



JUMP TRAJECTORIES FOR GENERAL QUANTUM MASTER EQUATIONS

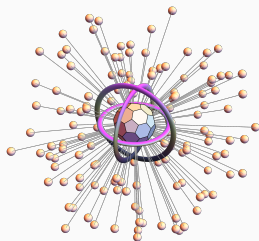
Paul Menczel, Theoretical Physics Laboratory, RIKEN

12 July 2024

Frontiers in Non-Equilibrium Physics 2024

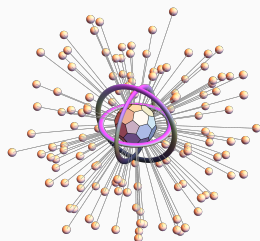
Setting

- OQS described by master equation
- Discuss quantum jump trajectories



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What master equation?

- Lindblad:

$$\partial_t \rho = -i[H, \rho] + \sum_{\alpha} \gamma_{\alpha} (L_{\alpha} \rho L_{\alpha}^{\dagger} - \{L_{\alpha}^{\dagger} L_{\alpha}, \rho\} / 2)$$

→ Markovian weak-coupling limit (cf. next talk)

→ Jump trajectory framework well known

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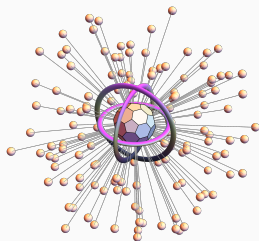
What master equation?

- Lindblad
- Non-Markovian Lindblad (negative rates):

$$\partial_t \rho = -i[H, \rho] + \sum_{\alpha} \gamma_{\alpha} (L_{\alpha} \rho L_{\alpha}^{\dagger} - \{L_{\alpha}^{\dagger} L_{\alpha}, \rho\} / 2)$$

→ TCL master equation, Redfield equation

→ May be completely positive (not CP-divisible)



Setting

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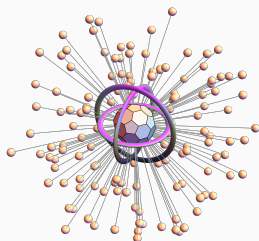
What master equation?

- Lindblad
- Non-Markovian Lindblad (negative rates)
- Non-Physical Lindblad (complex rates, non-Hermitian H):

$$\partial_t \rho = -i[H, \rho] + \sum_{\alpha} \gamma_{\alpha} (L_{\alpha} \rho L_{\alpha}^{\dagger} - \{L_{\alpha}^{\dagger} L_{\alpha}, \rho\} / 2)$$

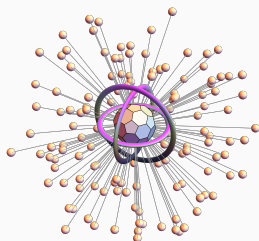
→ Embedding for physical non-Markovian dynamics
(pseudomodes)

→ Non-Hermitian state ρ



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What master equation?

- Lindblad
- Non-Markovian Lindblad (negative rates)
- Non-Physical Lindblad (complex rates, non-Hermitian H)
- Generic time-local master equation:

$$\partial_t \rho = A\rho + \rho B^\dagger + \sum_{\alpha} C_{\alpha} \rho D_{\alpha}^\dagger$$

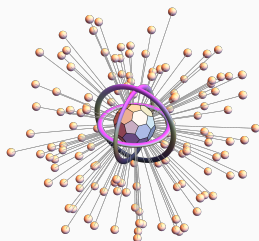
- Most general mathematical form
- Feedback, counting fields

Setting

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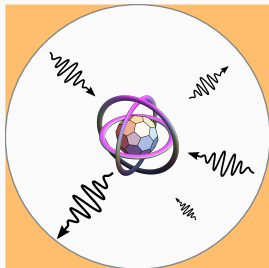
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Lindblad equation:

$$\partial_t \rho = -i[H, \rho] + \sum_{\alpha} \gamma_{\alpha} (L_{\alpha} \rho L_{\alpha}^{\dagger} - \{L_{\alpha}^{\dagger} L_{\alpha}, \rho\}/2)$$

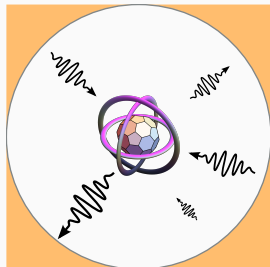


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Jump Trajectories:

- Stochastic evolution of pure state $|\psi\rangle$ such that $\rho = \mathbb{E}\{|\psi\rangle\langle\psi|\}$
- $\partial_t |\psi\rangle = \sum_{\alpha} \left[\frac{1}{\langle L_{\alpha}^{\dagger} L_{\alpha} \rangle^{1/2}} L_{\alpha} |\psi\rangle - |\psi\rangle \right] dN_{\alpha} + \left[-iH - \sum_{\alpha} \gamma_{\alpha} (L_{\alpha}^{\dagger} L_{\alpha} - \langle L_{\alpha}^{\dagger} L_{\alpha} \rangle) \right] |\psi\rangle dt$
- dN_{α} are Poisson increments, $\mathbb{E}\{dN_{\alpha}\} = \gamma_{\alpha} \langle L_{\alpha}^{\dagger} L_{\alpha} \rangle dt$
→ only works for positive rates

