

T violation at a future neutrino factory

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Ref. arXiv: 2407.05807 [hep-ph]

Introduction

Motivation

- Testing T violation in lepton sector has not been achieved.
- CP and T violation measurements as a non-trivial check of the CPT theorem in QFT.

Framework of standard 3 flavor ν oscillation in matter

Assumption: constant electron density

$$i \frac{d}{dt} \begin{pmatrix} \nu_e(\bar{\nu}_e) \\ \nu_\mu(\bar{\nu}_\mu) \\ \nu_\tau(\bar{\nu}_\tau) \end{pmatrix} = \left[U_{PMNS}^* \text{diag}(0, \Delta E_{21}, \Delta E_{31}) U_{PMNS}^{(T)} + \text{diag}(\pm A, 0, 0) \right] \begin{pmatrix} \nu_e(\bar{\nu}_e) \\ \nu_\mu(\bar{\nu}_\mu) \\ \nu_\tau(\bar{\nu}_\tau) \end{pmatrix}$$

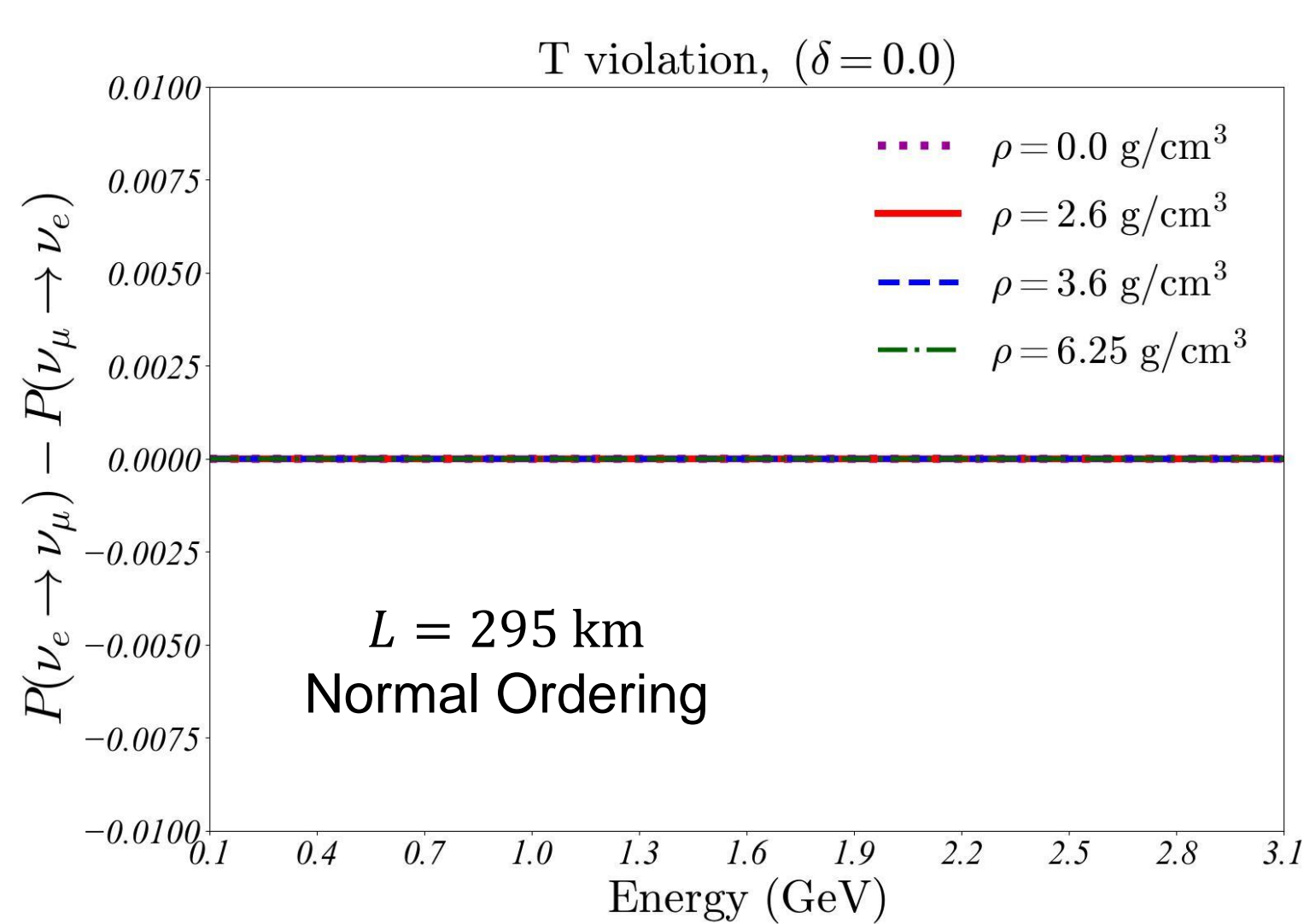
T violation : $P(\nu_\alpha \rightarrow \nu_\beta) - P(\nu_\beta \rightarrow \nu_\alpha) \propto J$

