

# New Vacuum Solutions and Black Strings in the 4D Standard Model

Yu Hamada<sup>1</sup>, Yuta Hamada<sup>2,3</sup>, Hayate Kimura<sup>3</sup>

<sup>1</sup>DESY, <sup>2</sup>KEK, <sup>3</sup>SOKENDAI

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## Short summary

String theory  $\xrightarrow{\text{compactify}}$  enormous lower dimensional vacua (String landscape)

The same is true for standard model : SM landscape

Ex : compactified on  $S^1$   $\rightarrow$  stabilized by Casimir energy of SM particles

On the other hand,

this  $S^1$  vacuum is reproduced near the horizon of “quantum” black string in our 4D world

Sec. 1~3 are the review of [Arkani-Hamed et. al. 07]

Compactified SM  
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Radion Stabilization  
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Interpolating BB  
ooooooo

Numerical Calculation  
ooo

Conclusion  
ooo

We construct this black string solution  
by numerical calculation

Compactified SM  
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Radion Stabilization  
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Interpolating BB  
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# Compactified SM

# The action for SM+GR

Our starting point is 4D SM+GR action :

$$S = \int d^4x \sqrt{-g^{(4)}} \left( \frac{1}{2} M_4^2 \mathcal{R}^{(4)} - \Lambda_4 + \dots \right)$$

$\mathcal{R}^{(4)}$  : 4D Ricci scalar     $M_4$  : 4D Planck mass

$\Lambda_4$  : 4D cosmological constant

... include other fermions, gauge fields in SM

# Kaluza Klein mechanism for 4D metric

Compactify one spacial dimension on  $S^1$  by Kaluza-Klein mechanism.

The metric components are decomposed :

$$g_{\mu\nu}^{(4)} \rightarrow g_{ij}^{(3)}, A_i, \textcolor{red}{R} \quad (i, j = 0, 1, 2)$$

$R$  :  $S^1$  radius in 4D  $\rightarrow$  scalar in 3D (radion)

integrate out the compact coordinate  $x_3 \rightarrow$  effective 3D action :

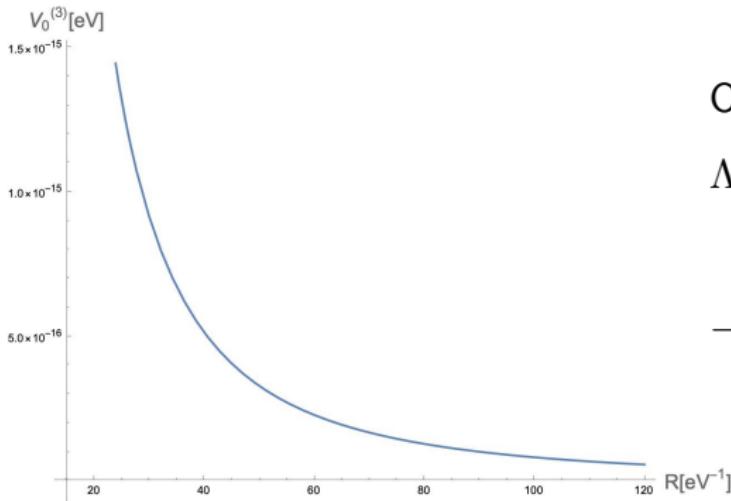
$$\begin{aligned} S = \int d^3x \sqrt{-g^{(3)}} & \left[ \frac{1}{2} r M_4^2 \mathcal{R}^{(3)} - r M_4^2 \left( \frac{\partial R}{R} \right)^2 \right. \\ & \left. - \frac{r M_4^2}{8} \left( \frac{2\pi R}{r} \right)^2 R^2 F_{ij} F^{ij} - \frac{r^3 \Lambda_4}{(2\pi R)^2} + \dots \right] \end{aligned}$$

# Radion potential ?

Classical potential for the radion is the cosmological const. term :

$$V_0^{(3)}(R) = \frac{r^3 \Lambda_4}{(2\pi R)^2}$$

The radion field (moduli of  $S^1$ ) can not be stabilized.



On the other hand ...

