

# Speed limits and optimal transport: General bounds and maximal speed of particle transport

Tan Van Vu  
YITP, Kyoto University

Vu and Saito, PRL 130, 010402 (2023)

Vu, Kuwahara, and Saito, Quantum 8, 1483 (2024)

Hydrodynamics of low-dimensional interacting systems  
June 2-13 2025, YITP, Kyoto



# Outline

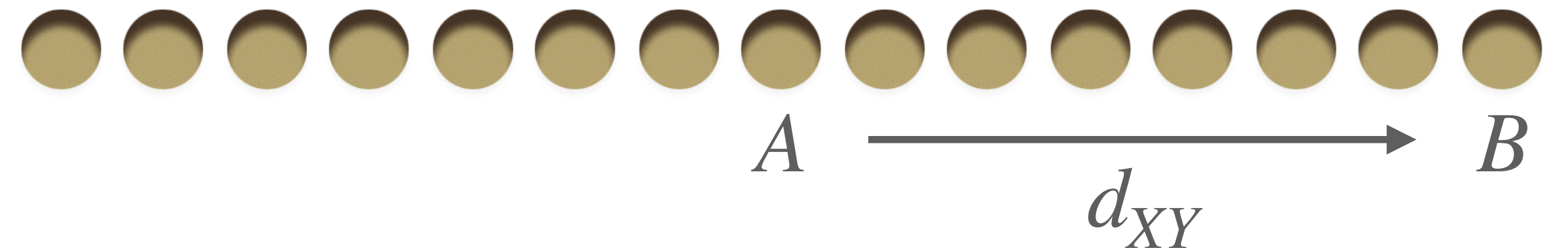
- Speed of information propagation and particle transport
- Optimal transport theory & quantum speed limit
- General speed limit based on optimal transport theory
- Results on bosonic transport in closed quantum systems

# Information propagation

- Upper limit on speed of information propagation

# Information propagation

- Upper limit on speed of information propagation



# Information propagation

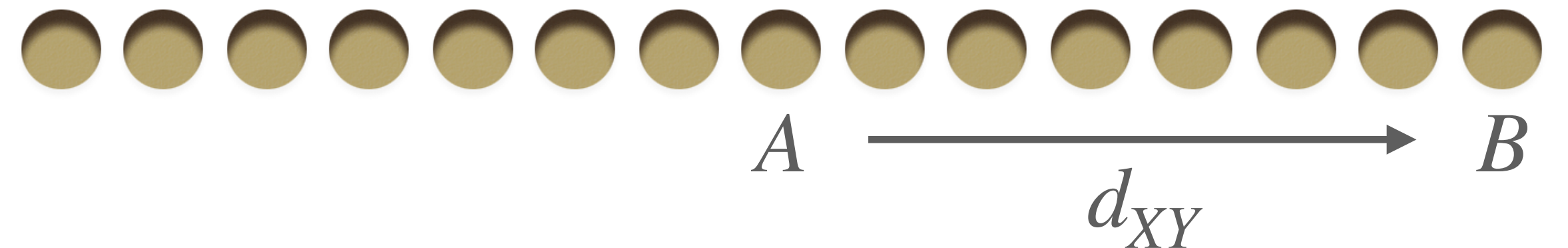
- Upper limit on speed of information propagation

$$\|[A(\tau), B]\| \leq c \exp[-a(d_{XY} - v\tau)]$$

Lieb and Robinson, Commun. Math. Phys. (1972)

$A, B$  : local operators with supports  $X$  and  $Y$

$A(t) := e^{iHt} A e^{-iHt}$  : Heisenberg operator



# Information propagation

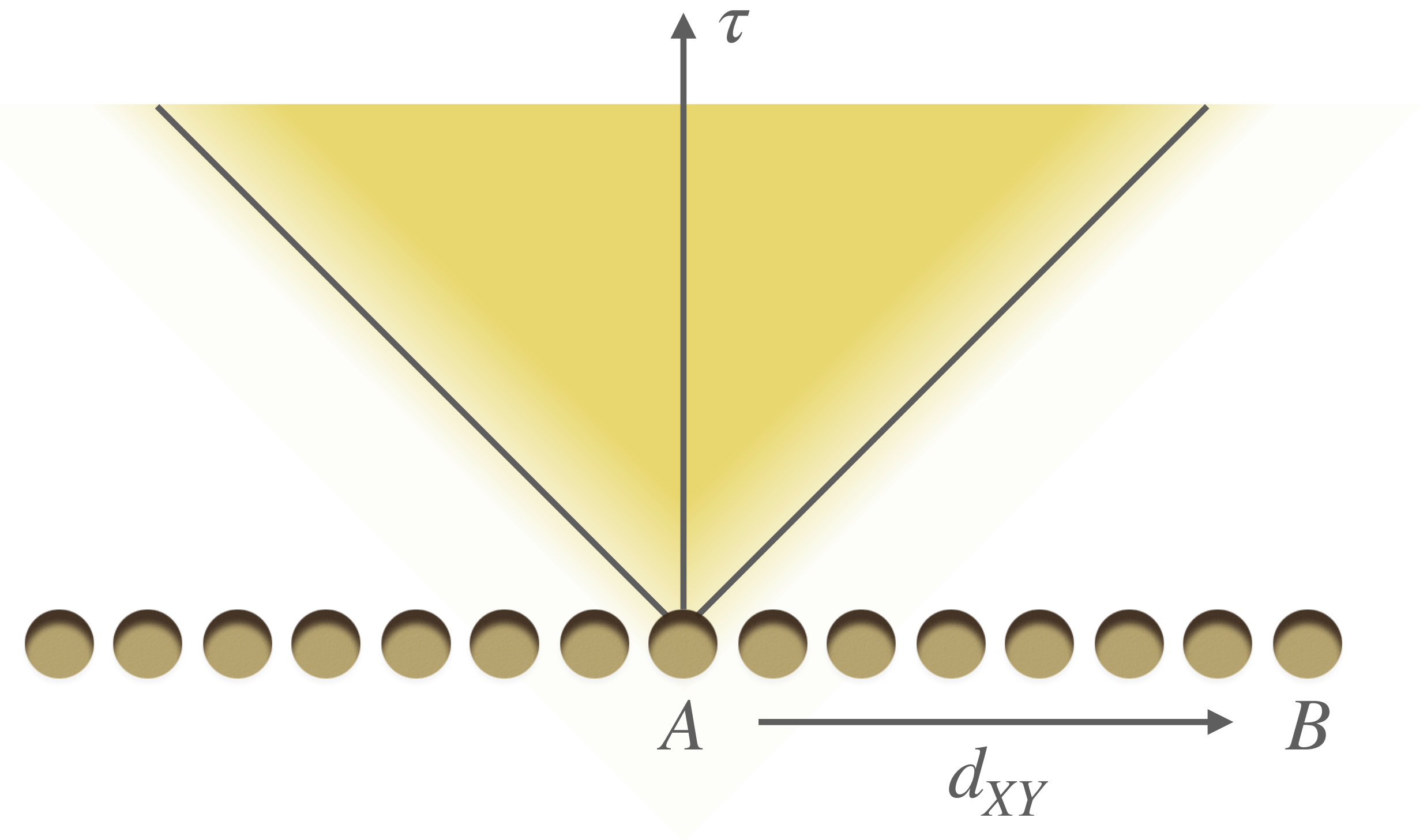
- Upper limit on speed of information propagation

$$\|[A(\tau), B]\| \leq c \exp[-a(d_{XY} - v\tau)]$$

Lieb and Robinson, Commun. Math. Phys. (1972)

$A, B$  : local operators with supports  $X$  and  $Y$

$A(t) := e^{iHt} A e^{-iHt}$  : Heisenberg operator



effective linear light cone  $\tau \gtrsim d_{XY}$

# Operator spreading & information propagation

- Quantum metrology perspective



Wysocki and Chwedenczuk, Phys. Rev. Lett. (2025)

# Operator spreading & information propagation

- Quantum metrology perspective
  - I. Encoding parameter  $Q(\theta) = e^{-i\theta A_X} Q e^{i\theta A_X}$



Wysocki and Chwedenczuk, Phys. Rev. Lett. (2025)

# Operator spreading & information propagation

- Quantum metrology perspective
  1. Encoding parameter  $Q(\theta) = e^{-i\theta A_X} Q e^{i\theta A_X}$
  2. Propagating information  $Q(\theta, t) = e^{-iHt} Q(\theta) e^{iHt}$



Wysocki and Chwedenczuk, Phys. Rev. Lett. (2025)

# Operator spreading & information propagation

- Quantum metrology perspective
  1. Encoding parameter  $\varrho(\theta) = e^{-i\theta A_X} \varrho e^{i\theta A_X}$
  2. Propagating information  $\varrho(\theta, t) = e^{-iHt} \varrho(\theta) e^{iHt}$
  3. Estimating  $\theta$  from  $\text{tr}_{Y^c}\{\varrho(\theta, t)\}$



Wysocki and Chwedenczuk, Phys. Rev. Lett. (2025)

# Operator spreading & information propagation

- Quantum metrology perspective
  1. Encoding parameter  $\varrho(\theta) = e^{-i\theta A_X} \varrho e^{i\theta A_X}$
  2. Propagating information  $\varrho(\theta, t) = e^{-iHt} \varrho(\theta) e^{iHt}$
  3. Estimating  $\theta$  from  $\text{tr}_{Y^c}\{\varrho(\theta, t)\}$
- Cramer-Rao inequality

$$\text{var}[\hat{\theta}] \geq \frac{1}{\mathcal{F}(\theta)} \quad \mathcal{F}(\theta) : \text{quantum Fisher information}$$



Wysocki and Chwedenczuk, Phys. Rev. Lett. (2025)

# Operator spreading & information propagation

- Quantum metrology perspective
  1. Encoding parameter  $\varrho(\theta) = e^{-i\theta A_X} \varrho e^{i\theta A_X}$
  2. Propagating information  $\varrho(\theta, t) = e^{-iHt} \varrho(\theta) e^{iHt}$
  3. Estimating  $\theta$  from  $\text{tr}_{Y^c}\{\varrho(\theta, t)\}$



Wysocki and Chwedenczuk, Phys. Rev. Lett. (2025)

- Cramer-Rao inequality

$$\text{var}[\hat{\theta}] \geq \frac{1}{\mathcal{F}(\theta)} \quad \mathcal{F}(\theta) : \text{quantum Fisher information}$$

- Quantum Fisher information is upper bounded by operator spreading

$$\mathcal{F}(\theta) \leq \sum_{i=1}^3 \|[O_{Y,i}(t), A_X]\|^2$$

# Lieb-Robinson bound

# Lieb-Robinson bound

- Completely resolved for spin and fermion systems

$$\tau \gtrsim d_{XY}^{\min(1, \alpha-2D)}$$

Kuwahara and Saito, Phys. Rev. X (2020); Tran+, Phys. Rev. Lett. (2021)

# Lieb-Robinson bound

- Completely resolved for spin and fermion systems

$$\tau \gtrsim d_{XY}^{\min(1, \alpha-2D)}$$

Kuwahara and Saito, Phys. Rev. X (2020); Tran+, Phys. Rev. Lett. (2021)

- Under investigation for boson systems    ▶ interaction strength is unbounded

# Lieb-Robinson bound

- Completely resolved for spin and fermion systems

$$\tau \gtrsim d_{XY}^{\min(1, \alpha-2D)}$$

Kuwahara and Saito, Phys. Rev. X (2020); Tran+, Phys. Rev. Lett. (2021)

- Under investigation for boson systems    ▶ interaction strength is unbounded
  - ✓ short-range hopping with low-boson density initial states

Yin and Lucas, Phys. Rev. X (2022); Kuwahara, Vu, and Saito, Nat. Commun. (2024)