$$J = \frac{c_{13}}{8} \sin 2\theta_{12} \sin 2\theta_{13} \sin 2\theta_{23} \sin \delta : \text{Jarlskog factor}$$

CP violation : $P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$

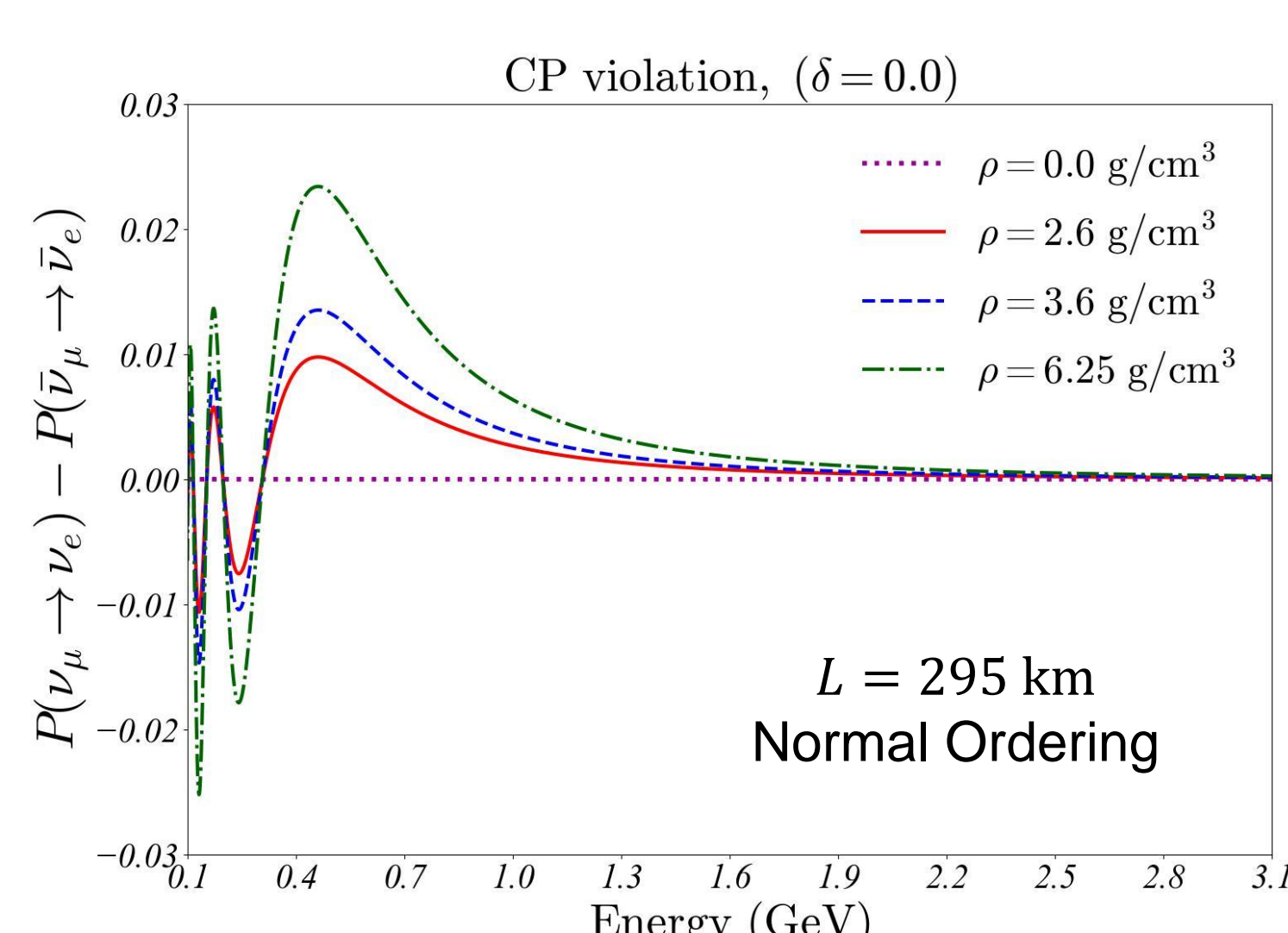
Non-negligible matter effects and δ_{CP} contribution

T violation



- Formally T-odd when we assume that the matter density is symmetric under the exchange of the neutrino source and the detector.
- No matter effect for $\delta = 0$.
- **Matter effects is not important.**

CP violation



- **Fake CP violation** due to the matter effect.
- Matter effects may mimic CP violation.

