

Route to turbulence of bacterial suspensions under confinement

Kazumasa A. Takeuchi (UTokyo)

with



Daiki Nishiguchi

(UTokyo) → ~~Tokyo Tech~~ from Oct

[experiment]



Sora Shiratani

(UTokyo)

[numerics]

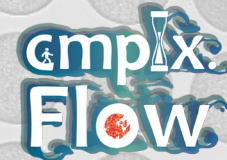


Igor S. Aranson

(Penn State & UTokyo)

[theory]

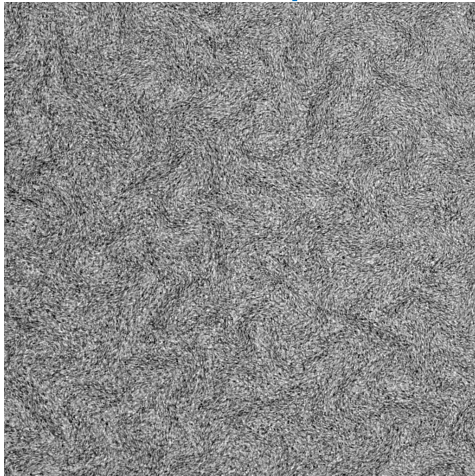
Nishiguchi et al., arXiv: 2407.05269



JST PRESTO (D.N.) Active Matter c2c

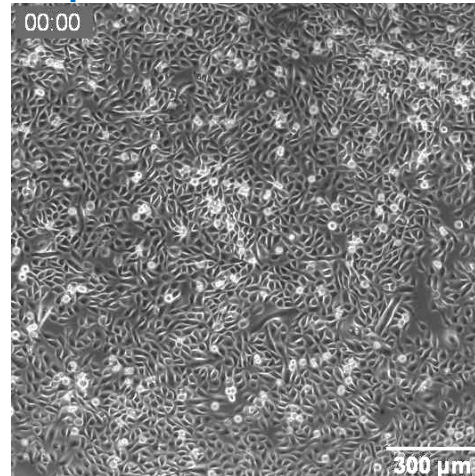
Collective motion is often chaotic

bacterial suspension



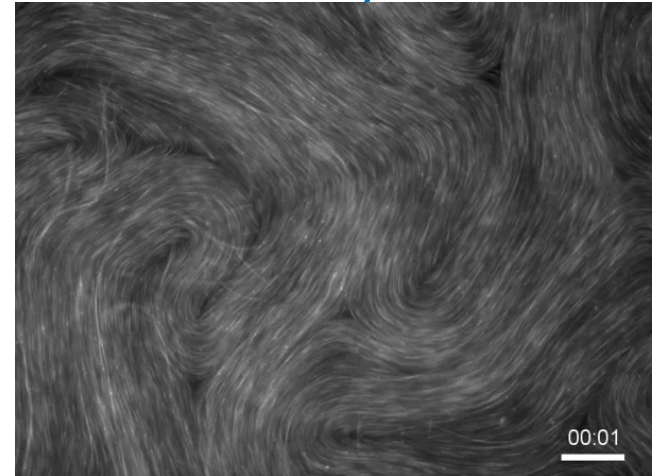
[Nishiguchi+, Nat Comm 2018]

epithelial cell sheet



[Blanch-Mercader+, PRL 2018]

reconstituted cytoskeletons



[Sanchez+, Nature 2012]

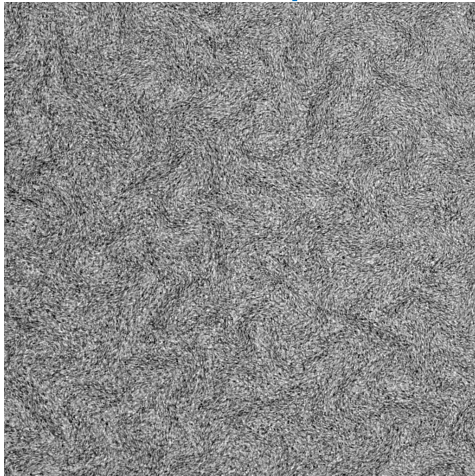
“active turbulence”

[review:Alert et al., Annu. Rev. Condens. Matter Phys. 13, 143 (2022)]

Vortices have a characteristic length scale.

Active turbulence self-organizes by confinement

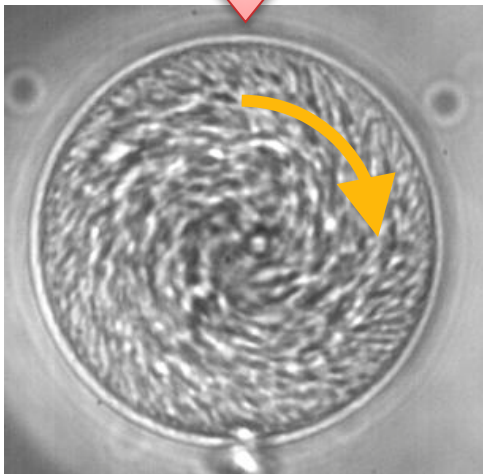
bacterial suspension



[Nishiguchi+, Nature Comm 2018]

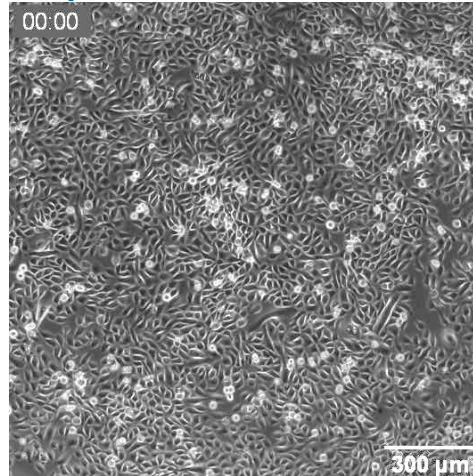


confine



[Wioland+, PRL 2013]

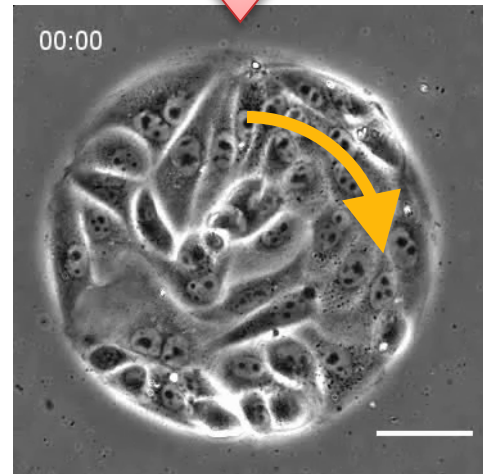
epithelial cell sheet



[Blanch-Mercader+, PRL 2018]

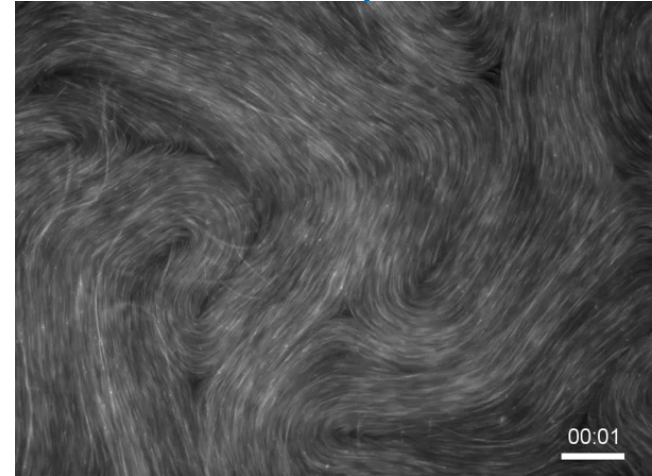


confine

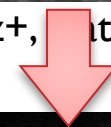


[Doxzen+, Integr. Biol. 2013]

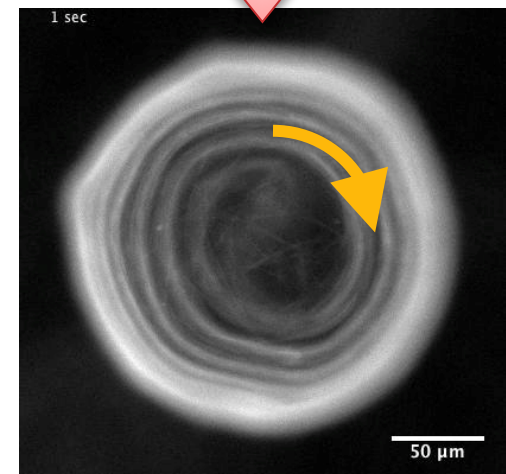
reconstituted cytoskeletons



[Sanchez+, Nature 2012]



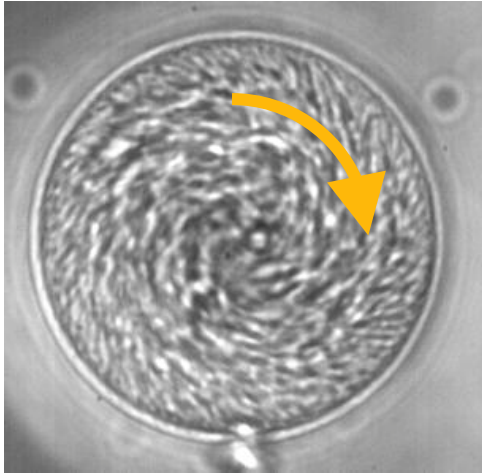
confine



[Opathalage+, PNAS 2019]

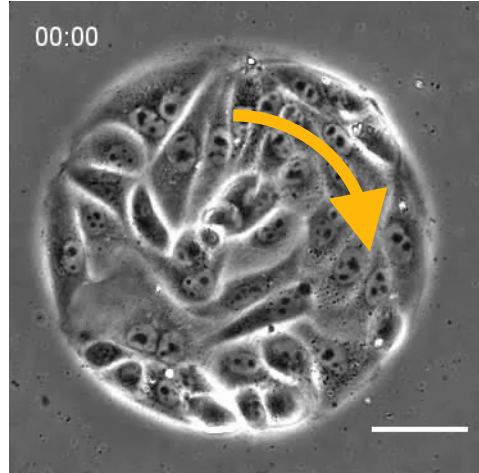
How does the vortex get destabilized?

bacterial suspension



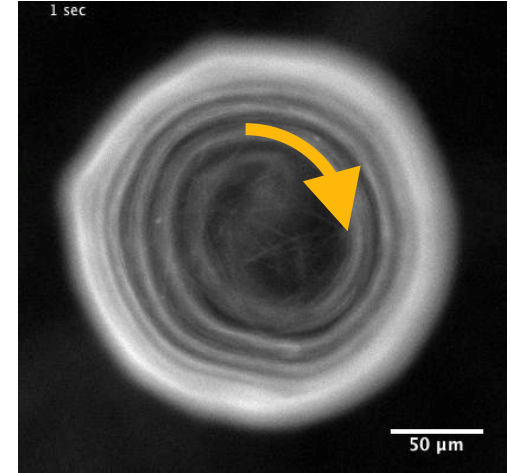
[Wioland+, PRL 2013]

epithelial cell sheet



[Doxzen+, Integr. Biol. 2013]

reconstituted cytoskeletons



[Opathalage+, PNAS 2019]

polar



nematic



What's the "route to turbulence"
& underlying mechanism
for this case?

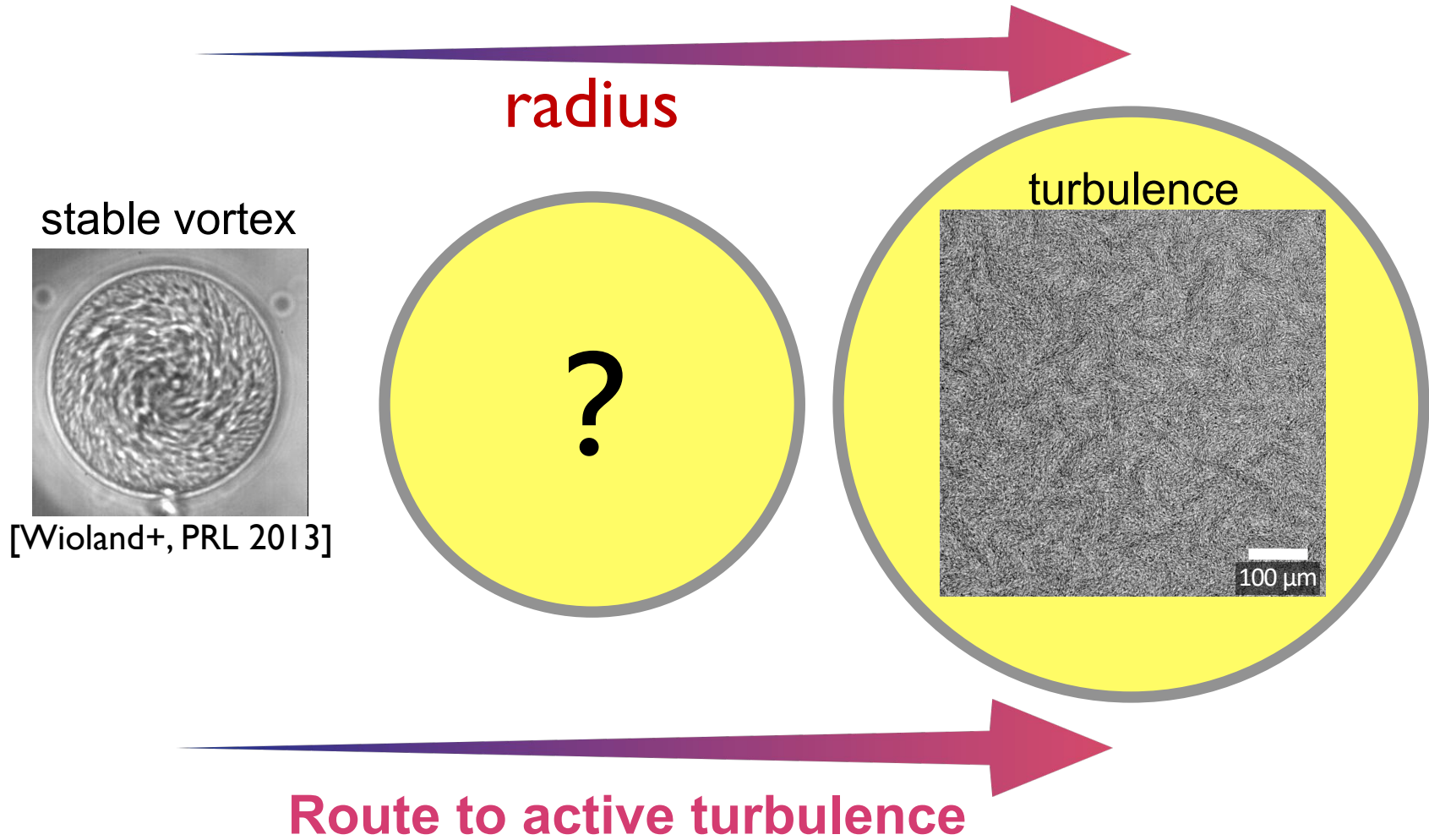
Today's question

unbinding of $\pm 1/2$ defects destabilizes vortices,
leading to active turbulence

[Opathalage+, PNAS 116, 4788 (2019)]

Let's change the degree of confinement!

