

Hypernuclei Studies: A Key to Resolving the Hyperon Puzzle in NSs

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Xiamen U.

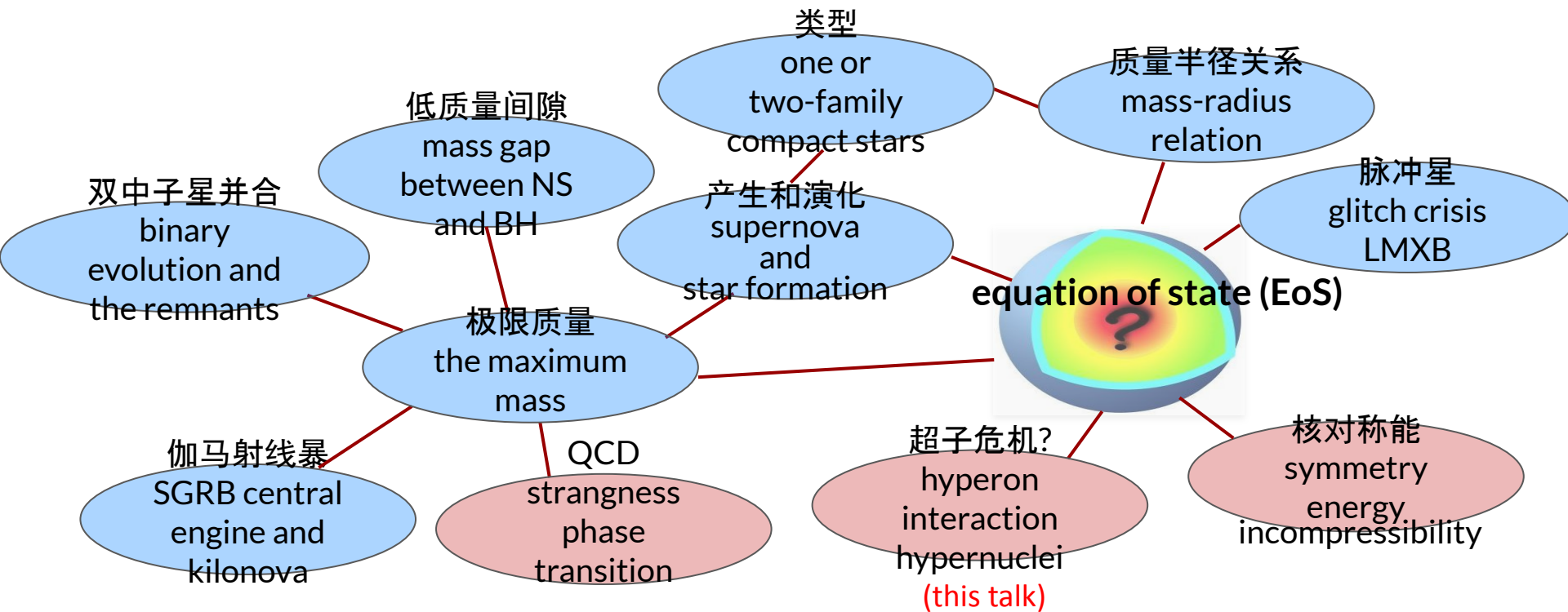
Based on arXiv: 2205.10631 ApJ
and ongoing works with
Shiyuan Ding, Baoyuan Sun (LZU, Lanzhou)
Jinniu Hu, Hong Shen (NKU, Tianjin)

Many thanks for
Invitation!







Nucleosynthesis and Evolution of Neutron Stars
27–30 Jan 2025 Kyoto University, Japan

Why is understanding the EoS important for **nuclear/astrophysicists**?





Astrophysical Implications on Hyperon Couplings and Hyperon Star Properties with Relativistic Equations of States

Xiangdong Sun¹ , Zhiqiang Miao¹ , Baoyuan Sun² , and Ang Li¹ 

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²Frontiers Science Center for Rare Isotopes, Lanzhou University, Lanzhou 730000, People's Republic of China

From Referee “The present article **addresses a long-standing issue** in neutron star physics, namely the hyperon puzzle. The authors **incorporate new information** from hypernuclei calculations and **treat the hyperon couplings in a more general way** than what exists in the present literature. This is an **interesting work that can have important future implications.**”

an RMF with density dependent couplings. The authors of Sun et al. (2023) have recently developed a Bayesian inference approach, in the framework of several nuclear RMF, to determine how GW and NICER measurements constrain the $\Lambda - \sigma$ and $\Lambda - \omega$ couplings, while fixing the Σ and Ξ couplings to reasonable values. A major advantage of this methodology is the possibility, once the inference is completed, to discuss the possible composition of matter or the nuclear properties. In the present study, we will base our approach

A major advantage of this methodology is the possibility, once the inference is completed, to discuss the possible composition of matter or the nuclear properties.

Huang, Raaijmakers, Watts,
Tolos, & Providência, 2303.17518
MNRAS

Outline

- **Basic** for neutron star structure and the EoS
- **Recent work** on connecting #consistently NS observations and (hyper)nuclear experiments
- **Summary and Exciting future**

What is the nature of the particles that make up neutron stars?



Lev Landau



James Chadwick



Walter Baade



Fritz Zwicky



J. Robert Oppenheimer



Antony Hewish



Dame Jocelyn Bell Burnell

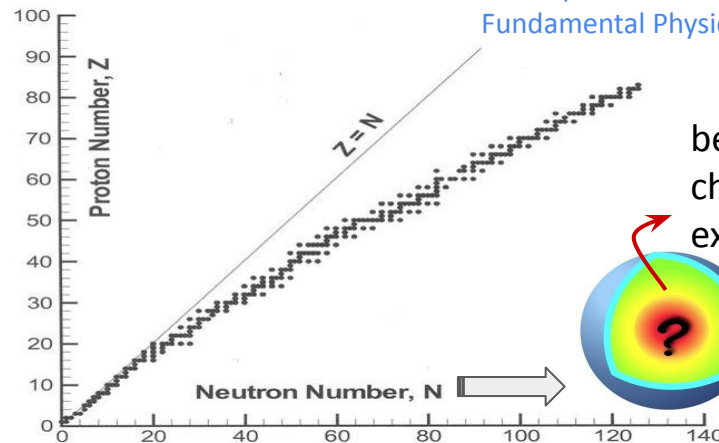
1932-

“the density of matter becomes so great that atomic nuclei come in close contact, forming one gigantic nucleus.”

$A \sim 10^{57}$; For $R=10$ km, $M=1.4 M_{\odot}$,
average density $\sim(2-3)$ nuclear density ρ_0

\sim several hundreds MeV (nonperturbative)

Extreme conditions make it impossible to attain by theo./exp. methods only.



beta equilibrium;
charge neutrality;
extremely neutron-rich

Nobel prize 1974
Fundamental Physics Breakthrough Prize 2018

1967

EoS uncertainty from QCD phase uncertainty and model uncertainty

- **Hyperon puzzle**; $\Delta(1232)$ /hyperon/Kaon/quark complication
- 1) Unified crust-core; 2) High-density extrapolation

Hyperon puzzle: Heavy pulsars larger than $2M_{\odot}$ is a pain

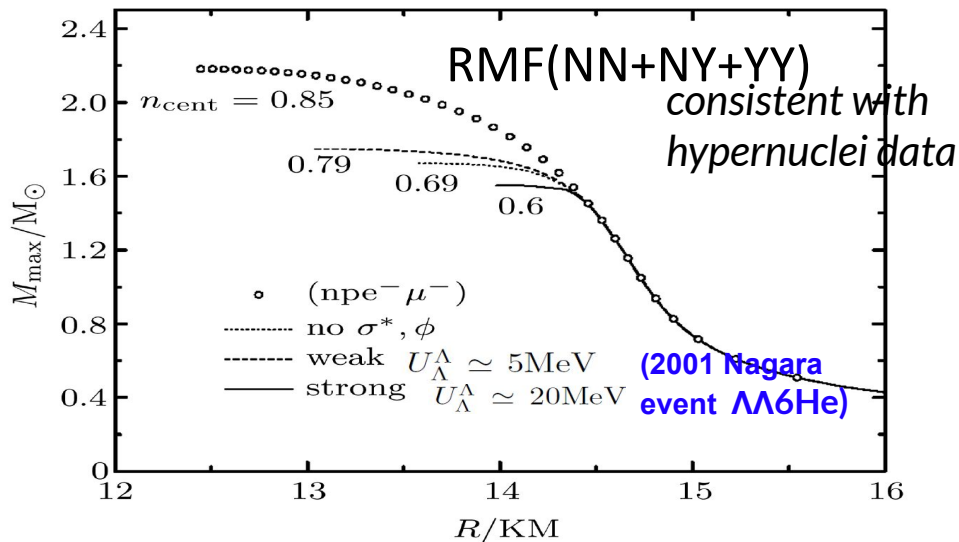
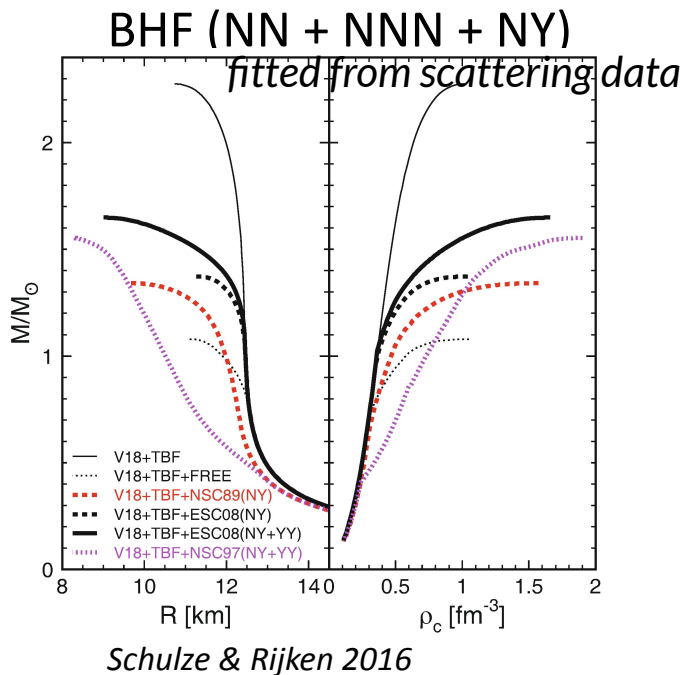


Fig.3. The mass–radius relation calculated for the strong and weak hyperon–hyperon interaction models, in comparison with the one without strange mesons (σ^* , ϕ), and also the results without hyperons. See text for details. *Li et al. 2007*

Is it possible to combine NS multi-messenger (2017-) observations with hypernuclei experiments to understand better the hypernuclear force? **How? What new?**