How to measure the T violation

- Utilize μ^+ beam from muon collider experiment at the J-PARC site (μ TRISTAN).
- Assume to point the beam toward the Hyper-Kamiokande (HK).
- Only ν_e and $\bar{\nu}_\mu$ beam can be used. Therefore, we consider combining it with the T2HK experiment.

$P(\nu_e \rightarrow \nu_\mu)$ from μ TRISTAN

$P(\nu_\mu \rightarrow \nu_e)$ from T2HK

combine

Can observe T violation $P(\nu_e \rightarrow \nu_\mu) - P(\nu_\mu \rightarrow \nu_e) \propto \sin \delta$
Robust test of CP phase independent of matter effects!

Statistical analysis

$$\chi_{TV}^2 \equiv \sum_j \frac{[P_j^{TV}(\delta_0, \rho_0) - P_j^{TV}(\delta^{\text{test}}, \rho^{\text{test}})]^2}{\Delta P_j^2}$$

$$P_j^{TV}(\delta, \rho) \equiv P_j(\nu_e \rightarrow \nu_\mu) - P_j(\nu_\mu \rightarrow \nu_e) \Big|_{T2HK}$$

- Measure $\nu_e \rightarrow \nu_\mu$ at the HK.
- The charges of the muons must be identified (generated by the CC process $\nu_\mu n \rightarrow \mu^- p$ and $\bar{\nu}_\mu p \rightarrow \mu^+ n$).
- At the HK, in principle, ν_μ and $\bar{\nu}_\mu$ are distinguished by neutron tagging method.

Oscillation probability $P(\nu_e \rightarrow \nu_\mu)$

$$P(\nu_e \rightarrow \nu_\mu) = \frac{1}{\kappa} \frac{(\kappa N_{\text{far}}^{\nu_e \rightarrow \nu_\mu} + (1 - \kappa) N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu}) - (1 - \kappa) N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu} \Big|_{T2HK}}{\kappa N_{\text{near}}^{\nu_e \rightarrow \nu_e}}$$

$$\kappa \equiv \frac{1 + C_{\text{id}}}{2}, C_{\text{id}}: \text{charge identification efficiency}$$