# Lieb-Robinson bound

- Completely resolved for spin and fermion systems

$$\tau \gtrsim d_{XY}^{\min(1, \alpha-2D)}$$

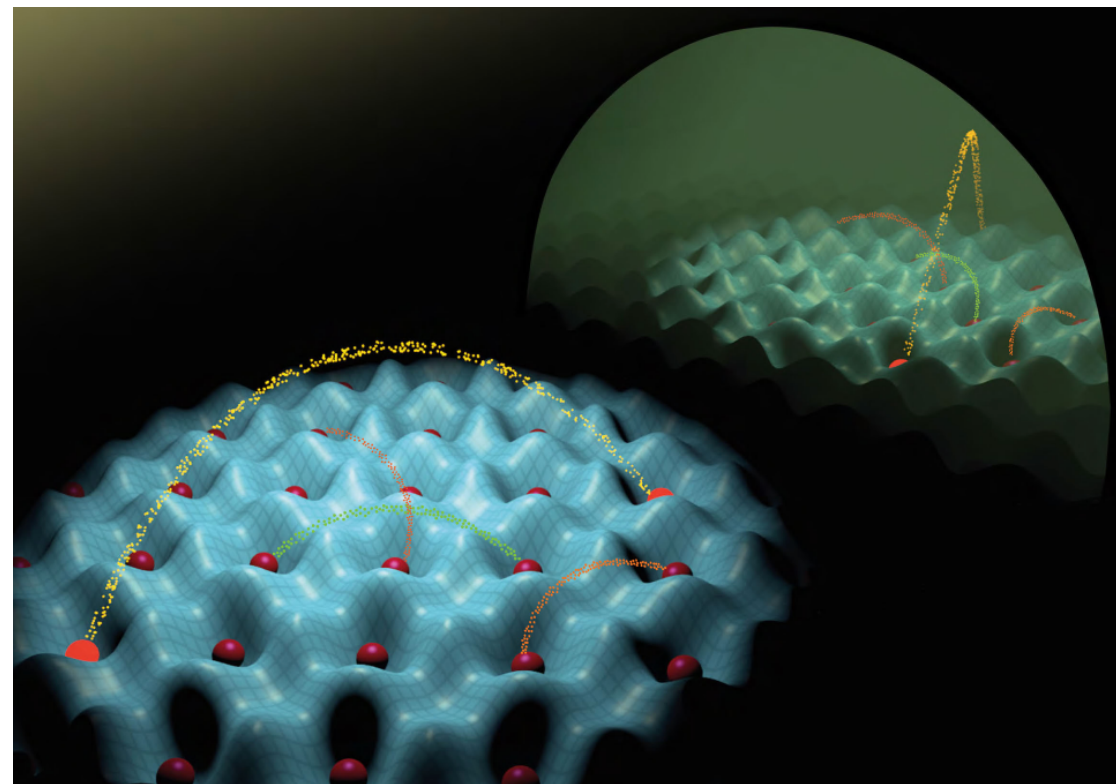
Kuwahara and Saito, Phys. Rev. X (2020); Tran+, Phys. Rev. Lett. (2021)

- Under investigation for boson systems    ▶ interaction strength is unbounded

✓ short-range hopping with low-boson density initial states

Yin and Lucas, Phys. Rev. X (2022); Kuwahara, Vu, and Saito, Nat. Commun. (2024)

× long-range hopping



Defenu+, Rev. Mod. Phys. (2023)

# Lieb-Robinson bound

- Completely resolved for spin and fermion systems

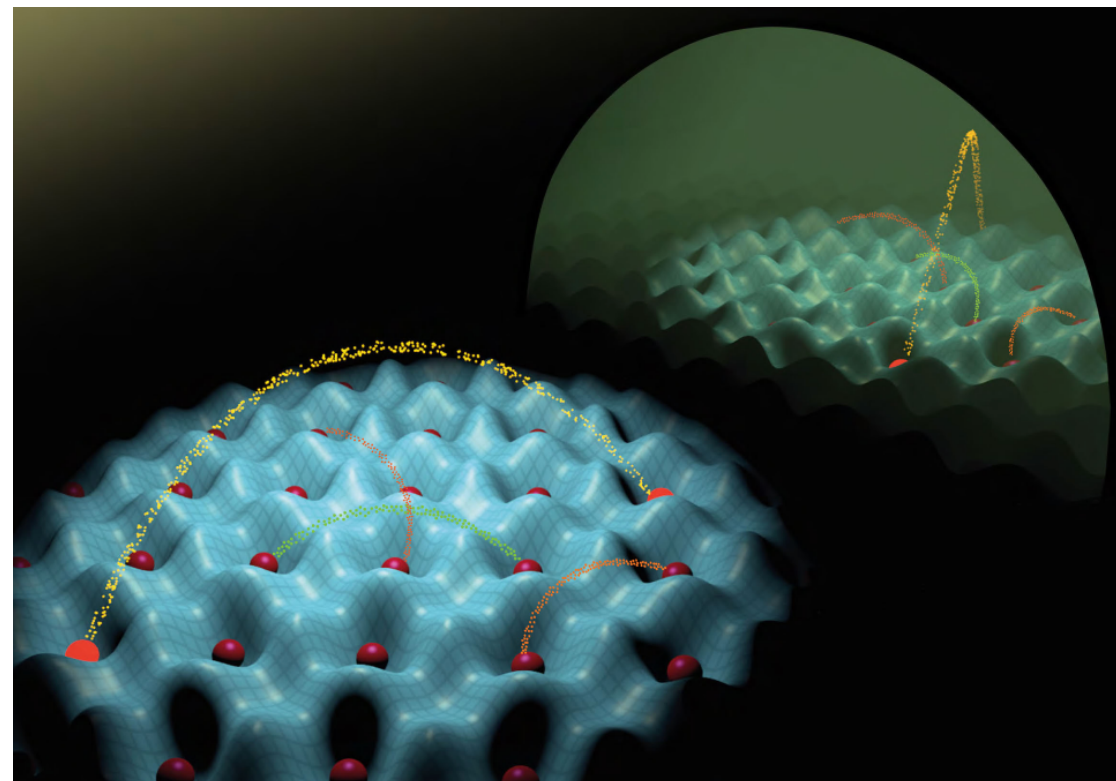
$$\tau \gtrsim d_{XY}^{\min(1, \alpha-2D)}$$

Kuwahara and Saito, Phys. Rev. X (2020); Tran+, Phys. Rev. Lett. (2021)

- Under investigation for boson systems    ▶ interaction strength is unbounded
  - ✓ short-range hopping with low-boson density initial states

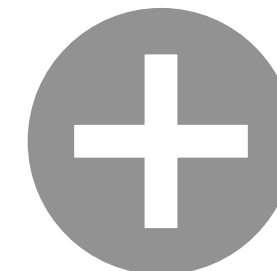
Yin and Lucas, Phys. Rev. X (2022); Kuwahara, Vu, and Saito, Nat. Commun. (2024)

× long-range hopping



Defenu+, Rev. Mod. Phys. (2023)

speed of particle  
transport



speed of information  
propagation

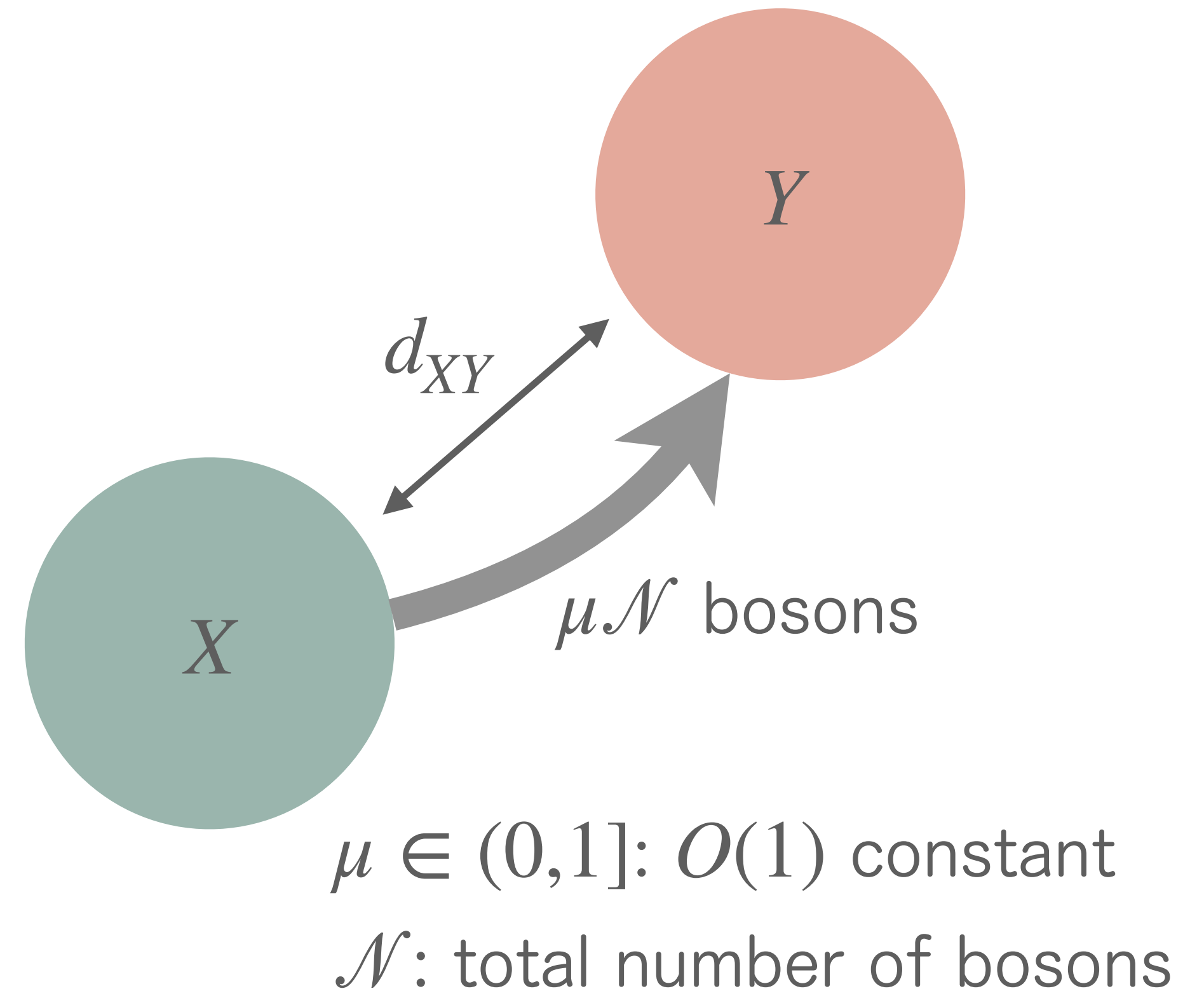
# Macroscopic particle transport

# Macroscopic particle transport

- Bosonic system on lattices
  - Time-dependent Hamiltonian
  - Long-range hopping & interactions

# Macroscopic particle transport

- Bosonic system on lattices
  - Time-dependent Hamiltonian
  - Long-range hopping & interactions
- **Macroscopic** particle transport

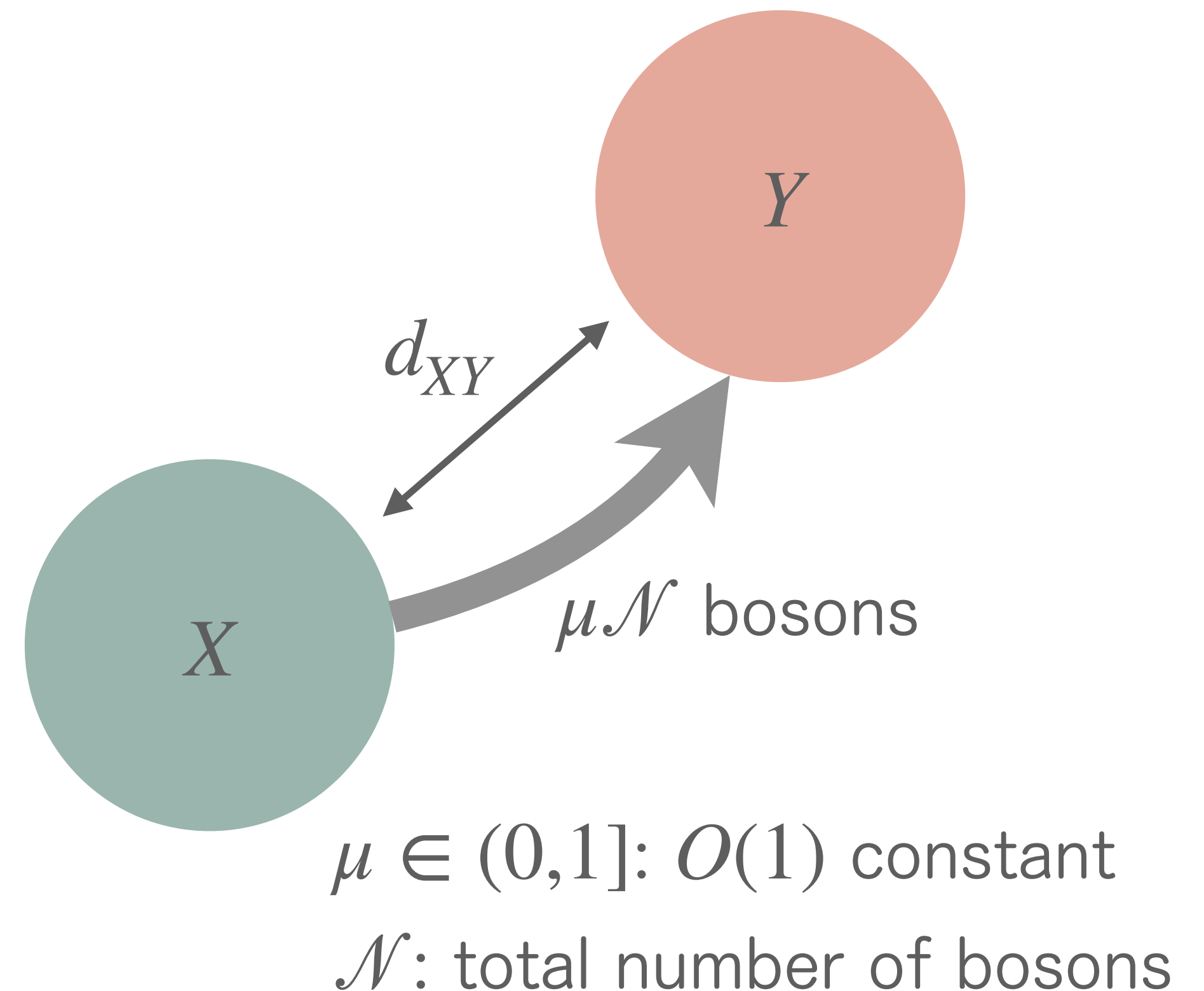


# Macroscopic particle transport

- Bosonic system on lattices
  - Time-dependent Hamiltonian
  - Long-range hopping & interactions