Neutron star group @XMU

[arXiv:2412.09219](#) Submitted

[arXiv:2408.15022](#) MNRAS

[arXiv:2402.02799](#) PRD

[arXiv:2312.17102](#) ApJ

[arXiv:2312.12185](#) ApJ

[arXiv:2312.04305](#) MNRAS

[arXiv:2305.16058](#) PRC

[arXiv:2305.08401](#) ApJ

[arXiv:2304.12050](#) PRD

[arXiv:2211.04978](#) PRD

[arXiv:2211.02007](#) ApJ

[arXiv:2205.10631](#) **ApJ**

[arXiv:2204.05560](#) ApJ

[arXiv:2203.04798](#) PRD

[arXiv:2201.12053](#) PRC

[arXiv:2108.00560](#) ApJ

[arXiv:2107.13997](#) ApJL

[arXiv:2107.07979](#) MNRAS

[arXiv:2103.15119](#) ApJ

[arXiv:2011.11934](#) ApJ

[arXiv:2009.12571](#) MNRAS

[arXiv:2007.05116](#) JHEAp (review)

[arXiv:2006.00839](#) ApJ

[arXiv:2005.12875](#) ApJS

[arXiv:2005.02677](#) PRD

[arXiv:2001.03859](#) PRC



Wenli Yuan 苑文莉



Quark matter;
QCD phase diagram

**Graduated in 2023;
postdoc in PKU**

Zhenyu Zhu 朱镇宇



Many-body theory;
Merger simulation
Numerical relativity

**Graduated in 2021;
postdoc in CCRG-RIT**

Peng Liu 刘鹏



Glitch;
Pulsar observation

Zhiqiang Miao 缪志强



Hybrid star;
Quark matter;
Bayesian analysis

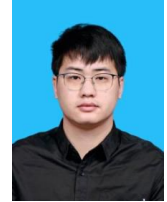
GR
**Graduated in 2023;
postdoc in TDLee inst.**

Xiangdong Sun 孙向东



Nuclear matter;
Hyperon matter;
Many-body theory

Zhonghao Tu 涂中豪



Superfluidity;
Neutron star cooling;
Nuclear pinning force

Shuochong Han 韩烁冲



Many-body theory;
Nuclear transport

with 4 undergraduate students

Outline

- **Basic** for neutron star structure and the EOS
- **Recent work** towards the determination of the $\#$ (hyper)nuclear force and NS properties from multimessenger astronomy
(Biased selected results; Highlighting work done by our group)
- **Summary and Exciting future**

Black box = EoS microphysics

Procedure used until now

$$\Gamma = (\rho + P)(dP/d\rho)/P$$

- ❑ Piecewise polytrope EoS: Logarithm of the adiabatic index (Γ) of the EoSs are treated as polynomial: **Little to no microphysics**;

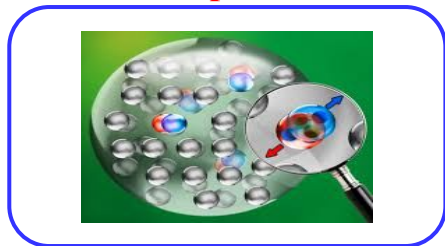
Better procedure

- ❑ More physics at the EoS modelling stage;
- ❑ Support **quantitative** studies of the EoS at different density regimes;
- ❑ With prior **explicitly** considering phase transition;
- ❑ Can discuss the **composition** of matter and the underlying strong **interaction**;
- ❑ Allow for straightforward **extensions** to higher-dimensional models, accommodating the inclusion of additional particles within NSs, even dark matter;
- ❑ Facilitate a **connection** with the ongoing research efforts in the field of relativistic HIC.



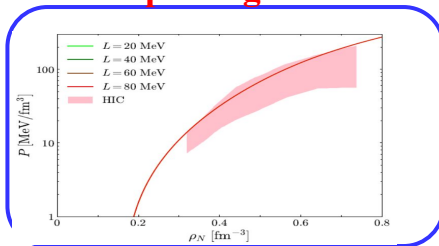
EoS: The roadmap

1. Model for interaction between particles



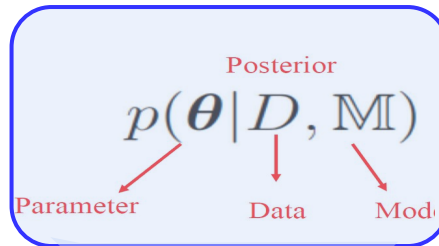
$$\hat{\mathcal{H}}\Psi = E\Psi$$

2. EoS prior pasting test

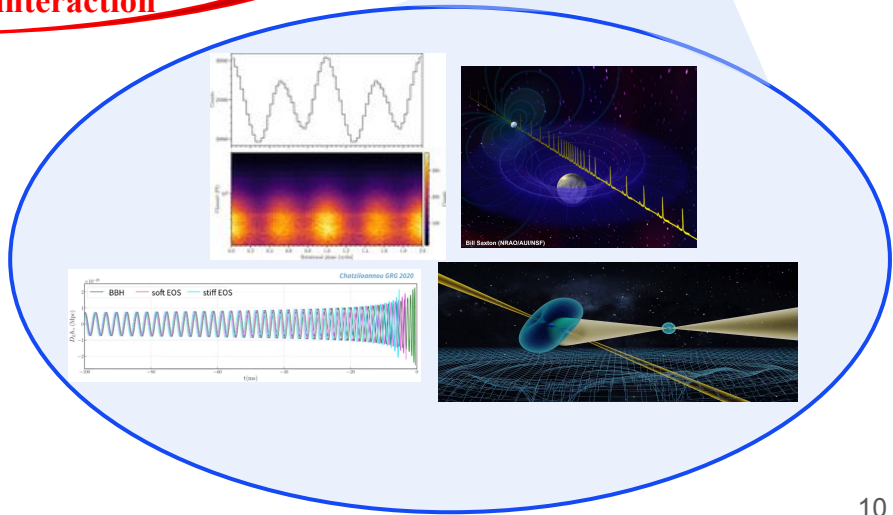


$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

3. EoS inference from NS obs. (GW, photons, neutrinos)



Physics on the EoS, the composition and the underlying strong interaction



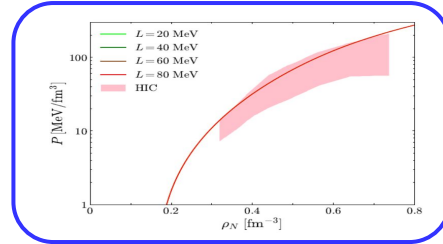
Prepare EoS prior starting for e.g., RMF parameter set

1. Model for interaction between particles



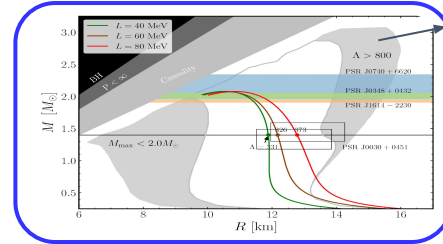
$$\hat{H}\Psi = E\Psi$$

2. the EoS



$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

3. NS observations (M-R relation, Lambda)



excluded by $\lambda(1.4) \leq 800$ of GW170817

$$\mathcal{L} = \bar{\psi} (i\gamma_\mu \partial^\mu - M_N^* - g_{\omega N} \omega \gamma^0 - g_{\rho N} \rho \tau_3 \gamma^0) \psi - \frac{1}{2} (\nabla \sigma)^2 - \frac{1}{2} m_\sigma^2 \sigma^2 - \frac{1}{3} g_2 \sigma^3 - \frac{1}{4} g_3 \sigma^4$$

$$+ \frac{1}{2} (\nabla \rho)^2 + \frac{1}{2} m_\rho^2 \rho^2 + \frac{1}{2} (\nabla \omega)^2 + \frac{1}{2} m_\omega^2 \omega^2 + \frac{1}{2} g_{\rho N}^2 \rho^2 \Lambda_v g_{\omega N}^2 \omega^2,$$

(P, I, J symmetry)

Eqs. of motion of baryons and mesons can be generated by the Euler-Lagrangian eq. from the Lagrangian:

$$\left[i\gamma^\mu \partial_\mu - M_B^* - \gamma^0 \left(g_{\omega B} \omega + g_{\phi B} \phi + \frac{g_{\rho B}}{2} \rho \tau_3 \right) \right] \psi_B = 0,$$

$$m_\sigma^2 \sigma + g_2 \sigma^2 + g_3 \sigma^3 = \sum_B g_{\sigma B} \rho_B^s,$$

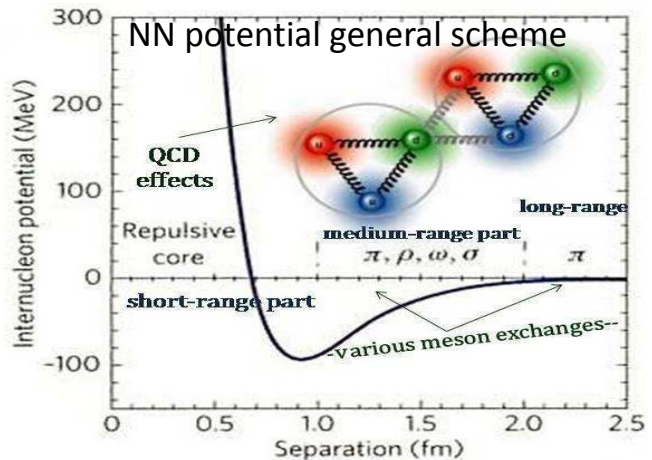
$$m_\omega^2 \omega + c_3 \omega^3 + 2\Lambda_v (g_{\omega N}^2 \omega) (g_{\rho N}^2 \rho^2) = \sum_B g_{\omega B} \rho_B^v,$$

$$m_\rho^2 \rho + 2\Lambda_v (g_{\omega N}^2 \omega^2) (g_{\rho N}^2 \rho) = \sum_B \frac{g_{\rho B}}{2} \rho_B^v,$$

Then, \mathbf{p} & $\boldsymbol{\varepsilon}$ (with arbitrary isospin asymmetry) from nuclear part can be generated by the energy-momentum tensor.