$\Lambda_4$  is sufficiently small

$$(\sim 3.25 \times 10^{-11} \text{ eV}^4)$$

→ quantum correction should be taken into account

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## Radion Stabilization

# Casimir energy as radion potential

(constant) expectation values for scalars with quantum correction

= the stationary point of the effective potential

1-loop free fields correction for effective potential is following :

$$V^{(4)} = \frac{\text{classical potential}}{\frac{V_0^{(4)}}{2} - \frac{i}{2} \sum_{\text{particle}} (-1)^{2s_p} n_p \int \frac{d^4 k}{(2\pi)^4} \ln(m_p^2 + k^2)} \quad \text{Casimir potential} \equiv V_1^{(4)}$$

Note that one dimension is compactified

⇒ the momentum is quantized as  $\frac{2\pi n}{R}$  and  $\int \rightarrow \sum$

⇒  $R$  appear in  $V_1^{(4)}$  : **radion potential**

# Boundary condition

The sum is evaluated as following :

$$V_1^{(4)} = - \sum_{\text{particle}} (-1)^{2s_p} n_p \frac{m_p^4}{2\pi^2} \sum_{n=1}^{\infty} \frac{\cos(2\pi n \theta_p)}{(2\pi R m_p n)^2} K_2(2\pi R m_p n)$$

$\theta_p$  is the parameter which characterize the boundary conditions of fields :

$$\psi_p(x_3 + 2\pi R) = e^{i2\pi\theta_p} \psi_p(x_3)$$

Possible  $\theta_p$  reflects symmetry of the field :

$\psi_p$  has  $U(1)$  symmetry  $\Rightarrow \theta_p$  can be arbitrary value

gauge boson, graviton  $\Rightarrow$  periodic ( $\theta_p = 0$ )

# Stabilized radion ; new SM vacuum solution

particle with mass  $m_p$

→ Casimir potential is

relevant for  $R < 1/m_p$

Casimir energy of

(massless Bosons :

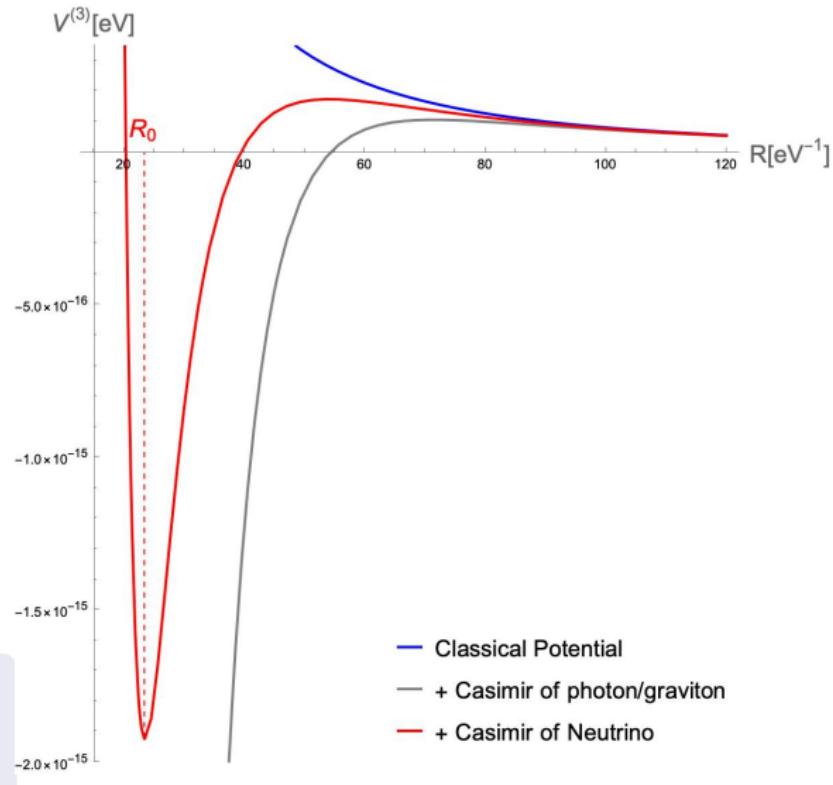
photon, graviton)

+ (Majorana Neutrinos)