- **Macroscopic** particle transport

requirement :  $n_Y(\tau) \geq n_{X^c}(0) + \mu\mathcal{N}$

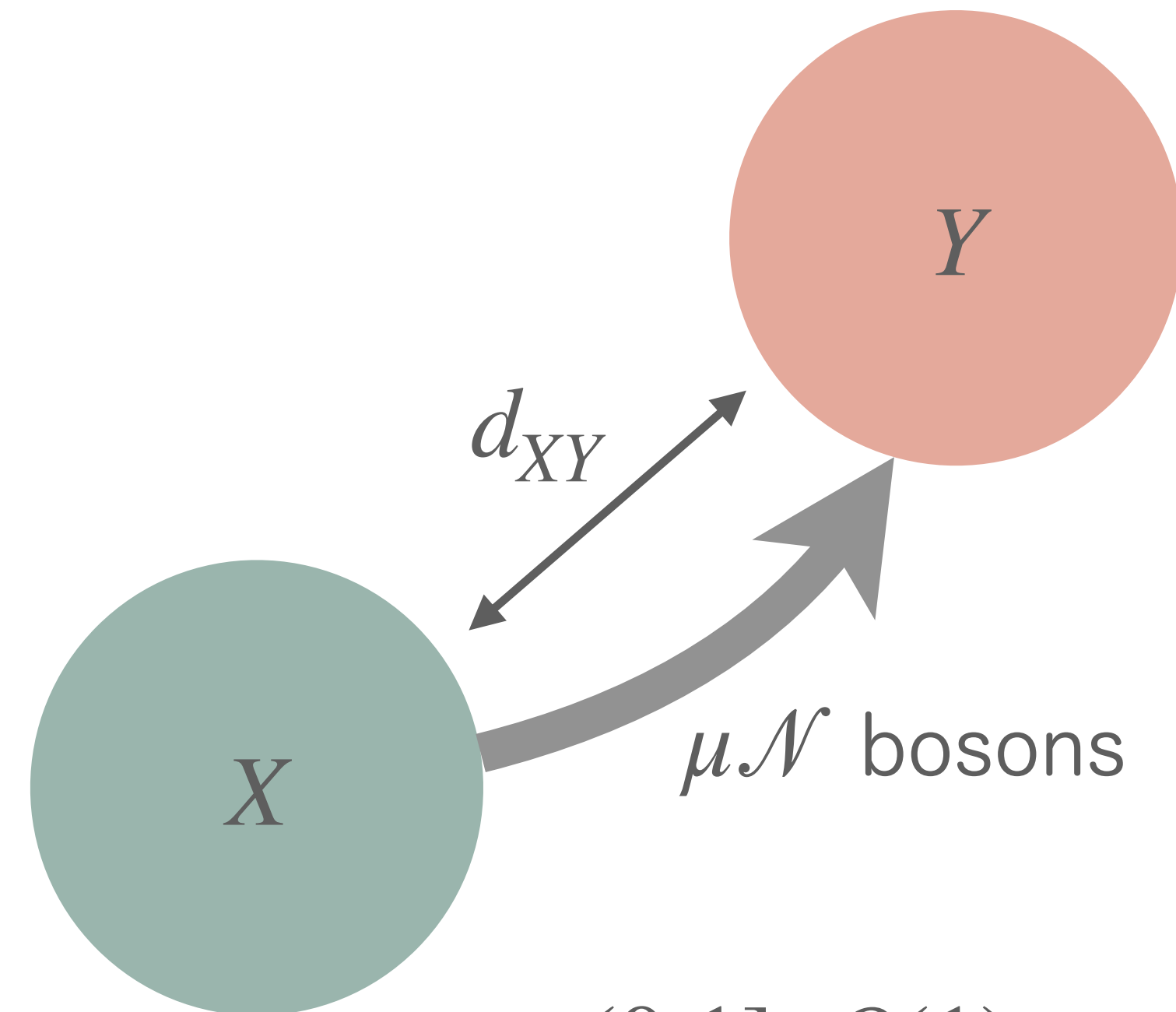


# Macroscopic particle transport

- Bosonic system on lattices
  - Time-dependent Hamiltonian
  - Long-range hopping & interactions

- **Macroscopic** particle transport

requirement :  $n_Y(\tau) \geq n_{X^c}(0) + \mu\mathcal{N}$



$\mu \in (0,1]$ :  $O(1)$  constant  
 $\mathcal{N}$ : total number of bosons

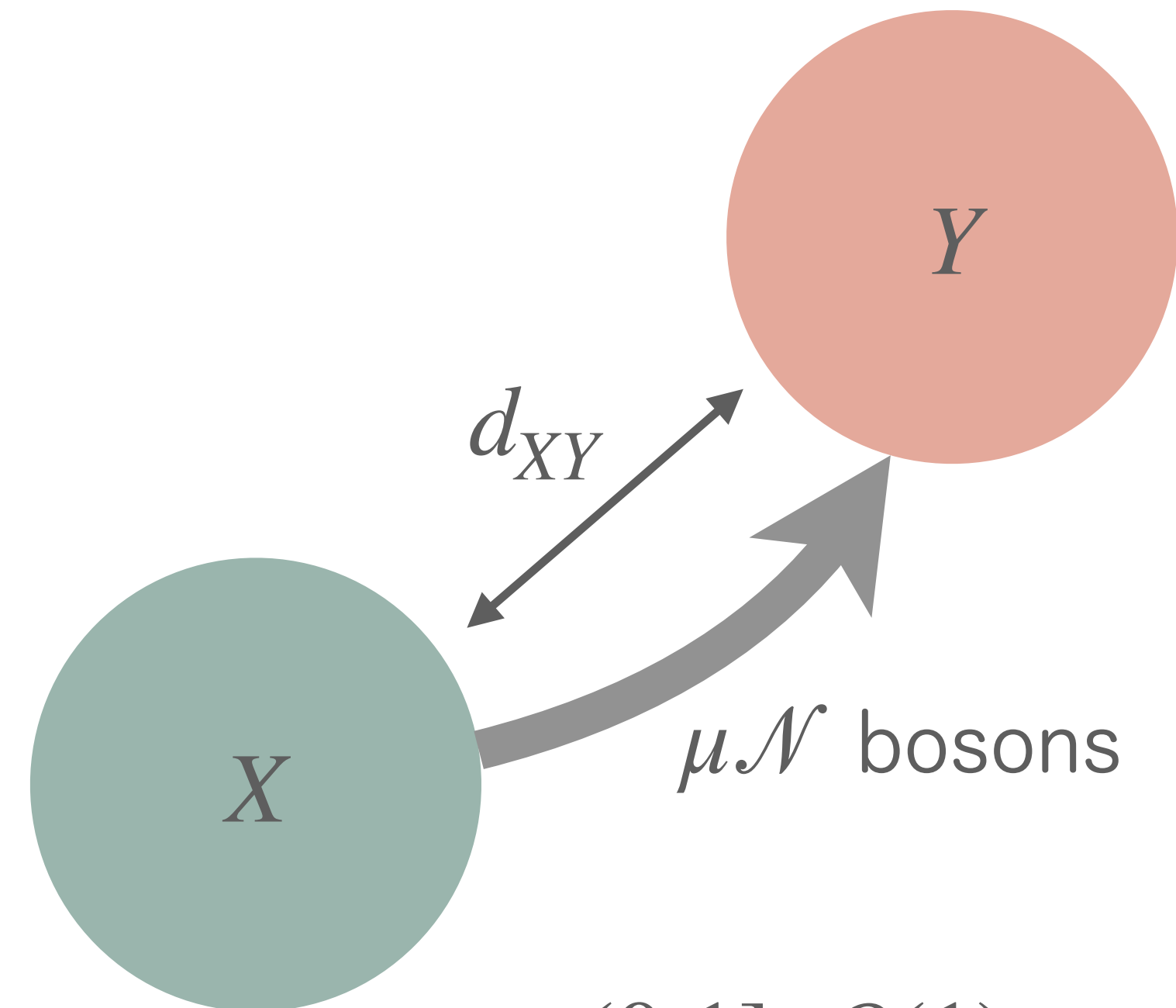
- Question: How fast can we transport particles?

# Macroscopic particle transport

- Bosonic system on lattices
  - Time-dependent Hamiltonian
  - Long-range hopping & interactions

- **Macroscopic** particle transport

requirement :  $n_Y(\tau) \geq n_{X^c}(0) + \mu \mathcal{N}$



$\mu \in (0,1]$ :  $O(1)$  constant  
 $\mathcal{N}$ : total number of bosons

- Question: How fast can we transport particles?

Our approach: **Optimal transport theory** & **Quantum speed limit**

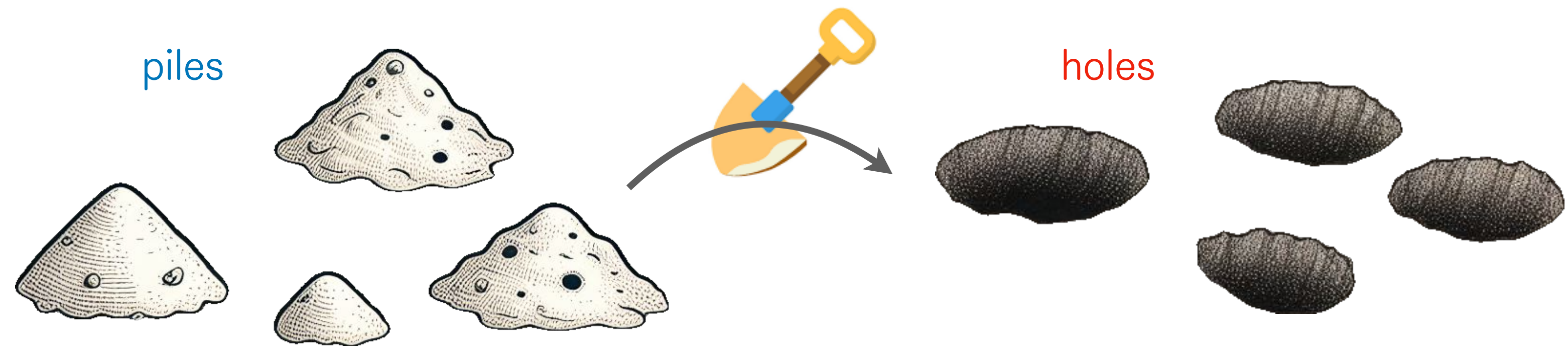
# Optimal transport

# Optimal transport

- Problem of transporting a source distribution to a target distribution



Gaspard Monge (1781)

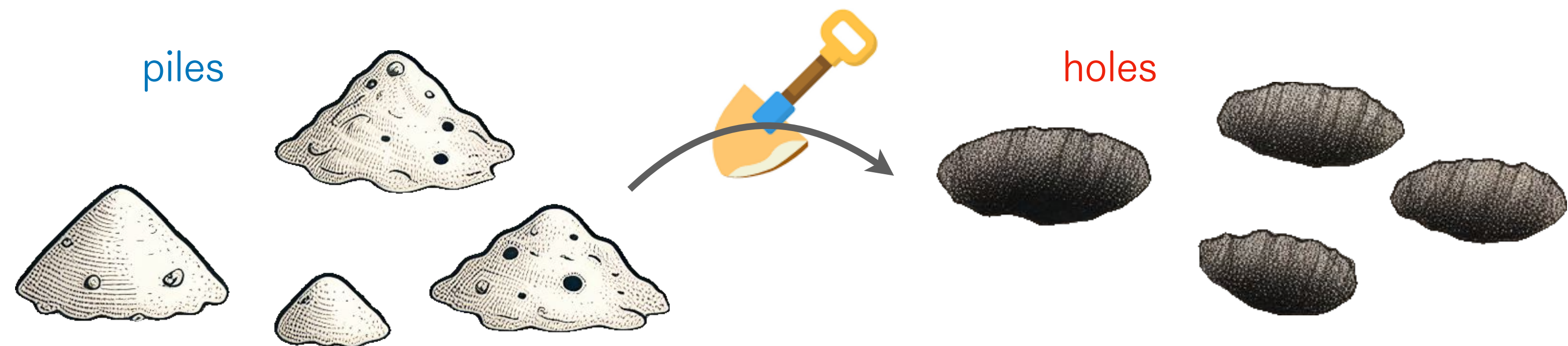


# Optimal transport

- Problem of transporting a source distribution to a target distribution



Gaspard Monge (1781)