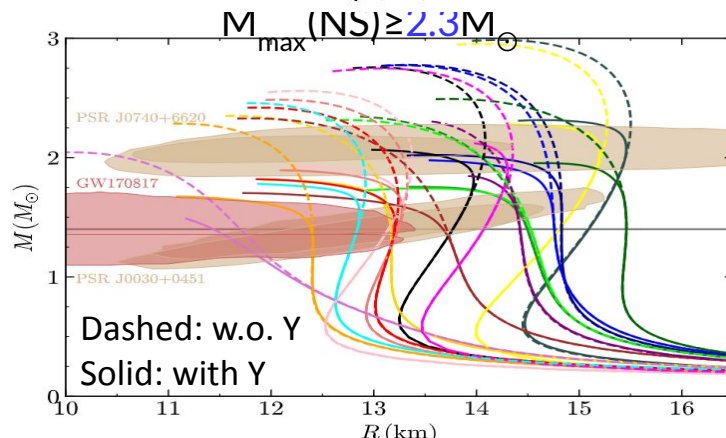
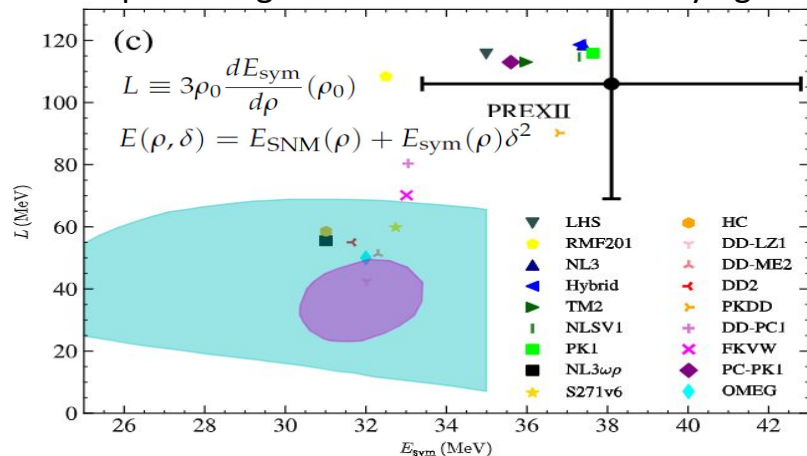
Static approximation considered in the Lagrangian on mesons so that their time components are neglected;
 Spatial part of $\boldsymbol{\omega}$ meson disappears for the **time reversal symmetry**;
Infinite nuclear matter has **translational invariance**, removing the partial part of the coordinate space.

Employed 18 stiff relativistic EoSs with 7 types of variation of the RMF effective interactions



- 1) Original linear Walecka model,
- 2) Nonlinear Walecka model with σ **self-interacting** mesons
- 3) Nonlinear Walecka model with σ, ω **self-interacting** mesons,
- 4) Nonlinear Walecka model with σ, ω **self-interacting** mesons and possible mesonic **cross** terms,
- 5) Models in which the parameters that couple the baryons with the mesons are **density-dependent**,
- 6) Models with the inclusion of δ mesons, i.e., $a_0(980)$,
- 7) **Point-coupling (PC)** models without exchanging mesons.

Reproducing nuclear saturation with varying stiffness



To combine (binary) NS observations with hypernuclei experiments

$$\rho(\theta, D) \propto \pi(\theta) \times L(D|\theta, \mathbb{M})$$

Assuming the sources are hyperon stars:

Formerly:
GW + X-ray



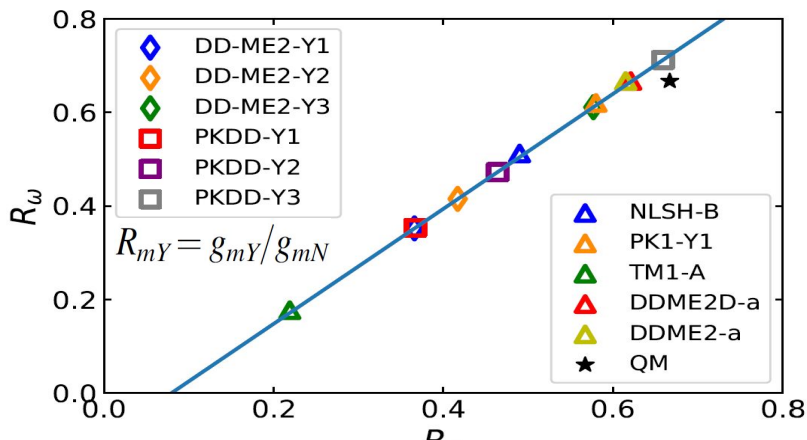
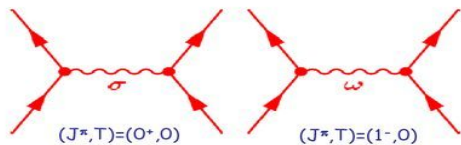
Presently: GW + X-ray + NUCL

$$\mathcal{L}_{\text{NUCL}}(d_{\text{NUCL}}|\theta_{\text{EOS}}) \propto \exp \left[-\frac{1}{2} \frac{(R_{\sigma\Lambda} - \bar{R}_{\sigma\Lambda})^2}{\sigma_{R_{\sigma\Lambda}}^2} \right]$$

$$\theta_{\text{EOS}} = \{R_{\sigma\Lambda}, R_{\omega\Lambda}\}$$

Strong linear $R_{\sigma\Lambda}$ - $R_{\omega\Lambda}$ relations from fitting (with some statistical error) calculated Λ separation energies of **eleven** $A \geq 12$ single Λ hypernuclei

(Rong, Tu & Zhou, 2021): $R_{\omega} = 1.228R_{\sigma} - 0.097$

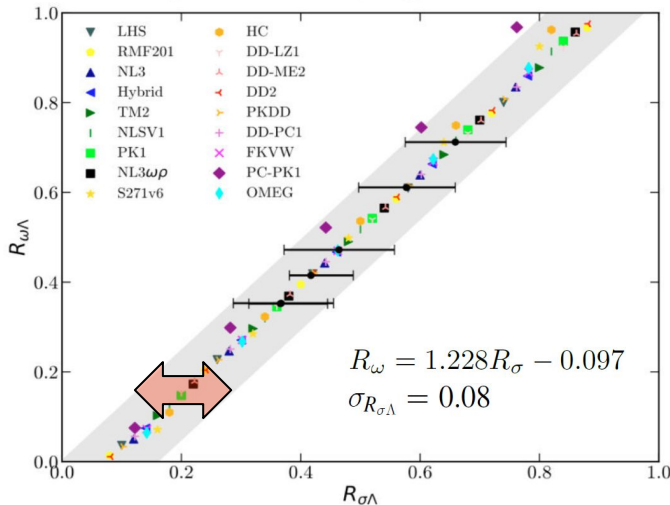


Hypernuclei	Exp.
$^{12}_{\Lambda}\text{C}$	11.36 ± 0.20
$^{13}_{\Lambda}\text{C}$	12.0 ± 0.2
$^{16}_{\Lambda}\text{O}$	13.0 ± 0.2
$^{28}_{\Lambda}\text{Si}$	17.2 ± 0.2
$^{32}_{\Lambda}\text{S}$	17.5 ± 0.5
$^{40}_{\Lambda}\text{Ca}$	18.7 ± 1.1
$^{51}_{\Lambda}\text{V}$	21.5 ± 0.6
$^{52}_{\Lambda}\text{V}$	21.8 ± 0.3
$^{89}_{\Lambda}\text{Y}$	23.6 ± 0.5
$^{139}_{\Lambda}\text{La}$	25.1 ± 1.2
$^{208}_{\Lambda}\text{Pb}$	26.9 ± 0.8



Linear $R_{\sigma\Lambda}$ - $R_{\omega\Lambda}$ relations consistent with both finite-range meson-exchange and zero-range PC models

$$U_N = -70 \text{ MeV}, U_\Lambda = -30 \text{ MeV}$$



- Meson-exchange

$$\mathcal{L} = \mathcal{L}_{\text{free}}^B + \mathcal{L}_{\text{int}}^B + \mathcal{L}_m + \mathcal{L}_{\text{NL}}$$

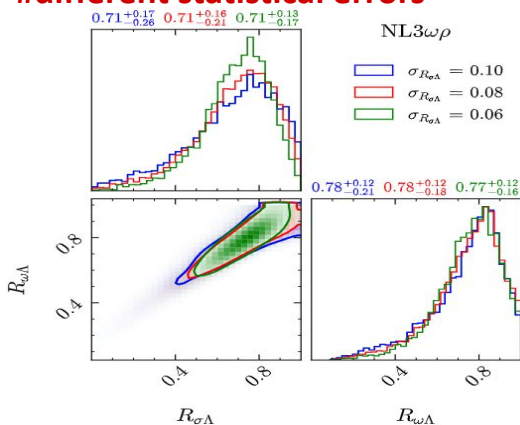
$$\begin{aligned} \mathcal{L}_{\text{free}}^B &= \sum_B \bar{\psi}_B [i\gamma_\mu \partial^\mu - m_B] \psi_B \\ \mathcal{L}_{\text{int}}^B &= \sum_B \bar{\psi}_B [-g_{\sigma B} \sigma - g_{\omega B} \gamma_\mu \omega^\mu - g_{\rho B} \gamma_\mu \vec{\rho}^\mu \cdot \vec{\tau}] \psi_B \\ \mathcal{L}_m &= -\frac{1}{2} m_\sigma^2 \sigma^2 + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu + \frac{1}{2} m_\rho^2 \vec{\rho}_\mu \cdot \vec{\rho}^\mu \\ &\quad + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - \frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} - \frac{1}{4} \vec{R}_{\mu\nu} \cdot \vec{R}^{\mu\nu} \\ \mathcal{L}_{\text{NL}} &= -\frac{1}{3} g_2 \sigma^3 - \frac{1}{4} g_3 \sigma^4 + \frac{1}{4} c_3 (\omega_\mu \omega^\mu)^2 \\ &\quad + \Lambda_\nu (g_{\omega B}^2 \omega_\mu \omega^\mu) (g_{\rho B}^2 \rho_\mu \rho^\mu); \end{aligned}$$

- Point-coupling

$$\mathcal{L}_{\text{PC}} = \mathcal{L}_{\text{free}}^B + \mathcal{L}_{4f}^B + \mathcal{L}_{\text{hot}}^B$$

$$\begin{aligned} \mathcal{L}_{\text{free}}^B &= \sum_B \bar{\psi}_B [i\gamma_\mu \partial^\mu - m_B] \psi_B \\ \mathcal{L}_{4f}^B &= -\frac{1}{2} \sum_B \alpha_S^{NB}(\rho) (\bar{\psi}_N \psi_N) (\bar{\psi}_B \psi_B) \\ &\quad - \frac{1}{2} \sum_B \alpha_V^{NB}(\rho) (\bar{\psi}_N \gamma_\mu \psi_N) (\bar{\psi}_B \gamma_\mu \psi_B) \\ &\quad - \frac{1}{2} \sum_B \alpha_{TS}^{NB}(\rho) (\bar{\psi}_N \vec{\tau} \psi_N) (\bar{\psi}_B \vec{\tau} \psi_B) \\ &\quad - \frac{1}{2} \sum_B \alpha_{TV}^{NB}(\rho) (\bar{\psi}_N \vec{\tau} \gamma_\mu \psi_N) (\bar{\psi}_B \vec{\tau} \gamma_\mu \psi_B) \\ \mathcal{L}_{\text{hot}}^B &= -\frac{1}{3} \beta_S^{NN} (\bar{\psi}_N \psi_N)^3 - \frac{1}{4} \gamma_S^{NN} (\bar{\psi}_N \psi_N)^4 \\ &\quad - \frac{1}{4} \gamma_V^{NN} [(\bar{\psi}_N \gamma_\mu \psi_N)] [(\bar{\psi}_N \gamma^\mu \psi_N)]^2, \end{aligned}$$