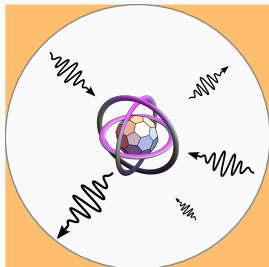


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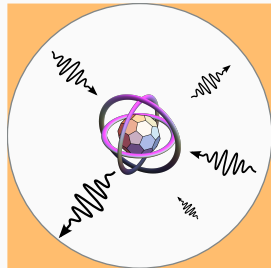


Focus on quantum jump unravellings (Poisson noise)

Quantum diffusion (Gaussian noise) → Landi *et al.*, PRXQ (2024)

Advantages:

- Monte-Carlo simulations:
 $\mathcal{O}(N^2) \rightarrow \mathcal{O}(N)$ in parallel
- Theoretical access to fluctuations



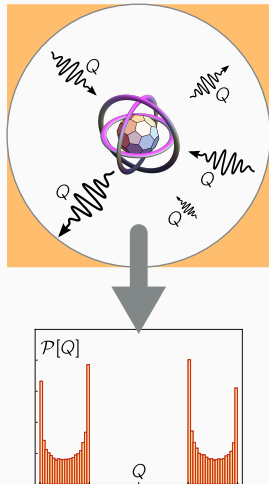
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Small advertisement:

- In cyclic operation,
$$\Delta S \geq 2A \lambda_Q \operatorname{artanh} \lambda_Q$$
 λ_Q characterizes energy frequency dist.
- Leads to power-efficiency trade-off etc.

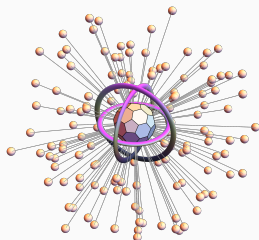
PM, C. Flindt, and K. Brandner, PRR (2020)



Negative rates (at some or all times):

$$\partial_t \rho = -i[H, \rho] + \sum_{\alpha} \gamma_{\alpha} (L_{\alpha} \rho L_{\alpha}^{\dagger} - \{L_{\alpha}^{\dagger} L_{\alpha}, \rho\} / 2)$$

- Time-convolutionless (TCL) master equation
- Redfield equation
- Classical noise models



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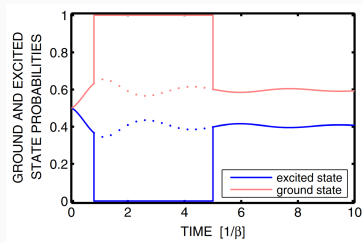
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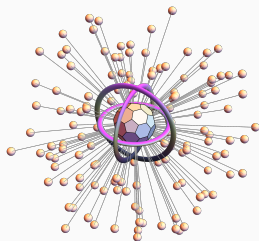
→ Redfield equation

→ Classical noise models

Approach 1: “backward jumps”



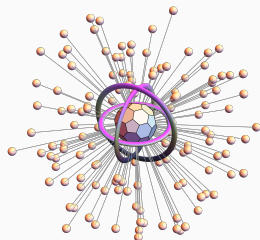
J. Piilo *et al.*, PRL (2008)



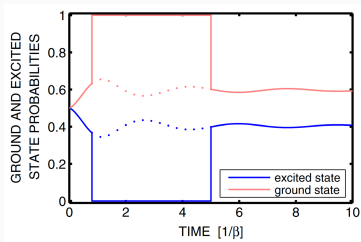
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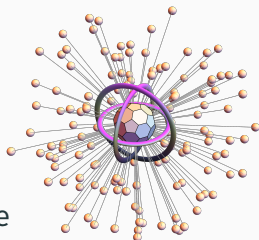


J. Piilo *et al.*, PRL (2008)

- Only works for completely positive dynamics
- Evolution of different trajectories is coupled

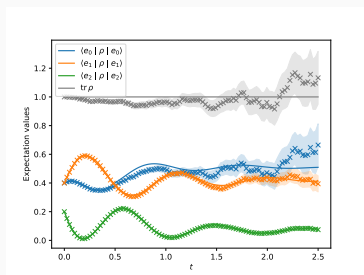
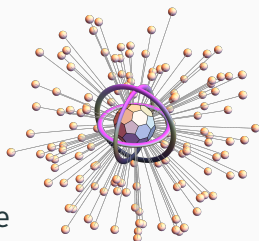
Approach 2: “influence martingale”

- Introduce scalar μ with $\rho = \mathbb{E}\{\mu|\psi\rangle\langle\psi|\}$
- $|\psi\rangle$ and μ satisfy coupled stochastic differential equations
- Method is part of QuTiP v5 as `nm_mcsolve`



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- Example: Redfield equation for two qubits coupled to same reservoir (10k trajectories)
- Instability at longer times (μ grows exponentially)

