

*Non-equilibrium physics*

*Many-body (equilibrium) physics*

# Non-reciprocal frustration

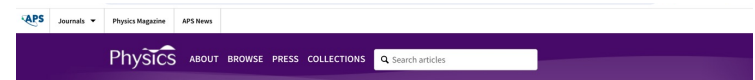
## physics

Ryo Hanai

*Yukawa Institute for Theoretical Physics (YITP), Kyoto University*



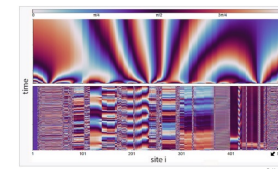
### Physics "Viewpoint"



#### Nonreciprocal Frustration Meets Geometrical Frustration

Peter Littlewood  
Department of Physics, University of Chicago, Chicago, IL, US  
February 26, 2024 • *Physics* 17, 32

New theoretical work establishes an analogy between systems that are dynamically frustrated, such as glasses, and thermodynamic systems whose members have conflicting goals, such as predator-prey ecosystems.



Nonreciprocal Frustration: Time Crystalline Order-by-Disorder Phenomenon and a Spin-Glass-like State  
Ryo Hanai  
*Phys. Rev. X* 14, 011029 (2024)  
Published February 26, 2024  
[Read PDF](#)

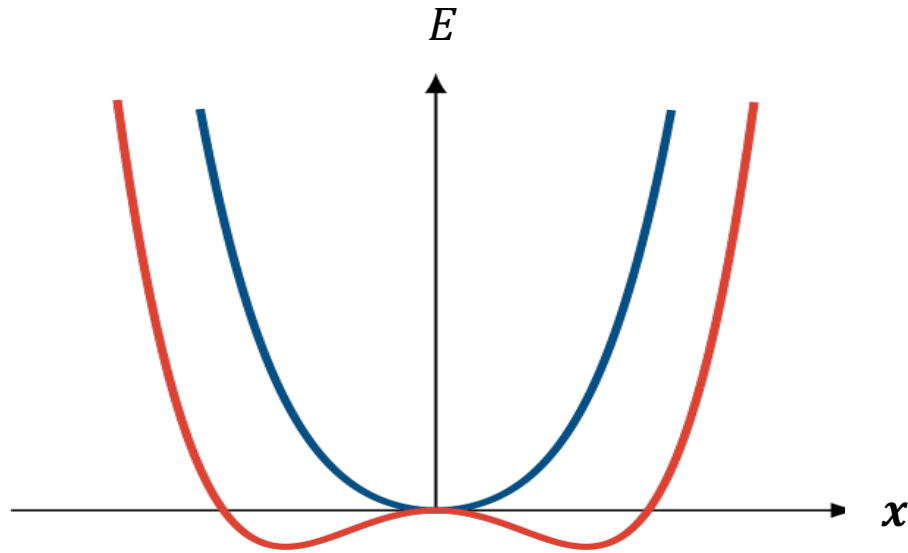
**Recent Articles**  
Fine Control of Ultracold Polar Molecules  
The ability to store molecules in reconfigurable optical traps could allow researchers to harness the rich physics of molecules in quantum applications.

A Chiral Crystal's Orbital Texture  
X-ray experiments reveal that a semimetal exhibits "orbital texture"—an exotic electronic structure resulting in spin-dependent electron transport.

Making Neutron-Deficient Nuclei

R. Hanai, *Phys. Rev. X* **14**, 011029 (2024).

# Equilibrium paradigm: (Free) energy minimization principle



Energy

$$E = - \sum_{i,j} J_{ij} \cos(x_j - x_i)$$

Ground state:  $x = \{x_i\}$  that minimizes the energy  $E$

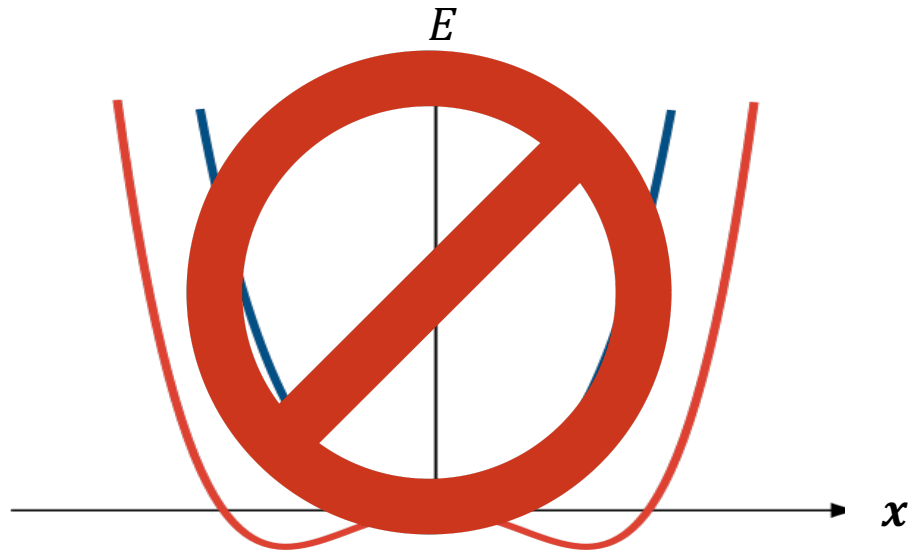
Reciprocal coupling

$$J_{ij} = J_{ji}$$

# ~~Equilibrium~~ paradigm:

(Free energy minimization principle)

## Nonequilibrium



~~Energy~~

$$E = - \sum_{i,j} J_{ij} \cos(x_j - x_i)$$

~~Ground state:  $x = \{x_i\}$  that minimizes the energy  $E$~~

Reciprocal coupling  
 $J_{ij} = J_{ji}$

~~Non-reciprocal coupling~~  
 $J_{ij} \neq J_{ji}$   
can arise!

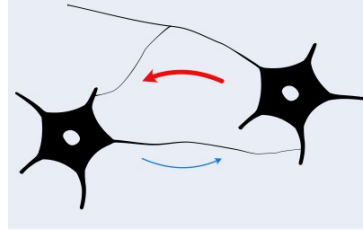
# Non-reciprocally interacting systems

## Ecological systems



<http://animal.memozee.com/view.php?tid=3&did=3513>

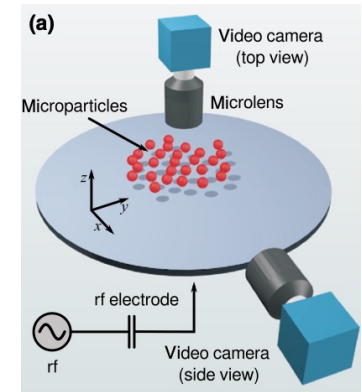
## Neural networks



e.g., J. W. Krakauer, et al., *Neuron* 93, 480 (2017).

## Driven systems

(Complex plasma)



Ivlev, et al., *PRX* 5, 011035 (2015).

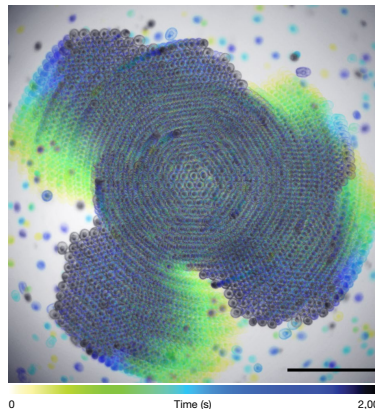
## Social network



e.g., A. Pluchino, et al., *Inter. J. Mod. Phys. C* 16, 515 (2005).

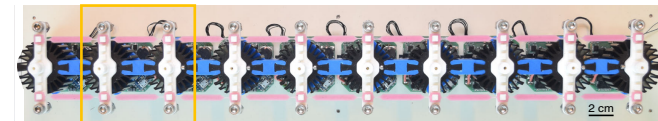
## Living matter

(Starfish embryos)



T. H. Tan, et al., *Nature* 2022

## Robotic metamaterial



Brandenbourger, et al., *Nat. Comm.* 2019

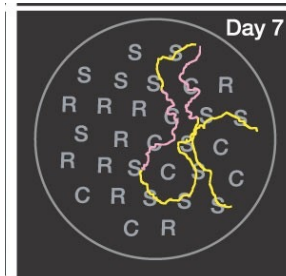
Ghatak, et al., *PNAS* 2020

Baconnier, et al., *Nat. Phys.* 2022



# Collective phenomena in non-reciprocal many-body systems

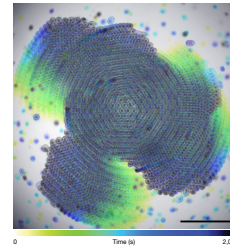
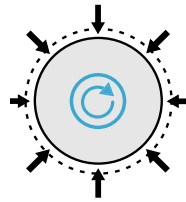
## Biodiversity in ecosystems