Optimal transport plan & optimal transport cost

# Monge formulation

Optimal transport cost with respect to a cost function  $c(x, y)$  :

$$M(p^A, p^B) = \min_T \int dx \, c(x, T(x)) p^A(x)$$

$T$  : one-to-one map that preserves the total probability

# Monge formulation

Optimal transport cost with respect to a cost function  $c(x, y)$  :

$$M(p^A, p^B) = \min_T \int dx c(x, T(x)) p^A(x)$$

$T$  : one-to-one map that preserves the total probability

- Non-existence of a valid transport map:  $T$  might not exist in discrete cases because no mass can be split

# Monge formulation

Optimal transport cost with respect to a cost function  $c(x, y)$  :

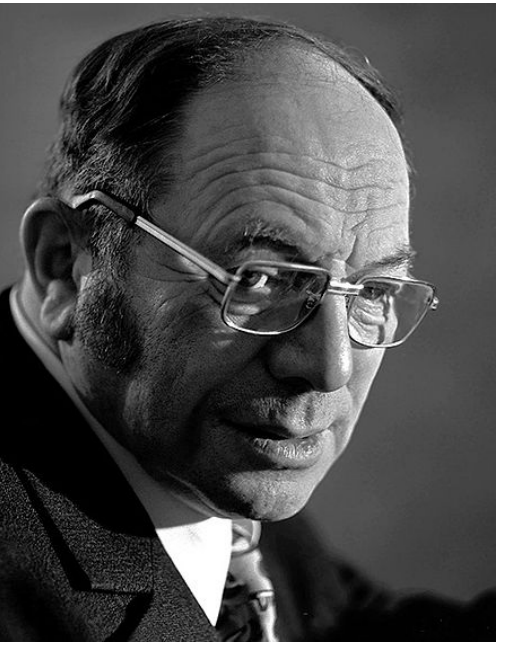
$$M(p^A, p^B) = \min_T \int dx c(x, T(x)) p^A(x)$$

$T$  : one-to-one map that preserves the total probability

- Non-existence of a valid transport map:  $T$  might not exist in discrete cases because no mass can be split
- Resolved by the relaxation of Kantorovich

# Kantorovich formulation

Leonid Kantorovich  
(1942)



# Kantorovich formulation

Leonid Kantorovich  
(1942)



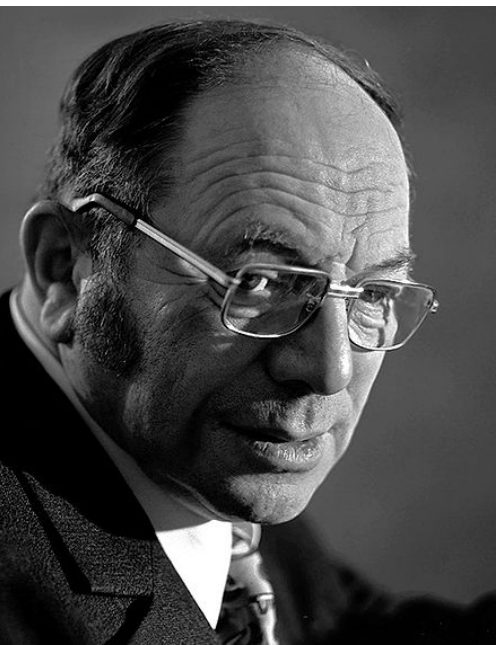
Optimal transport cost with respect to a cost function  $c(x, y)$  :

$$K(p^A, p^B) := \min_{\pi} \int dx dy c(x, y) \pi(x, y)$$

coupling  $\pi$  : joint probability distribution of  $p^A$  and  $p^B$

# Kantorovich formulation

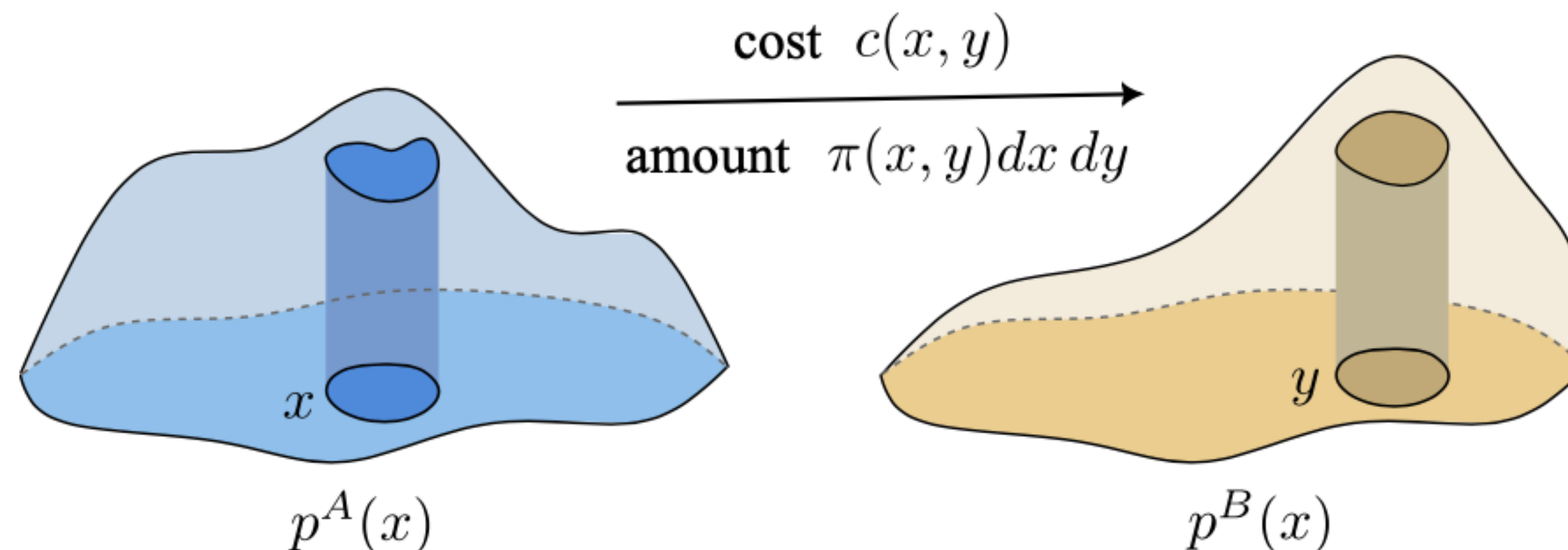
Leonid Kantorovich  
(1942)



Optimal transport cost with respect to a cost function  $c(x, y)$  :

$$K(p^A, p^B) := \min_{\pi} \int dx dy c(x, y) \pi(x, y)$$

coupling  $\pi$  : joint probability distribution of  $p^A$  and  $p^B$



# Kantorovich formulation

Leonid Kantorovich  
(1942)



Optimal transport cost with respect to a cost function  $c(x, y)$  :

$$K(p^A, p^B) := \min_{\pi} \int dx dy c(x, y) \pi(x, y)$$

coupling  $\pi$  : joint probability distribution of  $p^A$  and  $p^B$

# Kantorovich formulation

Leonid Kantorovich  
(1942)



Optimal transport cost with respect to a cost function  $c(x, y)$  :

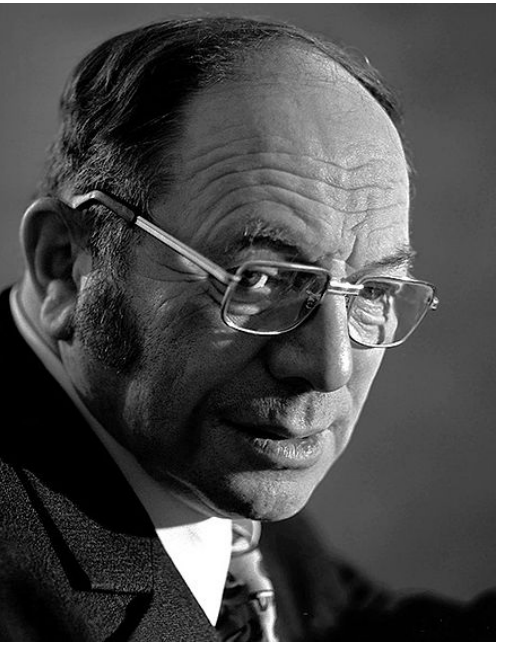
$$K(p^A, p^B) := \min_{\pi} \int dx dy c(x, y) \pi(x, y)$$

coupling  $\pi$  : joint probability distribution of  $p^A$  and  $p^B$

- In general,  $K(p^A, p^B) \leq M(p^A, p^B)$

# Kantorovich formulation

Leonid Kantorovich  
(1942)



Optimal transport cost with respect to a cost function  $c(x, y)$  :

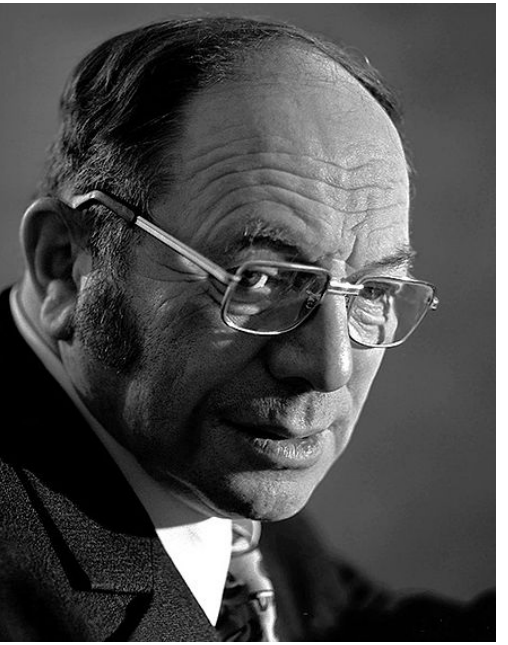
$$K(p^A, p^B) := \min_{\pi} \int dx dy c(x, y) \pi(x, y)$$

coupling  $\pi$  : joint probability distribution of  $p^A$  and  $p^B$

- In general,  $K(p^A, p^B) \leq M(p^A, p^B)$
- Optimal coupling of Kantorovich formulation always exists

# Kantorovich formulation

Leonid Kantorovich  
(1942)



Optimal transport cost with respect to a cost function  $c(x, y)$  :

$$K(p^A, p^B) := \min_{\pi} \int dx dy c(x, y) \pi(x, y)$$

coupling  $\pi$  : joint probability distribution of  $p^A$  and  $p^B$

- In general,  $K(p^A, p^B) \leq M(p^A, p^B)$
- Optimal coupling of Kantorovich formulation always exists
- Two formulations are equivalent if  $p^A$  has no atom and  $c$  is continuous