Generic time-local master equation:

$$\partial_t \rho = A\rho + \rho B^\dagger + \sum_{\alpha} C_{\alpha} \rho D_{\alpha}^\dagger$$

Hermiticity not preserved, $\rho = \mathbb{E}\{|\psi\rangle\langle\psi|\}$ clearly impossible

- Trajectories in double Hilbert space: $\rho = \mathbb{E}(|\psi_1\rangle\langle\psi_2|)$
→ coupled SDEs for $|\psi_1\rangle$ and $|\psi_2\rangle$
→ same instability at longer times

Breuer *et al.*, PRA (1999)

- Embedded in Lindblad equation for $\bar{\rho}$ on double space:

$$\rho = e^{\int_0^t d\tau \alpha(\tau)} \text{tr}_{\mathfrak{a}} [X_{\mathfrak{a}} \bar{\rho}]$$

→ same instability: $\alpha(\tau) \geq 0$ (zero iff Lindblad)

Hush *et al.*, PRA (2015)

Complex rates and non-Hermitian Hamiltonian:

$$\partial_t \rho = -i[H, \rho] + \sum_{\alpha} \gamma_{\alpha} (L_{\alpha} \rho L_{\alpha}^{\dagger} - \{L_{\alpha}^{\dagger} L_{\alpha}, \rho\}/2)$$

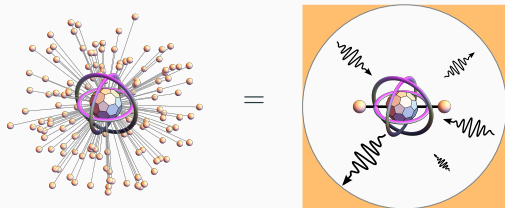
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Why?

Pseudomode framework provides embedding of non-Markovian dynamics (Gaussian bath) in such equation:

$\rho_{\text{sys}} = \text{tr}_{\text{pm}} \rho$ exactly equals correct system state



PM, K. Funo, M. Cirio, N. Lambert, and F. Nori, arXiv:2401.11830

Trajectories in double Hilbert space: $\rho = \mathbb{E}(|\psi_1\rangle\langle\psi_2|)$

$$\begin{aligned}
 d|\psi_1\rangle &= \left[-iH - \frac{1}{2} \sum_{\alpha} (\gamma_{\alpha} L_{\alpha}^{\dagger} L_{\alpha} - r_{\alpha}) \right] |\psi_1\rangle dt \\
 &\quad + \sum_{\alpha} \left[\sqrt{\frac{\gamma_{\alpha}}{r_{\alpha}}} L_{\alpha} |\psi_1\rangle - |\psi_1\rangle \right] dN_{\alpha}, \\
 d|\psi_2\rangle &= \left[-iH^{\dagger} - \frac{1}{2} \sum_{\alpha} (\gamma_{\alpha}^{*} L_{\alpha}^{\dagger} L_{\alpha} - r_{\alpha}) \right] |\psi_2\rangle dt \\
 &\quad + \sum_{\alpha} \left[\sqrt{\frac{\gamma_{\alpha}^{*}}{r_{\alpha}}} L_{\alpha} |\psi_2\rangle - |\psi_2\rangle \right] dN_{\alpha}
 \end{aligned}$$

→ jump rates $r_{\alpha} > 0$ can be chosen freely

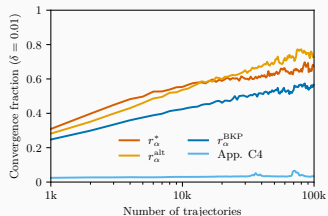
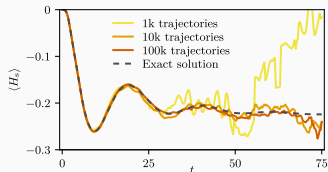
→ choose them to minimize fluctuations of $|\langle\psi_2 | \psi_1\rangle|$

→ exponential growth of fluctuations unavoidable

PM, K. Funo, M. Cirio, N. Lambert, and F. Nori, arXiv:2401.11830

FINAL EXAMPLE

- Spin-Boson Model:
 $H_s = \frac{\Delta}{2}\sigma_x$, coupling $Q = \sigma_z$
- Underdamped bath, finite temperature, neglect Matsubara
- Exponential growth of required trajectories
- Other stochastic methods have similar problems
 - Stockburger, EPL (2016)
 - Luo *et al.*, PRXQ (2023)



PM, K. Funo, M. Cirio, N. Lambert, and F. Nori, arXiv:2401.11830

- Quantum jump trajectories are a useful tool.
- For general master equations, one can do jump unravelings in the double Hilbert space.
- They appear to be unavoidably unstable at long times.
- We determined the optimal rates to minimize instability.
- Outlook:
 - Can we use unravelings like this to study fluctuations of physical quantities? (If not, why not?)
 - Is there a similar issue in classical stochastic dynamics?
 - Can the instability be used to characterize completely positive maps?
 - Does it tell us anything about Markovian embeddings?



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