B. Kerr, et al., Nature 2002

S. Allesina and S. Tang, Nature 2012

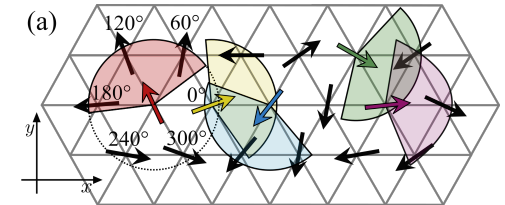
## Odd elasticity



C. Scheibner, et al., , Nat. Phys. 2020

T. H. Tan, et al., , Nature 2022

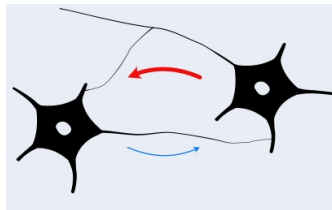
## Long-ranged order in 2D



Loos, Klapp, and Martyneć, arXiv:2206.10519

Dadhichi, et al., Phys. Rev. E 2020

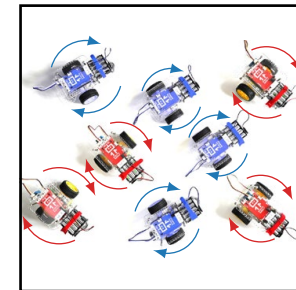
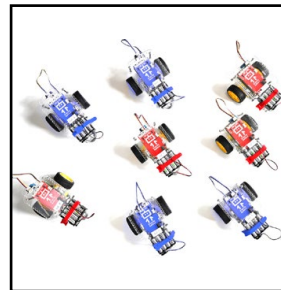
## Controlling neuron dynamics



H. Wilson and J. Cowan,  
Biophys. J. 1972

G. Parisi, J. Phys. A 1986

## Non-reciprocal phase transition

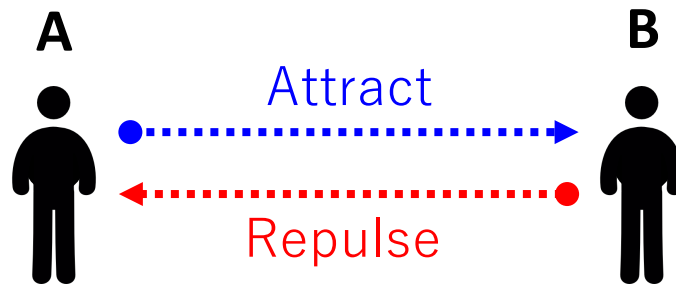


M. Fruchart\*, RH\*, P. B. Littlewood, and V. Vitelli, Nature 2021

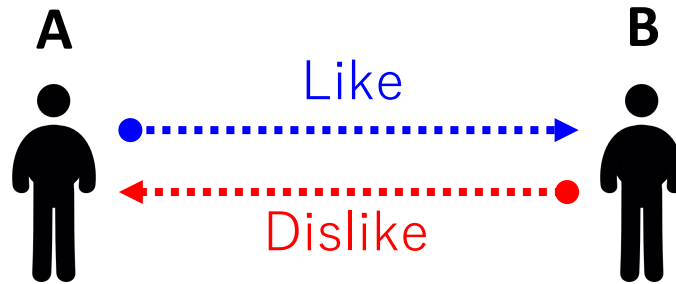
You, Baskaran, Marchetti, PNAS 2020

Saha, Agudo-Canalejo, Golestanian, PRX 2020

# Non-reciprocal interaction $\approx$ Frustration physics



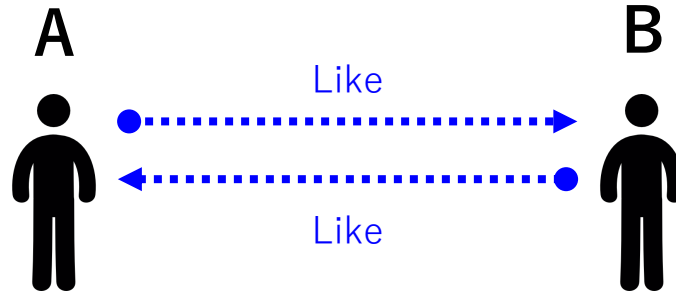
# Non-reciprocal friendship $\approx$ Frustration physics



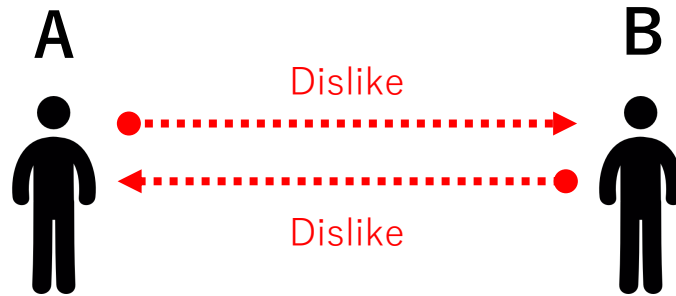


# Non-reciprocal friendship: source of frustration

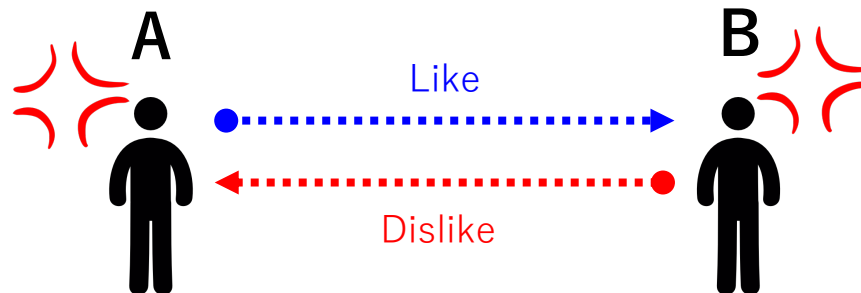
(a) Reciprocal relationship (like)



(b) Reciprocal relationship (dislike)



(c) Non-reciprocal relationship



No configuration that makes both happy

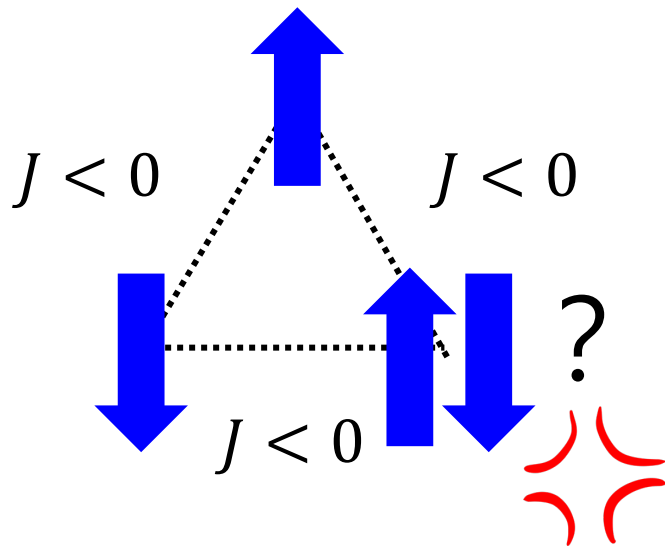


**FRUSTRATION!**

# Geometrical frustration

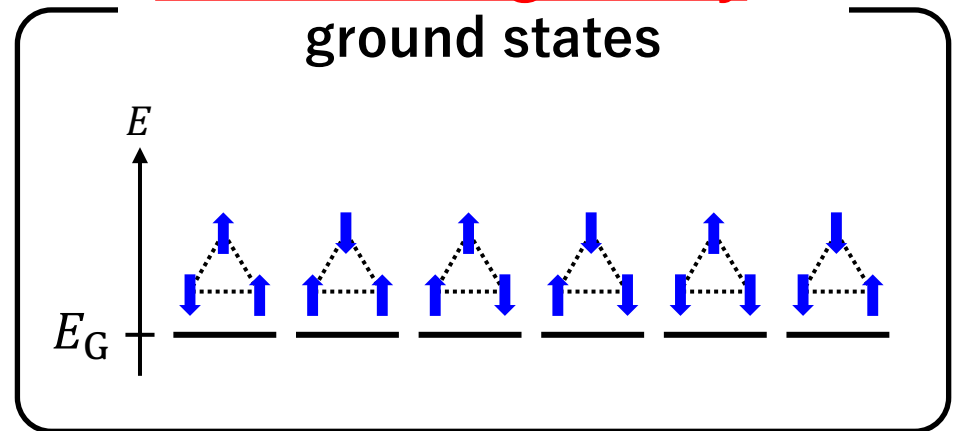
## Geometrically frustrated system

A system that cannot satisfy all the constituents' "desire" to minimize all interaction energies.