#different statistical errors



For finite-range interactions:

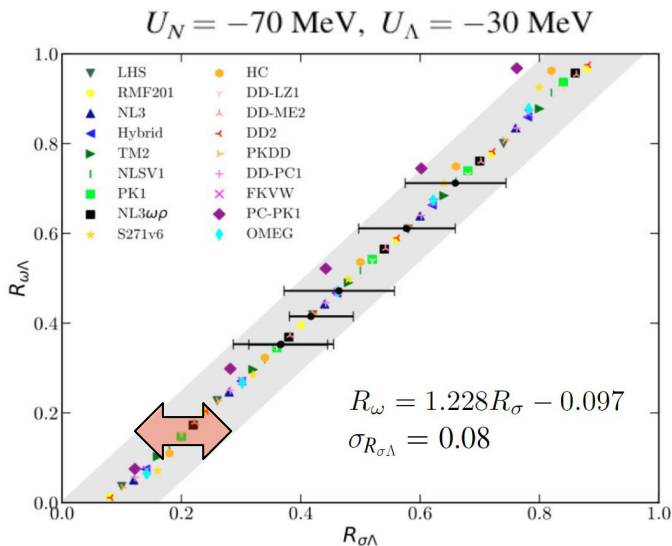
$$R_{\omega\Lambda} \approx \frac{-g_{\sigma N} \sigma}{U_N - g_{\sigma N} \sigma} R_{\sigma\Lambda} + \frac{U_\Lambda}{U_N - g_{\sigma N} \sigma}$$

For zero-range PC

$$\begin{aligned} R_{\omega\Lambda} &\approx \frac{-\alpha_S^{NN} \rho_S}{U_N - \alpha_S^{NN} \rho_S - \beta_S^{NN} \rho_S^2 - \gamma_S^{NN} \rho_S^3 - \gamma_V^{NN} \rho_V^3} R_{\sigma\Lambda} \\ &\quad + \frac{U_\Lambda}{U_N - \alpha_S^{NN} \rho_S - \beta_S^{NN} \rho_S^2 - \gamma_S^{NN} \rho_S^3 - \gamma_V^{NN} \rho_V^3}. \end{aligned}$$

Hyperon star properties do NOT rely sensitively on the choice of the statistical error.

Adding relativistic Fock diagram: Similar correlation within statistical error

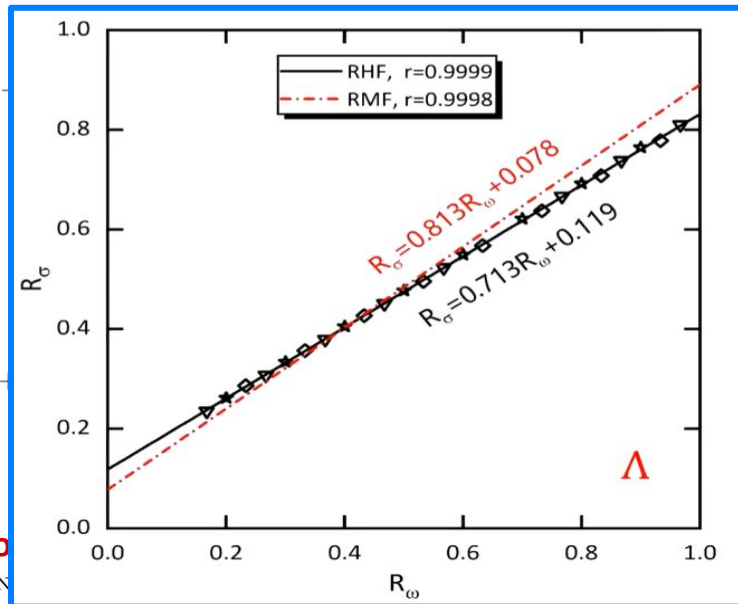


• Meson-exchange

$$\mathcal{L} = \mathcal{L}_{\text{free}}^B + \mathcal{L}_{\text{int}}^B + \mathcal{L}_m$$

• Point-coupling

$$\mathcal{L}_{PC} = \mathcal{L}_{\text{free}}^B + \mathcal{L}_{4f}^B$$



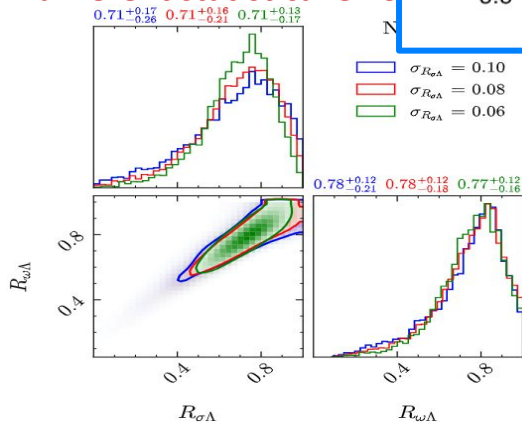
#different statistical error

For finite-range interactions:

$$R_{\omega\Lambda} \approx \frac{-g_{\sigma N\sigma}}{U_N - g_{\sigma N\sigma}} R_{\sigma\Lambda} + \frac{U_\Lambda}{U_N - g_{\sigma N\sigma}}$$

For zero-range PC

$$R_{\omega\Lambda} \approx \frac{-\alpha_S^{NN} \rho_S}{U_N - \alpha_S^{NN} \rho_S - \beta_S^{NN} \rho_S^2 - \gamma_S^{NN} \rho_S^3 - \gamma_V^{NN} \rho_V^3} R_{\sigma\Lambda} + \frac{U_\Lambda}{U_N - \alpha_S^{NN} \rho_S - \beta_S^{NN} \rho_S^2 - \gamma_S^{NN} \rho_S^3 - \gamma_V^{NN} \rho_V^3}$$



$$-\frac{1}{4} \gamma_V^{NN} [(\bar{\psi}_N \gamma_\mu \psi_N)] [(\bar{\psi}_N \gamma^\mu \psi_N)]^2$$

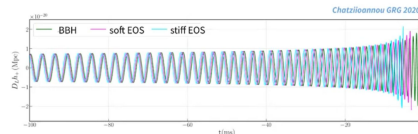
Hyperon star properties do NOT rely sensitively on the choice of the statistical error.

Bayesian inference of hyperon-nucleon interaction strengths from combining (binary) NS observations with hypernuclei experiments

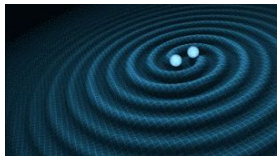
- Finite size effects of the merging star alter late inspiral GW signal:

$$\mathcal{L}_{\text{GW}}(d_{\text{GW}} | \boldsymbol{\theta}_{\text{GW}}, \mathbb{M}) \propto \exp \left[-2 \int_0^\infty \frac{|\tilde{d}(f) - \tilde{h}(f, \boldsymbol{\theta}_{\text{GW}})|^2}{S_n(f)} df \right],$$

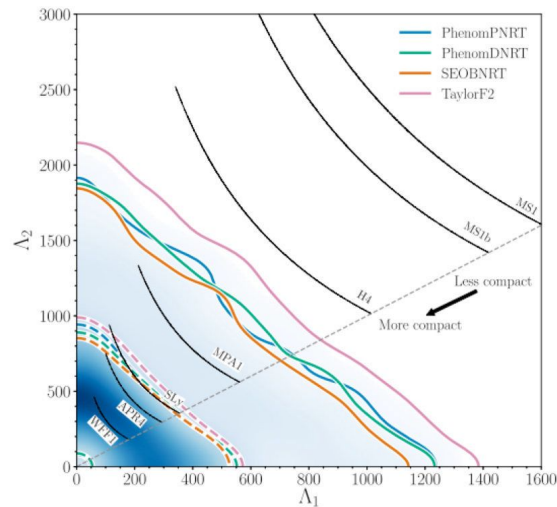
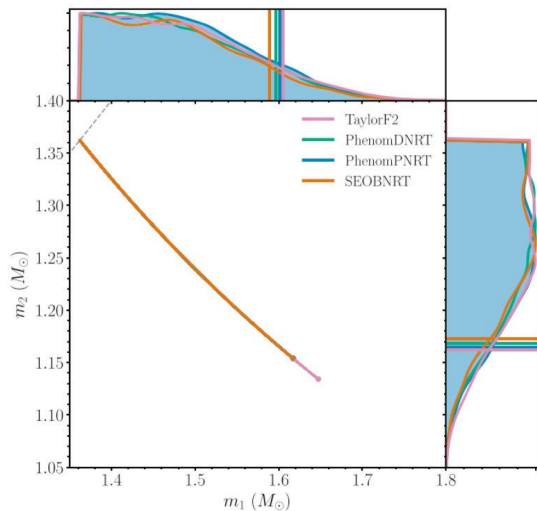
$$\boldsymbol{\theta}_{\text{GW}} = \{M_1, M_2, \Lambda_1, \Lambda_2, \chi_{1z}, \chi_{2z}, \varphi, \Psi, \theta_{\text{jn}}, t_c, d_L, \text{R.A.}, \text{Decl.}\}$$



waveform depending on 17(4) parameters



(1.36, 1.60) M_\odot
(1.16, 1.36) M_\odot



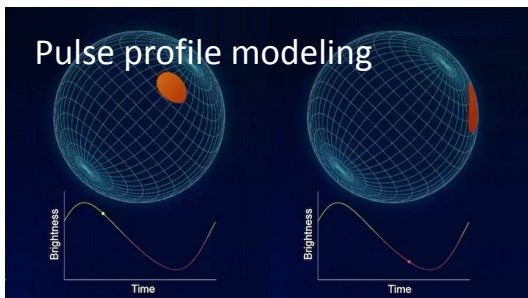
$$Q_{ij} = -\Lambda(\text{EOS}, m) m^5 \varepsilon_{ij}$$

Bayesian inference of hyperon-nucleon interaction strengths from combining (binary) NS observations with hypernuclei experiments

- $\mathcal{L}_{\text{NICER}}(M, R | \theta_{\text{EOS}} \cup \{\varepsilon_c\}, \mathbb{M}) = P_{\text{KDE}}(M(\theta_{\text{EOS}}; \varepsilon_c), R(\theta_{\text{EOS}}; \varepsilon_c)),$

Kernel density

Trace rays from hot spots on NS surface: estimation



PSR J0030+0451

Miller et al, ApJL, 2019

$$M = 1.44^{+0.15}_{-0.14} M_{\odot} \quad R = 13.02^{+1.24}_{-1.06} \text{ km}$$