Consider **varying radius of confining wells**

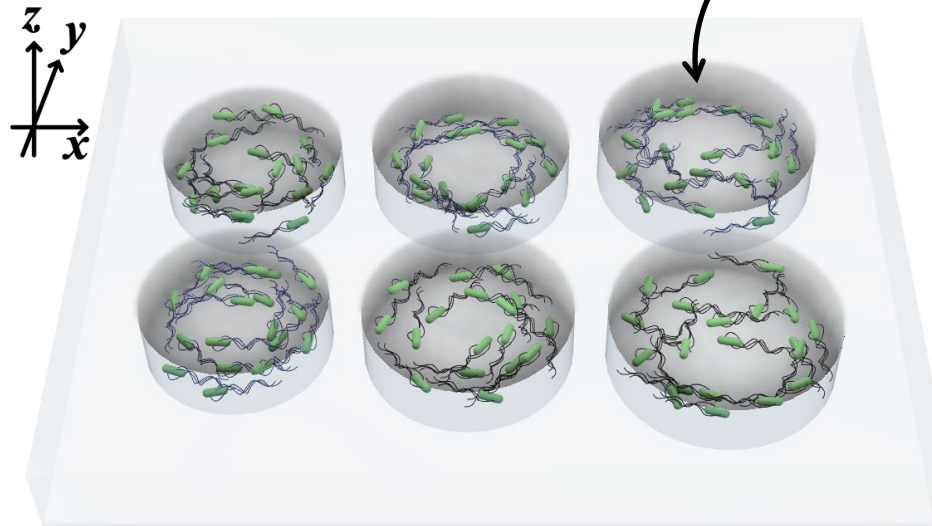


Experiment

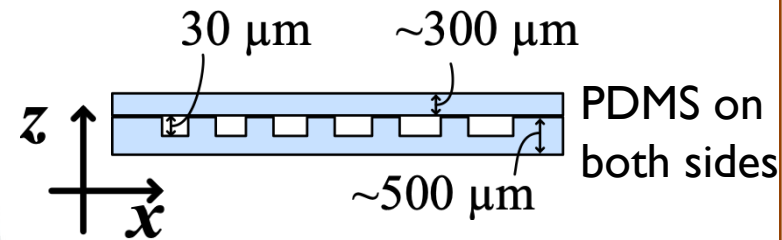


Daiki Nishiguchi

Our setup



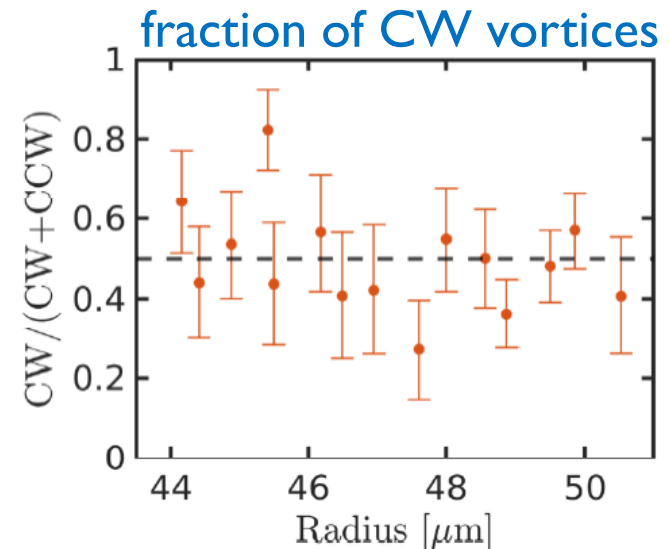
dense suspensions of *Bacillus subtilis*



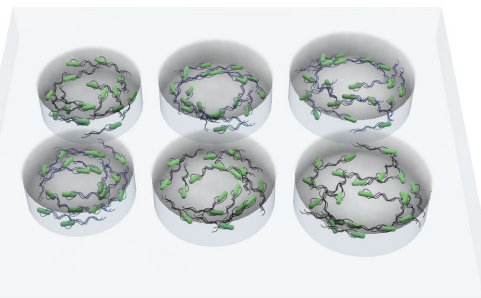
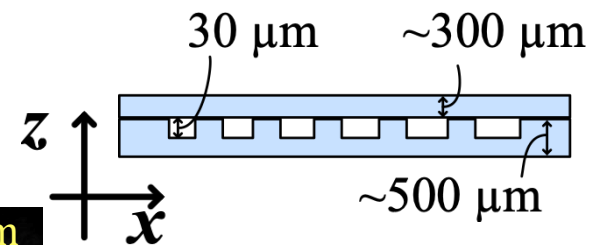
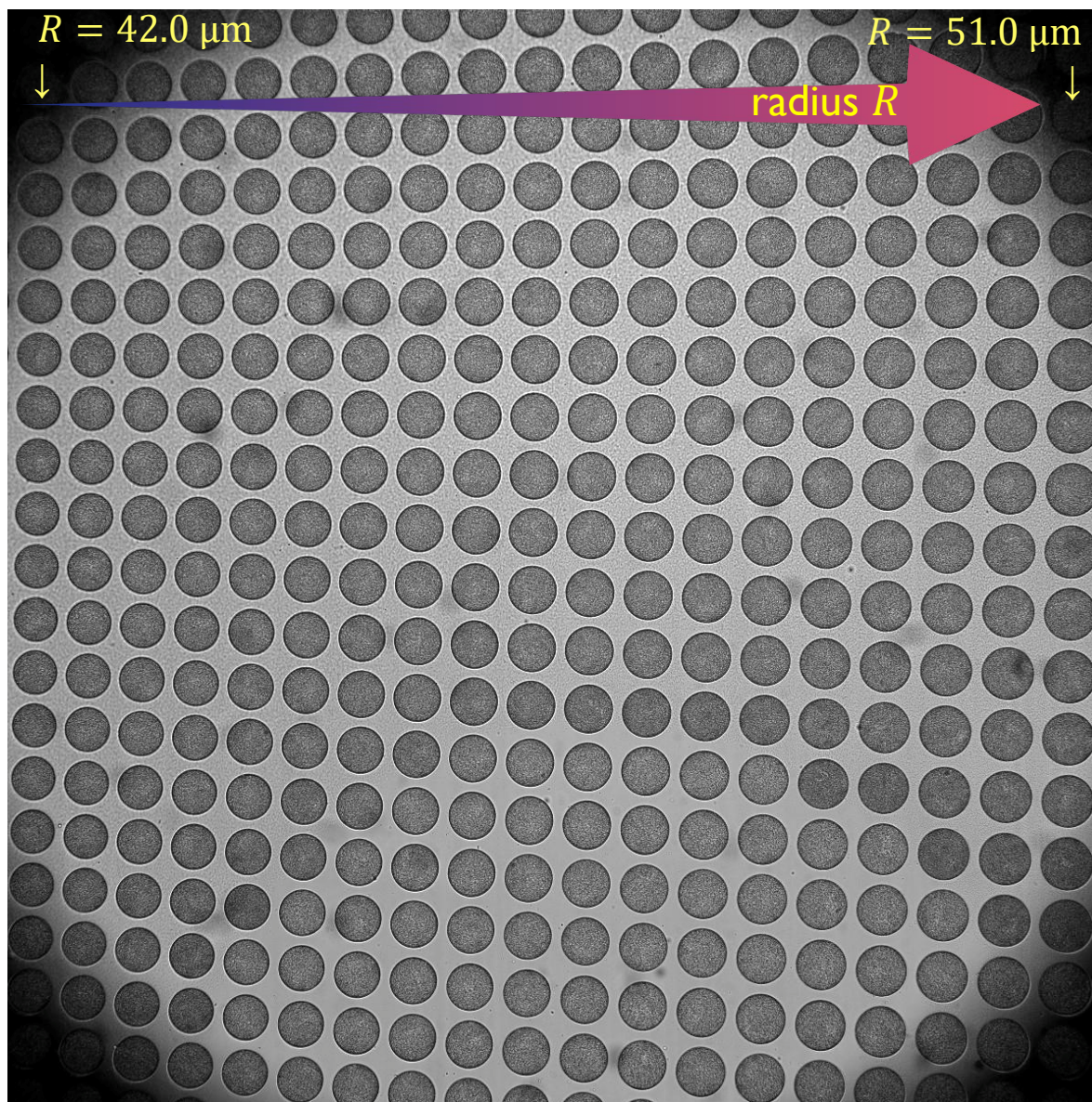
well radius: 42.0-51.0 μm ,
0.5 μm step

Key points

- Entirely covered by PDMS (high O_2 permeability)
→ Long observation (several minutes)
- Same solid surface on both sides
→ No chirality



Experiment



well radius: 42.0-51.0 μm ,
0.5 μm step

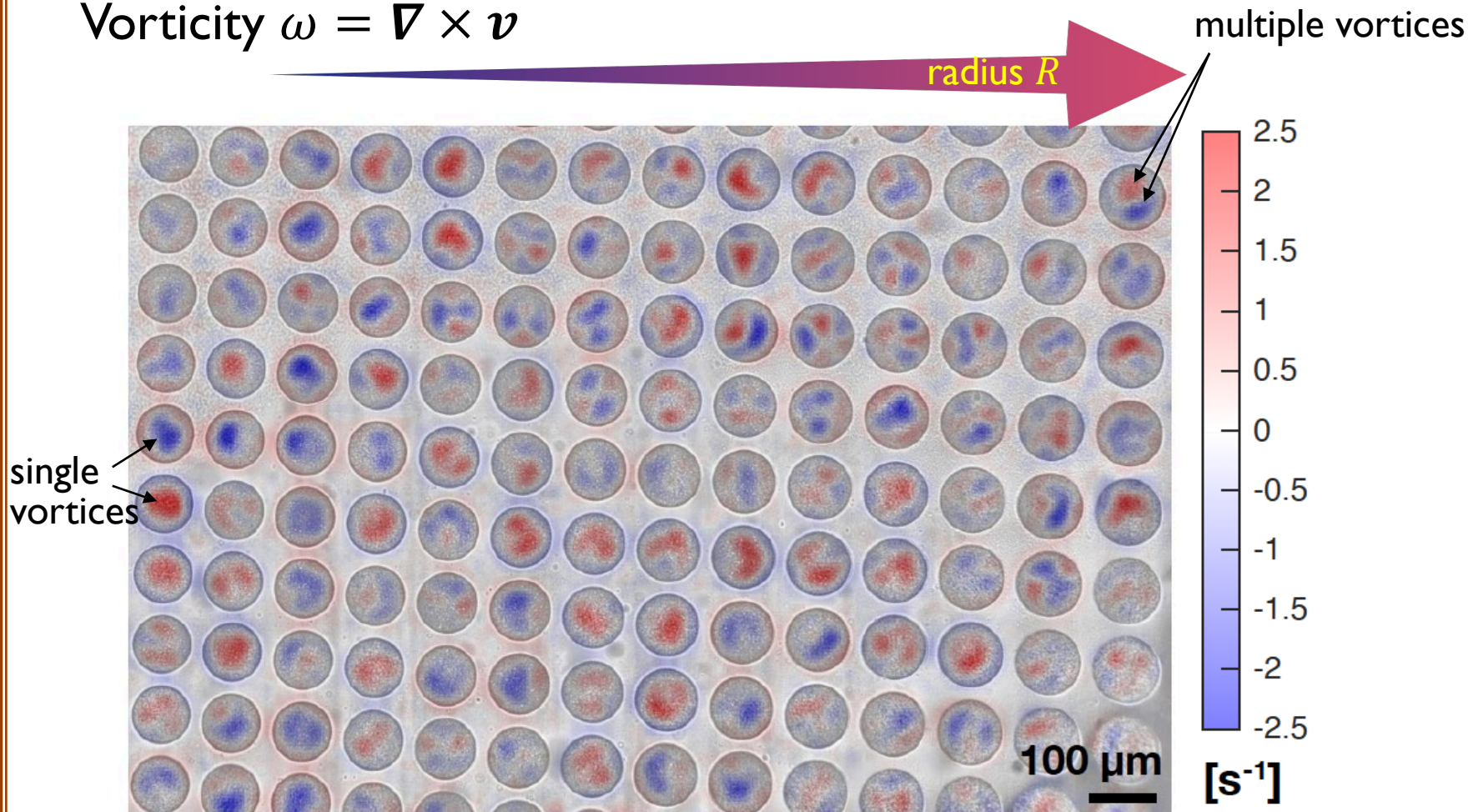
Key points

Simultaneous
observation of
numerous wells
→ allowed us to
capture transitions!