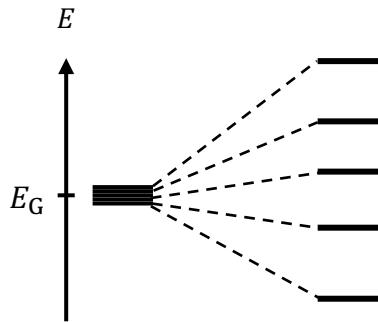
No configuration can make all spins happy

## Accidental degeneracy of ground states



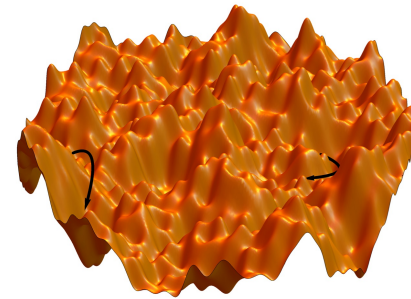
# Accidental degeneracy induces exotic phenomena in geometrical frustrated systems

## ➤ Order-by-disorder



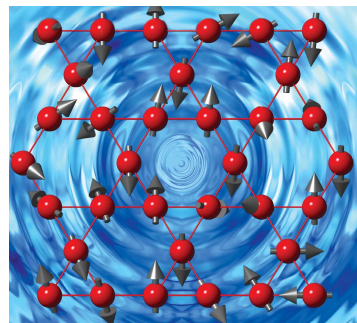
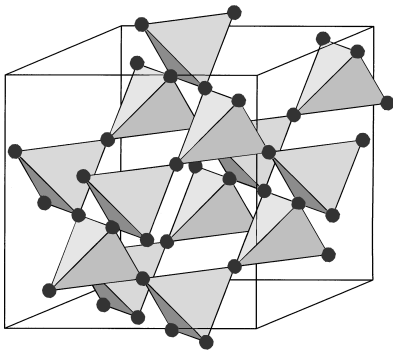
Villain, et al., J. Physique (1980)

## ➤ Spin glass



(Image from <https://scglass.uchicago.edu/>)

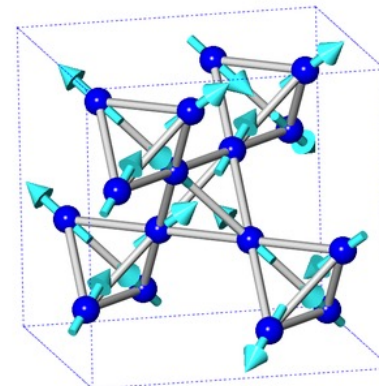
## ➤ Quantum/Classical spin liquid



R. Mossner and J. T. Chalker, PRL 1998

T. Imai and Y. Lee, Physics Today 2016

## ➤ Spin ice

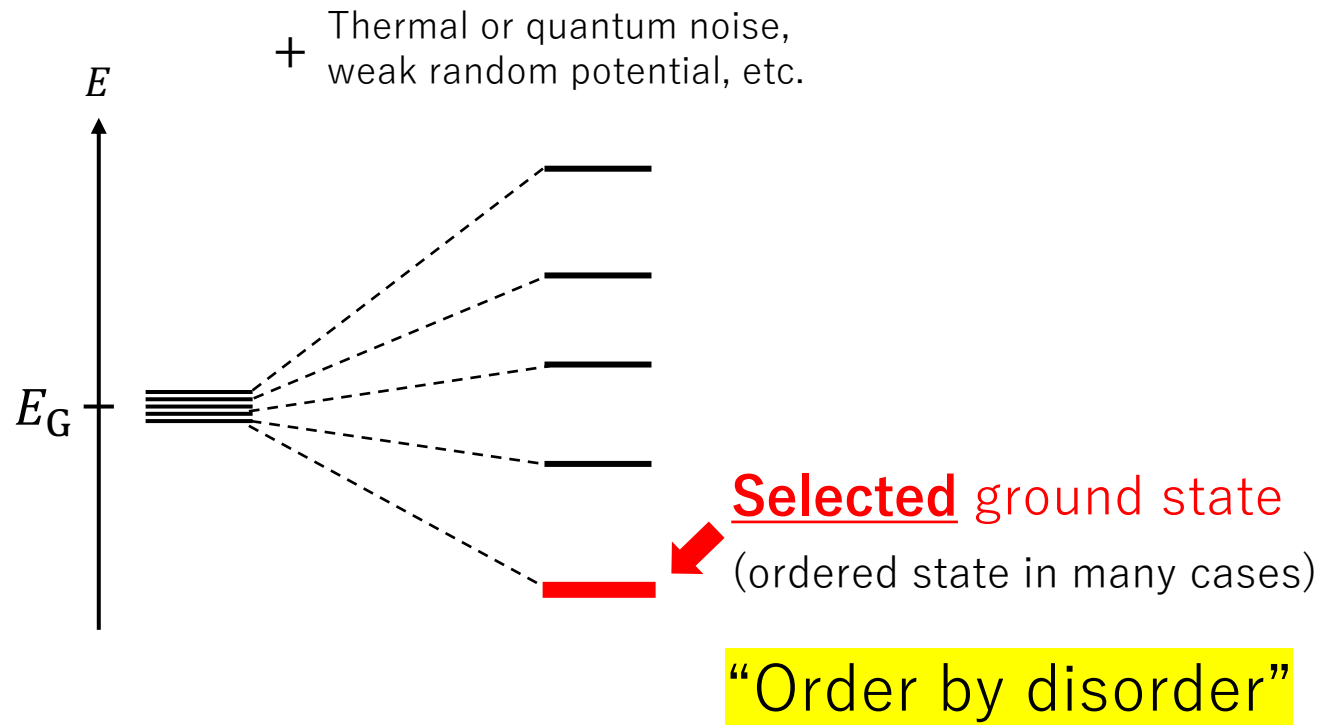


(Image from Wikipedia)

# Order by disorder phenomena

Villain, et al., J. Physique (1980)

**Accidental degeneracy:** Not protected by symmetry nor topology



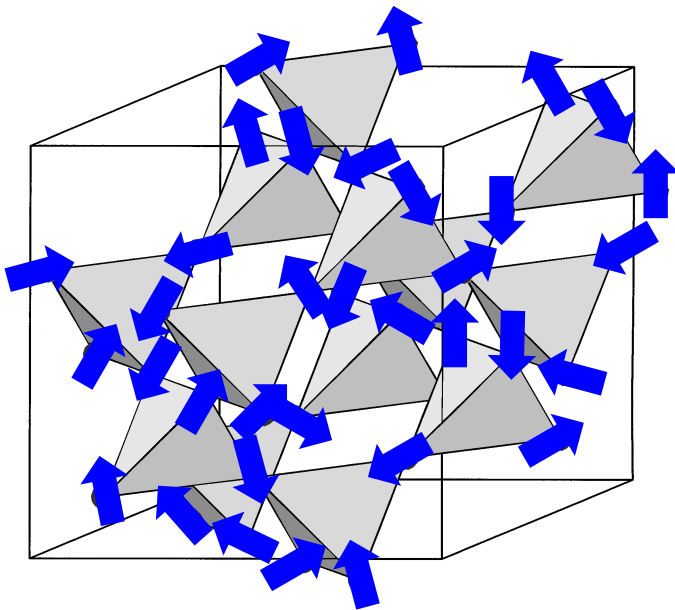
# Order by disorder phenomena

(Example)

Moessner and Chalker, PRL1998, PRB1998

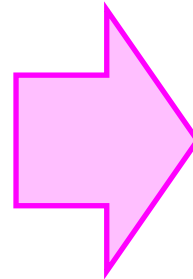
Antiferromagnetic XY spins on a pyrochlore lattice

Ground state

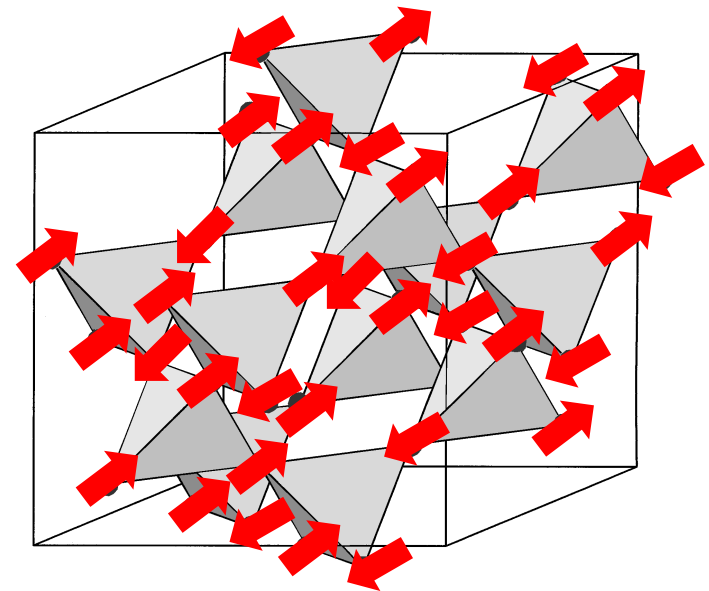


Disordered state

Increase  
temperature



Finite temperature

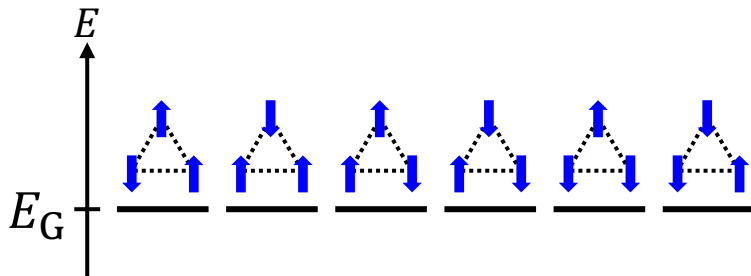


Long-ranged order

# Geometrical vs Non-reciprocal frustration

## Geometrical frustration

Accidental degeneracy of  
ground state



## Non-reciprocal frustration



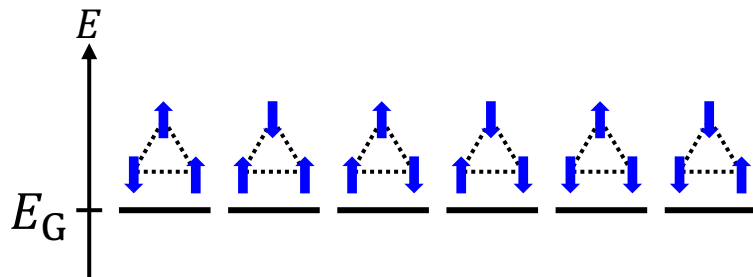
Energy cannot be defined ...