Riley et al, ApJL, 2019

$$M = 1.34^{+0.15}_{-0.16} M_{\odot} \quad R = 12.71^{+1.14}_{-1.19} \text{ km}$$

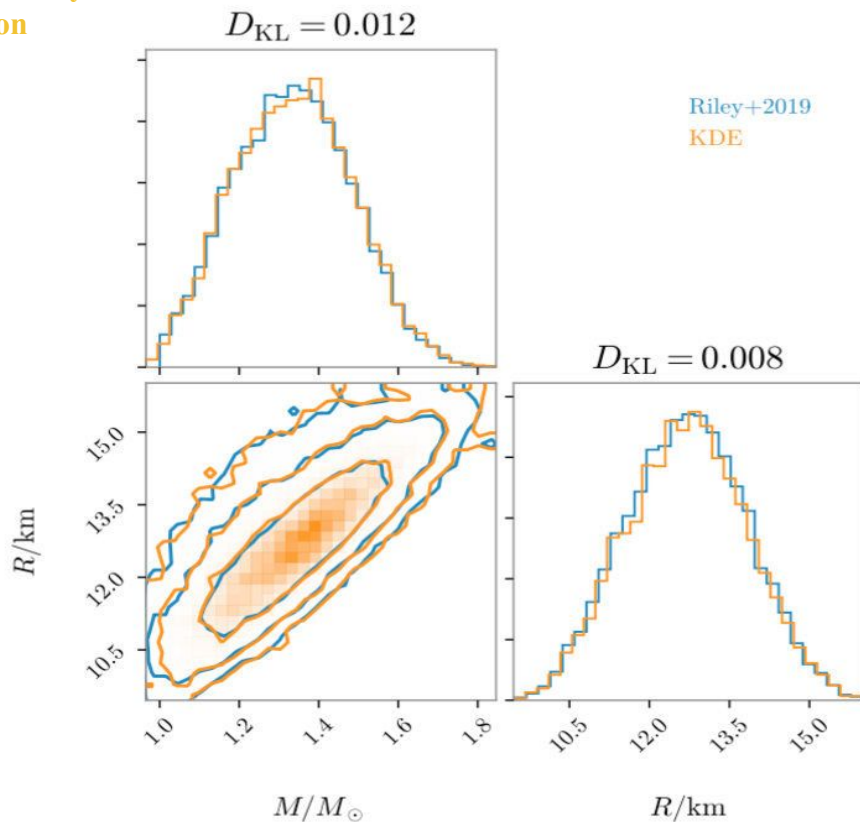
PSR J0740+6620

Miller et al, ApJL, 2021

$$M = 2.062^{+0.090}_{-0.091} M_{\odot} \quad R = 13.71^{+2.61}_{-1.50} \text{ km}$$

Riley et al, ApJL, 2021

$$M = 2.072^{+0.067}_{-0.066} M_{\odot} \quad R = 12.39^{+1.30}_{-0.98} \text{ km}$$



Bayesian inference of #hyperon-nucleon interaction strengths from combining (binary) NS observations with hypernuclei experiments

Likelihood

$$p(\boldsymbol{\theta} | \overset{\text{Posterior}}{\underset{\text{Data}}{D}}, \underset{\text{Model}}{M}) = \frac{\overset{\text{Likelihood}}{\mathcal{L}(D|\boldsymbol{\theta}, M)} \overset{\text{Prior}}{\pi(\boldsymbol{\theta})}}{\int p(D|\boldsymbol{\theta}, M)\pi(\boldsymbol{\theta})d\boldsymbol{\theta}}$$

Parameter Data Model

- $\mathcal{L}_{\text{NUCL}}(d_{\text{NUCL}}|\boldsymbol{\theta}_{\text{EOS}}) \propto \exp \left[-\frac{1}{2} \frac{(R_{\sigma\Lambda} - \bar{R}_{\sigma\Lambda})^2}{\sigma_{R_{\sigma\Lambda}}^2} \right]$
- $\mathcal{L}_{\text{GW}}(d_{\text{GW}}|\boldsymbol{\theta}_{\text{GW}}, M) \propto \exp \left[-2 \int_0^\infty \frac{|\tilde{d}(f) - \tilde{h}(f, \boldsymbol{\theta}_{\text{GW}})|^2}{S_n(f)} df \right],$
- $\mathcal{L}_{\text{NICER}}(M, R|\boldsymbol{\theta}_{\text{EOS}} \cup \{\varepsilon_c\}, M) = P_{\text{KDE}}(M(\boldsymbol{\theta}_{\text{EOS}}; \varepsilon_c), R(\boldsymbol{\theta}_{\text{EOS}}; \varepsilon_c)),$

Software:

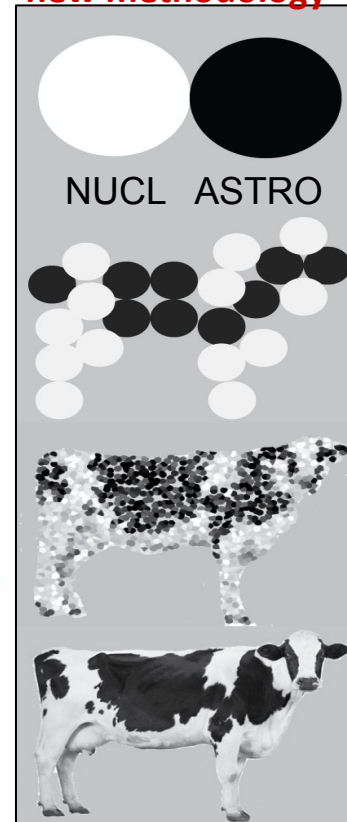
Bilby (Ashton et al. 2019, version 0.5.5, <https://git.ligo.org/lscsoft/bilby/>),

PyMultiNest (Buchner 2016, version 2.6, <https://github.com/JohannesBuchner/PyMultiNest>),

Toast (Hernandez Vivanco et al. 2020, <https://git.ligo.org/francisco.hernandez/toast>),

Corner (Foreman-Mackey 2016, <https://github.com/dfm/corner.py>).

new methodology



Span uncertainty in YN interaction

#parameters and priors

$$\theta_{\text{EOS}} = \{R_{\sigma\Lambda}, R_{\omega\Lambda}\}$$

$$R_{\sigma\Lambda} \sim U[0, 1]$$

and $R_{\omega\Lambda} \sim U[0, 1]$

- ONLY explore the couplings of Λ hyperons;
- Keep Σ, Ξ hyperon couplings fixed to their empirical values or based on SU(3) symmetry.

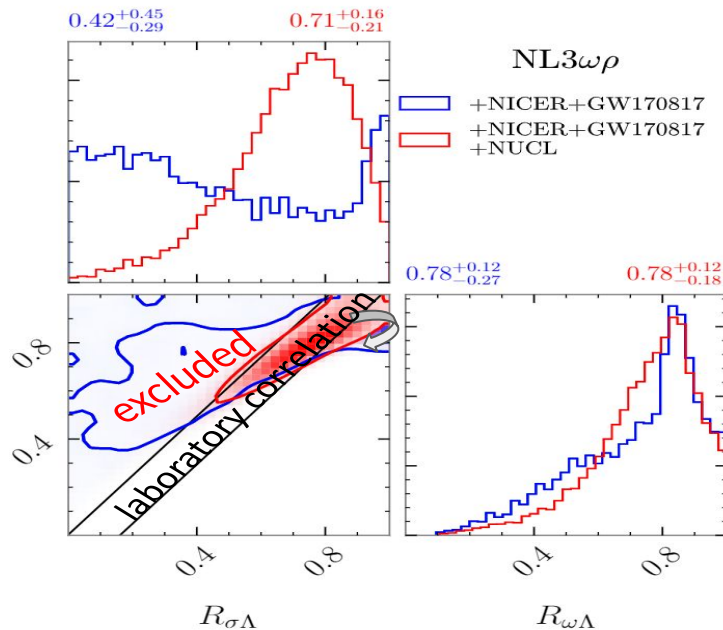
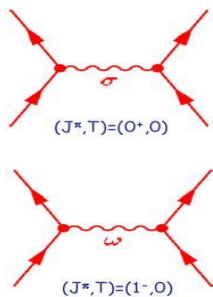
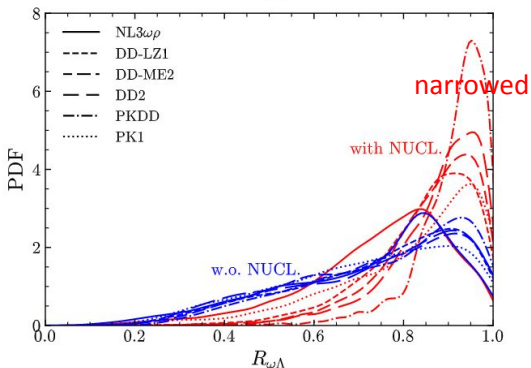
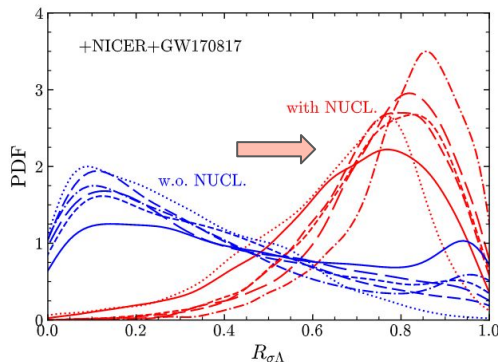
#posteriors

Most Probable Intervals of $R_{\sigma\Lambda}$ and $R_{\omega\Lambda}$ (68% Credible Intervals)