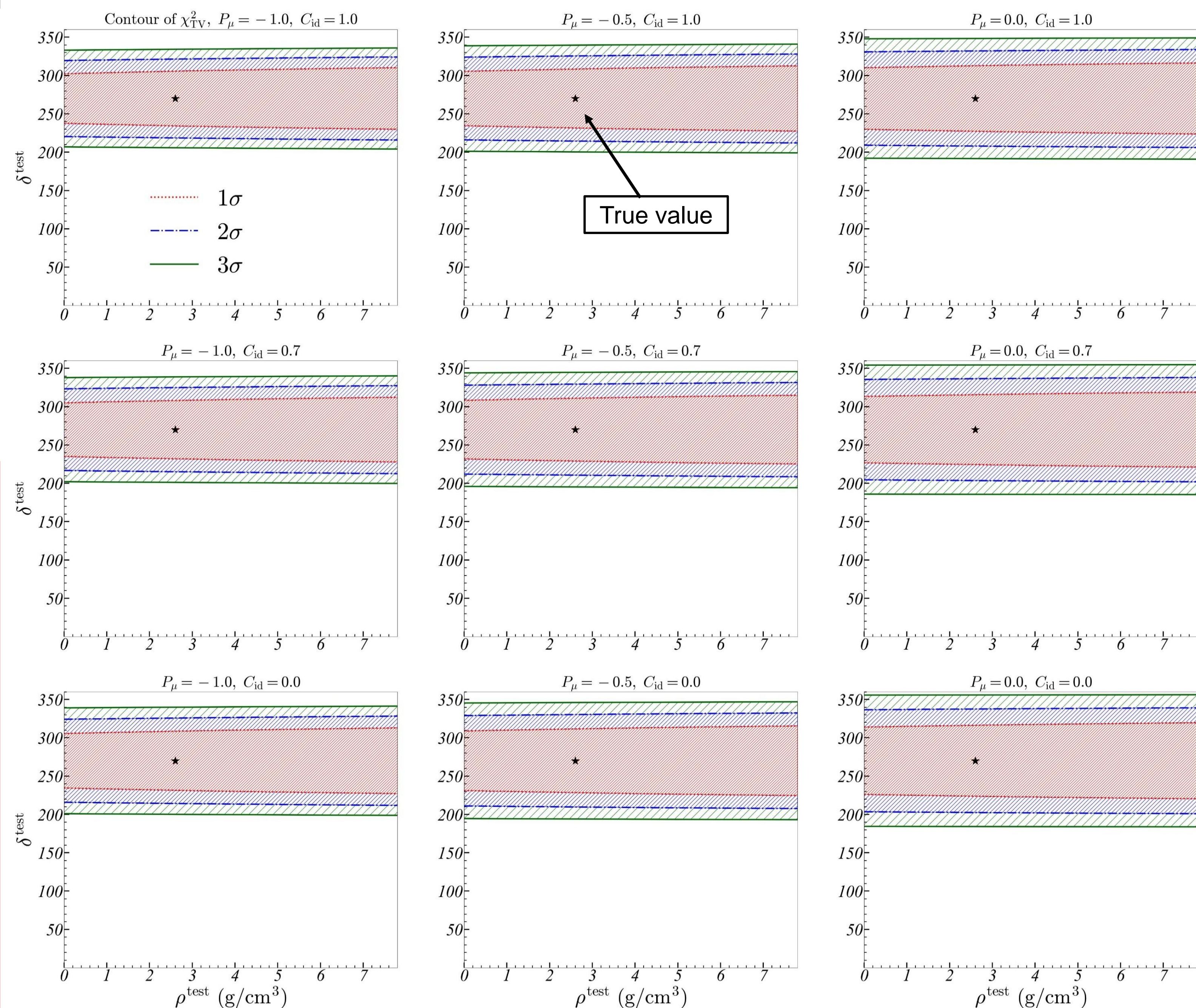
(probability to correctly identify ν_μ for an actual ν_μ event)

- Treatment of $C_{\text{id}} = 0.0$: adding the background events ($\kappa = 1$ and $1 - \kappa = 1$) without performing charge identification analysis.

Parameter determination

χ^2 analysis for T violation

- Set true values as $\delta_0 = -\pi/2$ and $\rho_0 = 2.6 \text{ g/cm}^3$.
- Take the bin size to be 50 MeV, and the muon polarization factor $P_\mu = -1.0, -0.5, 0.0$.
- The efficiency of the charge identification, $C_{\text{id}} = 1.0, 0.7, 0.0$.