May not even converge to  
a static state...

# Geometrical vs Non-reciprocal frustration

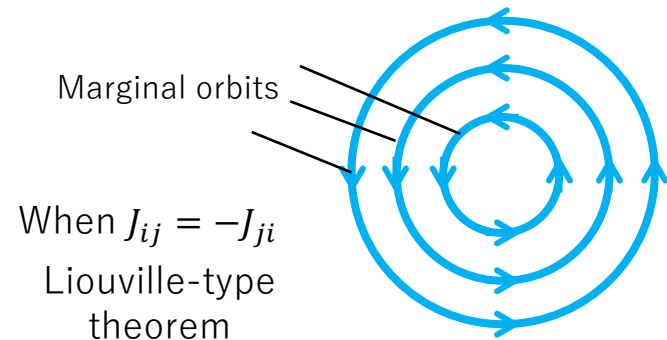
## Geometrical frustration

Accidental degeneracy of ground state



## Non-reciprocal frustration

"Accidental degeneracy" of orbits



➤ *Dynamical counterpart* of order-by-disorder and spin glass(-like) state occurs!

This talk: order by disorder

# Dissipative XY spin dynamics

$$\dot{\theta}_i = \sum_j J_{ij} \sin(\theta_j - \theta_i)$$

Here, couplings are **non-reciprocal** in general  $J_{ij} \neq J_{ji}$



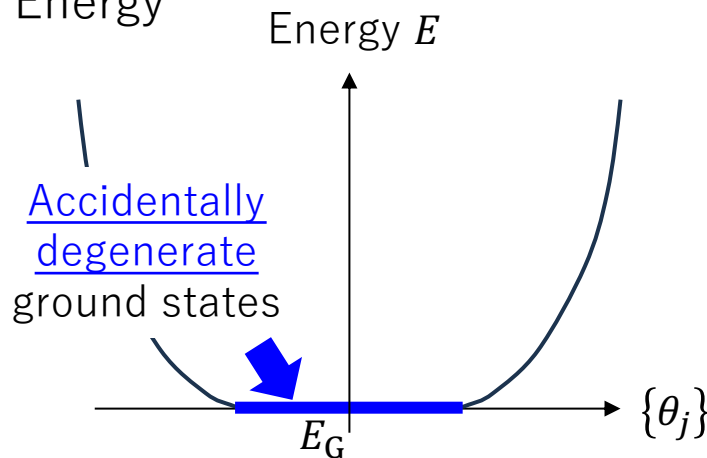
# Dissipative XY spin dynamics

- **Reciprocal** case ( $J_{ij} = J_{ji}$ ) : Energy minimization problem

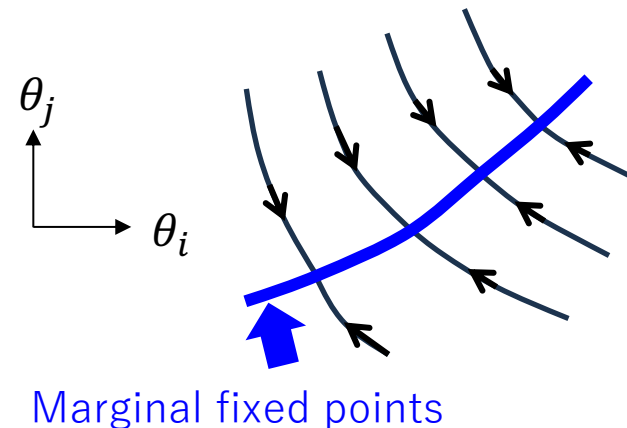
$$\dot{\theta}_i = -\frac{\partial E(\theta)}{\partial \theta_i} \quad \text{with} \quad E(\theta) = -\sum_{i,j} J_{ij} \cos(\theta_j - \theta_i)$$

Potential with **geometrically** frustration ➡ Accidentally degenerate ground states

- Energy



- Orbits



# "Accidental degeneracy" of *orbits*

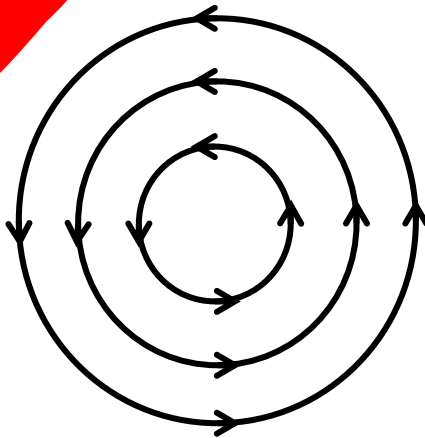
➤ **Anti-symmetric** case ( $J_{ij} = -J_{ji}$ )

Liouville-type theorem

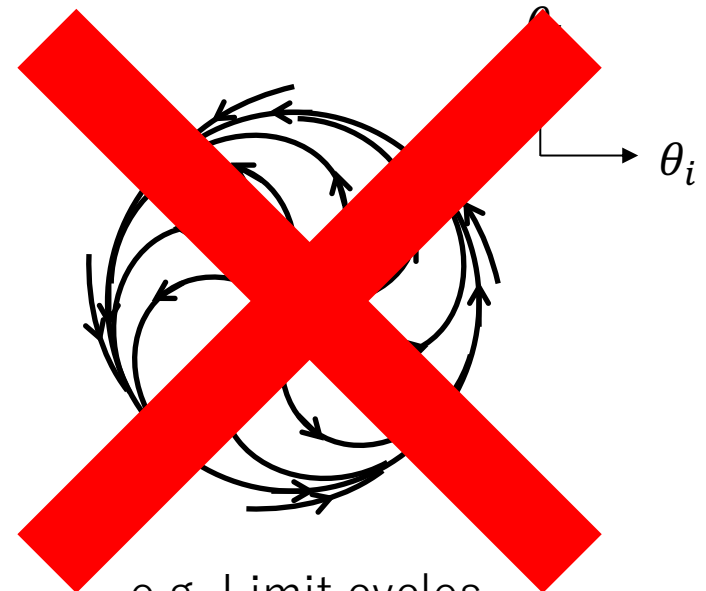
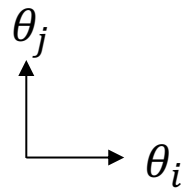
$$\frac{d\rho}{dt} = \frac{\partial \rho}{\partial t} + \sum_i \frac{\partial \rho}{\partial \theta_i} \dot{\theta}_i = 0$$