EOS	+NICER		+NICER +NUCL		+NICER +GW170817		+NICER +GW170817 +NUCL	
	$R_{\sigma\Lambda}$	$R_{\omega\Lambda}$	$R_{\sigma\Lambda}$	$R_{\omega\Lambda}$	$R_{\sigma\Lambda}$	$R_{\omega\Lambda}$	$R_{\sigma\Lambda}$	$R_{\omega\Lambda}$
LHS	0.821 ^{+0.125} _{-0.463}	0.755 ^{+0.073} _{-0.155}	0.865 ^{+0.074} _{-0.208}	0.658 ^{+0.130} _{-0.194}	0.941 ^{+0.035} _{-0.048}	0.763 ^{+0.034} _{-0.028}	0.658 ^{+0.163} _{-0.251}	0.752 ^{+0.049} _{-0.095}
RMF201	0.760 ^{+0.186} _{-0.520}	0.759 ^{+0.081} _{-0.224}	0.658 ^{+0.172} _{-0.249}	0.672 ^{+0.138} _{-0.215}	0.949 ^{+0.032} _{-0.056}	0.769 ^{+0.035} _{-0.028}	0.842 ^{+0.090} _{-0.250}	0.754 ^{+0.061} _{-0.136}
NL3	0.424 ^{+0.330} _{-0.293}	0.746 ^{+0.156} _{-0.261}	0.681 ^{+0.171} _{-0.247}	0.768 ^{+0.136} _{-0.214}	0.399 ^{+0.379} _{-0.291}	0.794 ^{+0.128} _{-0.216}	0.765 ^{+0.130} _{-0.191}	0.840 ^{+0.101} _{-0.163}
Hybrid	0.363 ^{+0.381} _{-0.265}	0.807 ^{+0.132} _{-0.276}	0.750 ^{+0.130} _{-0.179}	0.865 ^{+0.096} _{-0.157}	0.305 ^{+0.388} _{-0.217}	0.764 ^{+0.143} _{-0.254}	0.777 ^{+0.118} _{-0.181}	0.869 ^{+0.090} _{-0.147}
TM2	0.311 ^{+0.330} _{-0.221}	0.751 ^{+0.179} _{-0.494}	0.736 ^{+0.145} _{-0.201}	0.856 ^{+0.102} _{-0.193}	0.323 ^{+0.487} _{-0.237}	0.784 ^{+0.158} _{-0.300}	0.772 ^{+0.137} _{-0.239}	0.870 ^{+0.086} _{-0.204}
NLSV1	0.252 ^{+0.285} _{-0.183}	0.756 ^{+0.167} _{-0.281}	0.688 ^{+0.117} _{-0.227}	0.863 ^{+0.100} _{-0.199}	0.247 ^{+0.279} _{-0.177}	0.744 ^{+0.182} _{-0.259}	0.689 ^{+0.122} _{-0.225}	0.866 ^{+0.100} _{-0.206}
PK1	0.254 ^{+0.273} _{-0.185}	0.756 ^{+0.172} _{-0.250}	0.687 ^{+0.139} _{-0.222}	0.869 ^{+0.099} _{-0.216}	0.248 ^{+0.271} _{-0.170}	0.754 ^{+0.176} _{-0.247}	0.683 ^{+0.130} _{-0.220}	0.867 ^{+0.101} _{-0.222}
NL3 $\omega\rho$	0.384 ^{+0.393} _{-0.280}	0.773 ^{+0.147} _{-0.247}	0.690 ^{+0.163} _{-0.208}	0.759 ^{+0.131} _{-0.176}	0.420 ^{+0.448} _{-0.294}	0.777 ^{+0.127} _{-0.269}	0.712 ^{+0.157} _{-0.215}	0.778 ^{+0.121} _{-0.183}
S271v6	0.287 ^{+0.290} _{-0.207}	0.775 ^{+0.158} _{-0.232}	0.750 ^{+0.105} _{-0.144}	0.886 ^{+0.080} _{-0.123}	0.304 ^{+0.286} _{-0.183}	0.782 ^{+0.157} _{-0.147}	0.740 ^{+0.118} _{-0.161}	0.884 ^{+0.083} _{-0.134}
HC	0.266 ^{+0.253} _{-0.192}	0.517 ^{+0.316} _{-0.370}	0.733 ^{+0.110} _{-0.156}	0.902 ^{+0.070} _{-0.134}	0.266 ^{+0.304} _{-0.189}	0.783 ^{+0.157} _{-0.226}	0.737 ^{+0.106} _{-0.160}	0.902 ^{+0.072} _{-0.134}
DD- LZ1	0.298 ^{+0.321} _{-0.218}	0.775 ^{+0.152} _{-0.251}	0.769 ^{+0.122} _{-0.190}	0.871 ^{+0.083} _{-0.148}	0.327 ^{+0.381} _{-0.223}	0.792 ^{+0.142} _{-0.254}	0.772 ^{+0.128} _{-0.177}	0.870 ^{+0.087} _{-0.139}
DD-ME2	0.275 ^{+0.337} _{-0.192}	0.771 ^{+0.167} _{-0.299}	0.770 ^{+0.120} _{-0.172}	0.885 ^{+0.078} _{-0.137}	0.267 ^{+0.345} _{-0.188}	0.776 ^{+0.160} _{-0.237}	0.767 ^{+0.128} _{-0.168}	0.883 ^{+0.079} _{-0.124}
DD2	0.292 ^{+0.346} _{-0.205}	0.775 ^{+0.163} _{-0.252}	0.783 ^{+0.121} _{-0.173}	0.901 ^{+0.071} _{-0.135}	0.305 ^{+0.392} _{-0.221}	0.785 ^{+0.153} _{-0.276}	0.789 ^{+0.119} _{-0.157}	0.900 ^{+0.069} _{-0.120}
PKDD	0.267 ^{+0.347} _{-0.185}	0.806 ^{+0.140} _{-0.244}	0.820 ^{+0.095} _{-0.153}	0.930 ^{+0.051} _{-0.090}	0.282 ^{+0.420} _{-0.240}	0.813 ^{+0.136} _{-0.147}	0.835 ^{+0.102} _{-0.147}	0.932 ^{+0.047} _{-0.083}
FKVW	0.327 ^{+0.343} _{-0.236}	0.677 ^{+0.217} _{-0.260}	0.647 ^{+0.196} _{-0.250}	0.706 ^{+0.171} _{-0.211}	0.353 ^{+0.356} _{-0.240}	0.696 ^{+0.272} _{-0.203}	0.658 ^{+0.177} _{-0.254}	0.716 ^{+0.158} _{-0.217}
PC-PK1	0.283 ^{+0.310} _{-0.210}	0.701 ^{+0.215} _{-0.134}	0.650 ^{+0.150} _{-0.205}	0.770 ^{+0.147} _{-0.211}	0.282 ^{+0.319} _{-0.139}	0.703 ^{+0.212} _{-0.219}	0.651 ^{+0.148} _{-0.208}	0.771 ^{+0.146} _{-0.215}
OMEG	0.272 ^{+0.298} _{-0.194}	0.778 ^{+0.156} _{-0.244}	0.726 ^{+0.117} _{-0.171}	0.880 ^{+0.089} _{-0.153}	0.273 ^{+0.275} _{-0.188}	0.775 ^{+0.163} _{-0.242}	0.731 ^{+0.119} _{-0.167}	0.889 ^{+0.082} _{-0.152}

Hyperon-nucleon interactions in the relativistic Lagrangian

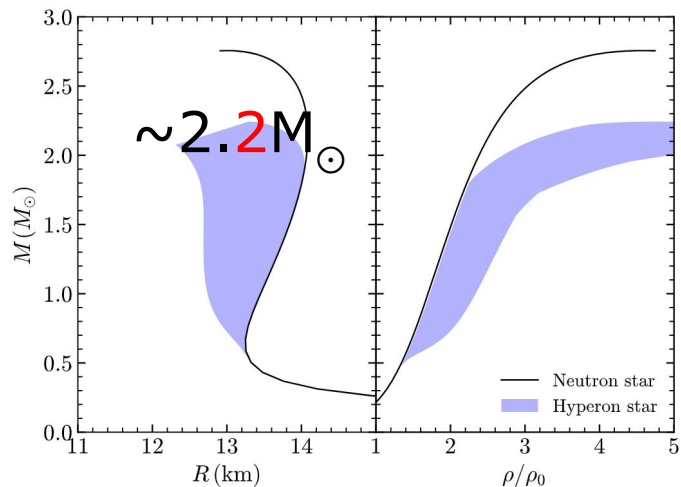
- Hypernuclei constraint favors large values of $R_{\sigma\Lambda}$ and $R_{\omega\Lambda}$ and disfavors small values of both couplings;
- The addition of astrophysical observational data on top of the laboratory $R_{\sigma\Lambda}$ - $R_{\omega\Lambda}$ correlation rotates the linear correlation **slightly** towards the direction of small values of $R_{\omega\Lambda}$.



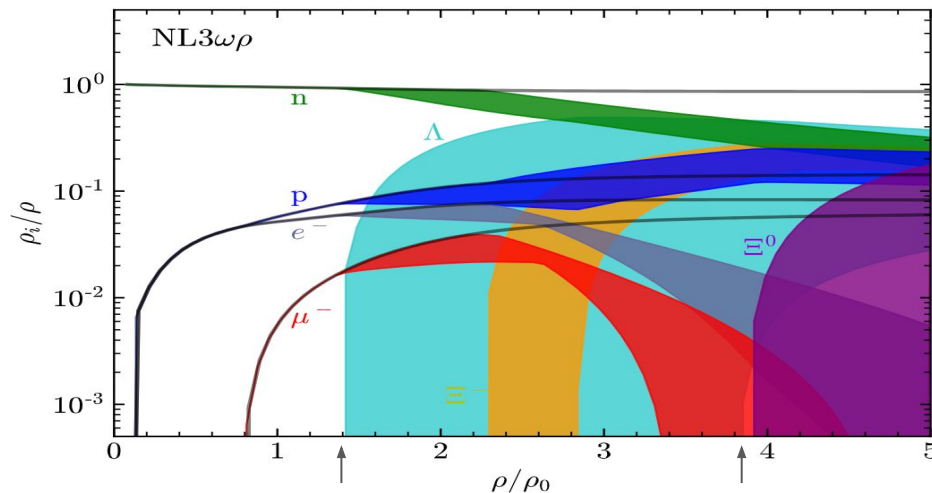
Hypernuclei results dominate !

Current status of the hypernuclear matter and hyperon star properties due to the uncertain YN+ interaction

- Taking the NL3 $\omega\rho$ one as an exemplary stiffest one;



2205.10631



Due to hyperons, the maximum mass is lowered by $\sim 20\%$: $M_{\max} = 2.176^{+0.085}_{-0.202} M_{\odot}$ (68% credible interval);

And the stellar radius is smaller above $\sim 0.5 M_{\odot}$ and grows with the stellar mass.

threshold density of Λ hyperons: $1.4\text{--}3.8\rho_0$

Unclear whether Λ or Ξ^- appear first.

