

## Summary

- We discuss if the two LFV decay modes of  $h(125)$  can be enhanced or not in generalized DFSZ axion models.
- In the case of non-chiral alignment Yukawa couplings ( $\Gamma_{\ell\ell'}^R \neq \Gamma_{\ell'\ell}^R$ ),  $h \rightarrow e\tau$  and  $h \rightarrow \mu\tau$  can be large simultaneously.
- The parameter region where hLFVs are enhanced can be explored by the searches of axion LFV couplings.

## Lepton Flavor violation (LFV)

### Representative process

- Lepton flavor violating (cLFV) decays:  $e \rightarrow \mu\gamma$ ,  $\mu \rightarrow 3e$ , ...
- Higgs LFV decays:  $h \rightarrow e\tau$ ,  $h \rightarrow \mu\tau$ ,  $h \rightarrow e\tau$

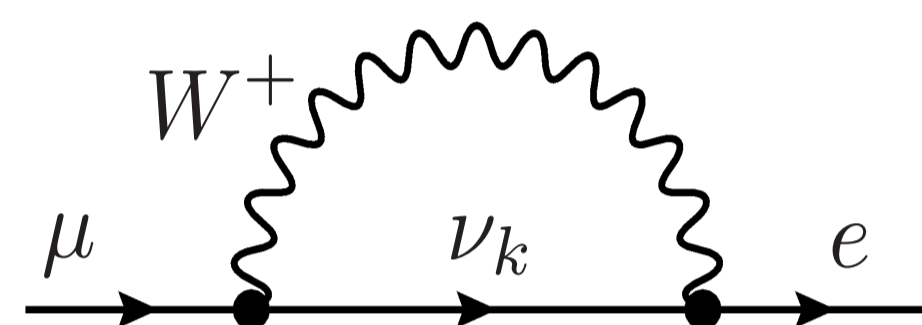
### In the framework of the SM + $m_\nu \bar{\nu}_L^c \nu_L$

- Lepton current appears in the charged current:

$$-\mathcal{L}_{cc} = \frac{g}{\sqrt{2}} (\bar{e}_L, \bar{\mu}_L, \bar{\tau}_L) \gamma^\mu U_{PMNS} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} W_\mu^+$$

- Due to the GIM mechanism, LFV appears at 1-loop level

e.g.)  $\mu \rightarrow e\gamma$



$$\text{BR} = \frac{3\alpha}{32\pi} \left| \sum_{k=1}^3 \frac{(U)_{\mu k} (U)_{ek}^* m_{\nu_k}^2}{m_W^2} \right|^2 = \mathcal{O}(10^{-54})$$

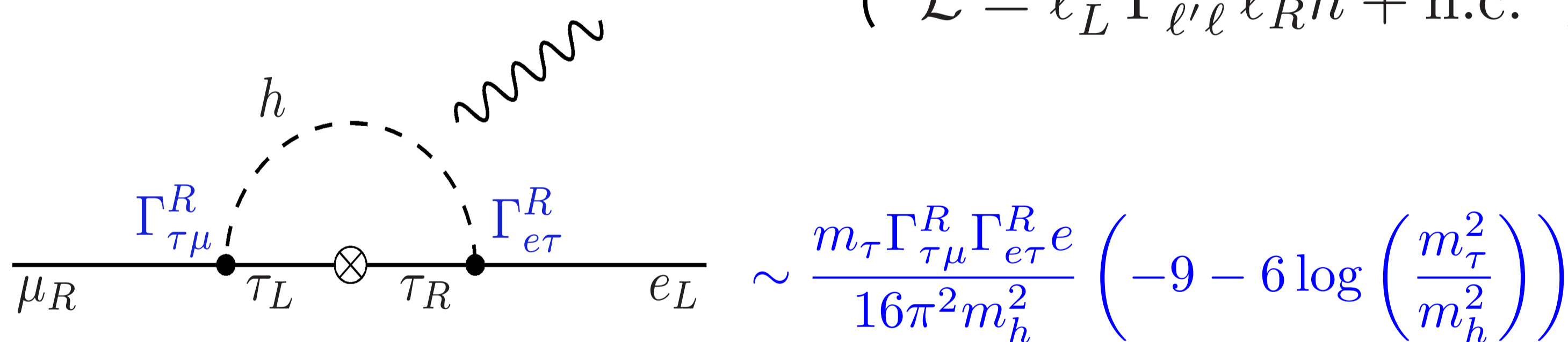
→ Clear evidence of new physics if LFV is discovered.