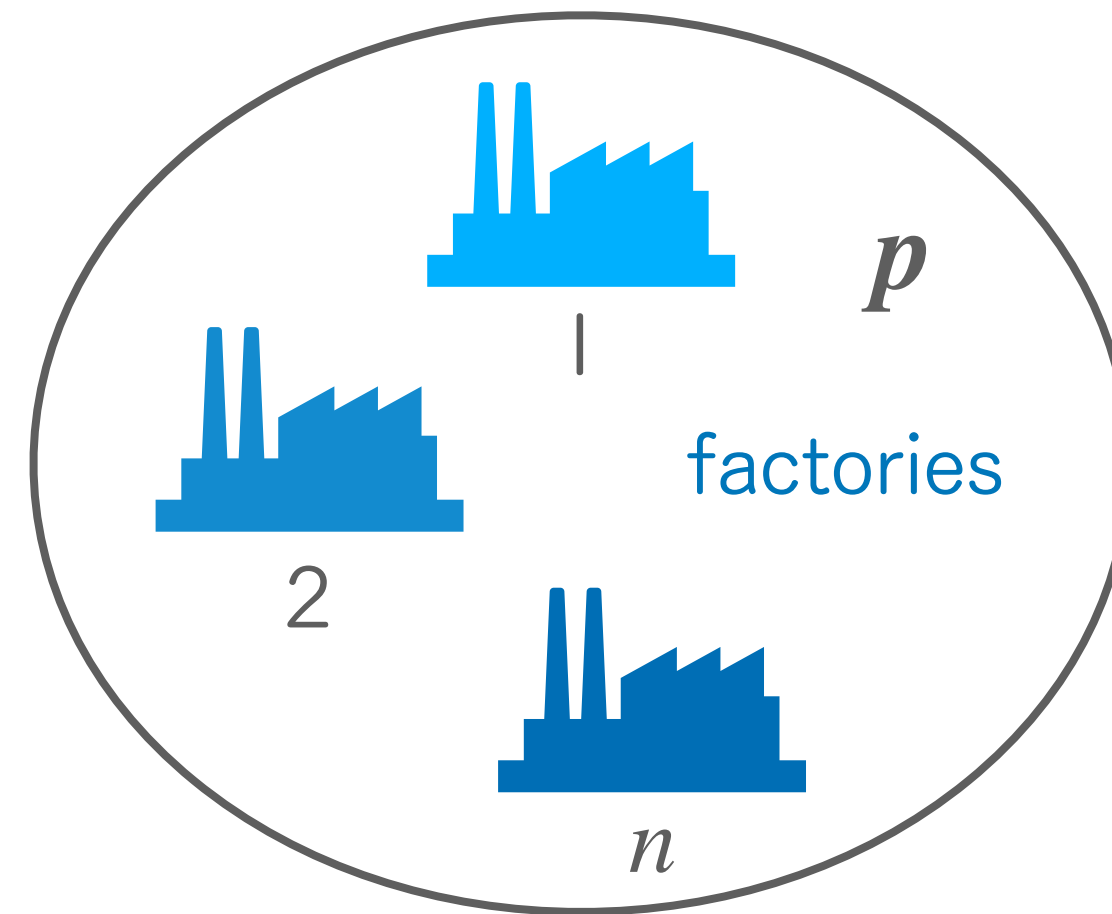
# Discrete optimal transport

# Discrete optimal transport

- Transport of goods

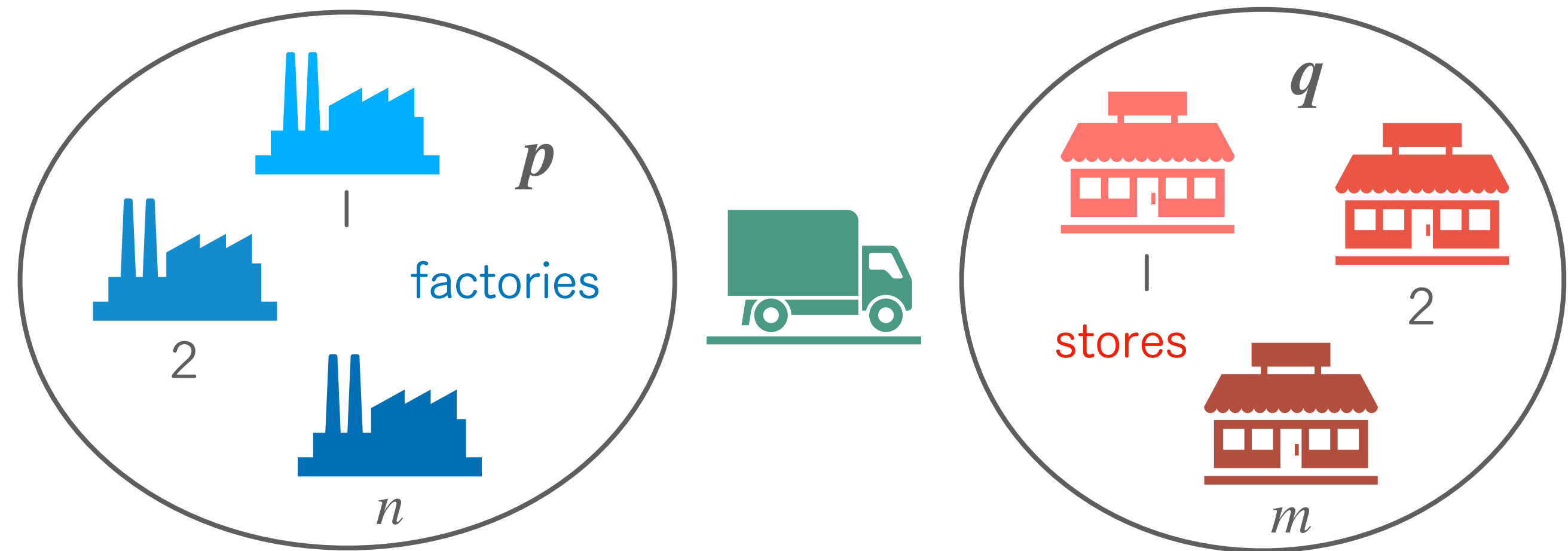
# Discrete optimal transport

- Transport of goods



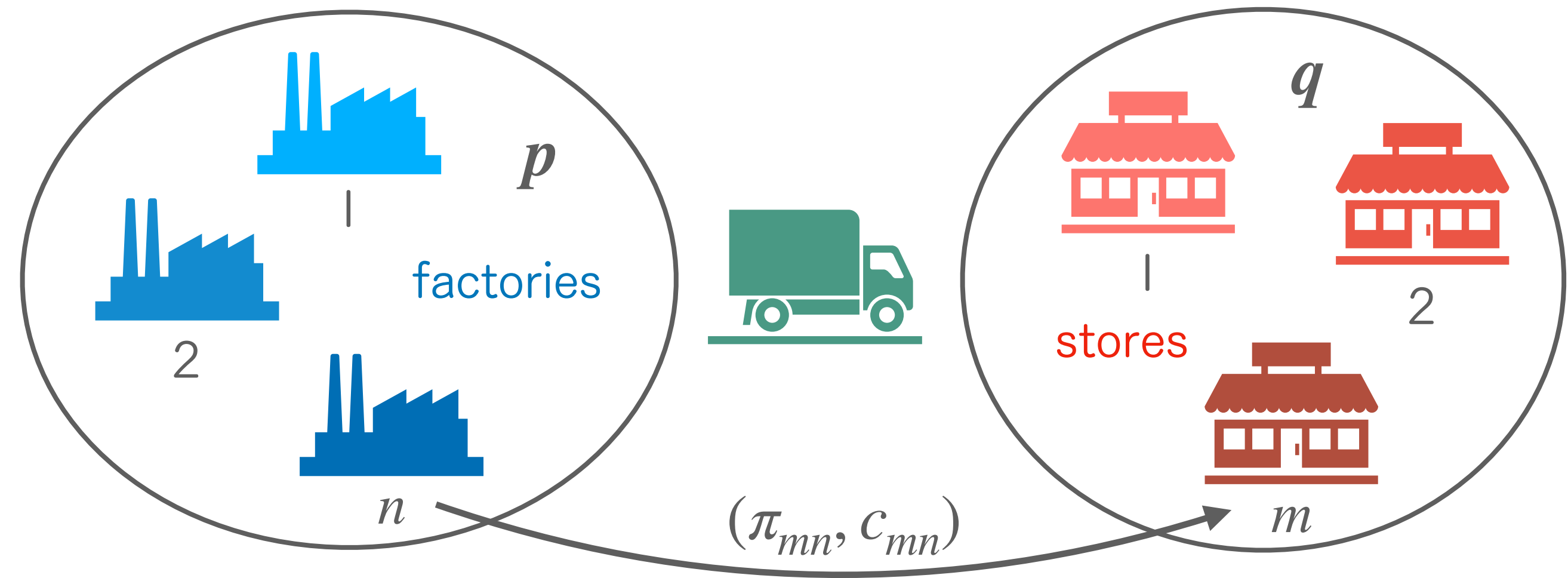
# Discrete optimal transport

- Transport of goods



# Discrete optimal transport

- Transport of goods



$\pi_{mn}$  : fraction of goods from factory  $n$  to store  $m$

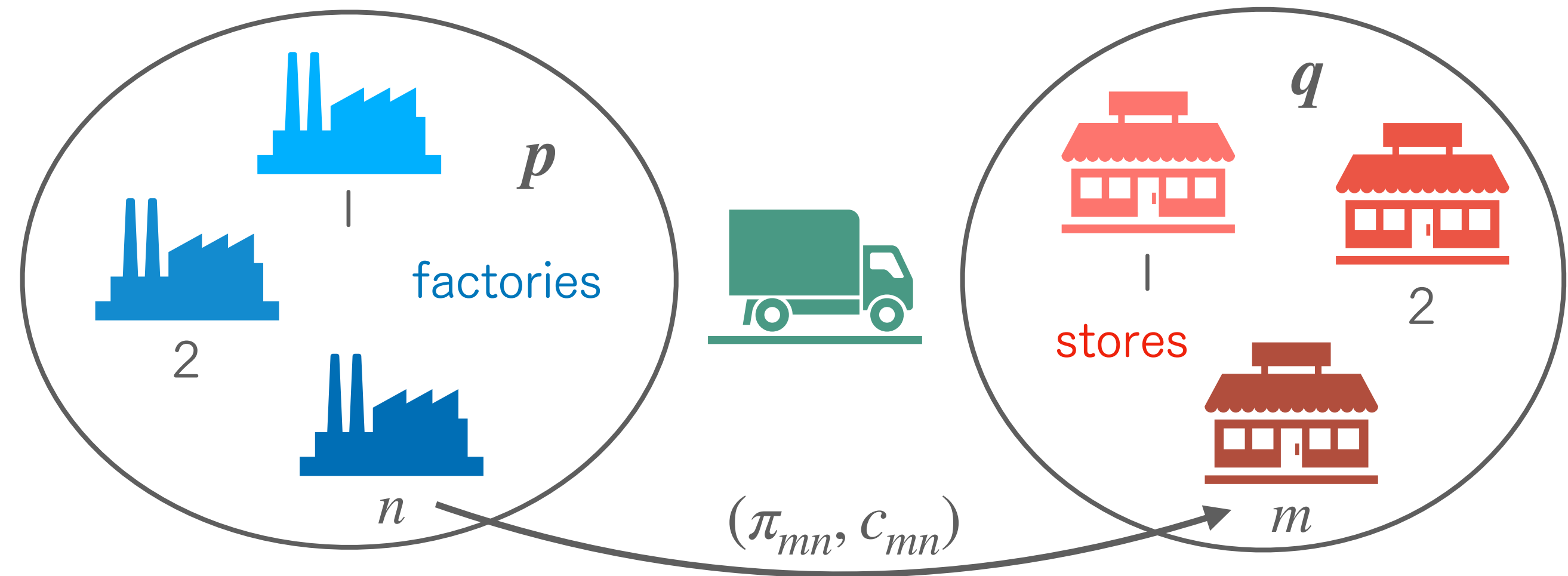
$c_{mn}$  : cost per unit of goods from factory  $n$  to store  $m$

# Discrete optimal transport

- Transport of goods

Total transport cost

$$\sum_{m,n} \pi_{mn} c_{mn}$$



$\pi_{mn}$  : fraction of goods from factory  $n$  to store  $m$

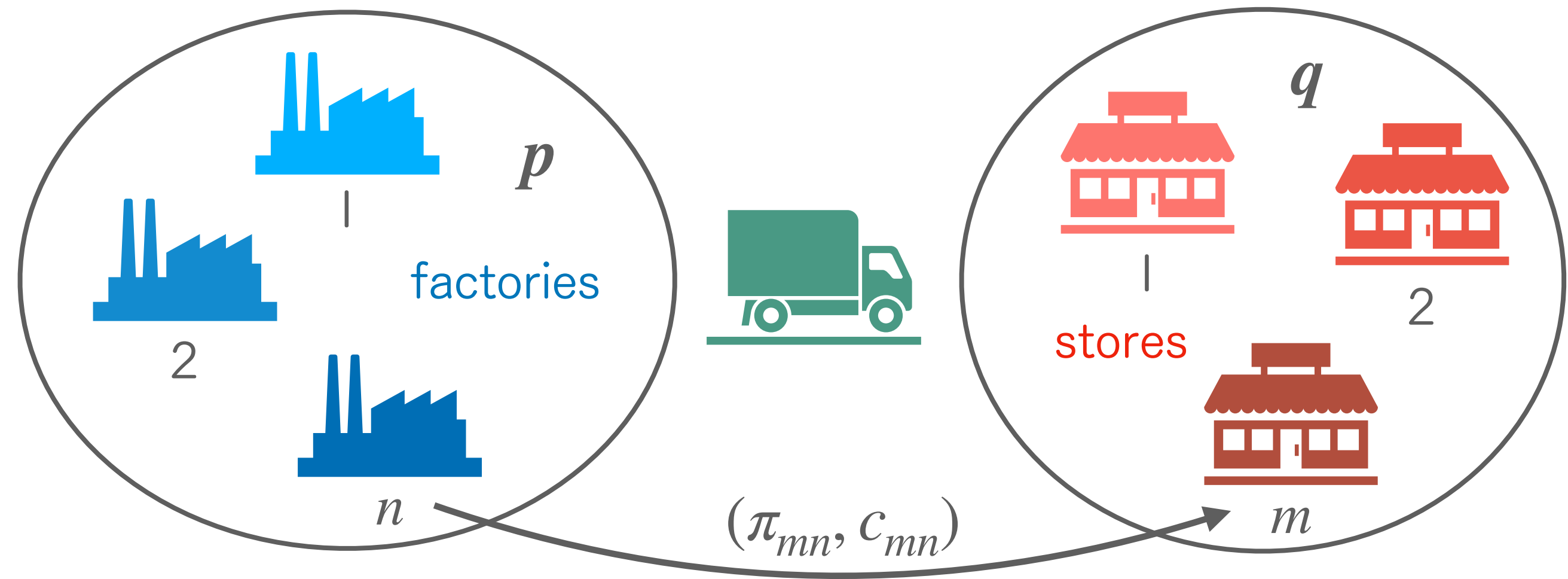
$c_{mn}$  : cost per unit of goods from factory  $n$  to store  $m$