200 μm

Vortices get destabilized as radius increases

$$\text{Vorticity } \omega = \nabla \times \mathbf{v}$$



Single vortex in small wells, multiple vortices in larger wells.

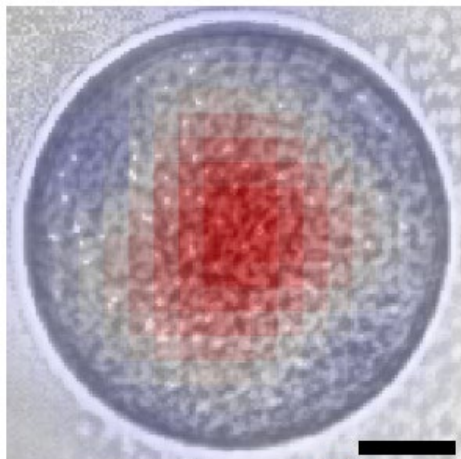
Vortices get destabilized as radius increases

$$\text{Vorticity } \omega = \nabla \times v$$

radius R

$R = 44.6 \mu\text{m}$

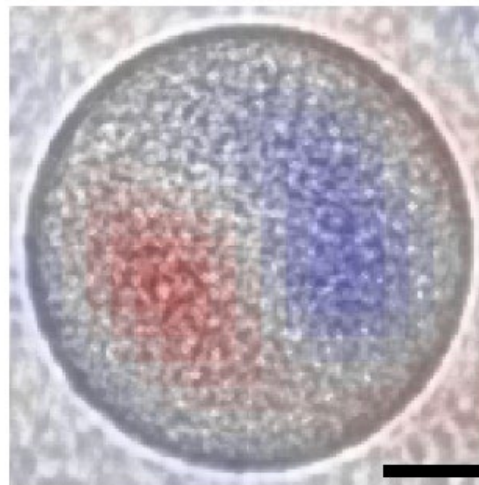
1 vortex



20 μm

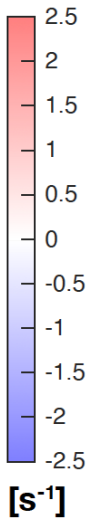
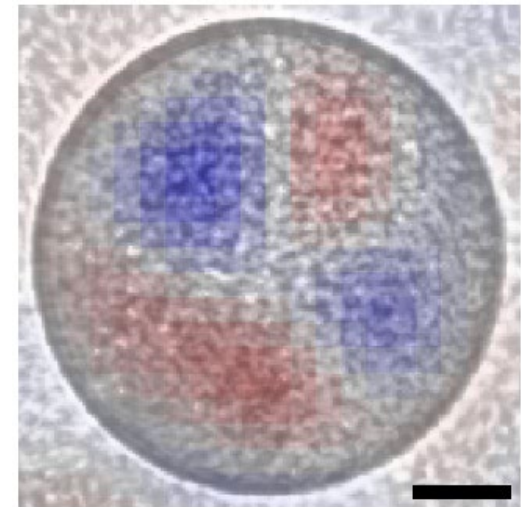
$R = 46.7 \mu\text{m}$

2 vortices



$R = 48.8 \mu\text{m}$

4 vortices



Single vortex in small wells, multiple vortices in larger wells.

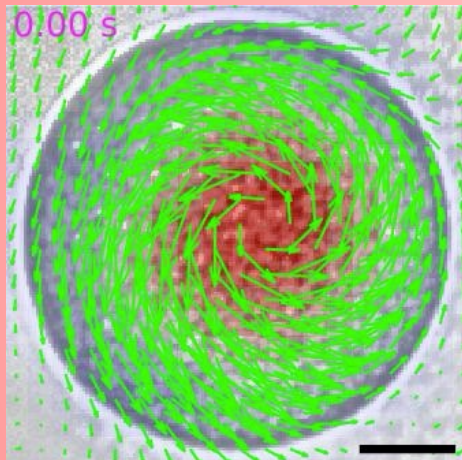
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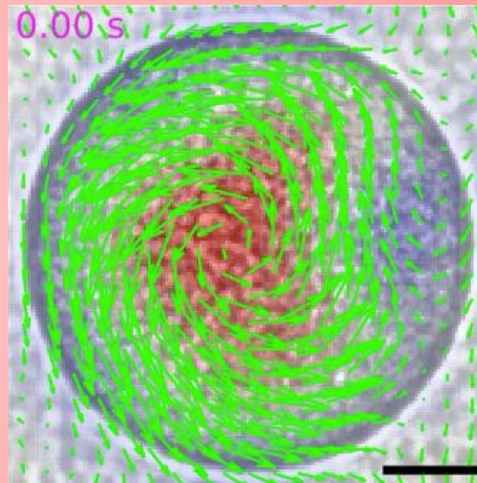
x4 fast

20 μm

stable for long time

$R = 46.7 \mu\text{m}$

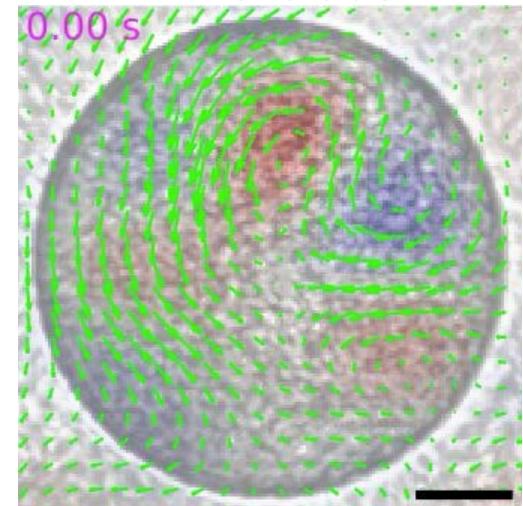
2 vortices



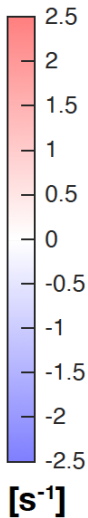
reversals

$R = 48.8 \mu\text{m}$

4 vortices



pulsations



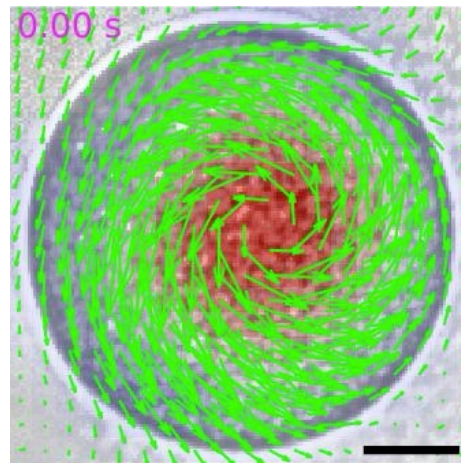
Let's focus on this transition!

Detecting the vortex reversals

Sign of vortex: “spin”

$$S_i(t) := \frac{\hat{z} \cdot \sum_{\mathbf{r} \in \text{well}} (\mathbf{r} - \mathbf{r}_{\text{center}}) \times \mathbf{v}(\mathbf{r}, t)}{\sum_{\mathbf{r} \in \text{well}} |\mathbf{r} - \mathbf{r}_{\text{center}}|}$$

stable single vortex

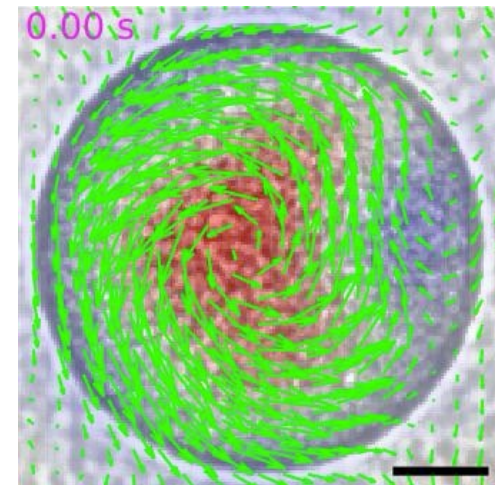


$R = 44.6 \mu\text{m}$

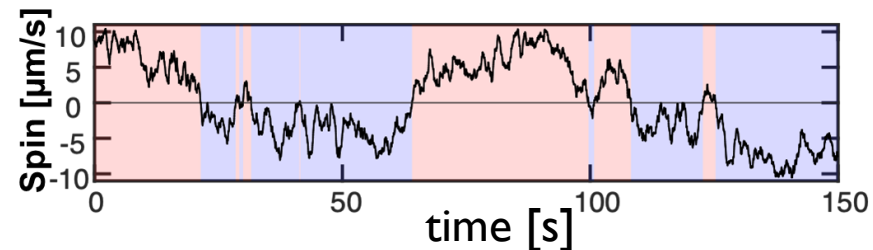
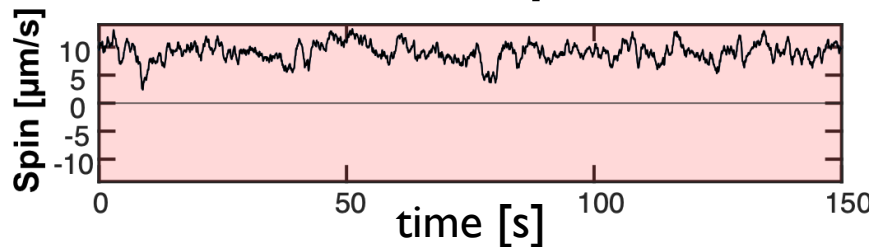
increasing radius



