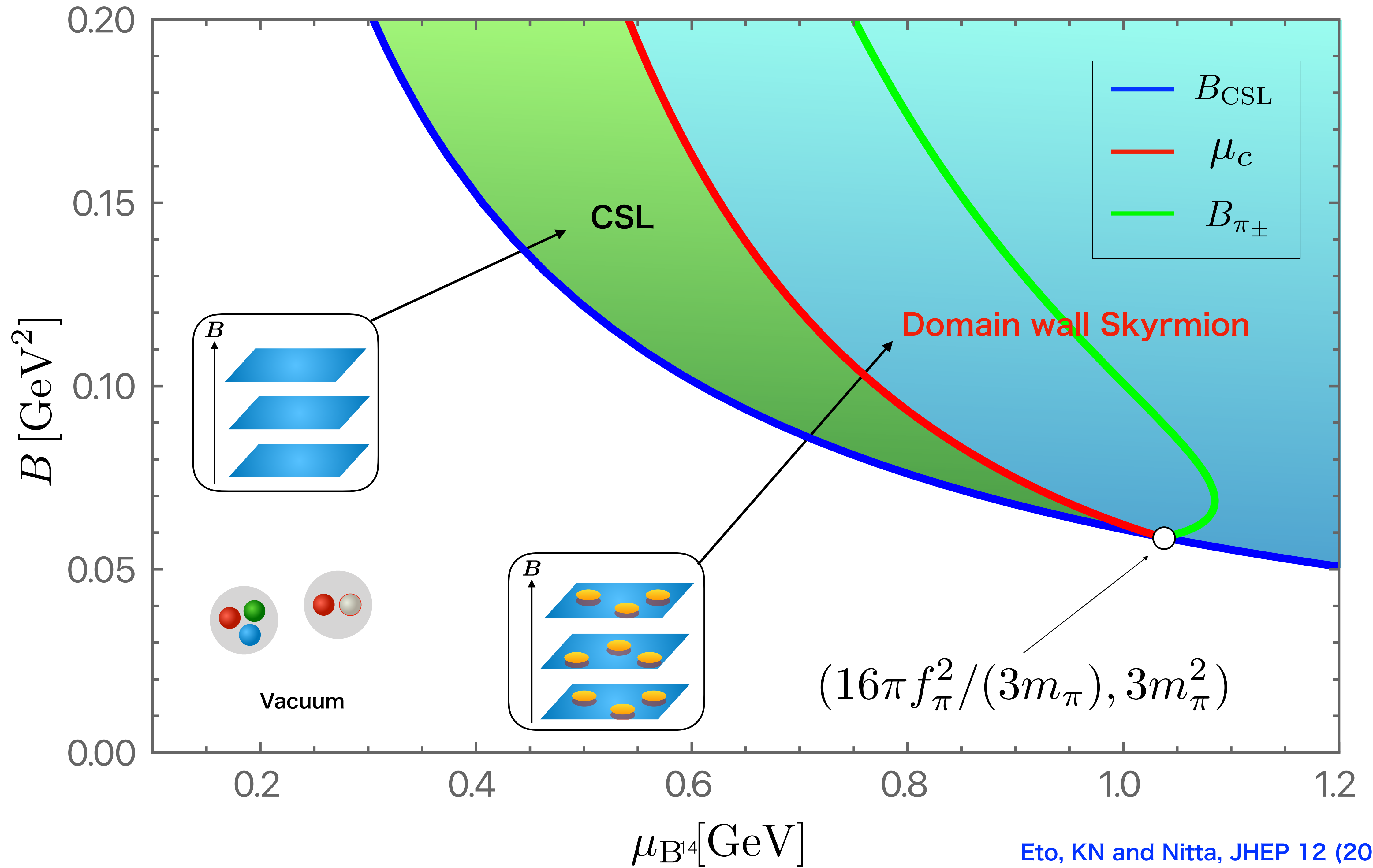
$$\mu_B > \pi\mathcal{C}(\kappa)$$

The total energy is negative $\mu > \mu_c$, and baby Skyrmion appears in the ground state! Some constraints on the lump

Baby Skyrmion

- Configuration on DW surrounding S^2 : $\uparrow = n_a \quad n^2 = 1$





Summary

- Pions couple to baryon as Skyrmion.
- At $B > B_c$, the stack of π_0 DWs is energetically stable.
- At $\mu > \mu_c$, baby Skyrmion appears on π_0 DWs.

Thank you for your attention!