RH, PRX2024

Conservation of phase volume = **Non-dissipative** dynamics



Marginal



e.g. Limit cycles

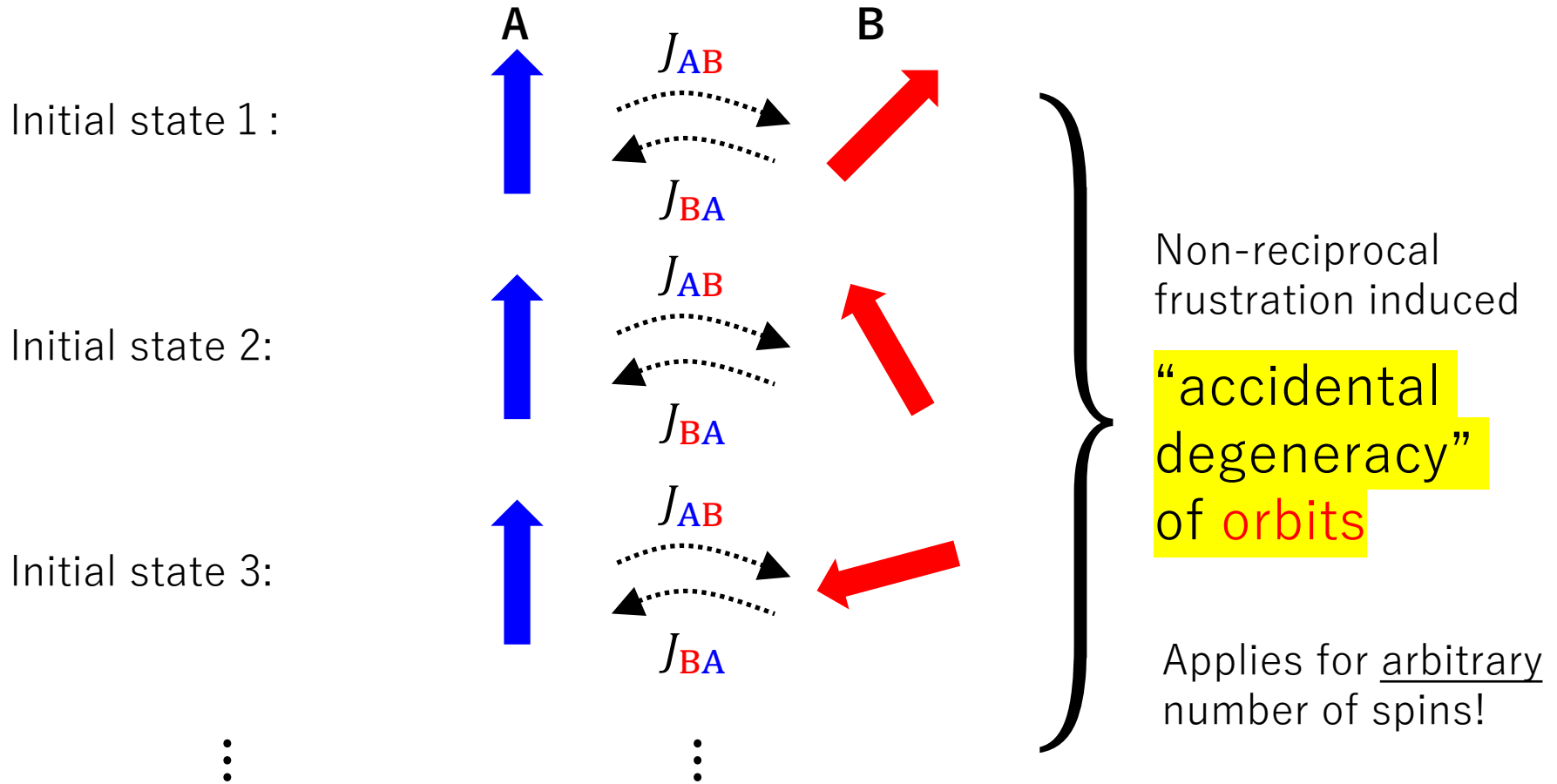
# "Accidental degeneracy" of *orbits*

(e.g., two XY spin system)

$$J_{AB} = -J_{BA}$$

$$\dot{\theta}_A = J_{AB} \sin(\theta_B - \theta_A)$$

$$\dot{\theta}_B = J_{BA} \sin(\theta_A - \theta_B)$$



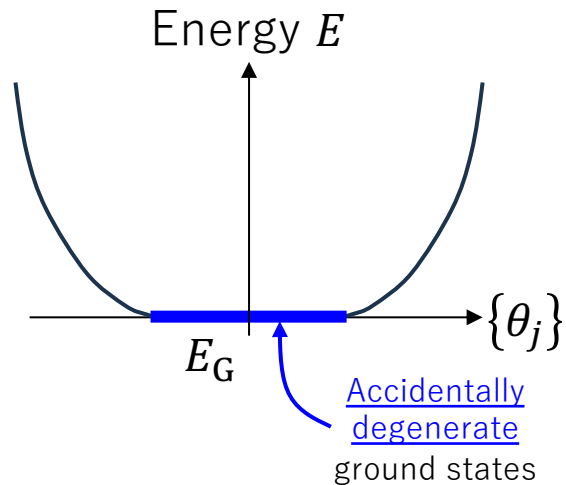
Common feature between **geometrically frustrated** and **non-reciprocal** system:

**Marginal orbits!**

# Order-by-disorder phenomena in geometrically frustrated systems

## ➤ Zero temperature

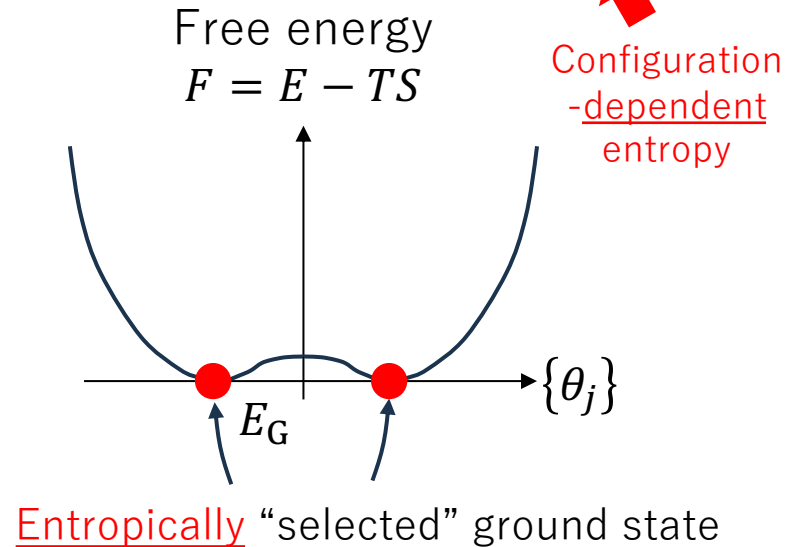
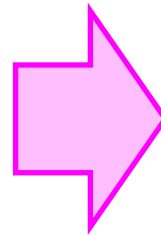
Minimize energy  $E$



## ➤ Finite temperature

Minimize free energy  $F = E - TS(\theta)$

Increase temperature  
 $T > 0$



## ➤ Dynamics ( $T > 0$ ):

$$\dot{\theta} = -\frac{\partial F}{\partial \theta} = -\cancel{\frac{\partial E}{\partial \theta}} + T \frac{\partial S}{\partial \theta}$$

Marginal

Entropic force

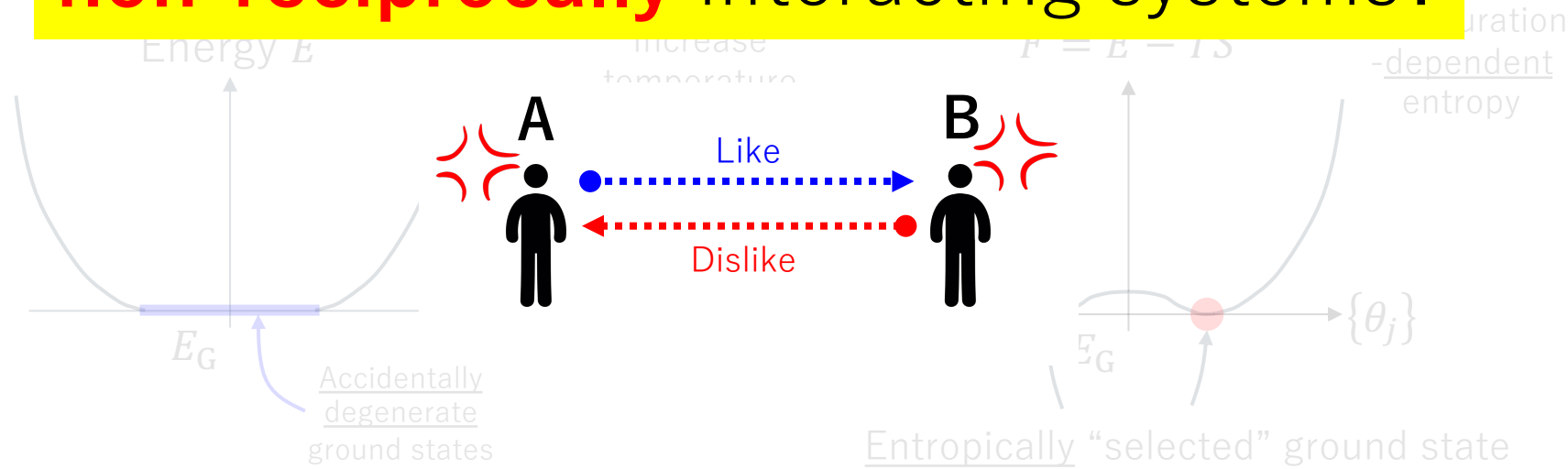
**Entropic force**  
selects the orbit!  
= order by disorder

# Order-by-disorder phenomena in geometrically frustrated systems

➤ Zero temperature

➤ Finite temperature

Q: Does the same physics occur in **non-reciprocally** interacting systems?