reversing vortex pair



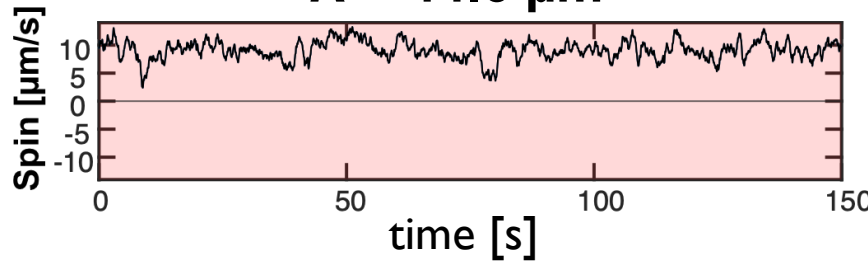
$R = 46.7 \mu\text{m}$



Detecting the vortex reversals

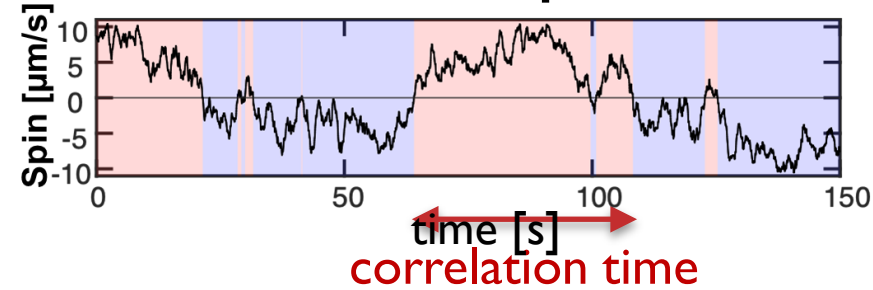
stable single vortex

$R = 44.6 \mu\text{m}$

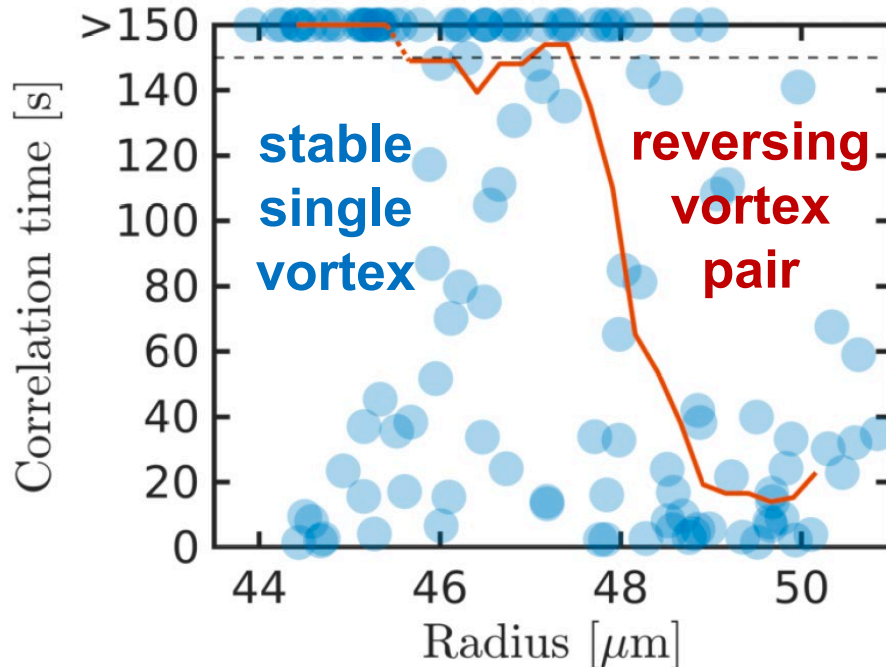


reversing vortex pair

$R = 46.7 \mu\text{m}$



correlation time vs radius



determined from $\overline{S_i(t)S_i(t + \tau)}$

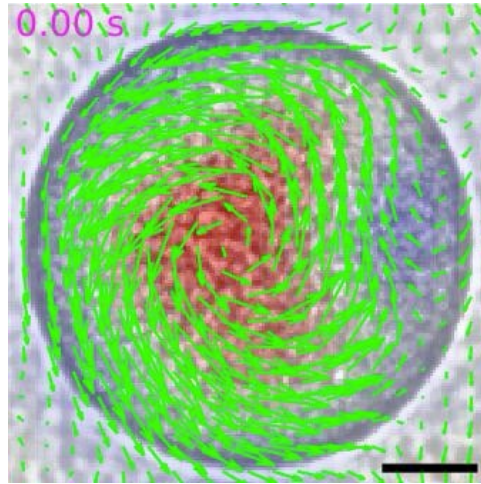
blue: each well

red: moving median

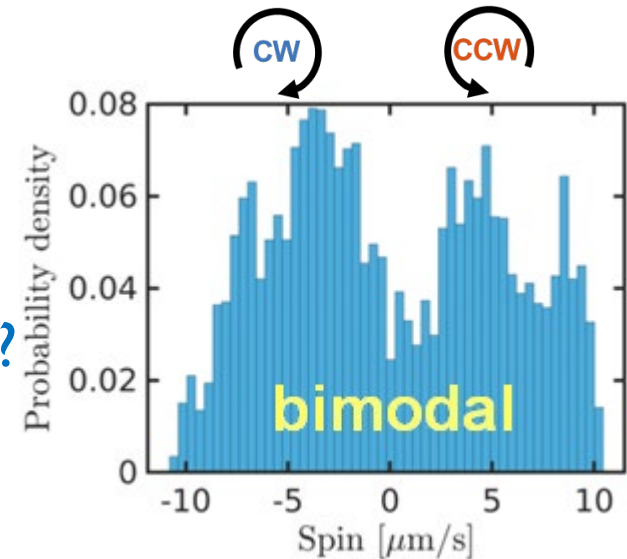
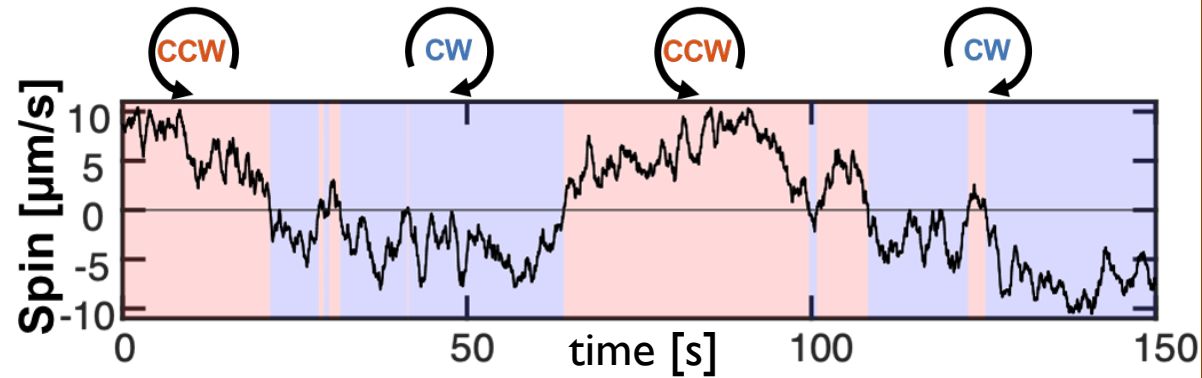
**Transition is
clearly detected**

Vortex reversals: regular or irregular?

reversing vortex pair



$R = 46.7 \mu\text{m}$



- Reversal looks rather regular & periodic...?
- Bimodal distribution of CW & CCW

Vortex reversal = periodic (?) switching of CW/CCW states

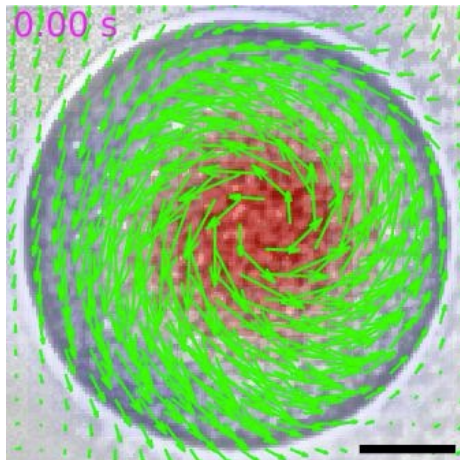
periodicity will be justified later

What about the 4-vortex state?



$R = 44.6 \mu\text{m}$

1 vortex



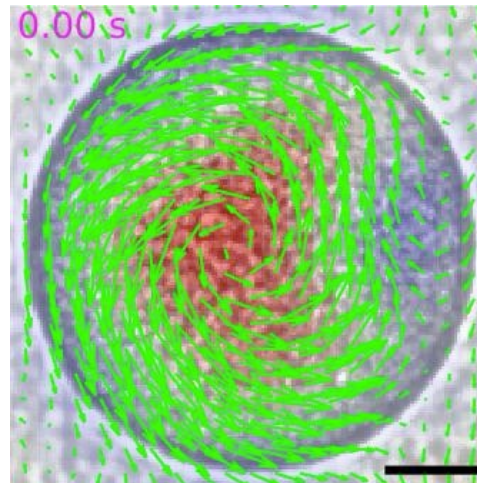
x4 fast

20 μm

stable for long time

$R = 46.7 \mu\text{m}$

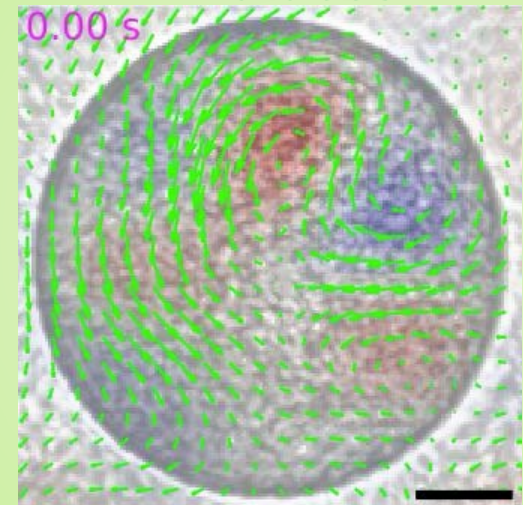
2 vortices



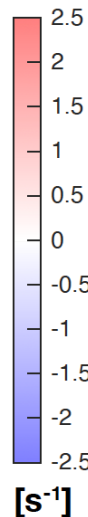
reversals

$R = 48.8 \mu\text{m}$

4 vortices



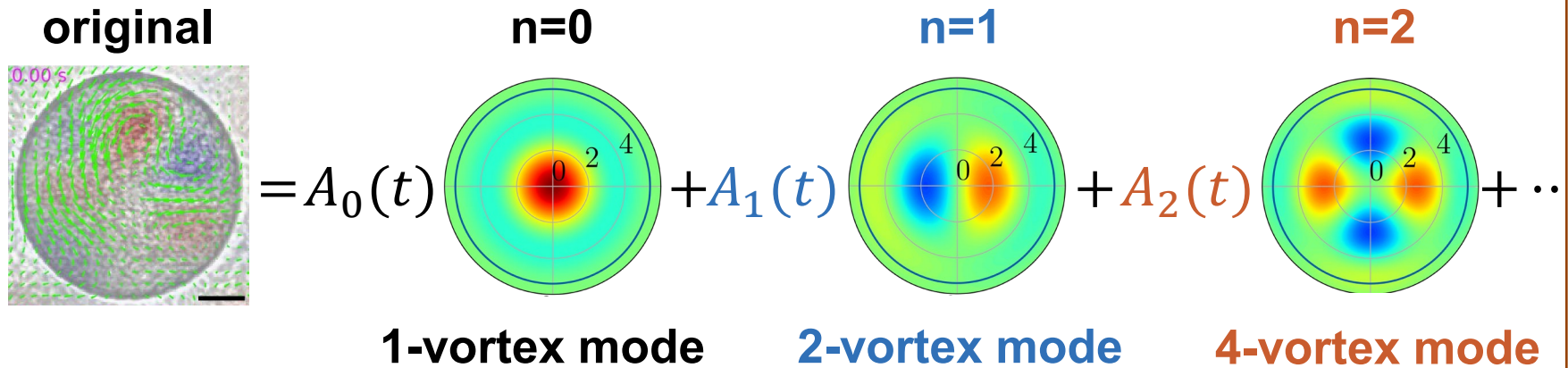
pulsations



Now let's characterize this!

Multipole expansion of 4-vortex state

Decompose $v(r, \theta)$ into angular Fourier modes: $\int d\theta e^{-in\theta} v(r, \theta) + \text{c. c.}$



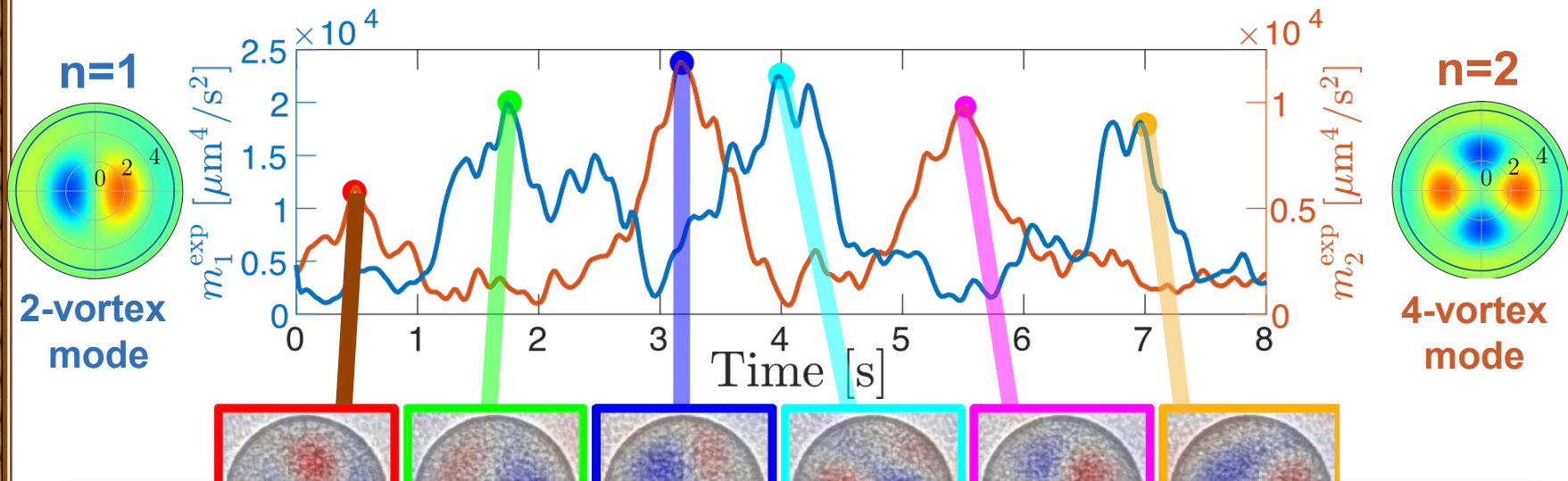
Calculate kinetic energy of each mode

$$m_n^{\text{exp}}(t) = \int_0^R dr r \left| \frac{1}{2\pi} \int d\theta e^{-in\theta} v(r, \theta; t) \right|^2$$

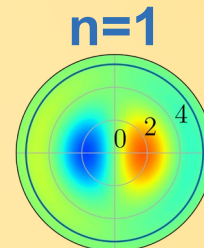
to evaluate its contribution

Multipole expansion of 4-vortex state

Kinetic energy of n th mode:
$$m_n^{\text{exp}} = \int_0^R dr r \left| \frac{1}{2\pi} \int d\theta e^{-in\theta} v(r, \theta) \right|^2$$

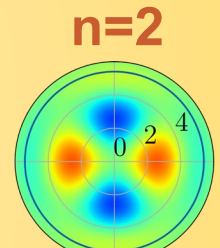


4-vortex state = anti-phase oscillation of



2-vortex mode

&

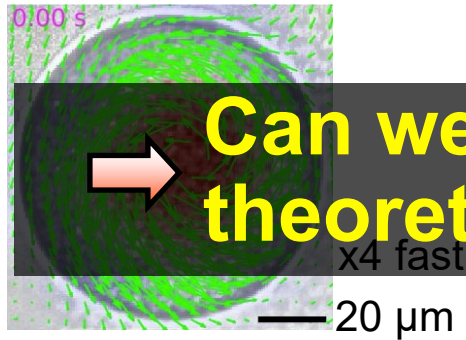


4-vortex mode

Summary so far: Route to turbulence (exp't)