# Discrete optimal transport

- Transport of goods

Total transport cost

$$\sum_{m,n} \pi_{mn} c_{mn}$$



$\pi_{mn}$  : fraction of goods from factory  $n$  to store  $m$

$c_{mn}$  : cost per unit of goods from factory  $n$  to store  $m$

- $L^1$ -Wasserstein distance

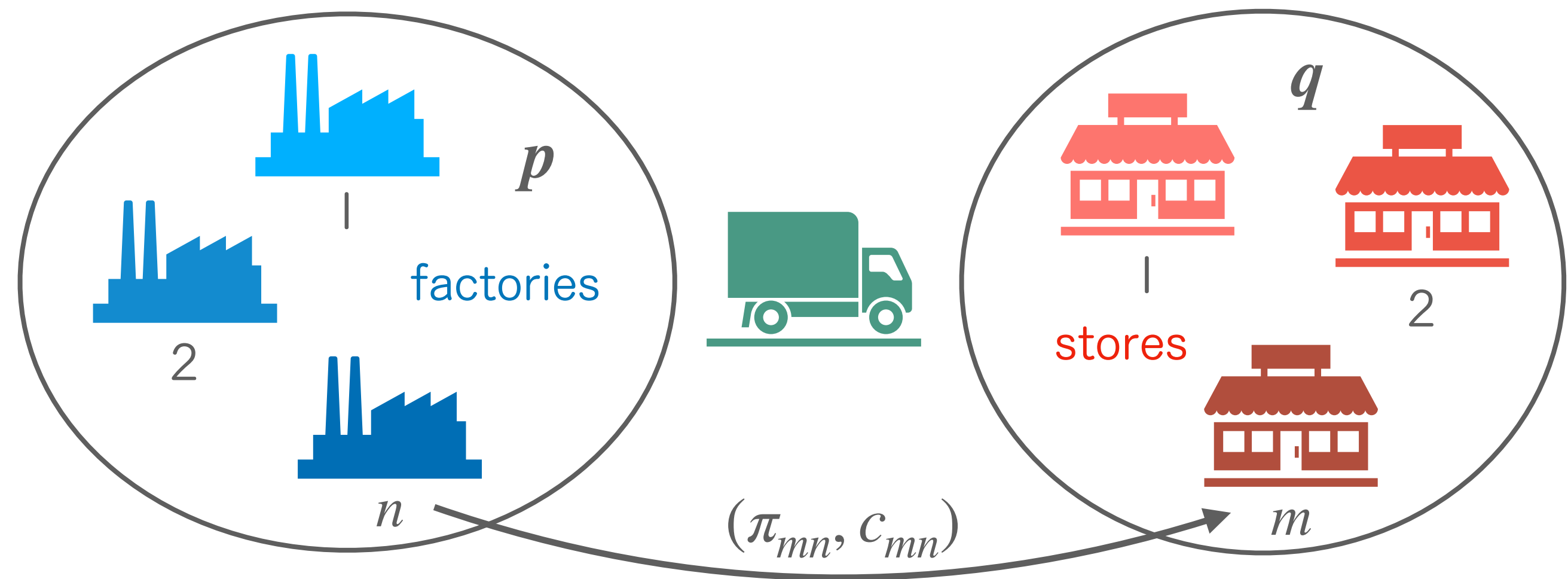
$$\mathcal{W}(p, q) := \min_{\pi} \sum_{m,n} \pi_{mn} c_{mn}$$

# Discrete optimal transport

- Transport of goods

Total transport cost

$$\sum_{m,n} \pi_{mn} c_{mn}$$



$\pi_{mn}$  : fraction of goods from factory  $n$  to store  $m$

$c_{mn}$  : cost per unit of goods from factory  $n$  to store  $m$

- $L^1$ -Wasserstein distance

$$\mathcal{W}(p, q) := \min_{\pi} \sum_{m,n} \pi_{mn} c_{mn}$$

$$\text{transport plan } \pi : \sum_m \pi_{mn} = p_n, \quad \sum_n \pi_{mn} = q_m$$

$$\text{cost matrix } [c_{mn}] : c_{mn} + c_{nk} \geq c_{mk}$$

# Kantorovich-Rubinstein duality

# Kantorovich-Rubinstein duality

- For symmetric cost matrix  $c_{mn} = c_{nm}$

$$\mathcal{W}(\boldsymbol{p}, \boldsymbol{q}) = \max_{\boldsymbol{\phi}} \boldsymbol{\phi}^\top (\boldsymbol{p} - \boldsymbol{q})$$

maximum is over all vectors  $\boldsymbol{\phi}$  satisfying  $|\phi_m - \phi_n| \leq c_{mn}$

# Kantorovich-Rubinstein duality

- For symmetric cost matrix  $c_{mn} = c_{nm}$

$$\mathcal{W}(\mathbf{p}, \mathbf{q}) = \max_{\boldsymbol{\phi}} \boldsymbol{\phi}^\top (\mathbf{p} - \mathbf{q})$$

maximum is over all vectors  $\boldsymbol{\phi}$  satisfying  $|\phi_m - \phi_n| \leq c_{mn}$

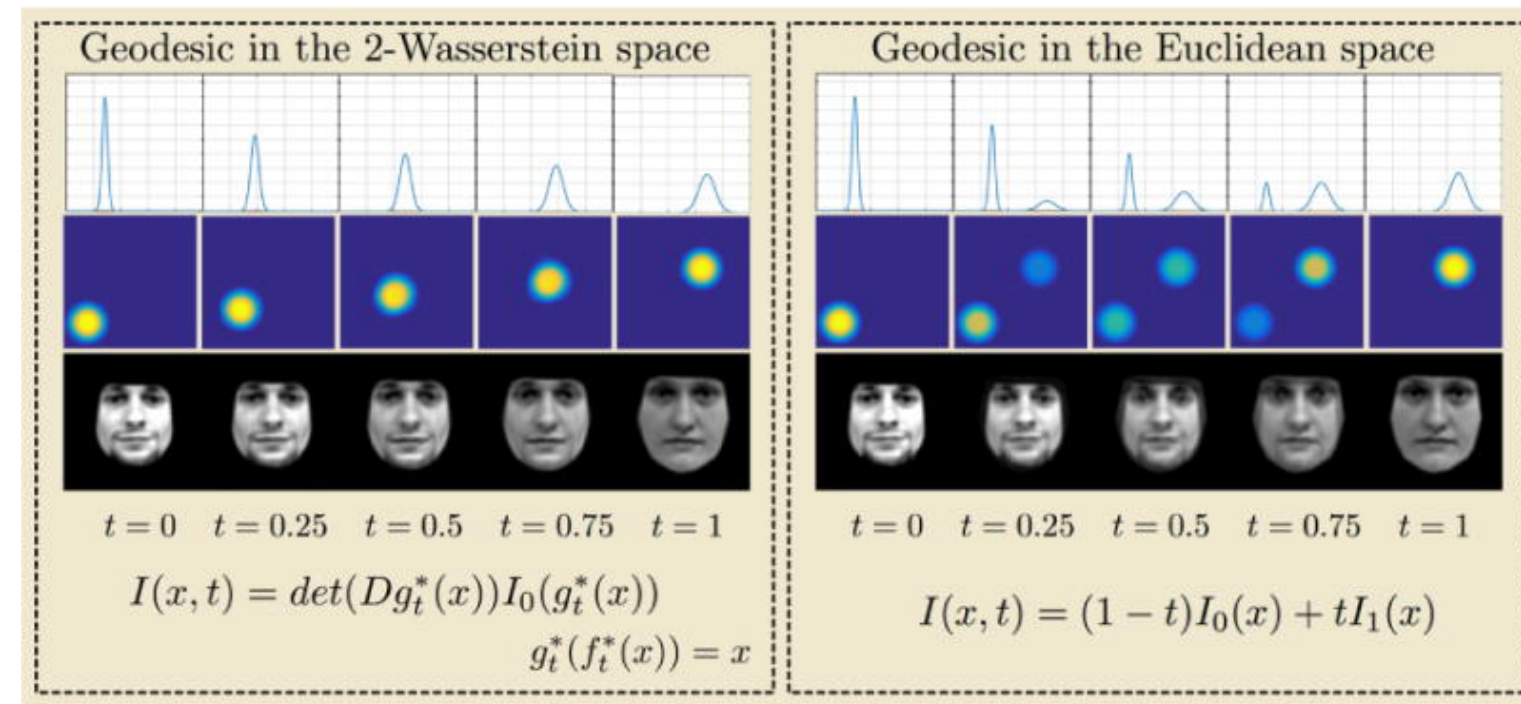
- For general cost matrix

$$\mathcal{W}(\mathbf{p}, \mathbf{q}) \leq \max_{\boldsymbol{\phi}} \boldsymbol{\phi}^\top (\mathbf{p} - \mathbf{q})$$

maximum is over all vectors  $\boldsymbol{\phi}$  satisfying  $\phi_m - \phi_n \leq c_{mn}$