➤ Dynamics ( $T > 0$ ):

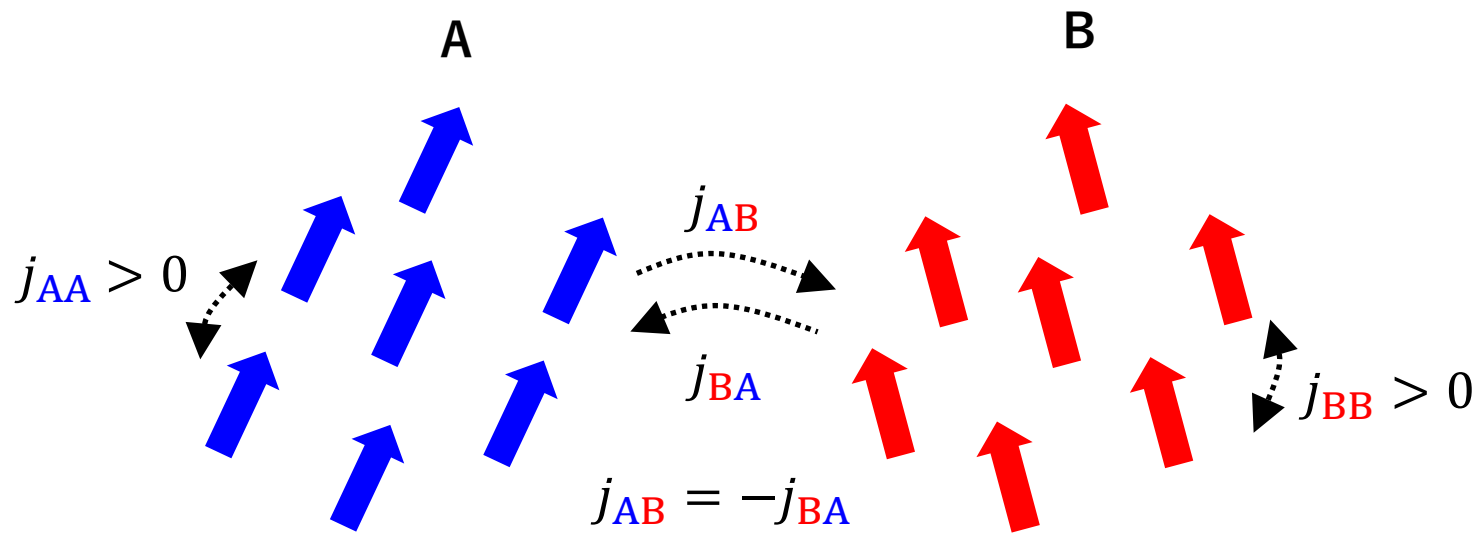
$$\dot{\theta} = -\frac{\partial F}{\partial \theta} = -\cancel{\frac{\partial E}{\partial \theta}} + T \frac{\partial S}{\partial \theta}$$

Entropic force

**Entropic force**  
selects the orbit!  
= order by disorder

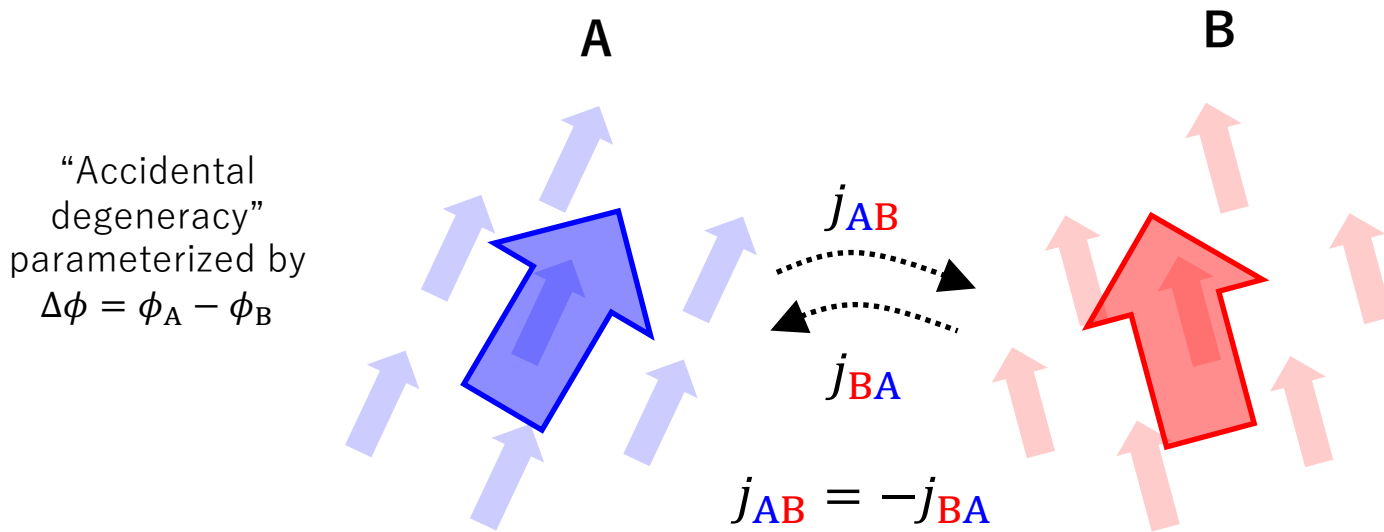
# Order-by-disorder phenomena

$$\dot{\theta}_i^a = \sum_b \sum_{i=1}^{N_b} \frac{j_{ab}}{N_b} \sin(\theta_j^b - \theta_i^a)$$



# Order-by-disorder phenomena

$$\dot{\theta}_i^a = \sum_b \sum_{i=1}^{N_b} \frac{j_{ab}}{N_b} \sin(\theta_j^b - \theta_i^a)$$



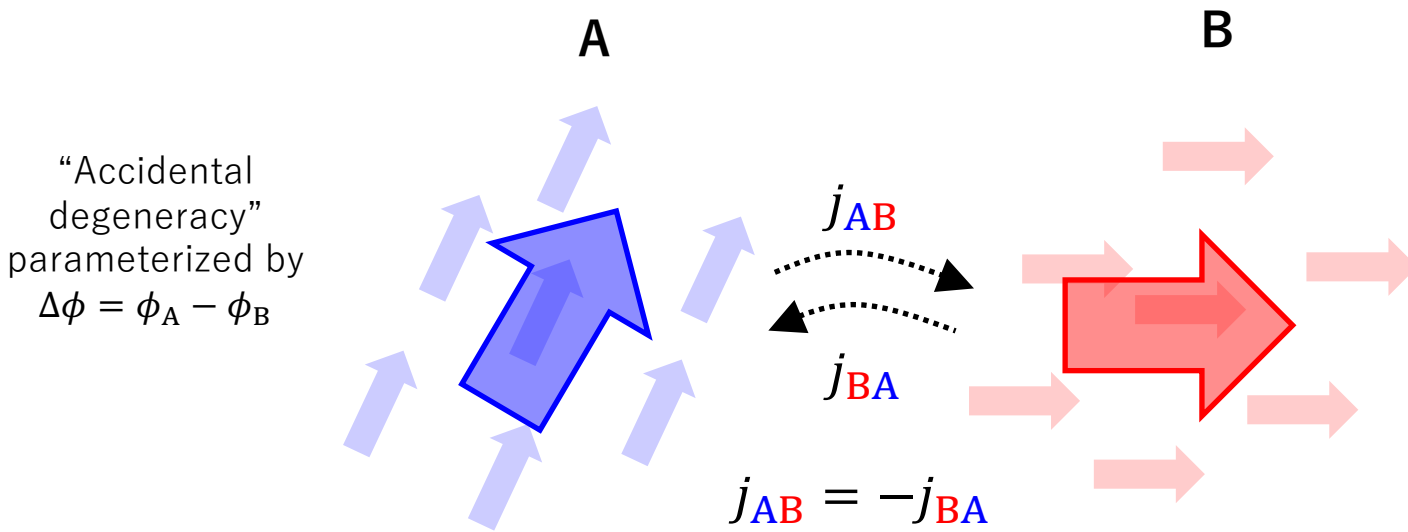
Macroscopic spin (No noise)

$$\dot{\phi}_a = \sum_b j_{ab} \sin(\phi_b - \phi_a)$$

# Order-by-disorder phenomena

$$\dot{\theta}_i^a = \sum_b \sum_{i=1}^{N_b} \frac{j_{ab}}{N_b} \sin(\theta_j^b - \theta_i^a) + \eta_i^a$$

Noise  
 $+\eta_i^a$



Macroscopic spin (No noise)