$R = 44.6 \mu\text{m}$
1 vortex



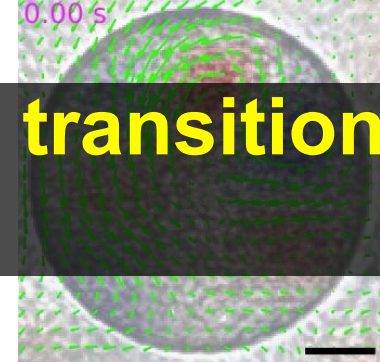
stable for long time

$R = 46.7 \mu\text{m}$
2 vortices



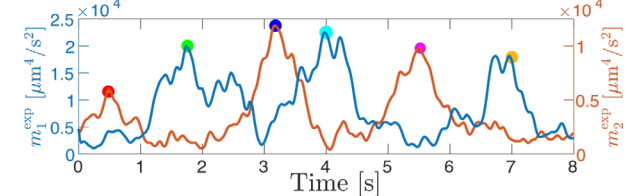
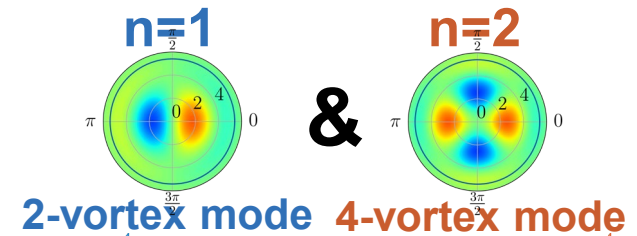
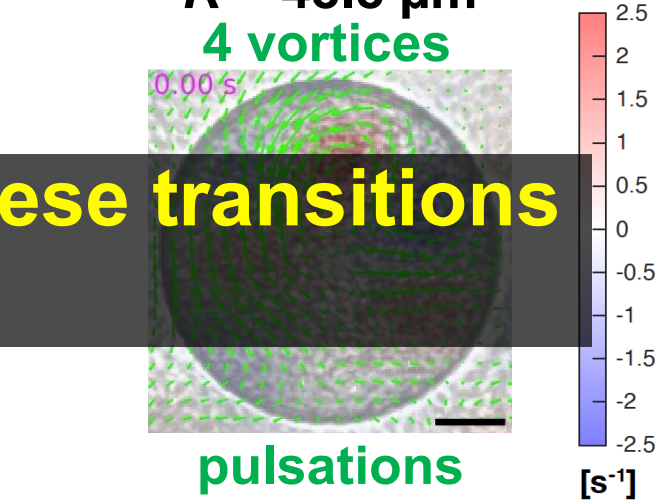
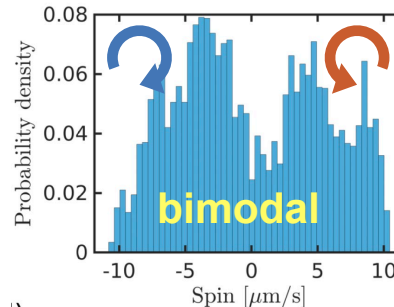
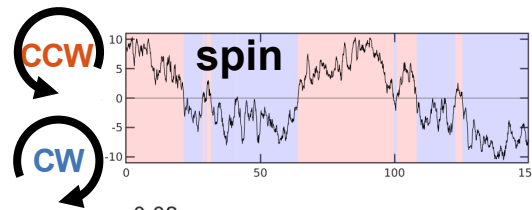
regular reversals

$R = 48.8 \mu\text{m}$
4 vortices



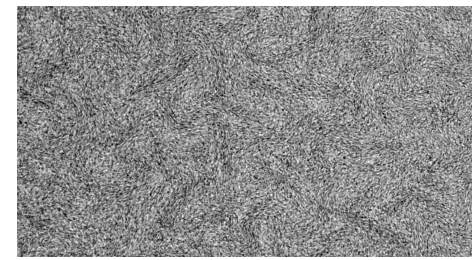
pulsations

Anti-phase oscillation



Can we understand these transitions theoretically?

Model for bacterial turbulence?



Hydro descriptions of bacterial turbulence w/o confinement.

Two phenomenological models in the literature.

Toner-Tu-Swift-Hohenberg (TTSH) equation

Wensink+, PNAS 2012
Dunkel+, New J Phys 2013

$$\begin{cases} \nabla \cdot \mathbf{v} = 0 \\ \partial_t \mathbf{v} + \lambda(\mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + a\mathbf{v} - b|\mathbf{v}|^2 \mathbf{v} - \Gamma_0 \nabla^2 \mathbf{v} - \Gamma_2 \nabla^4 \mathbf{v} \end{cases}$$

active stress propulsion vortex formation

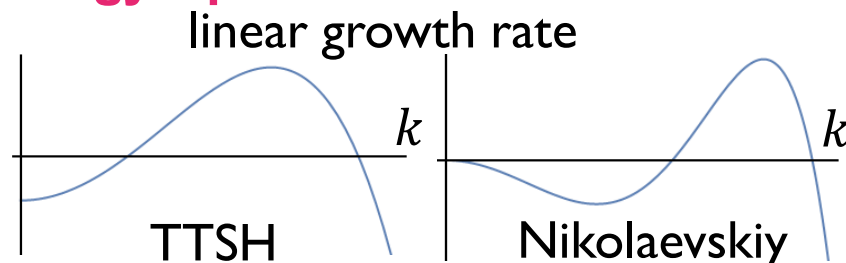
Nikolaevskiy equation

Beresnev & Nikolaevskiy, Physica D 1993
Słomka & Dunkel, PNAS 2017

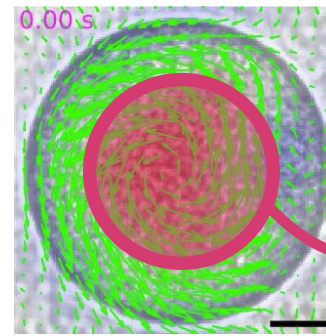
$$\begin{cases} \nabla \cdot \mathbf{v} = 0 \\ \partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + \Gamma_0 \nabla^2 \mathbf{v} + \Gamma_2 \nabla^4 \mathbf{v} + \Gamma_4 \nabla^6 \mathbf{v} \end{cases}$$

energy input

In both models, vortices have a characteristic size.



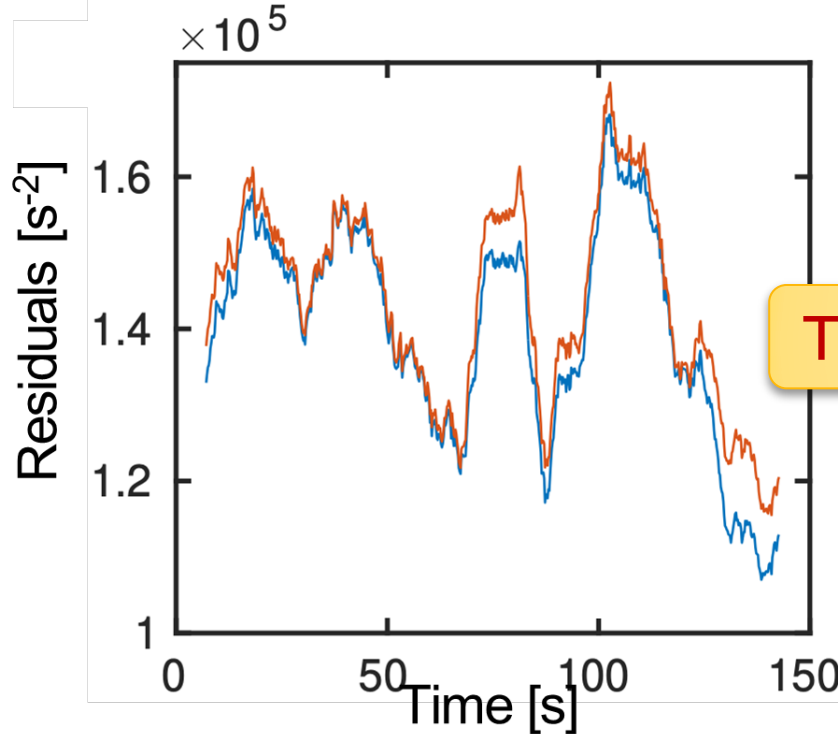
Experimental test



We only used
area far from boundary

For each equation, we

- determined all coefficients by least squares from exp't data
- evaluated the residual for each time t



Always,
 $\text{residual}(\text{Nikolaevskiy}) > \text{residual}(\text{TTSH})$

TTSH outperforms Nikolaevskiy!

In the following, we use TTSH only.

Boundary conditions

TTSH equation (non-dimensionalized)

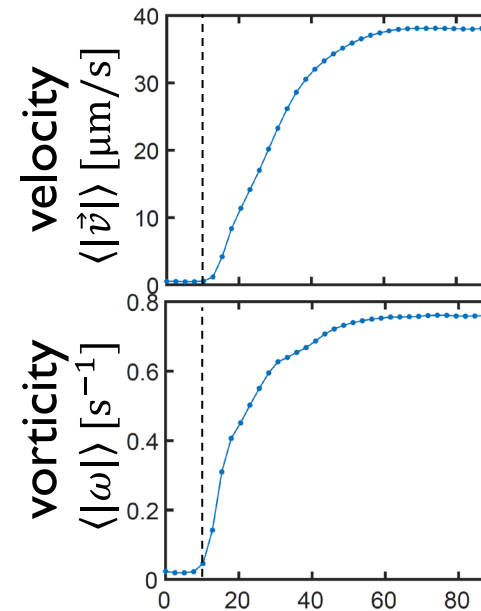
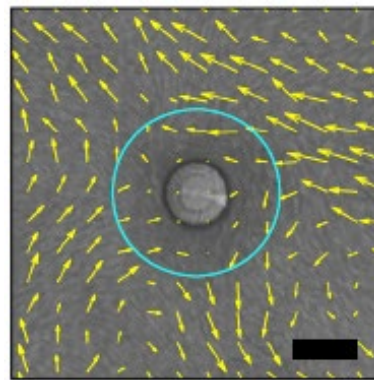
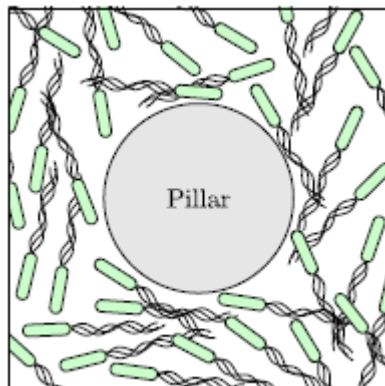
$$\partial_t \vec{v} + \lambda \vec{v} \cdot \nabla \vec{v} = -\nabla q + a \vec{v} - b |\vec{v}|^2 \vec{v} - (1 + \nabla^2)^2 \vec{v}$$

higher-order derivative

→ more boundary conditions needed

Boundary conditions were experimentally inferred.

[Reinken, Nishiguchi, ... Comm Phys 3, 76 (2020)]



$\vec{v} = 0$ & $\omega = 0$
at boundary



distance from pillar center [μm]

Daiki Nishiguchi

Boundary conditions

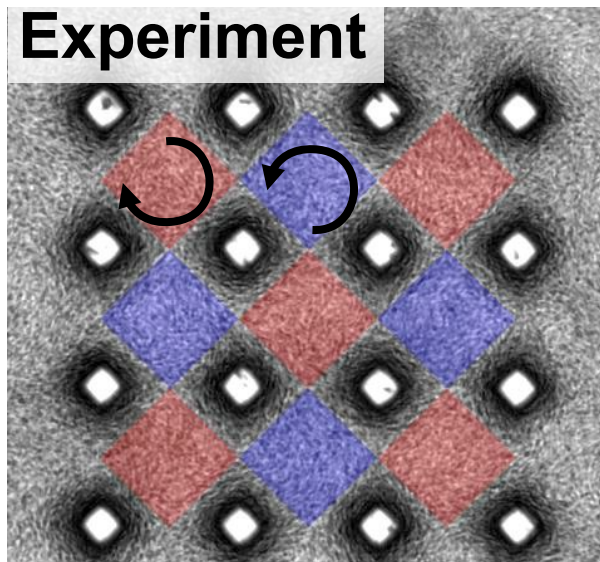
TTSH equation (non-dimensionalized)

$$\partial_t \vec{v} + \lambda \vec{v} \cdot \nabla \vec{v} = -\nabla q + a \vec{v} - b |\vec{v}|^2 \vec{v} - (1 + \nabla^2)^2 \vec{v}$$

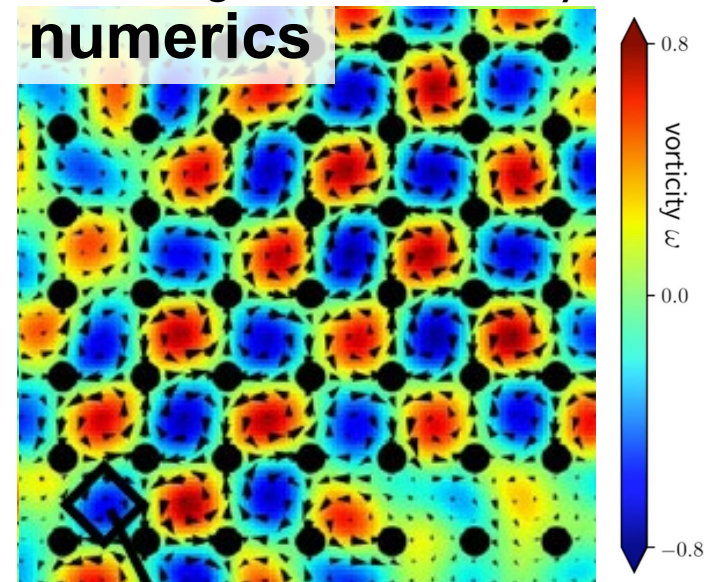
with $\vec{v} = 0$ & $\omega = 0$ at boundary.

This reproduced non-trivial vortex lattice order observed with pillar array.

[Reinken, Nishiguchi, ... Comm Phys 3, 76 (2020)]



red: ↻ blue: ↺ x0.5 speed

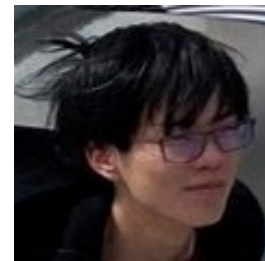


Let's use this for circular wells & inspect the route to turbulence!