## Motivations

- In Ref. [1], it is shown that due to the bound of cLFV, only one of  $h \rightarrow e\tau$  or  $h \rightarrow \mu\tau$  can be large:

$$\Gamma_{\tau\mu}^R \Gamma_{e\tau}^R, \Gamma_{\mu\tau}^R \Gamma_{\tau e}^R \lesssim 10^{-8} \quad (\text{from } \text{BR}_{\mu \rightarrow e\gamma}^{\text{exp}} < 4.2 \times 10^{-13})$$

$$(\mathcal{L} = \bar{\ell}'_L \Gamma_{\ell'\ell}^R \ell_R h + \text{h.c.})$$



- They assumed that  $\Gamma_{\ell\ell'}^R = \Gamma_{\ell'\ell}^R$ . However, if it's not the case,  $\mu \rightarrow e\gamma$  is not so severe.

- Can two Higgs LFV decay modes be enhanced at the same time?

- This possibility isn't studied much in the framework of general structure of the Yukawa couplings.
- Excess of  $2.1\sigma$  in  $h \rightarrow \mu\tau$  and  $h \rightarrow e\tau$  in ATLAS [2].

## Generalized DFSZ axion models

- The original motivation of the model is the realization of "nucleophobic" axion [3].

- Bounds from Neutron star cooling and SN 1987A for  $m_a < 20\text{meV}$  [4]

$$|g_{an}|, |g_{ap}| \lesssim 10^{-9}$$

- These couplings can be suppressed by specific conditions for  $g_{au}$  and  $g_{ad}$ .

$$C_p + C_n = 0.50(5)(C_u + C_d - 1) - 2\delta, \quad C_u + C_d = 1,$$

$$C_p - C_n = 1.273(2) \left( C_u - C_d - \frac{1-z}{1+z} \right), \quad \Rightarrow \quad C_u = \frac{1}{1+z} \sim 2/3.$$

$$C_u + C_d = N_1 / \left( \sum_i^3 N_i \right) = 1 \quad N_i: \text{Anomaly coefficient for the } i\text{th generation quarks}$$

$$\Rightarrow N_2 + N_3 = 0 \quad N_1 = (X_{uL} + X_{uR} - X_{dL} - X_{dR})/2$$

### → Nucleophobic axion induces Flavor violation.

- (2+1) flavor space  $Q(f_1) = Q(f_2) = X_{fa}$ ,  $Q(f_3) = X_{f3}$

The quark couplings are given by  $C_u = c_\beta^2, C_d = s_\beta^2$

Nucleophobia is realized  $\tan\beta \sim 0.7$ .

- Nonalignment Yukawa coupling [ $\xi_{ii}^{eL/R} \equiv (V_{eL/R})_{3i}^* (V_{eL/R})_{3i}$ ]

$$(\Gamma_{heie_i}^R) = - \left( \frac{m_{e_i}}{v} [s_{\beta-\alpha} + (-t_\beta + \frac{1}{c_\beta s_\beta}) c_{\beta-\alpha}] \delta_{ij} - \frac{c_{\beta-\alpha} m_\tau}{c_\beta s_\beta v} \sqrt{\xi_{ii}^{eL} \xi_{jj}^{eR}} \right),$$

( $ii = 11, 22, 33$ )

$$\xi_{11}^{eR} = \xi_{22}^{eR} = 0$$



$$\Gamma_{heie_j}^R = \begin{pmatrix} \Gamma_{hee}^R & 0 & \Gamma_{het}^R \\ 0 & \Gamma_{h\mu\mu}^R & \Gamma_{h\mu\tau}^R \\ 0 & 0 & \Gamma_{h\tau\tau}^R \end{pmatrix}$$

- ✓ cLFV bounds  $\mu \rightarrow e\gamma$ :  $\Gamma_{\tau\mu}^R \Gamma_{e\tau}^R, \Gamma_{\mu\tau}^R \Gamma_{\tau e}^R \lesssim 10^{-8}$

$$\tau \rightarrow \ell\gamma: |\Gamma_{\tau\mu,\mu\tau}^R|^2, |\Gamma_{\tau e,e\tau}^R|^2 \lesssim 10^{-4}$$

- ✓ Enhancement of  $h \rightarrow \mu\tau$  and  $h \rightarrow e\tau$

## Numerical calculations