$$\dot{\phi}_a = \sum_b j_{ab} \sin(\phi_b - \phi_a)$$

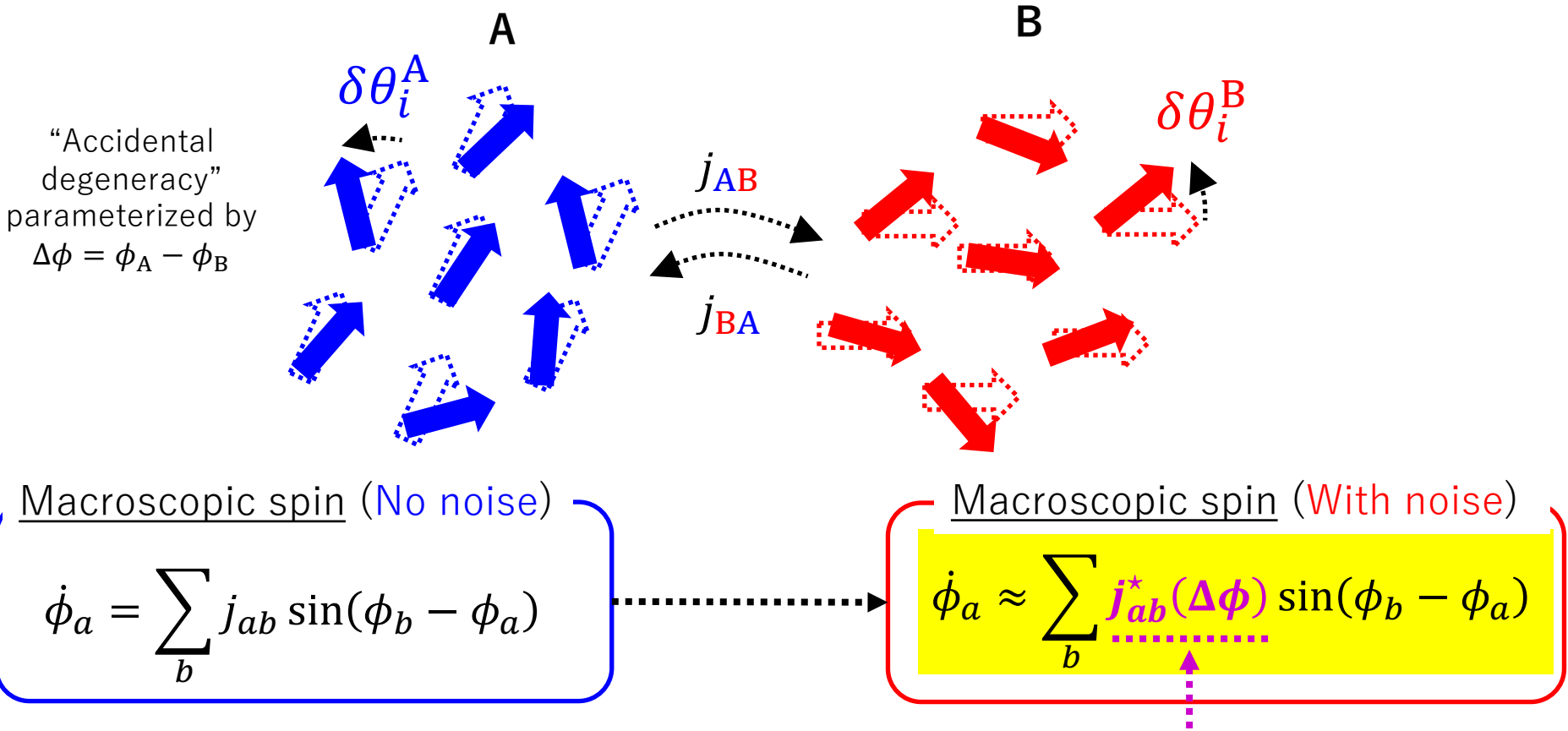


# Order-by-disorder phenomena

$$\dot{\theta}_i^a = \sum_b \sum_{i=1}^{N_b} \frac{j_{ab}}{N_b} \sin(\theta_j^b - \theta_i^a) + \eta_i^a$$

Noise  
+ $\eta_i^a$

Probability distribution of  $\delta\theta_i^a$  is  $\Delta\phi$ -dependent



$\Delta\phi$ -dependent renormalized coupling

# Order-by-disorder phenomena

## No noise

$$\Delta\dot{\phi} = -(j_{AB} + j_{BA}) \sin \Delta\phi = 0$$

Marginal when  $j_{AB} = -j_{BA}$

$$\Delta\phi(t) = \Delta\phi(0)$$

## With noise

$$\Delta\dot{\phi} = -(j_{AB}^*(\phi) + j_{BA}^*(\phi)) \sin \Delta\phi$$

$$\approx \frac{j_0 j_-^2 \sigma^2}{2} \frac{\cos \Delta\phi}{(j_0^2 - j_-^2 \cos^2 \Delta\phi)^2} \sin \Delta\phi$$

when  $j_{AB} = -j_{BA} = j_-$ ,  $j_{AA} = j_{BB} = j_0$

Fixed point at

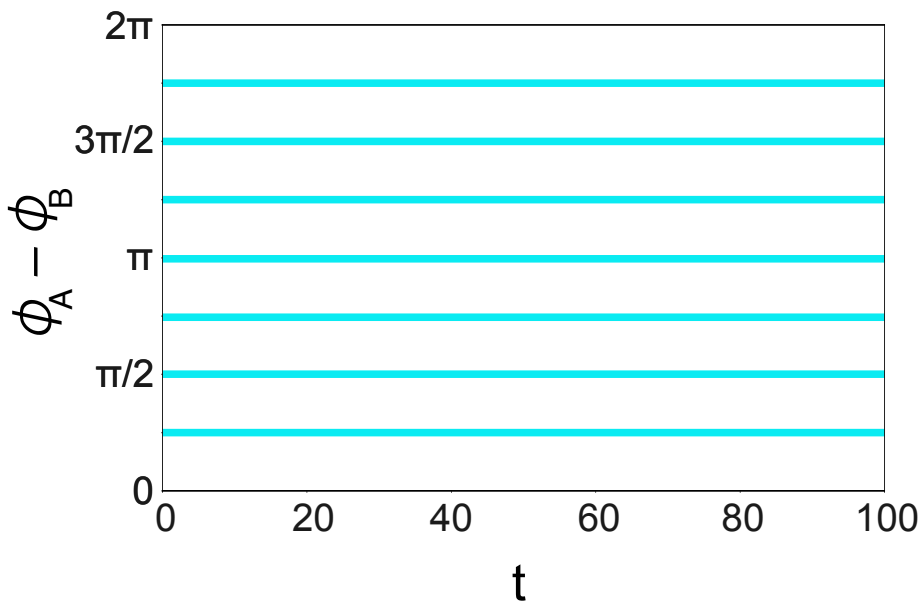
$$\Delta\phi_* = \pm \frac{\pi}{2}$$

“Entropic” force  
favoring  $\Delta\phi = \pm \frac{\pi}{2}$

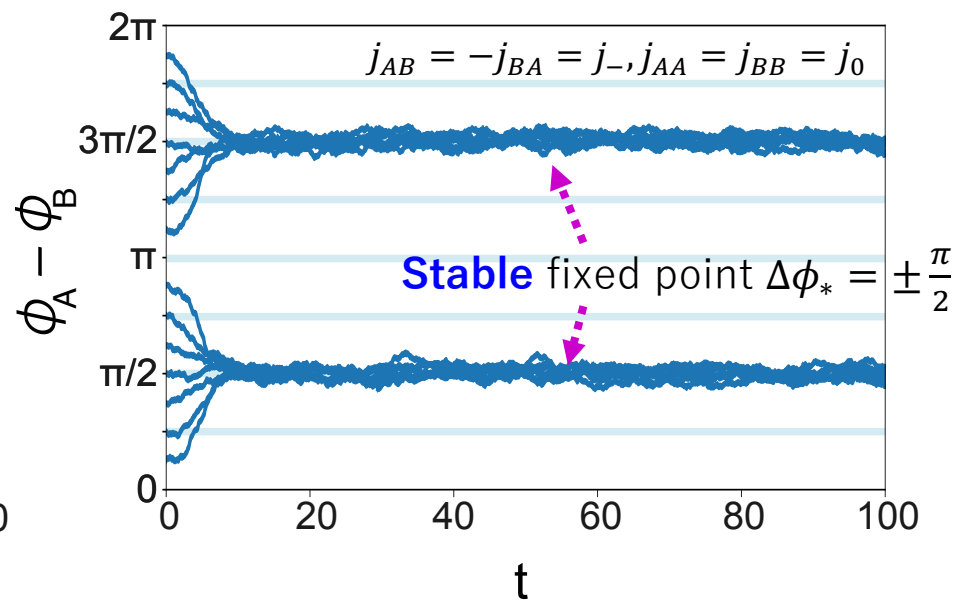
# Order-by-disorder phenomena

$$j_{AB} = -j_{BA}$$

No noise



With noise



Order-by-disorder!

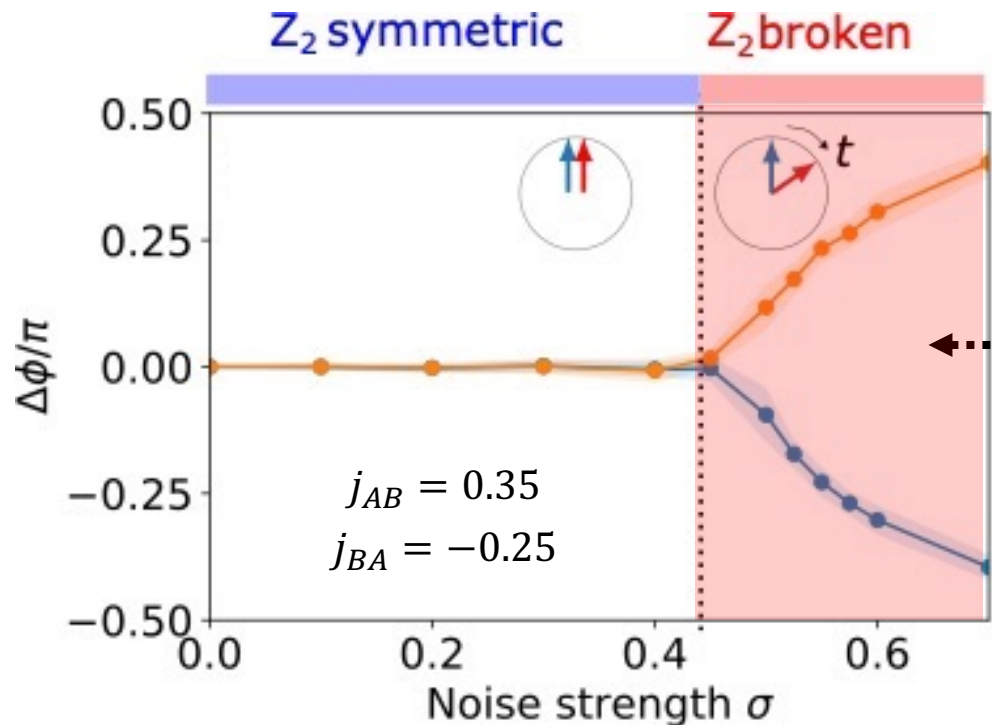
# Noise-induced non-reciprocal phase transition

$$j_{AB} > -j_{BA}$$

$$\Delta\phi = \left[ \underbrace{-2j_+}_{\text{“Energetics”}} + \underbrace{\frac{j_0 j_-^2 \sigma^2}{2} \frac{\cos \Delta\phi}{(j_0^2 - j_-^2 \cos^2 \Delta\phi)^2}}_{\text{“Entropic” force}} \right] \sin \Delta\phi \quad j_+ = j_{AB} + j_{BA} > 0$$

