Complex-valued problem on FRG analysis of relativistic BEC #14

Kazuya Mameda ^{1,2} Fumio Terazaki¹ Katsuhiko Suzuki¹

1. Tokyo University of Science 2. RIKEN iTHEMS

FT, K. Mameda and K. Suzuki Prog. Theor. Exp. Phys. 2024, 123B02

Nucleosynthesis and Evolution of Neutron Stars at YITP, Kyoto University, on Jan. 28, 2025

Introduction - relativistic Bose-Einstein Condensates (BEC) -



Electron chemical potential $\mu_{\rm e}$ increases with density ρ .

In the outer core, ($\rho \approx \rho_{\text{saturation}}$) $\mu_e \approx 100 \mathrm{MeV} \approx \mathrm{free \ pion \ mass}$

Pion-BEC affects key properties of neutron stars:

- softening equation of state
- enhancing cooling rates via neutrino emission

Analyses of relativistic BEC in the complex scalar theory $\mathcal{L} = \left|\partial_{\mu}\varphi\right|^{2} + m^{2}|\varphi|^{2} + \lambda|\varphi|^{4}$

- mean-field approximation (MFA) J. I. Kapusta (1981)
 including no effects of quantum fluctuations
- Functional Renormalization Group
 Complex-valued **(FRG)** L. F. Palfares (2012)

problem E. E. Svanes et al. (2011), O(N) model

Purpose: FRG analysis of the complex scalar theory

as a basic step for studying pion-BEC

Formalism - Functional Renormalization Group (FRG) -

• FRG calculates effective actions $\Gamma[\langle \phi \rangle]$ non-perturbatively.

$$\Gamma[\langle \phi \rangle] \coloneqq -\log \int \mathcal{D}\phi \, e^{-S + \int d^4 x \, j \phi} + \int d^4 x \, j \langle \phi \rangle$$
$$\int \mathcal{D}\phi = \prod_{q=0}^{\Lambda} \int_{-\infty}^{+\infty} d\phi_q \quad \text{integrations over} \\ \text{momentum modes } \phi_q$$

- An IR-regulator R_k is introduced in Γ . $\rightarrow \Gamma_k$
 - Γ_k follows the **Wetterich flow equation**:

$$\frac{\partial \Gamma_k}{\partial k} = \frac{1}{2} \operatorname{Tr} \left[\frac{\partial R_k}{\partial k} \frac{1}{\Gamma_k^{(2)} + R_k} \right] = \frac{1}{2} \bigotimes^{\text{C. Wetterich (1993)}}$$

 $\partial \Gamma_{k}^{(2)}$ $\partial_{\tau} \Gamma^{(n)} = f \left(\Gamma^{(n+1)} \Gamma^{(n+2)} \right)$ Model - the complex scalar theory -

$$f_k = \int_0^\beta \mathrm{d}\tau \int \mathrm{d}^3x \left[\nabla \varphi \cdot \nabla \varphi^* + (\partial_\tau - \mu) \varphi (\partial_\tau + \mu) \varphi^* + V_k(\rho) \right] \quad \rho := |\varphi|^2$$

 $\Gamma_k \to \beta V \cdot U_k(\rho)$: effective potential ($\varphi(x) \to \varphi$)

 $\rho_0 \neq 0$ $ho_0 = 0$ thermodynamics in FRG BEC normal order parameter (BEC²): $\rho_0(T,\mu)$: = argmin $U_{k\to 0}(\rho)$ phase phase

$$\Omega(T,\mu) = U_{k\to 0}(\rho_0) \quad \text{pressure: } p = -\Omega, \quad \text{number density: } n = -\partial\Omega/\partial\mu,$$

grand-canonical potential entropy: $s = -\partial\Omega/\partial\mu$, energy density: $\varepsilon = -p + \mu n + Ts$

flow equation in
$$R_k(p) = (k^2 - |\vec{p}|^2) \cdot \theta(k^2 - |\vec{p}|^2)$$
 D.F. Litim (2001)

$$\frac{k}{\partial k} = \int_{k}^{3} \int_{k}^{3} \frac{1}{2} \int_{k}^{3} \frac{1}{2} \int_{k}^{3} \frac{1}{2} \int_{k}^{4} \int_{k}^{4$$

(quantum). Solution

Derivative Expansion

$$\Gamma_{k} = \sum_{n=0}^{\infty} g_{2n} \left[\langle \phi \rangle, \left(\partial_{\mu} \langle \phi \rangle \right)^{2n} \right] \to g_{0} + g_{2}$$

Local Potential Approximation

 $Z_k(\partial_\mu \langle \phi \rangle)^2 \to 1 \cdot (\partial_\mu \langle \phi \rangle)^2$ in g_2

$$\frac{\partial U_k}{\partial k} = \frac{k^4}{6\pi^2} \sum_{s=\pm} \frac{E_s^2 + \mu^2 - E_k^2}{E_s^2 - E_{-s}^2} \frac{1 + 2n_{\rm B}(E_s)}{E_s} \qquad E_k^2 = k^2 + V_k' + \rho V_k'', n_{\rm B}(E) = \left(e^{E/T} - 1\right)^{-1} E_k + \left[E_k^2 + \mu^2 \pm \left\{4E_k^2 \mu^2 + \left(\rho V_k''\right)^2\right\}^{1/2}\right]^{1/2}\right]^{1/2}$$

E. E. Svanes et al. (2011), consider only real part E_{-} takes imaginary values at some ρ -regions. L. F. Palfares (2012), restrict parameters

We require U_k to be real and finite. cf. condition in free theory consider ρ -range $\mu^2 < m_{\mathrm{eff},k>0}^2 \coloneqq k^2 + V'_{k>0}$ where $E_-^2(\rho) > 0 \xrightarrow{\mu^2} \mu^2 \leq m_{\mathrm{eff},k\to0}^2 \leftarrow$ $\Omega_{\rm free}$ is defined only if $\mu^2 \leq m_{\rm hare}^2$. Correspondence

Results

BEC phase diagram $V_{k=\Lambda} = \overline{m}^2 \rho + \overline{\lambda} \rho^2, \overline{m} / \Lambda = 0.1$ $\overline{\lambda} = 0.50$ 2.5 $ar{\lambda}=1.0, T=0, \mu/ar{m}=2.0$ 2.0 $k/\Lambda = 1.0$ 0.8 ${}^{+}\bar{m}^{4}$ \bar{m}

- For values of ρ smaller than the potential minimum "", these ρ -points violate the condition, $\mu^2 > m_{\text{eff},k}^2(\rho)$.
- In a strong coupling regime ($\overline{\lambda} = 1.00$), quantum fluctuations strongly influence the behavior of the relativistic BEC.



• Regulator dependence...

The Litim R_k keeps poles away from relevant ρ -regions as much as possible. D.F. Litim (2001)

> The complex-valued problem would occur in other regulators.

Summary

The complex scalar theory in FRG suffers from the complex-valued problem.

- We developed the novel method and obtained real valued BEC.
- We found in a strong coupling regime, quantum fluctuations are important.

Outlooks:

 $\overline{\lambda} = 1.00$

- Application for models including pions
- Analyses of parameters that have not been explored before