In the following: + a few Ξ hypernuclei, $\Lambda\Lambda$ hypernuclei

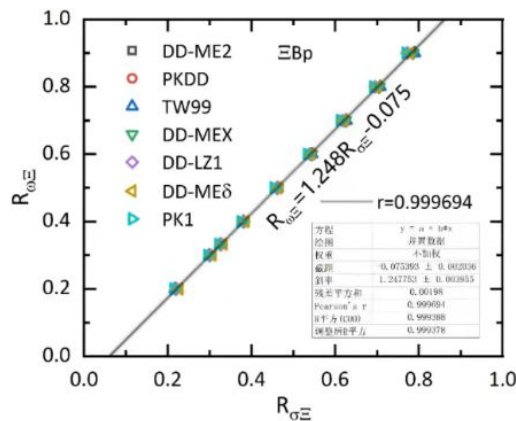
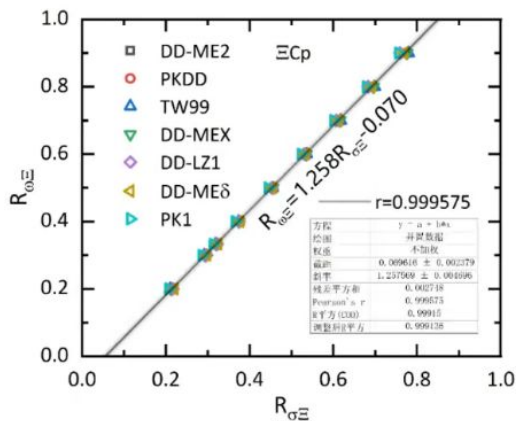
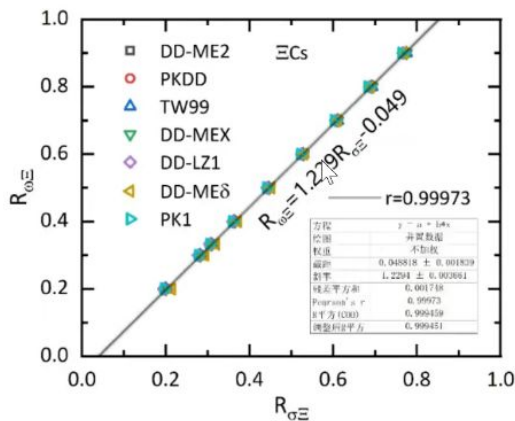
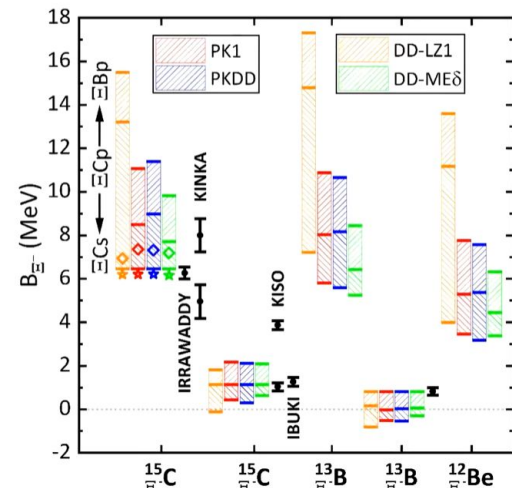
To include Ξ hypernuclei data

- Compared with Λ hypernuclei, Ξ hypernuclei are more difficult to produce and observe: Shorter lifetimes; Smaller cross sections;
- Fitting three s.p. separation energies of Ξ hypernuclei,

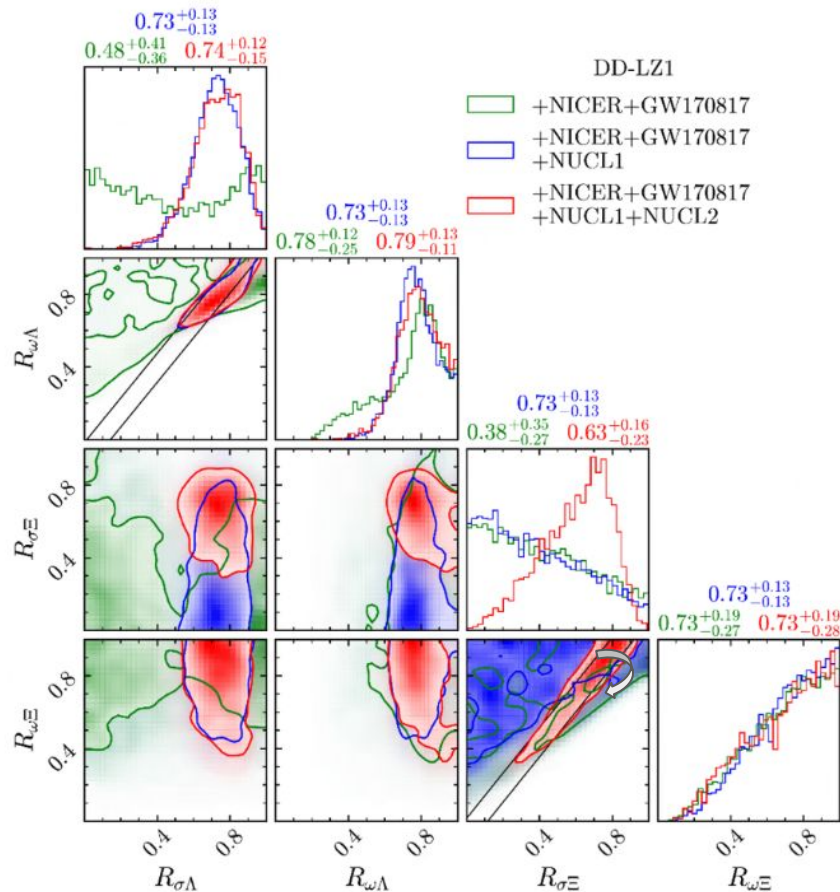
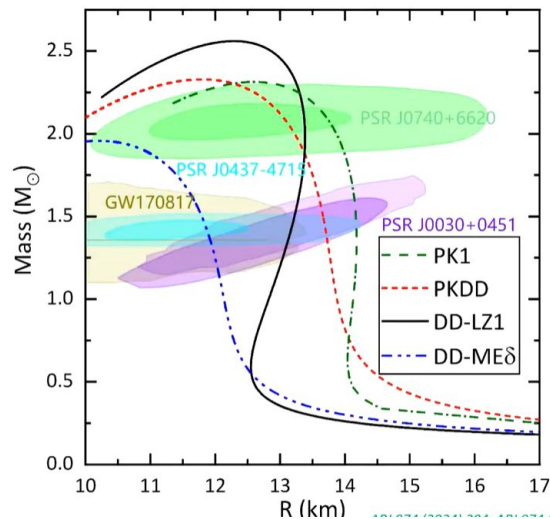
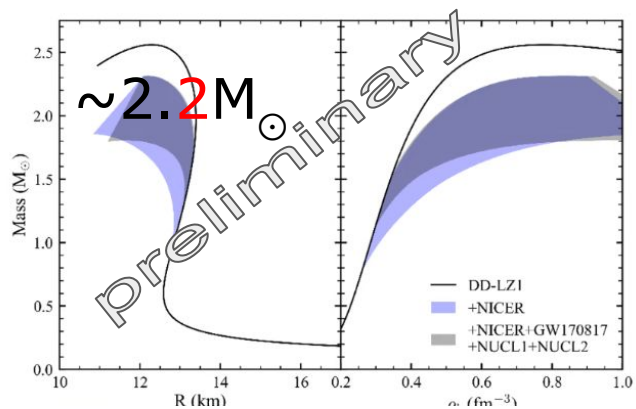
PRC 111 (2025) 014301

TABLE I. The ratio of σ - Ξ coupling strengths $g_{\sigma\Xi}/g_{\sigma N}$ for various RMF effective interactions, which are determined by fitting to the possible experimental values of the Ξ^- separation energy of $^{15}_{\Xi}\text{C}$ in the $1s$ state [11] (denoted as ΞCs), in the $1p$ state [10] (denoted as ΞCp), and of $^{13}_{\Xi}\text{B}$ in the $1p$ state [21] (denoted as ΞBp); see text for details. For other meson-hyperon coupling channels, the ratio of coupling strengths are fixed to be $g_{\omega\Xi}/g_{\omega N} = 0.333$, $g_{\rho\Xi}/g_{\rho N} = 1.000$, $g_{\delta\Xi}/g_{\delta N} = 1.000$, and, additionally, the ω - Ξ tensor coupling $f_{\omega\Xi} = -0.400g_{\omega\Xi}$.

	PK1	TW99	PKDD	DD-ME2	DD-MEX	DD-ME δ	DD-LZ1
ΞCs	0.304666	0.309145	0.312701	0.313264	0.309712	0.319533	0.305429
ΞCp	0.312236	0.318984	0.321078	0.322175	0.320552	0.324708	0.322607
ΞBp	0.320842	0.326105	0.328357	0.329127	0.326959	0.332777	0.327859



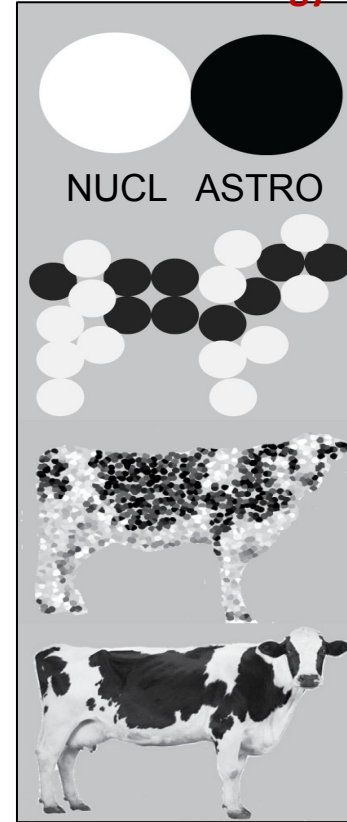
Adding likelihood of Ξ hypernuclei



Short summary and Exciting future

- We incorporate new information from hypernuclei calculations to address a long-standing issue in NS physics;
- We find that the strong correlation between the scalar and vector channel of YN interactions indicated by s.p. separation energy of available Λ , Ξ hypernuclei **ENSURE** that there is **sufficient (vector) repulsion** and a prediction of hyperon stars with $M_{\max} \sim 2.2M_{\odot}$;
- Comprehensive analysis of multi-messenger, multi-wavelength data ongoing to probe the EoS at **different** density regimes ->

new methodology



Connect consistently nuclear physics and GW+EM observations to probe the EoS at different density regimes (2021-)

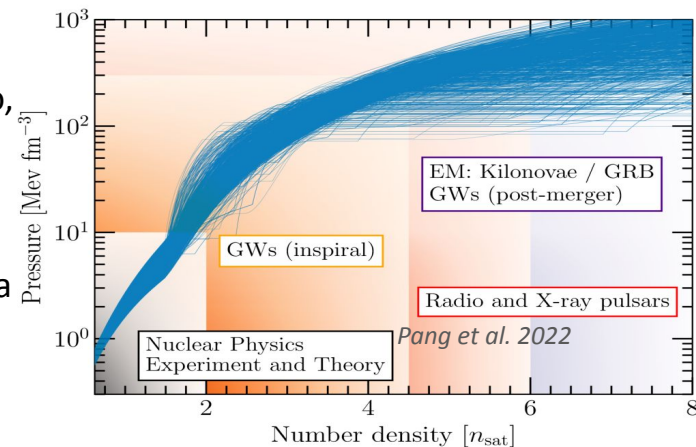
Data:

GW event of GW170817(+GW190425) & **kilonova** light curve of AT2017gfo,

NICER×**XMM-Newton**'s measurement of mass and radius of 2 PSRs,

(Mocked) **SKA**'s moment of inertia measurement on PSR J0737-3039,

Neutron-skin from PREX-II, CREX and the ab initio predictions on ^{208}Pb , ^{40}Ca



2021

2022

2023

2024

Massive PSRs (radio)

+ GW (static tide)

+ X-ray (NICER)

[2103.15119](#)

+mocked MOI (radio)

+DM

[2107.07979](#)

[2204.05560](#)

+n skin

[2305.16058](#)

+hypernuclei

[2205.10631](#)

+kilonova
(optical+)

[2211.02007](#)

+GW (dynamic tide)

+X-ray

(XMM-Newton)

[2402.02799](#)

+ΛΛ hypernuclei

(in preparation)

Constraint on the pressure at densities $\sim 1-3n_0$ effectively tightened

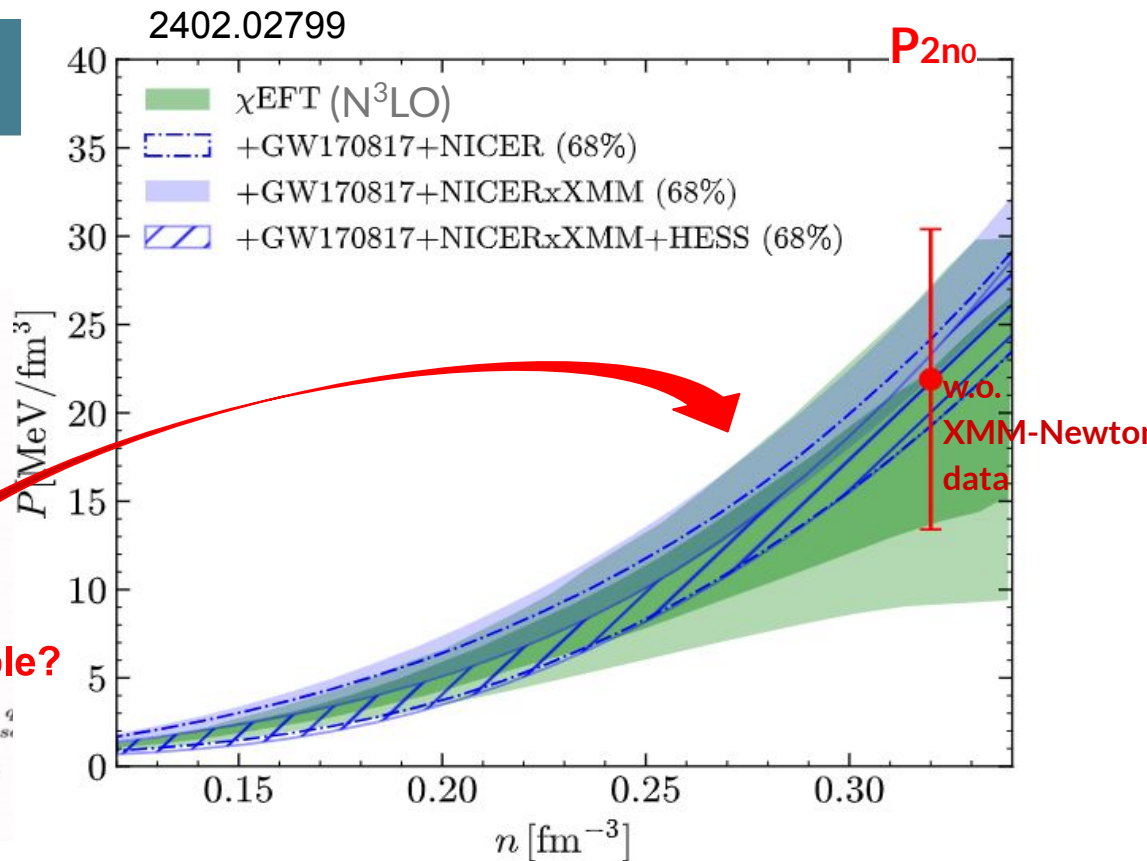
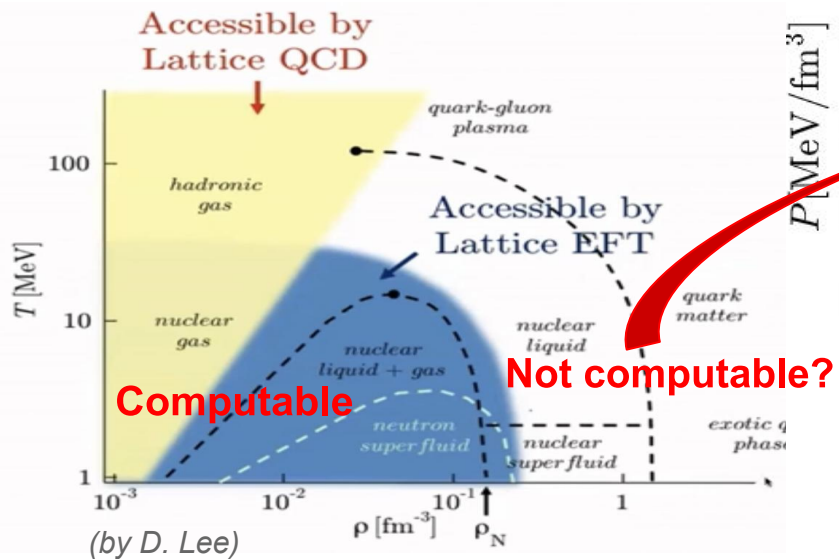
PHYSICAL REVIEW D

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Thermal x-ray studies of neutron stars and the equation of state

Zhiqiang Miao, Liqiang Qi, Juan Zhang, Ang Li, and Mingyu Ge
Phys. Rev. D **109**, 123005 – Published 3 June 2024



NICER view of PSR J0030+0451: MSP parameter estimation

Riley et al. 2019

Follow the propagation of light from the NS surface to the observer through the curved spacetime around the star;

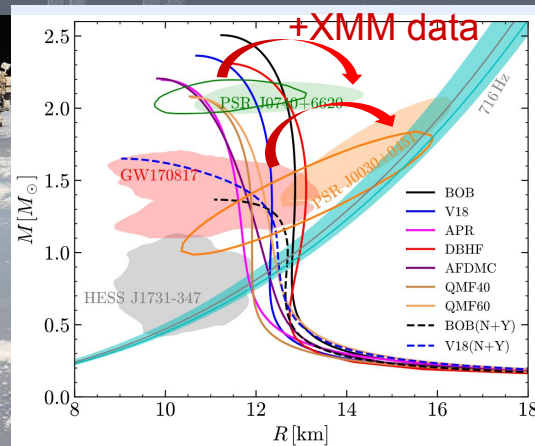
Main difficulty lies in the degeneracy between **MANY** parameters it employs: (in addition to M, R) source D ; H column density, parameters needed to describe the surface T map, angles encoding the orientation of this map with respect to the rotational axis and the observer's LOS;

Significant uncertainties in M, R estimate may be from **low counts** in the detected lightcurve or **noisy data** can lead to degeneracies between the model parameters or multi-modal posteriors:

X-ray telescopes with **large effective area, high energy resolution, and high time resolution** in the soft X-ray band, as well as adequate imaging capability to discriminate the **background**:
NICER (2017), EP-FXT (2024), eXTP (~2030!), ATHENA (~2030).

XMM-Newton data can provide (indirect) NICER background constraint:

The introduction of the XMM-Newton data (EPIC MOS1 and MOS2) led to a reduction in background estimates from NICER-only data, resulting in higher compactness value:



2402.02799

“new precision era”



SKAO



CSST

Electromagnetic waves

FAST



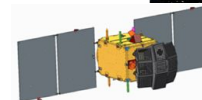
JCMT



TMT



Insight/
HXMT



Einstein
Probe (2024)

Radio

mm/sub-mm

FIR

MIR

NIR

NUV

UV

Soft X-ray

Hard X-ray

γ -ray

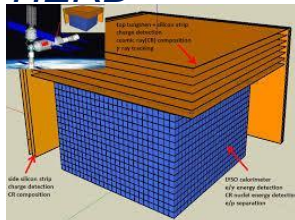
Neutrinos

JUNO

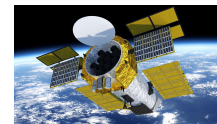


Cosmic Rays

HERD



HUBS



eXTP
(~2030)

enhanced X-ray Timing and Polarimetry mission

web: <https://astro.xmu.edu.cn/>



18 faculty members
11 postdocs
70 graduate students

magnetic waves

SKAO



FAST



R

Neut

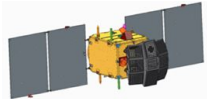
JUNO



Soft X-ray Hard X-ray γ -ray



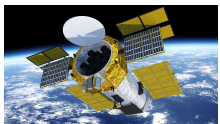
Insight/
HXMT



Einstein
Probe (2024)



HUBS



eXTP
(~2030)

enhanced X-ray Timing and Polarimetry mission

web: <https://astro.xmu.edu.cn/>

What we learn so far

- ❑ Produce self-consistent framework for EoS modeling
- ❑ Check if constraints from all available constrains are fulfilled
Demonstrate the consistency between laboratory and astrophysical nuclear matter in neutron stars by considering low-density nuclear physics constraints (from ^{208}Pb neutron-skin thickness) and high-density astrophysical constraints (from neutron star global properties).
- ❑ Prepare priors for a set of EoS parameters, also incorporating phase transitions
- ❑ Statistically establish the effective stiffness of neutron star EoS
General requirements adopted (e.g., causality) indicate the EoS is moderately stiff, with sound speed squared peaked at $\sim 0.8c_s^2$ in NS matter, not subject to the type of phase transitions.
- ❑ Examine whether current data favour (strong) 1st-order phase transition inside NS cores
Current data compatible with both possibilities; Evidence of a phase transition strengthened for stiff hadronic EoSs (like DD2); Perform preliminary test of the consistency of the parameters in the metastable NS matter and in the nearly symmetric nuclear matter formed in HIC
- ❑ **Confront Λ , Ξ hypernuclei data with the neutron star observational data**
With the relaxation of the commonly-assumed SU(3) symmetry, the data of single Λ hypernuclei ensures a large enough scalar hyperon coupling to match the large vector hyperon coupling
- ❑ Suggest possibility to distinguish neutron stars with quark stars from simultaneous measurement of the stellar radius and the moment of inertia
If the MOI is measured large for PSR J0737-3039 A, $I_A \geq 1.4 \times 10^{45} \text{ g cm}^2$, it is most likely a quark star rather than a neutron star with or without a quark core, provided that the accuracy of the radius measurement is at least $\sim 1 \text{ km}$
- ❑ Comprehensive analysis of multi-messenger, multi-wavelength data **ongoing** to probe the EoS at different density regimes...

Thank you!