“Energetics”  
favoring  $\Delta\phi = 0$

“Entropic” force  
favoring  $\Delta\phi = \pm \frac{\pi}{2}$



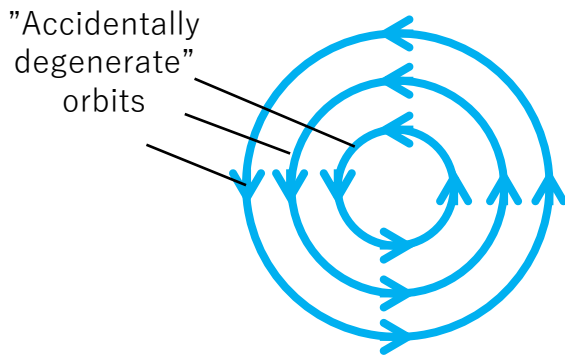
Noise induced  
symmetry breaking!

See also,  
Fruchart\*, RH\*, Littlewood, Vitelli, Nature 2021

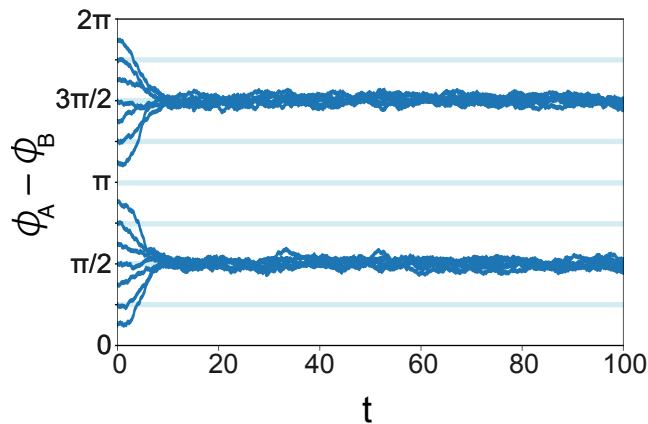
# Summary

- Pointed out a direct analogy between **geometrical** and **non-reciprocal** frustration

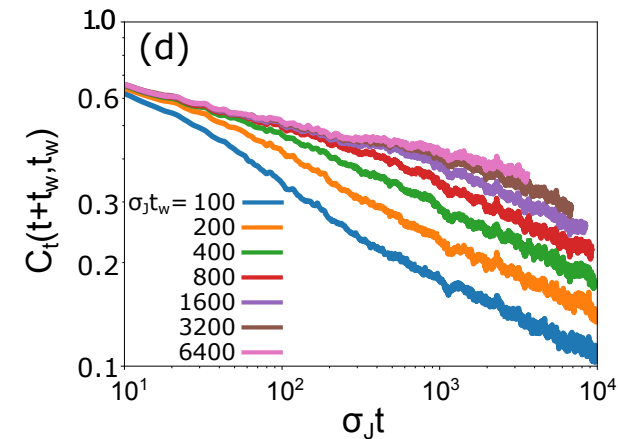
## "Accidental degeneracy" of orbits



## Order-by-disorder



## Spin-glass-like state



RH, Phys. Rev. X **14**, 011029 (2024).  
RH and Weis, in preparation

# Derivation of renormalized coupling

## Equation of motion

$$\dot{\theta}_i^a = - \sum_b \frac{j_{ab}}{N_b} \sum_{j=1}^{N_b} \sin(\theta_i^a - \theta_j^b) + \eta_i^a, \quad \xrightarrow{\text{dotted arrow}} \quad \dot{\theta}_i^a = - \sum_b j_{ab} r_b \sin(\theta_i^a - \phi_b) + \eta_i^a,$$

Order parameter  $\psi_a = r_a e^{i\phi_a} = \frac{1}{N_a} \sum_i^{N_a} e^{i\theta_i^a}$

## Order parameter dynamics in the absence of noise

$$\dot{\phi}_a = \sum_b j_{ab} \sin(\phi_b - \phi_a), \quad \leftarrow \text{Marginal orbit (i.e. initial state dependent orbits) emerges when } j_{ab} = -j_{ba}$$

Fluctuation dynamics  $\delta\theta_i^a = \theta_i^a - \phi_a \quad \delta\dot{\theta}_i^a \approx - \sum_b j_{ab} \cos(\phi_a(t) - \phi_b(t)) \delta\theta_i^a + \eta_i^a.$

## Distribution function

$$\rightarrow \rho_i^a(t, \delta\theta_i^a; \phi(t)) = \frac{1}{\sqrt{\pi} w_a(t; \phi(t))} e^{-(\delta\theta_i^a)^2 / w_a^2(t; \phi(t))}$$

with **configuration dependent** width  $w_a^2(t; \phi(t)) = 2\sigma \int_0^t d\tau e^{-2 \int_\tau^t d\tau' \sum_b j_{ab} \cos(\phi_a(\tau') - \phi_b(\tau'))}$

Especially when  $\Delta\phi_{ab} = \phi_a - \phi_b = \text{const.}$ ,  $w_a^2(\phi) = \frac{\sigma}{\sum_b j_{ab} \cos \Delta\phi_{ab}}$

# Derivation of renormalized coupling

Order parameter dynamics  $\psi_a = r_a e^{i\phi_a}$

$$\dot{\psi}_a = (\dot{r}_a + r_a i \dot{\phi}_a) e^{i\phi_a} = \frac{i}{N_a} \sum_{i=1}^{N_a} \dot{\theta}_i^a e^{i\theta_i^a}$$



$$\begin{aligned} \dot{\phi}_a &= - \sum_b \frac{j_{ab}}{N_a} \sum_{i=1}^{N_a} \frac{r_b}{r_a} \sin(\theta_i^a - \phi_b) \cos(\theta_i^a - \phi_a) + \bar{\eta}_a \\ &= - \sum_b j_{ab}^*(\phi(t)) \sin(\phi_a - \phi_b) + \bar{\eta}_a \end{aligned}$$

with renormalized coupling  $j_{ab}^*(\phi(t)) = j_{ab} \frac{r_b(\phi(t))}{r_a(\phi(t))} \langle \cos^2 \delta\theta_i^a \rangle_{\phi(t)}$

where  $\langle \dots \rangle_{\phi(t)} = \int d\delta\theta_i^a \rho_i^a(t, \delta\theta_i^a; \phi(t)) (\dots)$

Distribution function

$$\rho_i^a(t, \delta\theta_i^a; \phi(t)) = \frac{1}{\sqrt{\pi} w_a(t; \phi(t))} e^{-(\delta\theta_i^a)^2 / w_a^2(t; \phi(t))}$$

Here, self-averaging  $\frac{1}{N_a} \sum_i^{N_a} (\dots) = \langle \dots \rangle_{\phi}$  is assumed.

Macroscopic noise strength

$$\langle \bar{\eta}_a(t) \bar{\eta}_b(t') \rangle \approx \frac{\sigma}{N_a} \delta_{ab} \delta(t - t') \rightarrow 0 \quad (N_a \rightarrow \infty)$$

# "Accidental degeneracy" of *orbits*

## [Proof]

Continuity equation:  $\frac{\partial \rho}{\partial t} = - \sum_i \frac{\partial(\rho \dot{\theta}_i)}{\partial \theta_i} = - \sum_i \left[ \frac{\partial \rho}{\partial \theta_i} \dot{\theta}_i + \rho \frac{\partial \dot{\theta}_i}{\partial \theta_i} \right]$

$$\sum_i \frac{\partial \dot{\theta}_i}{\partial \theta_i} = \sum_{ij} [J_{ij} \cos(\theta_j - \theta_i)] = 0$$

$J_{ij} = -J_{ji}$

$$\dot{\theta}_i = \sum_j J_{ij} \sin(\theta_j - \theta_i)$$

Therefore,  $\frac{\partial \rho}{\partial t} + \sum_i \frac{\partial \rho}{\partial \theta_i} \dot{\theta}_i = 0.$  ■

✘A similar theorem holds for non-reciprocally interacting particles and Heisenberg spins.