New Vacuum Solution

$(S^1 \times AdS_3) !$

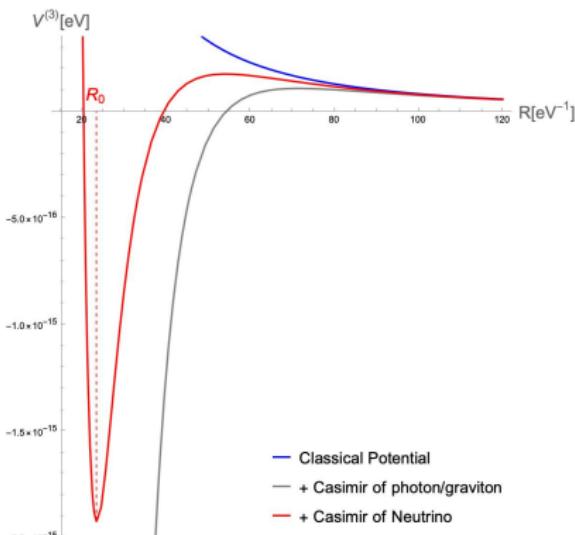


# Assumptions for stabilization

In previous solution, the assumption :

- ① Neutrino spectrum is normal hierarchy and lightest Neutrino is massless

- The existence and type ( $dS, AdS$ ) of solution is sensitive to Neutrino mass and dof
- Casimir of the other particle is not relevant in this scale



- ② The boundary condition of Majorana fields is **periodic** instead of anti-periodic

⇒ fields feel  $\mathbb{Z}_2$  flux on  $S^1$  (cf. Aharonov-Bhom effect)

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# Interpolating Black Brane

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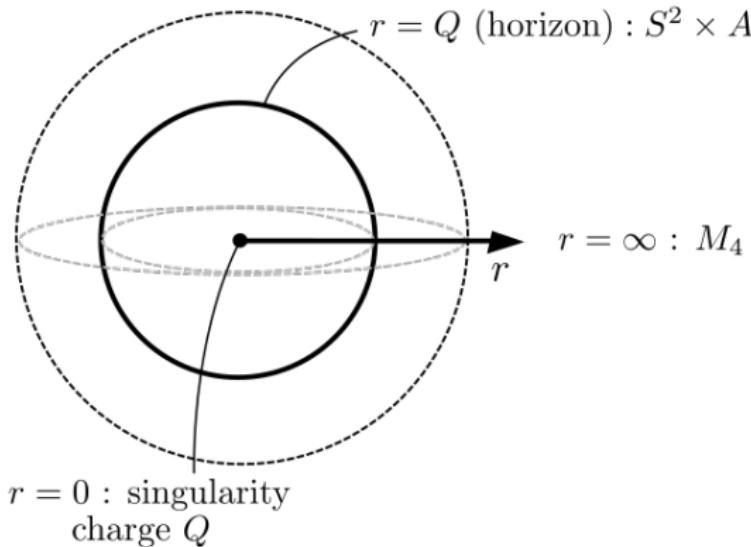
The important perspective :

Extremal charged black brane interpolate between non-compact spacetime and compact spacetime which is compactified by its charge flux

Then, is there any black brane which interpolate with previous vacuum?



# Example : Reissner-Nordström BH



$S^2 \times AdS_2$  : Vacuum solution stabilized by electric flux  $Q$

(potential for moduli :  $F_{\mu\nu}F^{\mu\nu}$ )

$$ds^2 = - \left(1 - \frac{Q}{r}\right)^2 dt^2 + \left(1 - \frac{Q}{r}\right)^{-2} dr^2 + r^2 d\Omega^2$$

$$\downarrow \xi = r - Q \ll Q$$

$$ds^2 \sim - \left(\frac{\xi}{Q}\right)^2 dt^2 + \left(\frac{Q}{\xi}\right)^2 d\xi^2 + Q^2 d\Omega^2$$

$: S^2 \times AdS_2$

# Analogy to the new vacuum

Our new vacuum is only one spacial dimension is compactified

⇒ consider an object extended 1D from the BH : **Black String**

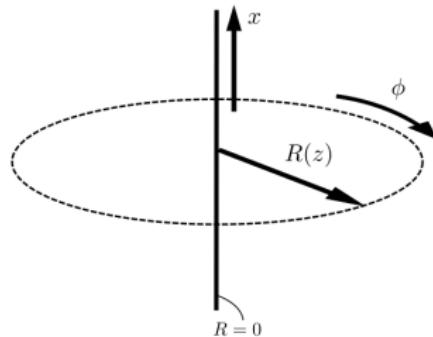
**$\mathbb{Z}_2$  charged Black String should reproduce the new  $S^1 \times AdS_3$  vacuum near its horizon**

metric ansatz :

$$ds^2 = A^2(z)(-dt^2 + dx^2) + dz^2 + R^2(z)d\phi^2$$

What we have to do :

$$\text{Solve Einstein eq. with } T^{\mu\nu} = \frac{-2\pi}{\sqrt{-g^{(4)}}} \frac{\delta}{\delta g_{\mu\nu}^{(4)}} \int dx^4 \sqrt{-g^{(4)}} (-\Lambda_4 - V_1^{(4)})$$