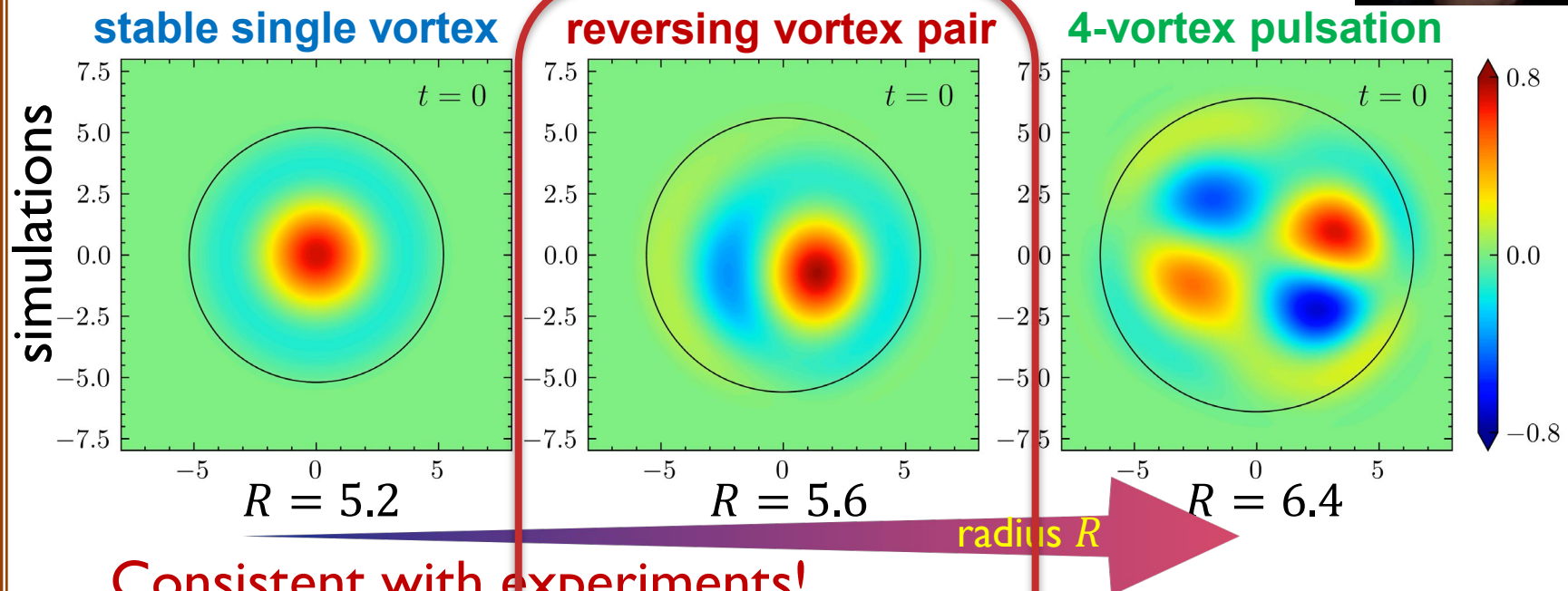
Circular well simulations

Sora Shiratani

(see also arXiv:2304.03306
for computation method)

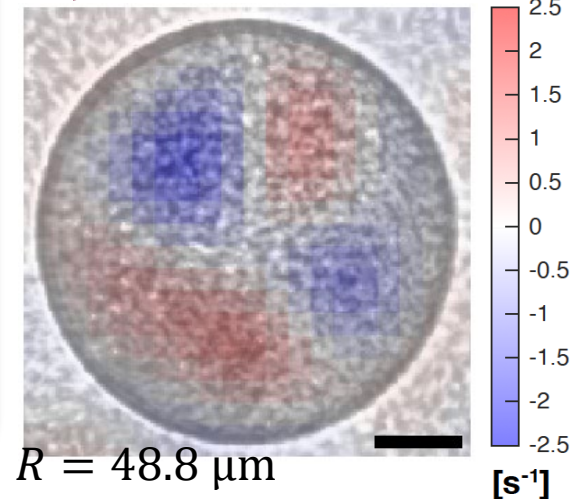
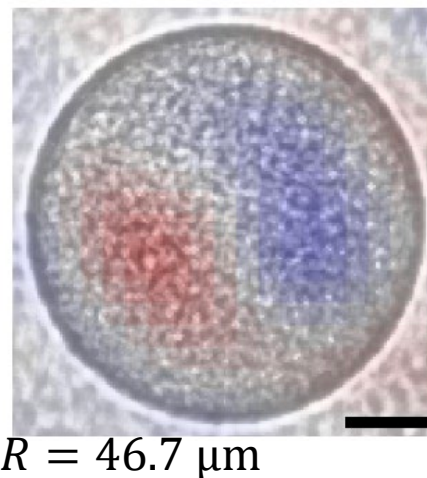
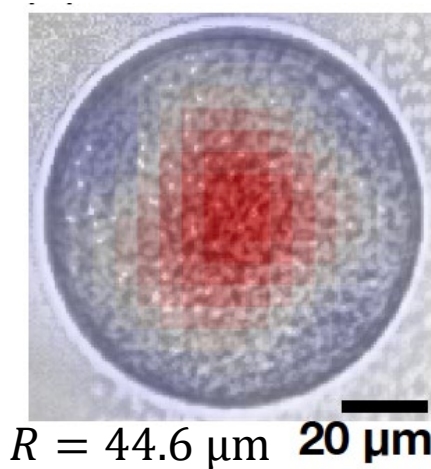


Let's analyze



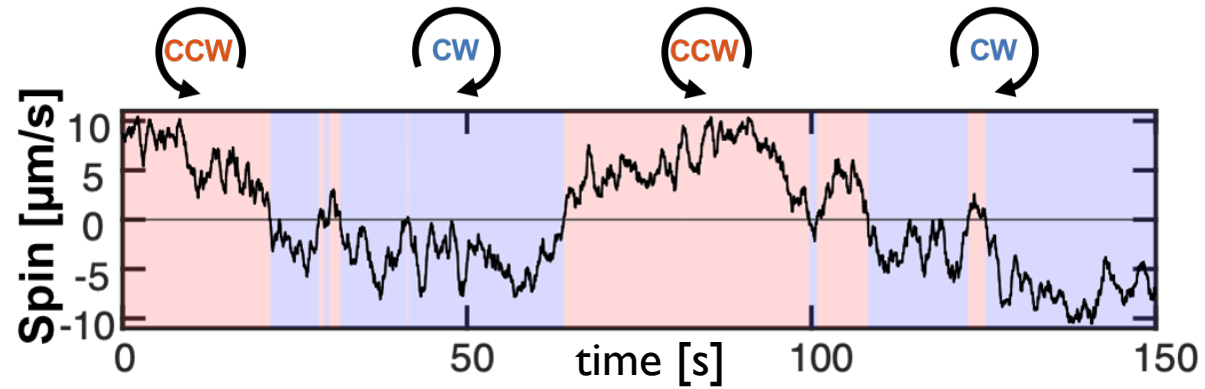
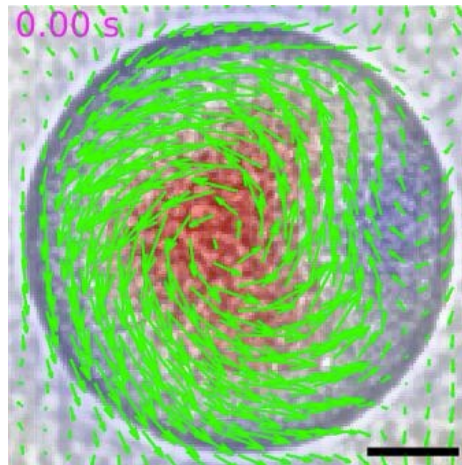
Consistent with experiments!

experiment



Reversing vortex pair

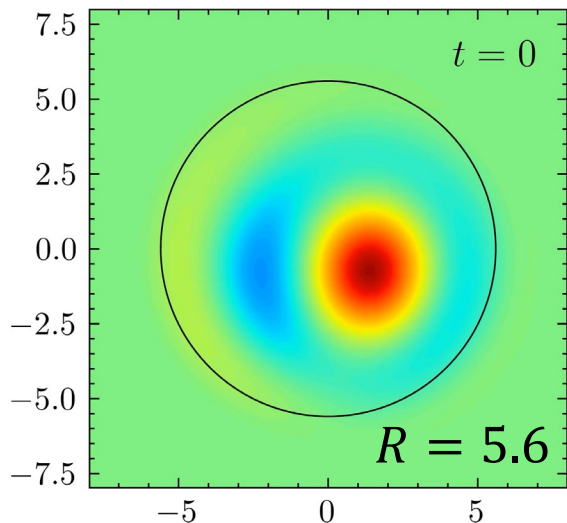
Experiment



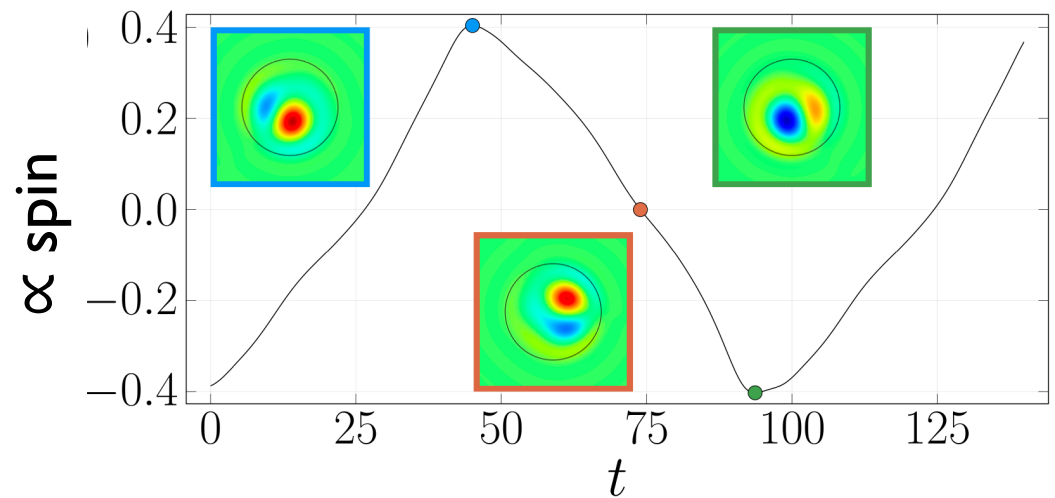
Seemingly regular reversals in experiment.



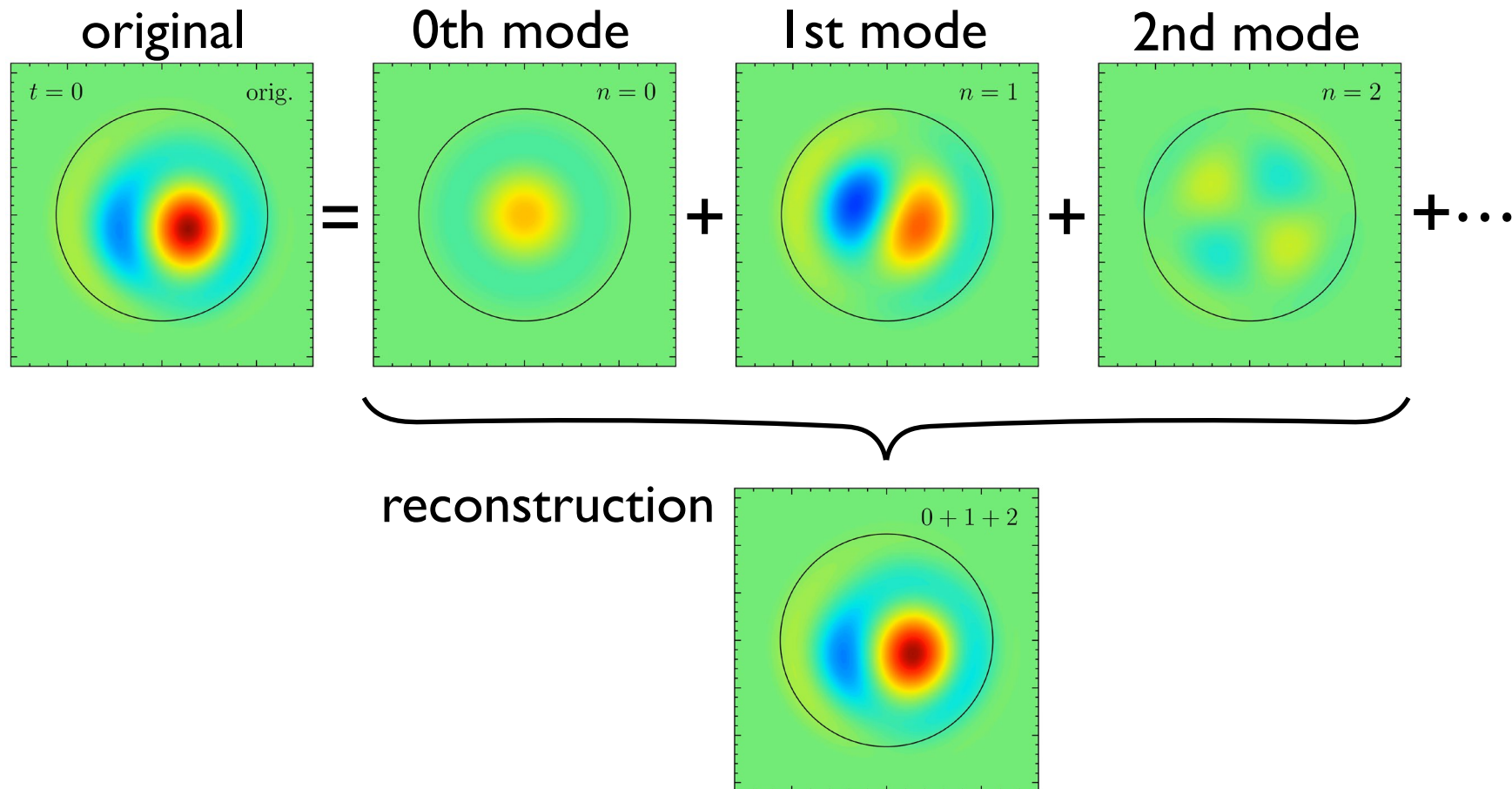
Simulation



Limit cycle was obtained in simulations!



Multipole expansion of reversing vortex pair



- **First 3 modes are relevant** (having comparable amplitudes)
- Can we describe them analytically?