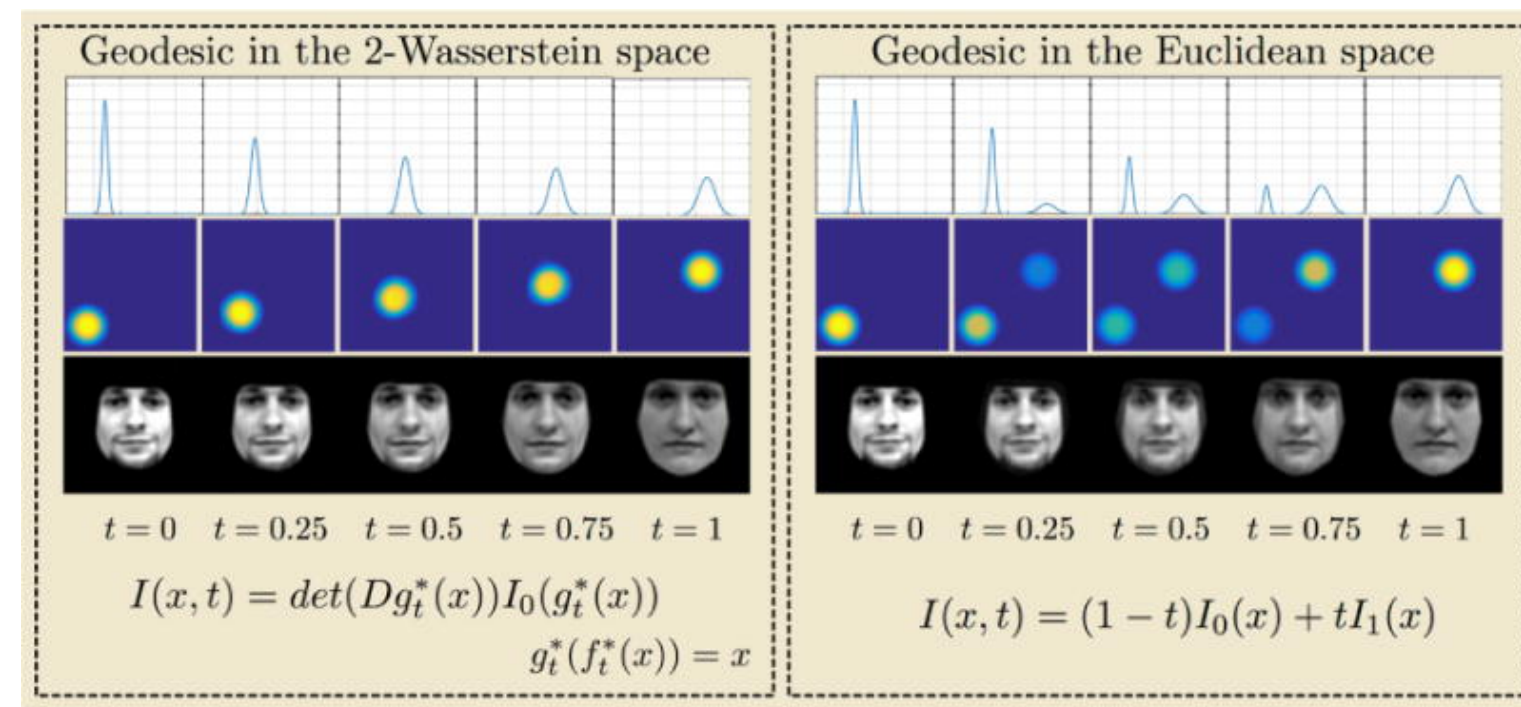
# Applications of optimal transport

# Applications of optimal transport

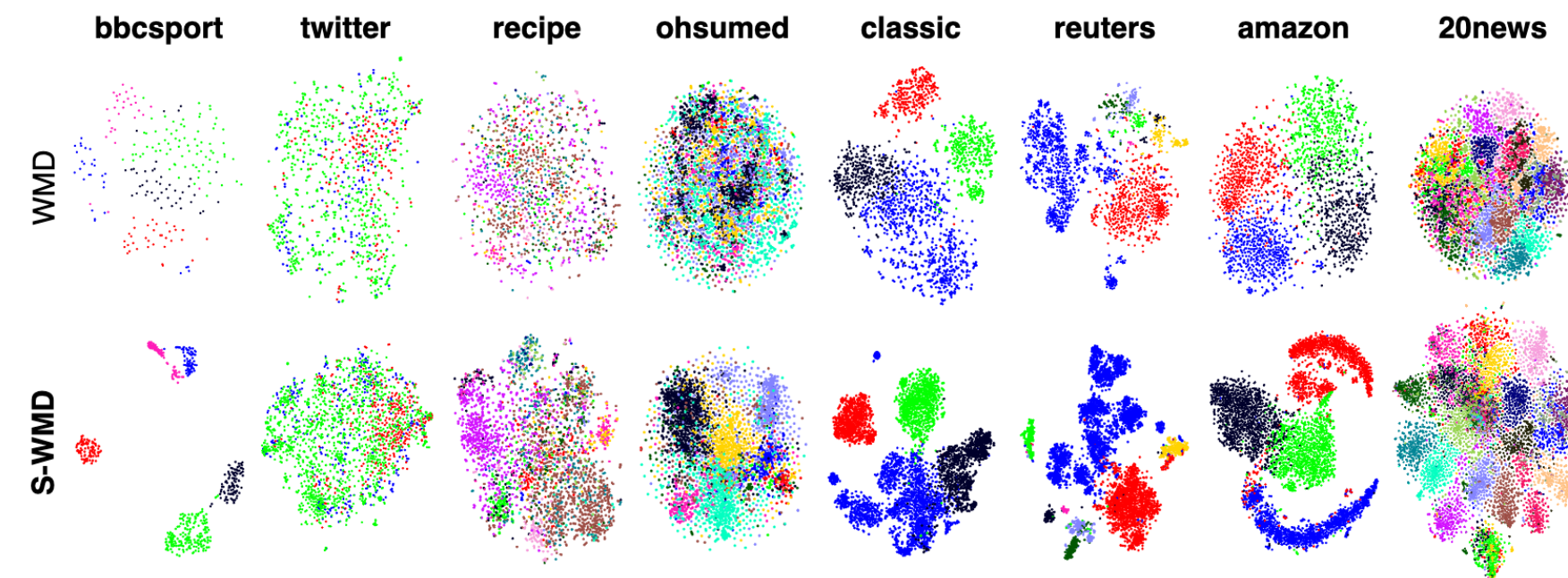


Kolouri+, IEEE Signal Process. Mag. (2017)

# Applications of optimal transport

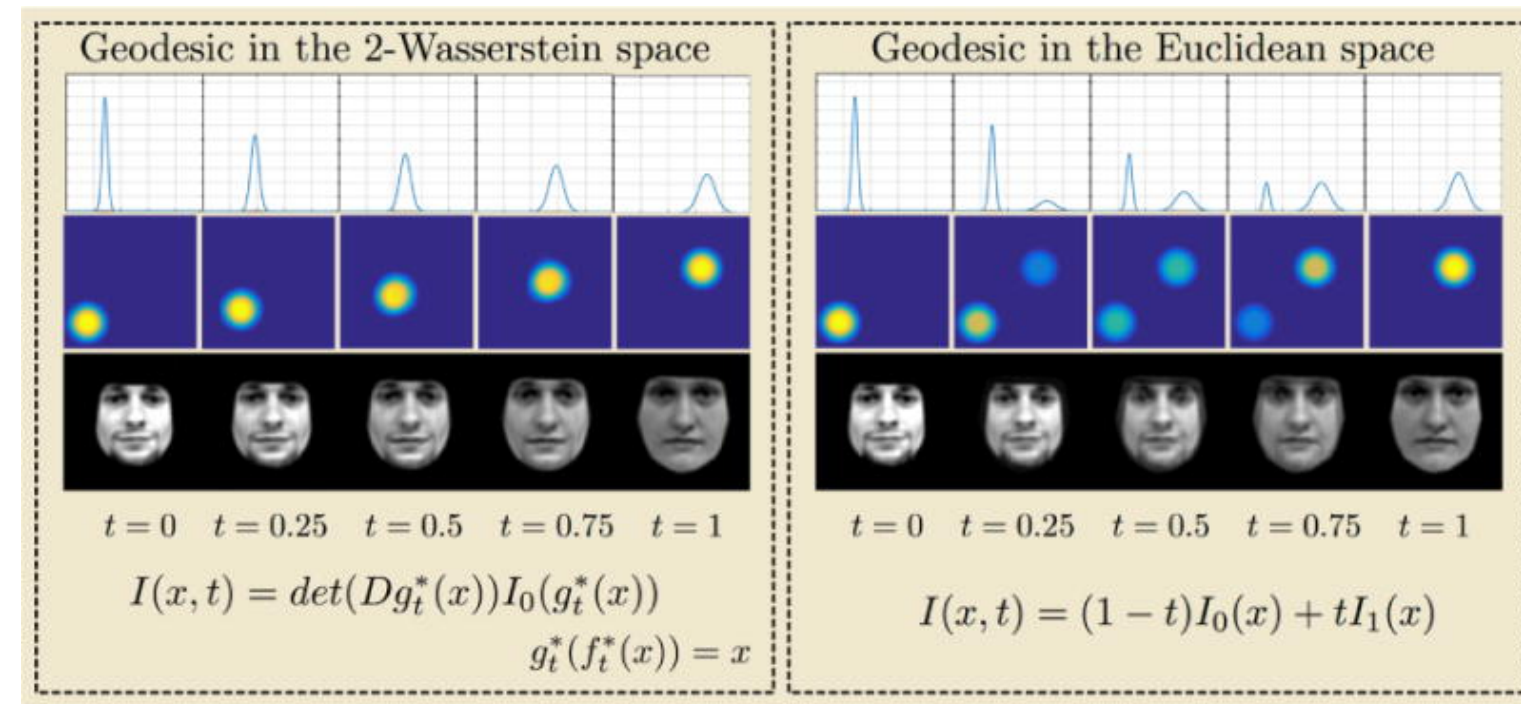


Kolouri+, IEEE Signal Process. Mag. (2017)

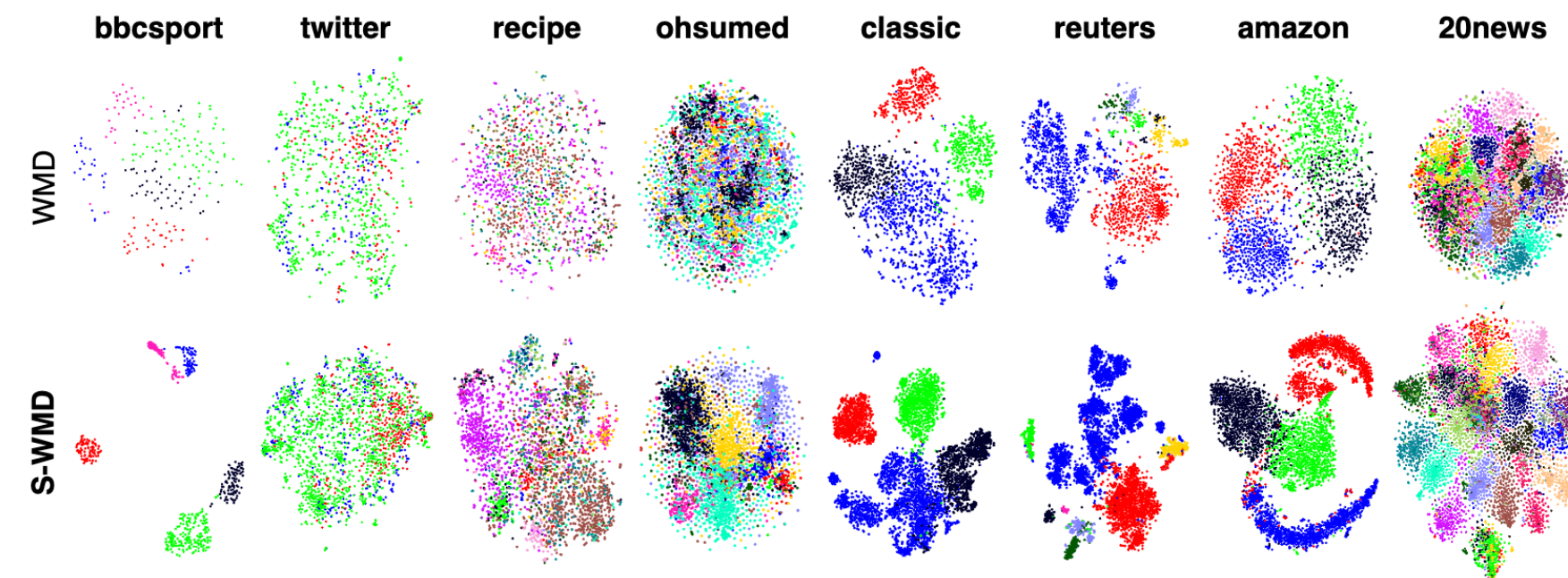


Huang+, NIPS (2016)

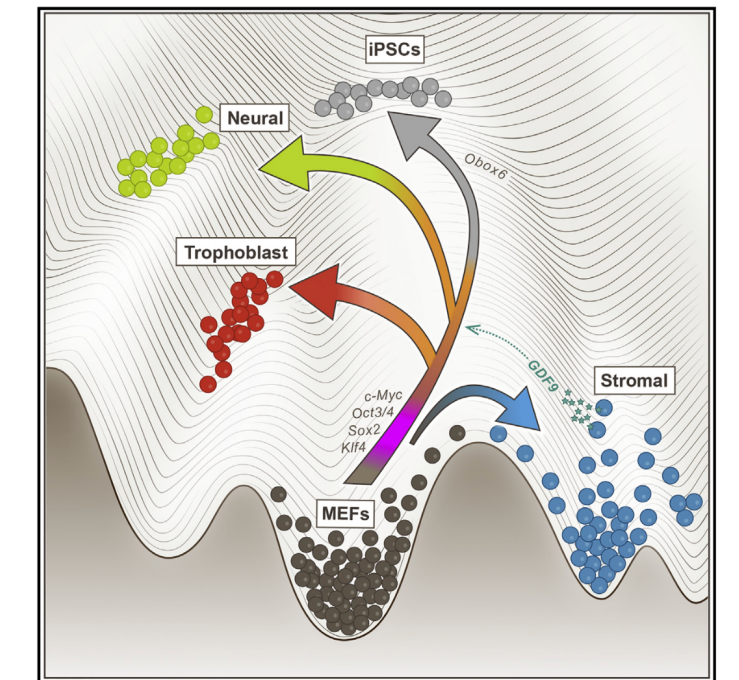
# Applications of optimal transport



Kolouri+, IEEE Signal Process. Mag. (2017)