# Einstein equation

Einstein equations are written as following :

$$\frac{A'}{A} = -\frac{R'}{R} - \sqrt{\frac{R'^2}{R^2} - \frac{V_1^{(4)} + \Lambda_4}{M_4^2}}, \quad R'' + \gamma R' = -\partial_R U(R)$$

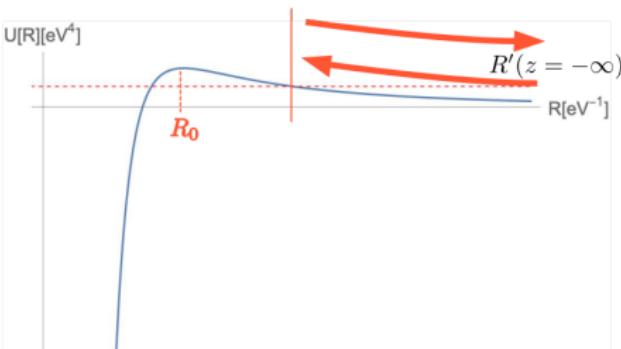
$$\left( \partial_R U(R) = \frac{R}{M_4^2} \left( V_1^{(4)} + \Lambda_4 - \frac{1}{2} R \partial_R V_1^{(4)} \right), \quad \gamma = -2 \left( \frac{R'}{R} + \sqrt{\frac{R'^2}{R^2} - \frac{V_1^{(4)} + \Lambda_4}{M_4^2}} \right) \right)$$

\* For simplicity, discuss the  $\Lambda_4 = 0$  case

$R$  EOM : classical mechanical EOM

with friction  $\gamma$  and potential  $U(R)$

⇒ shooting problem



# Extremal black string solution ?

For  $A$ ,

- consistent with expected

boundary condition at

$z \rightarrow -\infty$  :

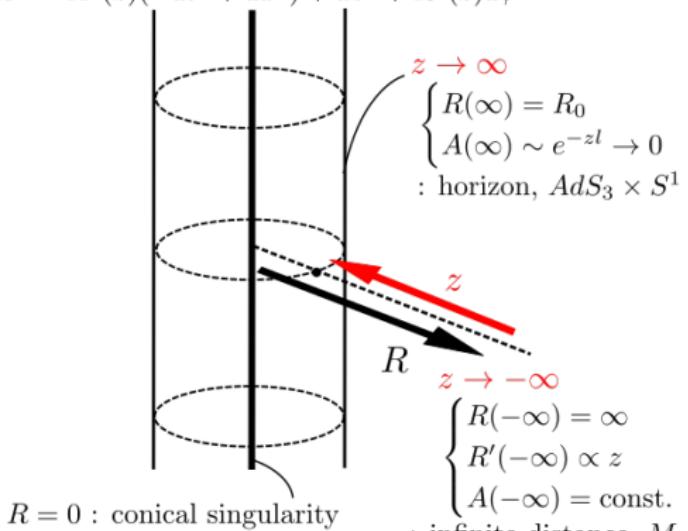
$$A'|_{z=-\infty} = 0, \quad R|_{z=-\infty} \propto z$$