Theory

TTSH vorticity eq.

$$\partial_t \omega + \lambda \vec{v} \cdot \nabla \omega = a\omega - b[\nabla \times (|\vec{v}|^2 \vec{v})]_z - (1 + \nabla^2)^2 \omega$$

↓ linearized

$$\partial_t \omega = a\omega - (1 + \nabla^2)^2 \omega$$

↓ general solutions

$$\omega = \sum_{n=-\infty}^{\infty} \omega_n e^{\lambda_n t}$$

with $\omega_n = (C_n^+ J_n(k_n^+ r) + C_n^- J_n(k_n^- r)) e^{in\theta} + \text{c.c.}$

↙ Bessel function

$$k_n^\pm = \sqrt{1 \pm \sqrt{a - \lambda_n}}$$

$$v_n^\theta = - \left(\frac{C_n^+}{2k_n^+} \Delta J_n(k_n^+ r) + \frac{C_n^-}{2k_n^-} \Delta J_n(k_n^- r) + n C_n^0 r^{n-1} \right) e^{in\theta} + \text{c.c.}$$

$$(\Delta J_n(x) \equiv J_{n-1}(x) - J_{n+1}(x))$$

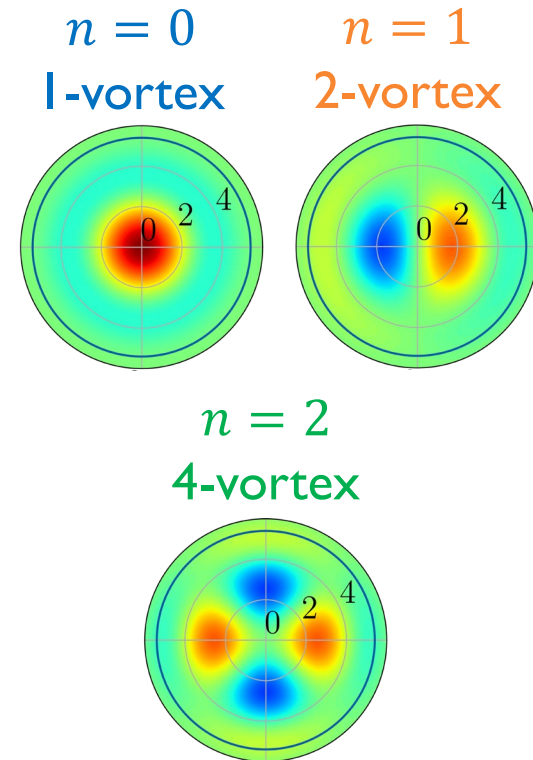
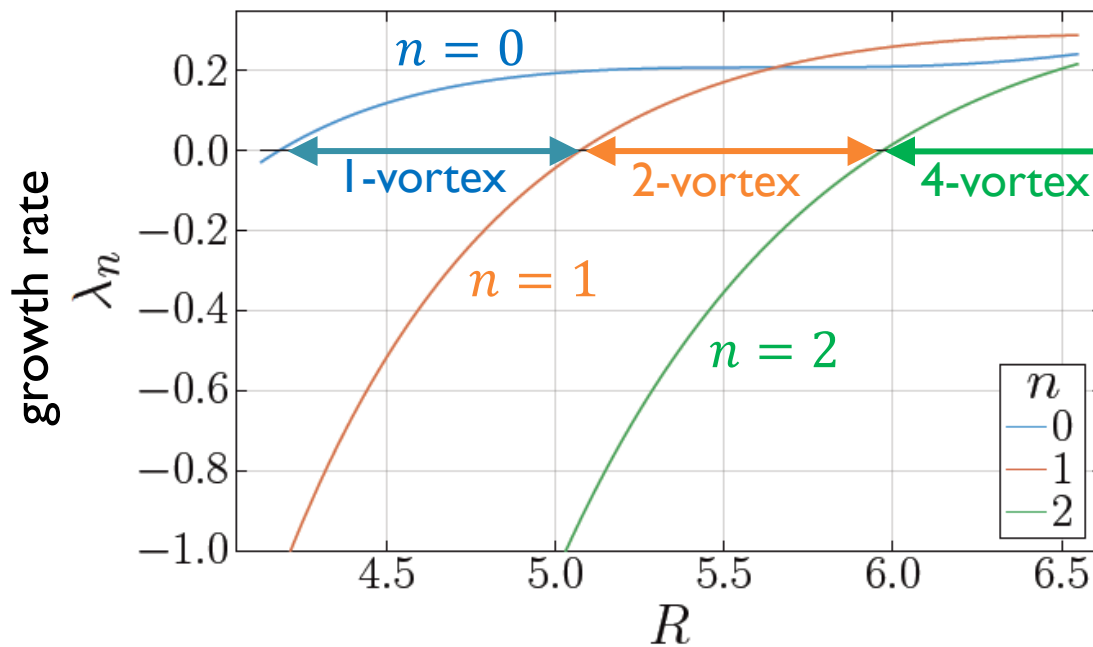
$$v_n^r = \frac{in}{r} \left(\frac{C_n^+}{(k_n^+)^2} J_n(k_n^+ r) + \frac{C_n^-}{(k_n^-)^2} J_n(k_n^- r) + C_n^0 r^n \right) e^{in\theta} + \text{c.c.}$$

BC: $\vec{v}_n = 0$ & $\omega_n = 0$ at $r = R \rightarrow \lambda_n, k_n^\pm, C_n^\pm$ determined

Linear growth rate λ_n vs radius R

Transitions qualitatively accounted for

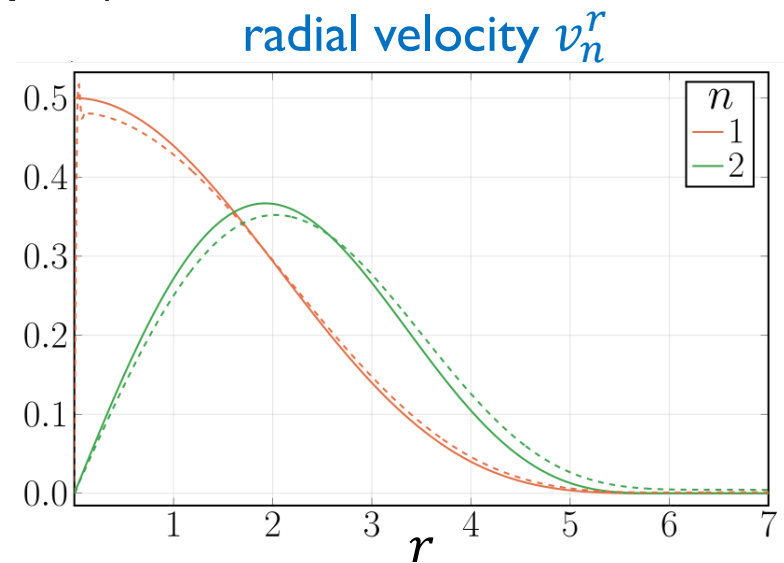
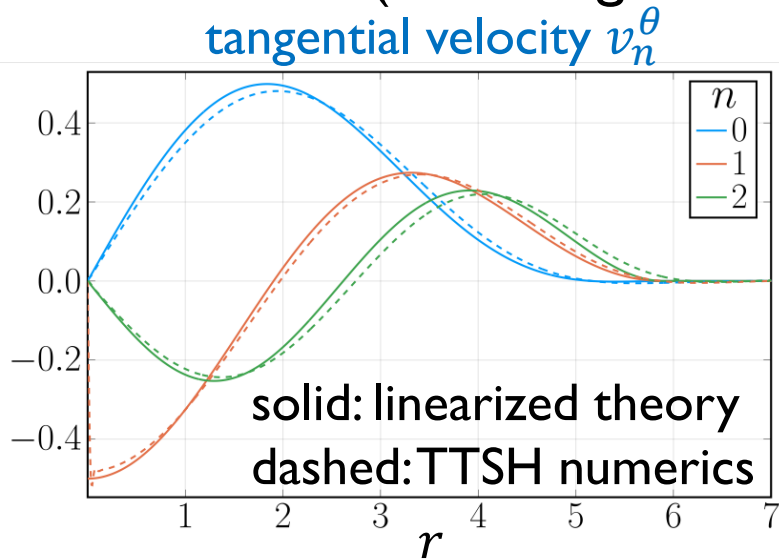
by linear instability of $n = 0, 1, 2$ modes



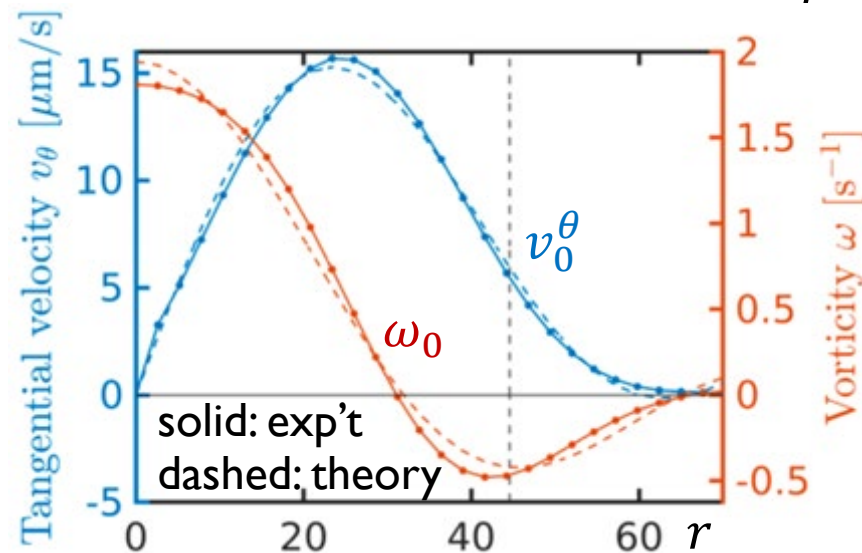
NB) Actual transition points are altered by nonlinear effects (described later)

Spatial structure of vortices

Simulations (reversing vortex pair)



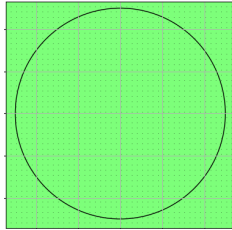
Experiment
(single vortex)



remarkable
agreement!

Dynamics

Mode decomposition



$$\omega(r, \theta, t) \simeq \underbrace{C(t)\omega_0(r)}_{\text{1-vortex}} + \underbrace{[A_1(t)e^{i\theta}\omega_1(r) + A_2(t)e^{2i\theta}\omega_2(r)]}_{\text{2-vortex}} + \text{c. c.}]$$

Bessel (eigenfunctions of linearized equation)

↓ plugging it into TTSH

Time evolution equations (ODEs) for mode coefficients

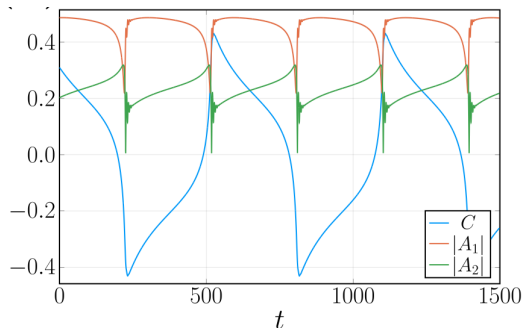
$$\partial_t C = \lambda_0 C - c_1 C^3 - c_2 C |A_1|^2 - c_3 C |A_2|^2 - 2c_4 \text{Re} A_2 A_1^{2*}$$

$$\partial_t A_1 = \lambda_1 A_1 - b_1 A_1 |A_1|^2 - b_2 A_1 C^2 - b_3 A_1 |A_2|^2 - b_4 C A_2 A_1^* + \delta_1 A_1 C + \gamma_1 A_2 A_1^*$$

$$\partial_t A_2 = \lambda_2 A_2 - a_1 A_2 |A_2|^2 - a_2 A_2 C^2 - a_3 A_2 |A_1|^2 - a_4 C A_1^2 + \delta_2 A_2 C + \gamma_2 A_1^2$$

↓ solving numerically

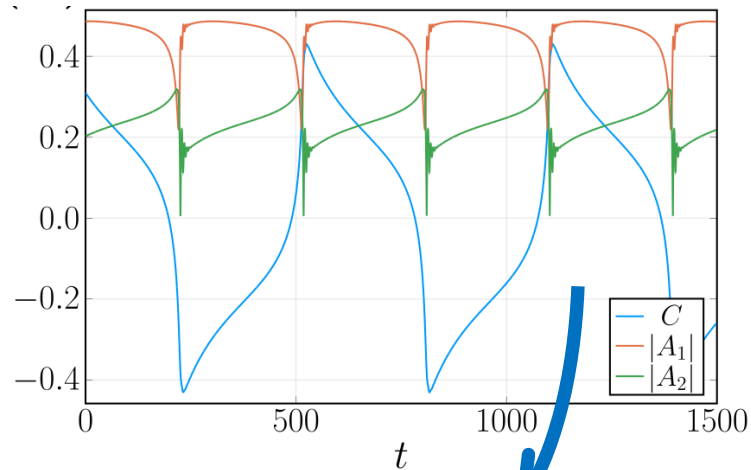
Time evolution of the coefficients in 2-vortex state



- Limit cycle solution!