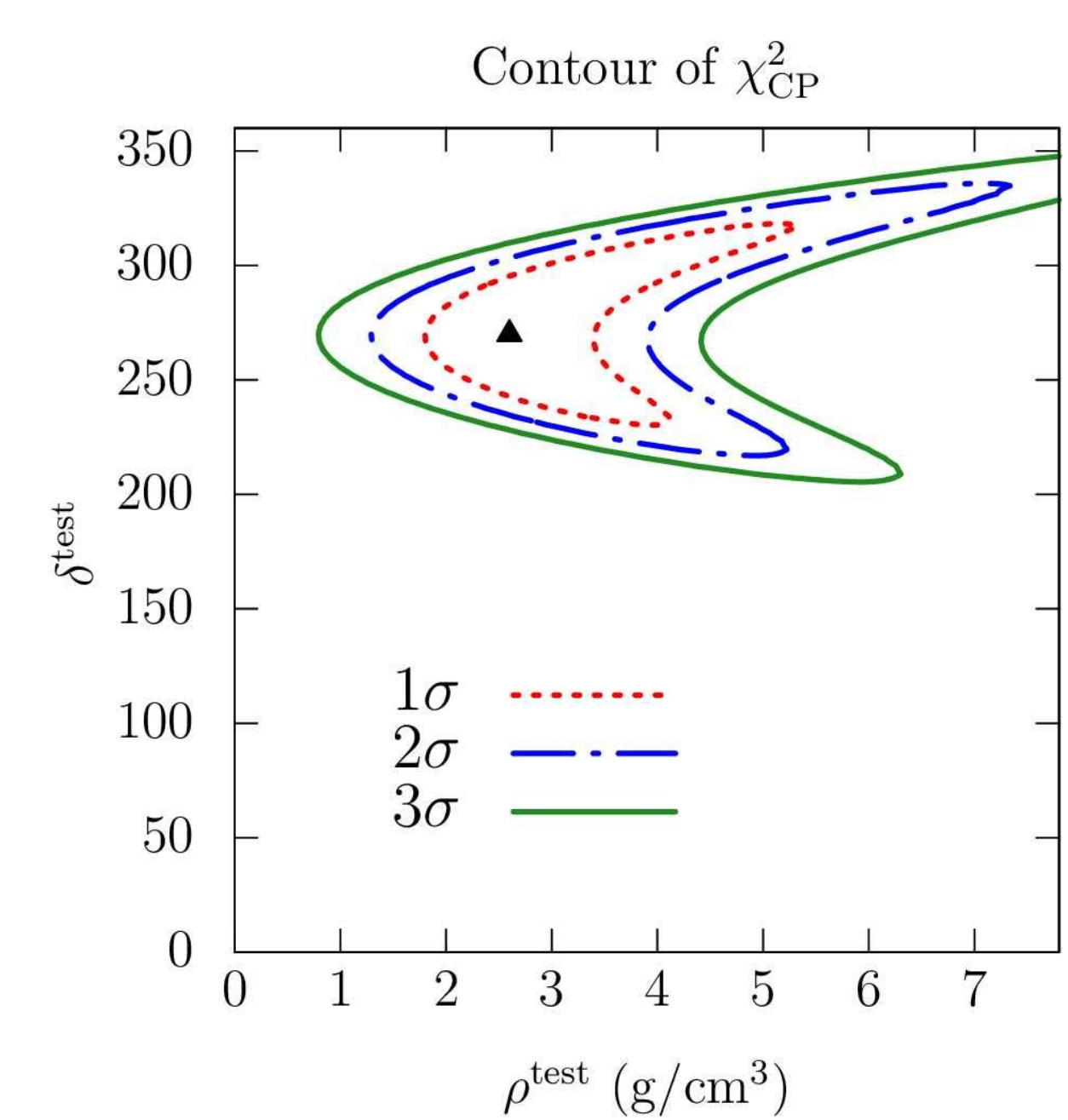
- CP (or T) conserving point, $\delta = 0, \pi$ can be excluded at the level of 3σ even for the unpolarized muon beam.
- For the best case, $C_{\text{id}} = 1.0$ and $P_\mu = -1.0$, the CP angle δ is determined with an accuracy of about $\pm 30^\circ$.
- χ_{TV}^2 only depends on the parameters, δ .
- 1σ confidence level is determined from $\chi^2 = 1$, i.e. for a single degree of freedom.

χ^2 analysis for CP violation

$\chi_{CP}^2(\delta^{\text{test}}, \rho^{\text{test}})$ is similarly analyzed

$$P_j^{CP} = P_j(\nu_\mu \rightarrow \nu_e) \Big|_{T2HK} - P_j(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \Big|_{T2HK}$$

- **Non-trivial dependence** of the allowed region on ρ^{test} appears.
- Information of ρ is critically important.
- The measurement of T violation will be an important additional information for the measurement of the CP angle δ .
- Additional measurement of T violation helps to discriminate the matter effect.



Possible CPT test?

$$P_j^{\text{CPT}} = P_j(\nu_e \rightarrow \nu_\mu) - P_j(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \Big|_{T2HK} = P_j^{TV} + P_j^{CP}$$

Under our assumptions, this quantity would not measure anything as it should be trivially vanishing up to matter effects. The analysis of this kind will be a quite important fundamental test of symmetry in physical laws of the Universe.

Summary

- Testing T violation in the lepton sector remains as an important task in particle physics.
- We studied the possibility of measuring T violation, $P(\nu_e \rightarrow \nu_\mu) - P(\nu_\mu \rightarrow \nu_e)$ considering the baseline from J-PARC to Hyper-Kamiokande.
- By combining $P(\nu_e \rightarrow \nu_\mu)$ from μ TRISTAN with $P(\nu_\mu \rightarrow \nu_e)$ from T2HK, the T violation can be defined.
- For the case of maximum CP violation, $\delta = -\pi/2$, CP invariant theories ($\delta = 0$ or π) can be excluded at more than 3σ level.
- T violation can be a good measure which is not sensitive to the detailed density profile of the Earth.
- Finally, A more complete analysis will be necessary to establish the feasibility.