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## The triple-alpha reaction at low temperatures by an exact three-body model

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The triple-alpha reaction plays a significant role in nucleosynthesis heavier than  $^{12}\text{C}$  and concomitant stellar evolution [1]. The reaction rates of this reaction at the helium-burning temperatures,  $T_9 > 0.1$ , are dominated by the sequential process via two narrow resonances:  $\alpha+\alpha \rightarrow ^8\text{Be}(0^+_{-1}; \text{g.s.})$ ,  $^8\text{Be}+\alpha \rightarrow ^{12}\text{C}(0^+_{-2}; E=0.379 \text{ MeV})$  [2,3], and they have been understood relatively well through the studies of the Hoyle state.  $T_9$  is temperature in unit of  $10^9 \text{ K}$ ;  $E$  is the center-of-mass energy to the 3 alpha threshold in  $^{12}\text{C}$ . In contrast, the direct triple-alpha process from ternary continuum states,  $\alpha+\alpha+\alpha \rightarrow ^{12}\text{C}$ , for  $T_9 < 0.1$  still seems to remain in an open question. The direct process corresponds to the non-resonant component in [4], and it is thought to be important in the astrophysical sites of novae, X-ray bursts, and Type-Ia SNe, leading to the nucleosynthesis of the hot CNO cycle and rp-process [5].

In NACRE [2],  $^8\text{Be}$  is assumed to be bound as a particle, and the reaction rates have been estimated by an improved model based on the pioneering works of [4]. To determine the rates more dynamically, the methods with hyper-spherical coordinates are used in [6-9], and the Coulomb modified Faddeev method is adopted in [10]. Whereas  $^8\text{Be}$  continuum states are treated adiabatically in [8-10], the direct process is calculated non-adiabatically in [6,7].

In this presentation, I discuss the direct triple-alpha process by using a non-adiabatic Faddeev hyper-spherical harmonics and R-matrix (HHR\*) expansion method [6]. I illustrate that the calculated photo-disintegration cross sections of  $^{12}\text{C}(2^+_{-1}(E=-2.835 \text{ MeV}) \rightarrow 0^+)$  of HHR\* are much smaller than those of the recent adiabatic models [9,10] for  $0.15 < E < 0.35 \text{ MeV}$ . The resultant rates have the strong temperature dependence at  $T_9 = 0.1$ , as well as NACRE, and their numerical values are expressed in a simple analytic form. From the comparison between the calculations, I find that the current standard rates [2,3] can be reduced by about  $10^{-4}$  at  $T_9 \sim 0.05$ , because of the accurate description of  $^8\text{Be}$  break-up. As an example of new rates, I also examine the ignition critical density [4] of helium burning in accreting white dwarfs. The present model with [11] leads to the ignition density of about  $(3 \times 10^8) \text{ g cm}^{-3}$  at  $T_9 = 0.01$ , which seems to be consistent with [4]. Due to the reduction of the rate at  $T_9 = 0.05$ , the derived ignition density appears to be insensitive to the temperatures in  $0.01 < T_9 < 0.05$ .

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