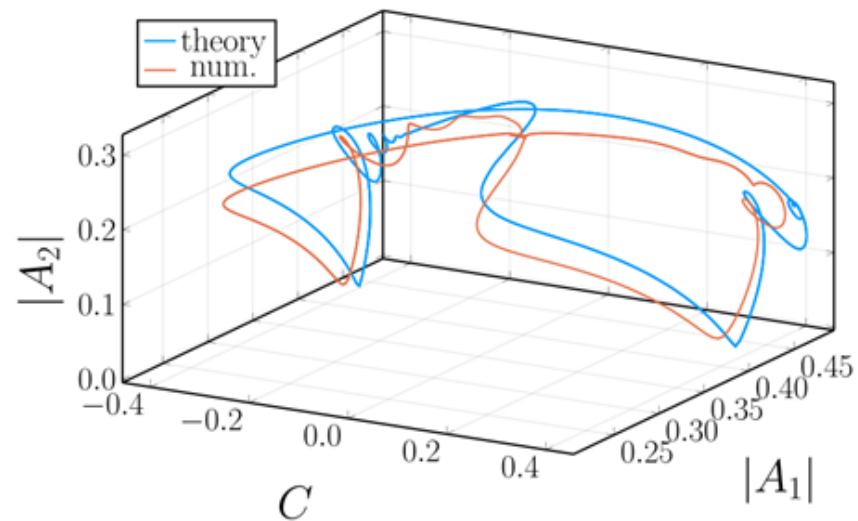
- $\begin{cases} C > 0 & \text{CCW} \\ C < 0 & \text{CW} \end{cases} \therefore \text{Periodic reversal of vortex pair!}$

Reversing vortex pair

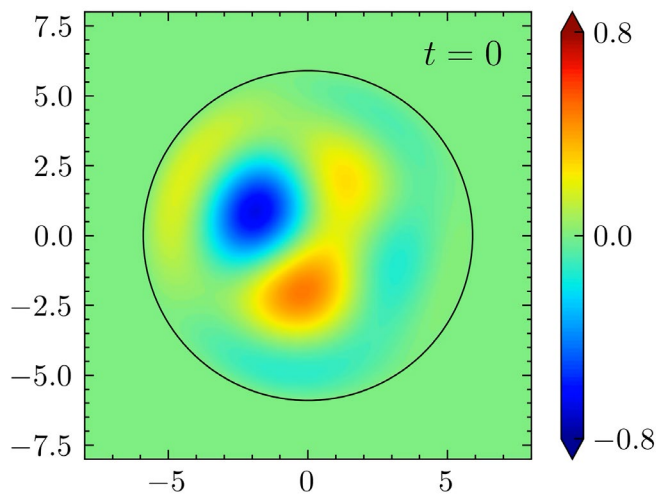


theory

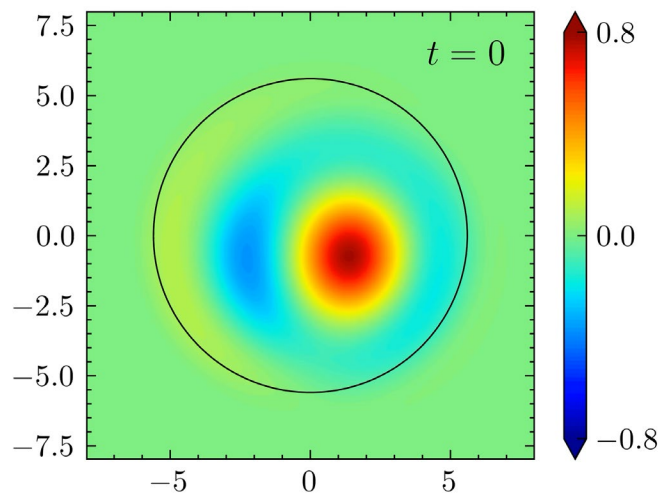
trajectories in phase space



numerics

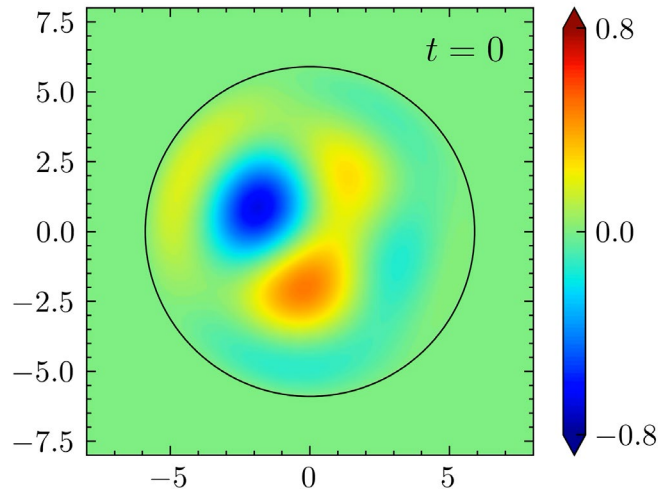


\approx

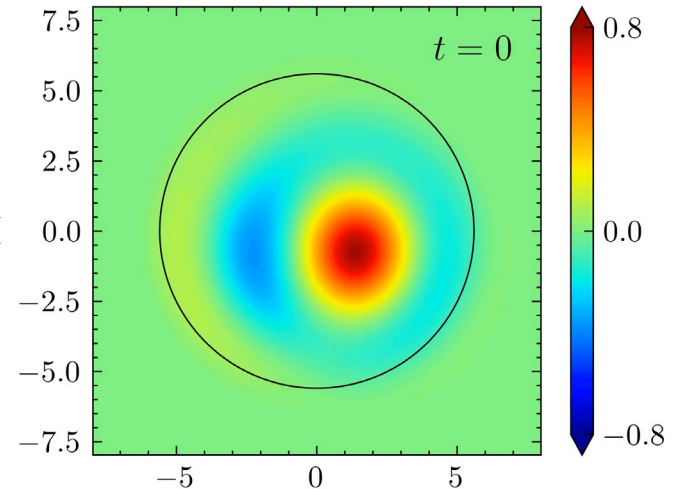


Reversing vortex pair to 4-vortex pulsations

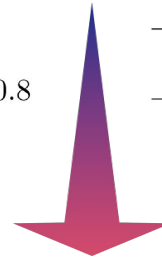
theory



numerics

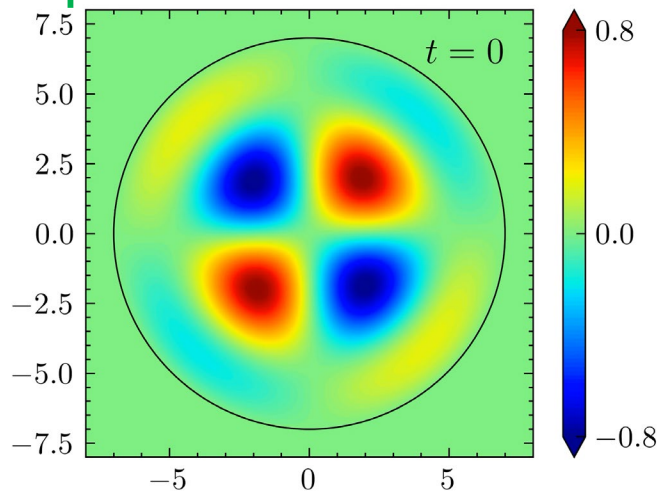


\approx

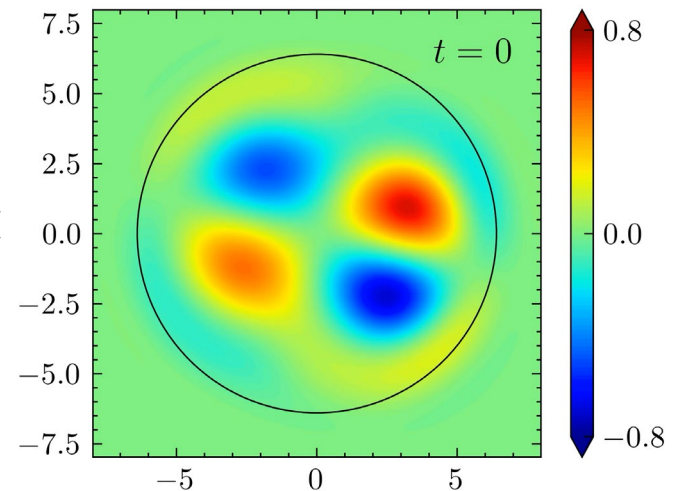


increasing radius

4-vortex pulsations
also reproduced!

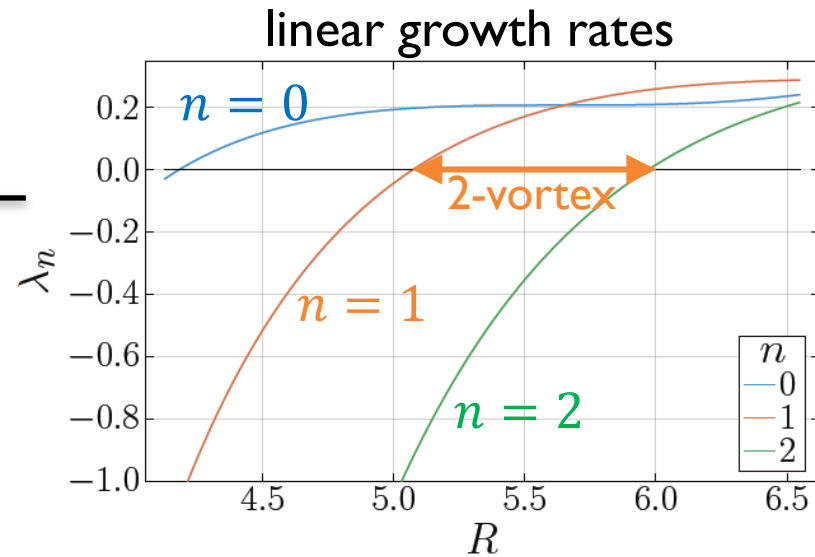


\approx



Coupling of the 3 modes

but we showed $n = 2$ mode is essential for vortex reversing.
What activates this mode?



Time evolution equations (ODEs) for mode coefficients

$$\partial_t C = \lambda_0 C - c_1 C^3 - c_2 C |A_1|^2 - c_3 C |A_2|^2 - 2c_4 \text{Re} A_2 A_1^{2*}$$

$$\partial_t A_1 = \lambda_1 A_1 - b_1 A_1 |A_1|^2 - b_2 A_1 C^2 - b_3 A_1 |A_2|^2 - b_4 C A_2 A_1^* + \delta_1 A_1 C + \gamma_1 A_2 A_1^*$$

$$\partial_t A_2 = \lambda_2 A_2 - a_1 A_2 |A_2|^2 - a_2 A_2 C^2 - a_3 A_2 |A_1|^2 - a_4 C A_1^2 + \delta_2 A_2 C + \gamma_2 A_1^2$$

damping

The rest can be activating!

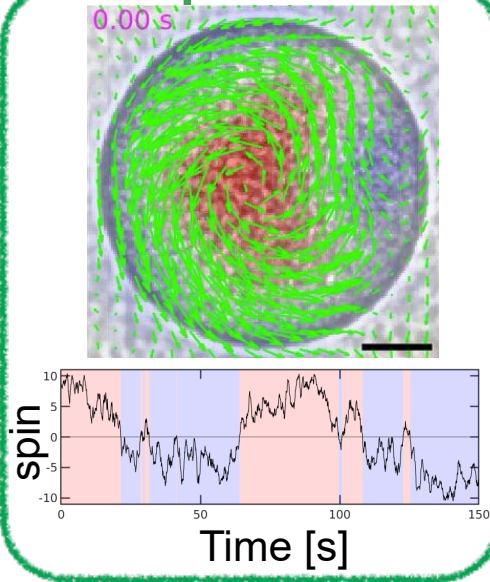
- We obtain $\delta_2, \gamma_2 \propto \lambda$ of $\lambda(\vec{v} \cdot \vec{\nabla}) \vec{v}$

→ active stress is crucial for the reversing vortex state!

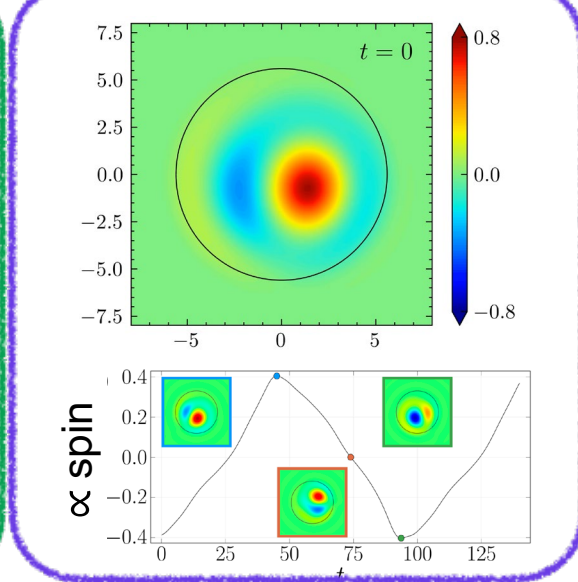
- In our theory, limit cycle was obtained only for $\lambda \gtrsim 3.75$.
- Experiment: $\lambda \approx 4.2$. Consistent!

Summary: Reversal is a precursor of active turbulence

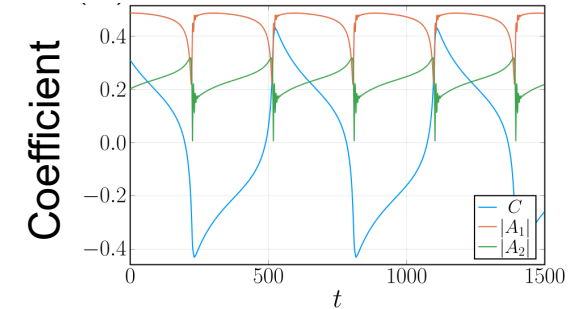
bacteria
experiment



TTSH
numerics



analytical theory

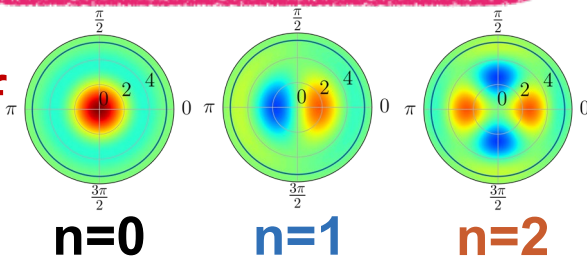


ODEs for the first 3
azimuthal modes obtained
& account for the instability!

Underlying mechanism = orchestration of

- Active stress brings this orchestration.
- Bifurcation is generic.

→ Expected to arise in other active turbulence systems too!



Ref: **arXiv: 2407.05269** (main) & 2304.03306 (numerical detail)