- When  $R \equiv R_0$ ,

$$A \sim e^{-z} \sqrt{V_1^{(4)}(R_0)/M_4} :$$

$$AdS_3 \times S^1$$

$$ds^2 = A^2(z)(-dt^2 + dx^2) + dz^2 + R^2(z)d\phi^2$$



“Stopping solution” of  $R$  EOM should be extremal solution

Compactified SM  
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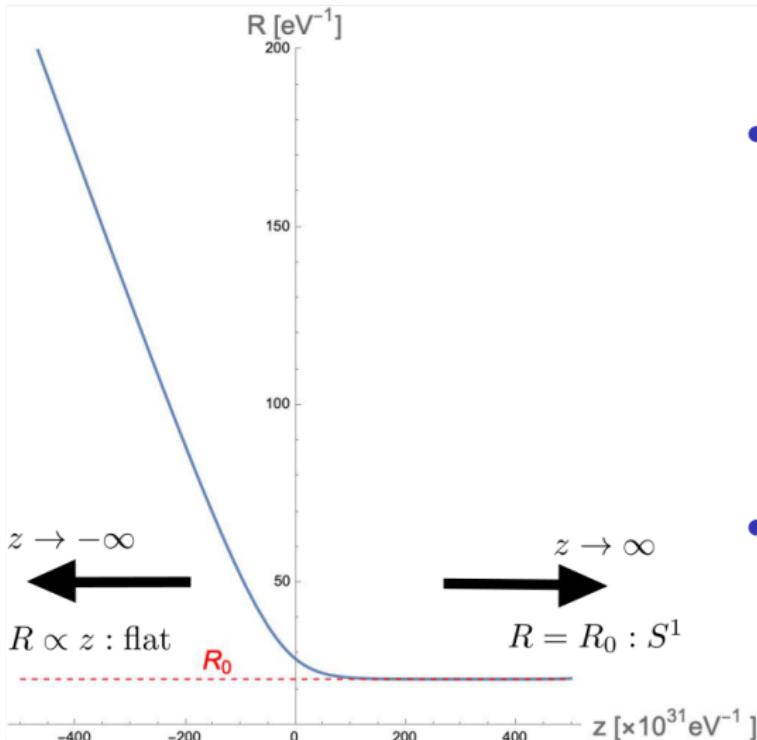
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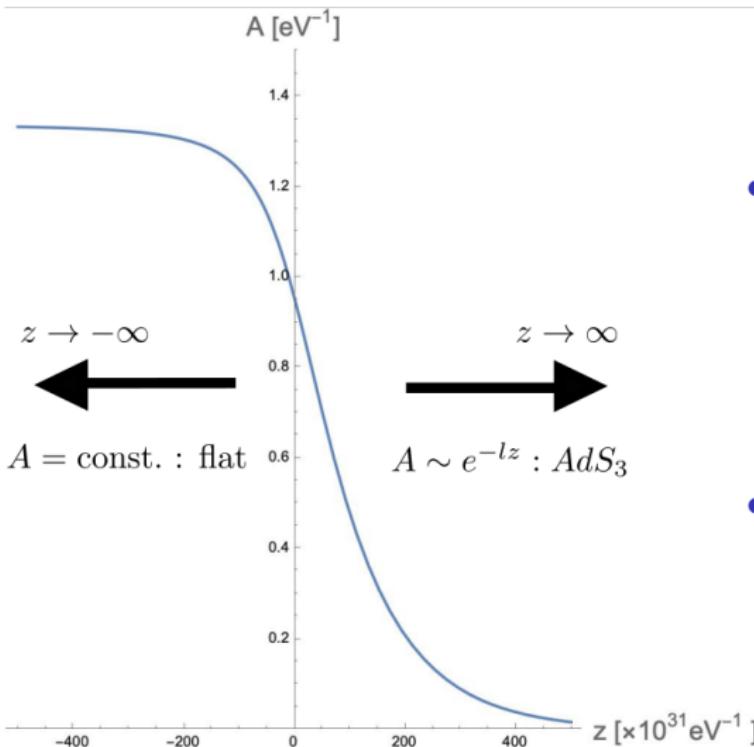
Conclusion  
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## Numerical Calculation

$\Lambda_4 = 0$  case :  $R(z)$ 

- This solution indeed interpolates the flat part ( $z \rightarrow -\infty$ ) and the part  $R$  is constant ( :  $S^1$  is decoupled from  $ds^2$ ) ( $z \rightarrow \infty$ )

- $|R'(-\infty)| \sim 4.3 \times 10^{-32}$   
 $\sim \frac{m_\nu}{M_4}$

$\Lambda_4 = 0$  case :  $A(z)$ 

- This solution indeed interpolates flat part and  $AdS$  part : when  $z \rightarrow \infty$ ,  $A \sim e^{-lz}$ ,  $l \sim 8.4 \times 10^{-34} [\text{eV}]$
- $A \rightarrow 0$  when  $z \rightarrow \infty$  :  $z \rightarrow \infty$  corresponds to the horizon

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## Conclusion

# Conclusion

- By using of Casimir energy, we can obtain the new vacuum solution in SM : SM vacuum is not unique, but there is SM Landscape
- The compact SM vacuum is reproduced near the horizon of a 4D black string

# Future direction

- asymptotically  $dS_4$
- non-extremal black string
- cosmological argument ?
- $E_8 \times E_8$  heterotic SUGRA