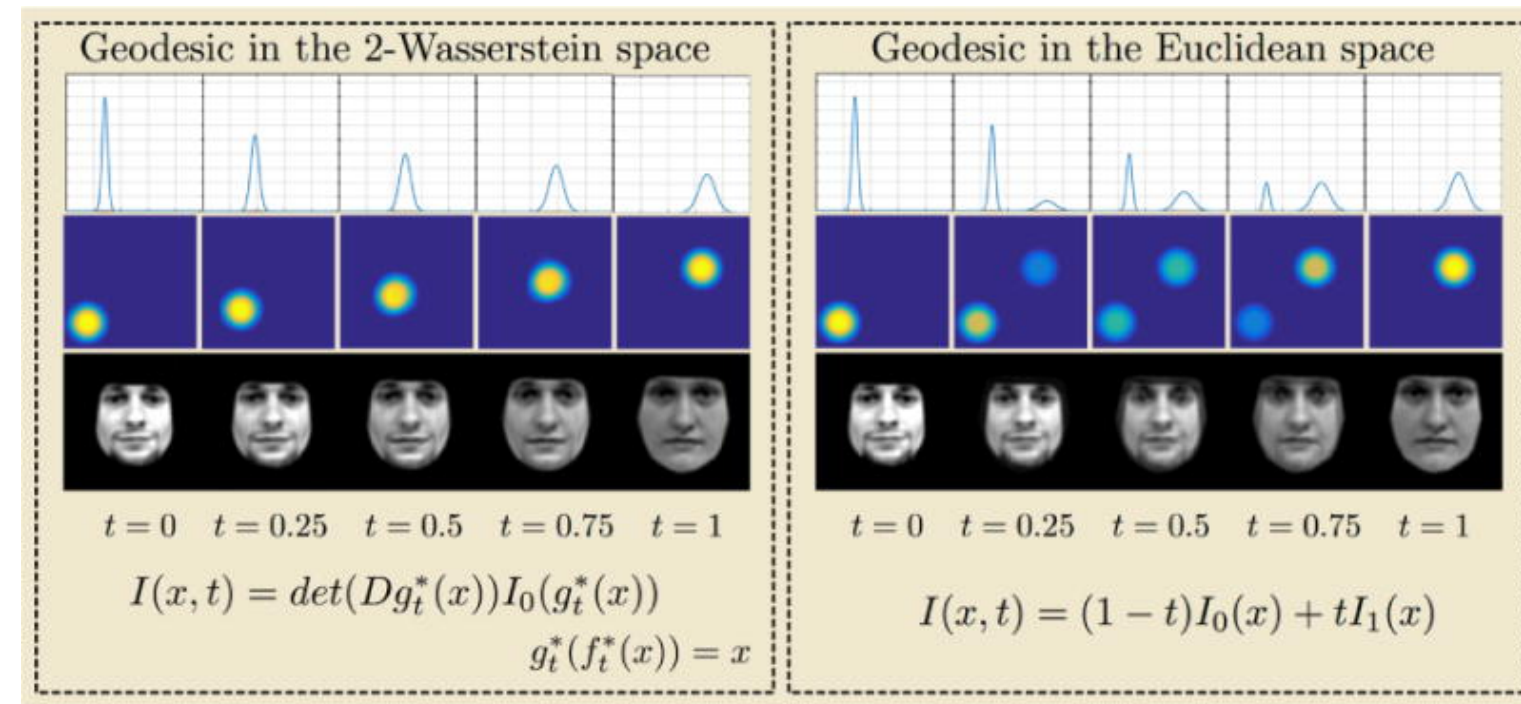


Huang+, NIPS (2016)

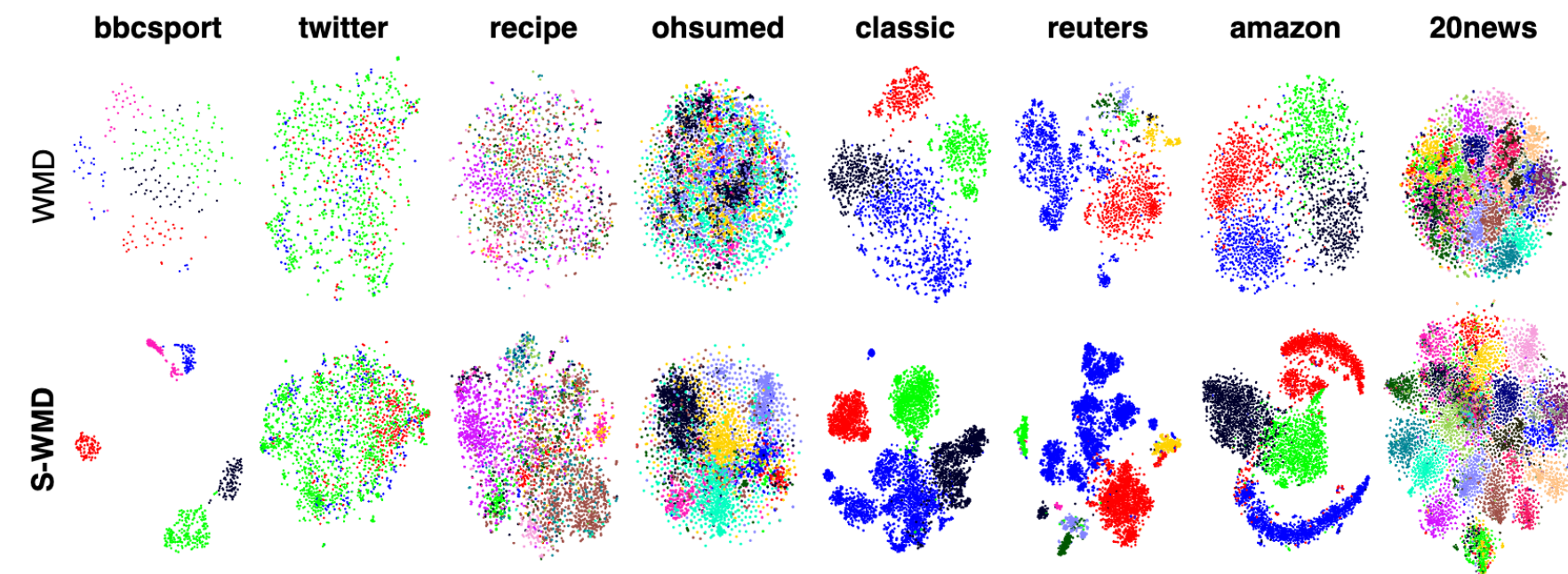


Schiebinger+, Cell (2019)

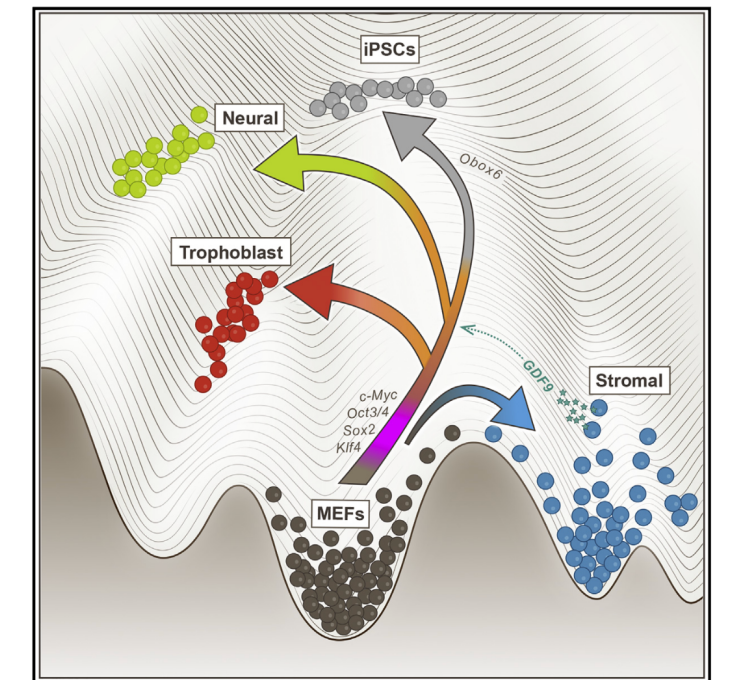
# Applications of optimal transport



Kolouri+, IEEE Signal Process. Mag. (2017)



Huang+, NIPS (2016)



Schiebinger+, Cell (2019)

PHYSICAL REVIEW LETTERS **128**, 201302 (2022)

Featured in Physics

## Accurate Baryon Acoustic Oscillations Reconstruction via Semidiscrete Optimal Transport

Sebastian von Hausegger<sup>1,2,3,\*</sup>, Bruno Lévy<sup>2</sup>, and Roya Mohayaee<sup>1,3</sup>

<sup>1</sup>Rudolf Peierls Centre for Theoretical Physics, University of Oxford, Parks Road, Oxford OX1 3PU, United Kingdom

<sup>2</sup>Université de Lorraine, CNRS, Inria, LORIA, F-54000 Nancy, France

<sup>3</sup>Institut d'Astrophysique de Paris, CNRS, Sorbonne Université, 98bis Bld Arago, 75014 Paris, France

# Optimal transport and thermodynamics

PHYSICAL REVIEW X **13**, 011013 (2023)

## Thermodynamic Unification of Optimal Transport: Thermodynamic Uncertainty Relation, Minimum Dissipation, and Thermodynamic Speed Limits

Tan Van Vu<sup>\*</sup> and Keiji Saito<sup>†</sup>

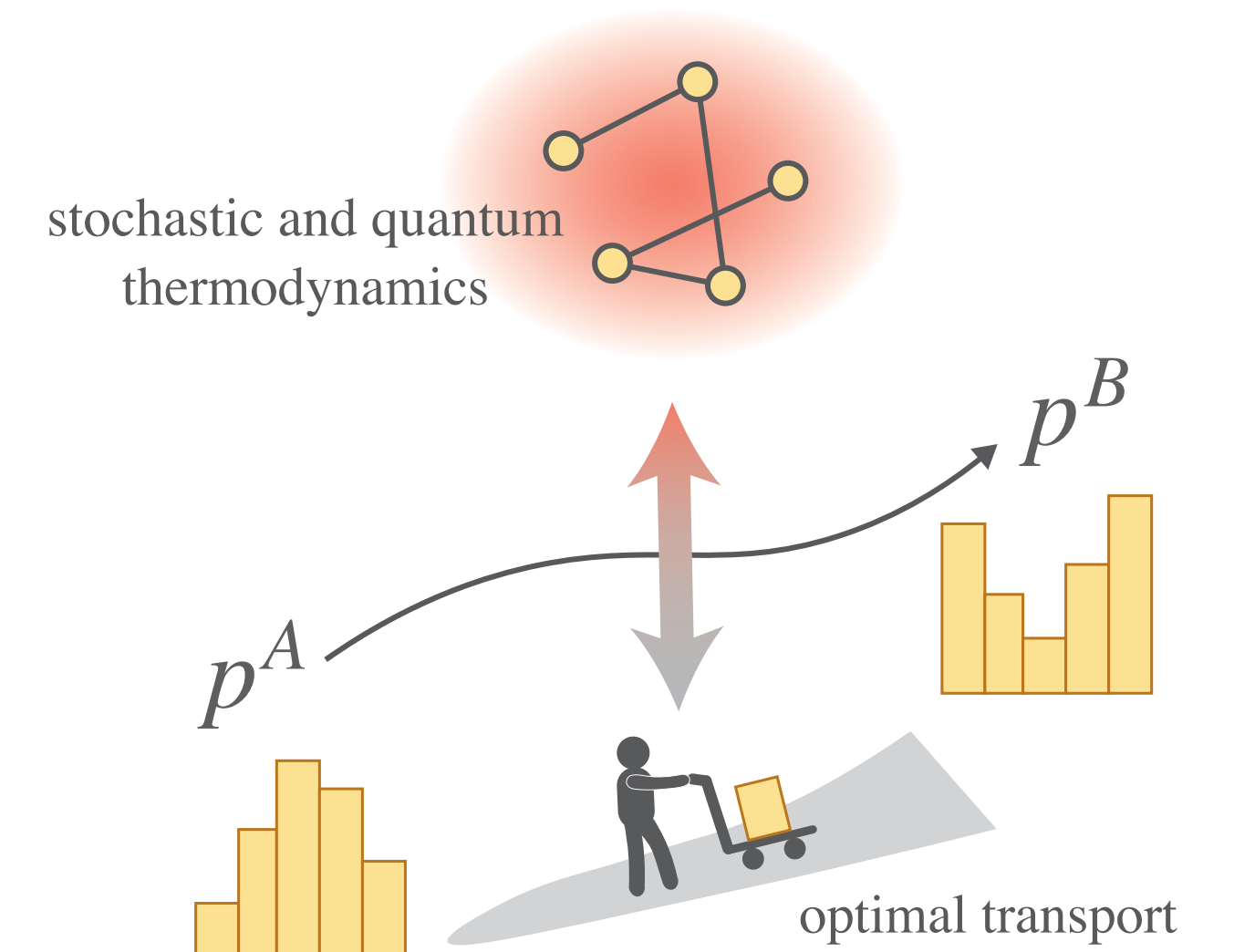
*Department of Physics, Keio University, 3-14-1 Hiyoshi, Kohoku-ku, Yokohama 223-8522, Japan*

 (Received 26 July 2022; revised 1 December 2022; accepted 13 December 2022; published 3 February 2023)

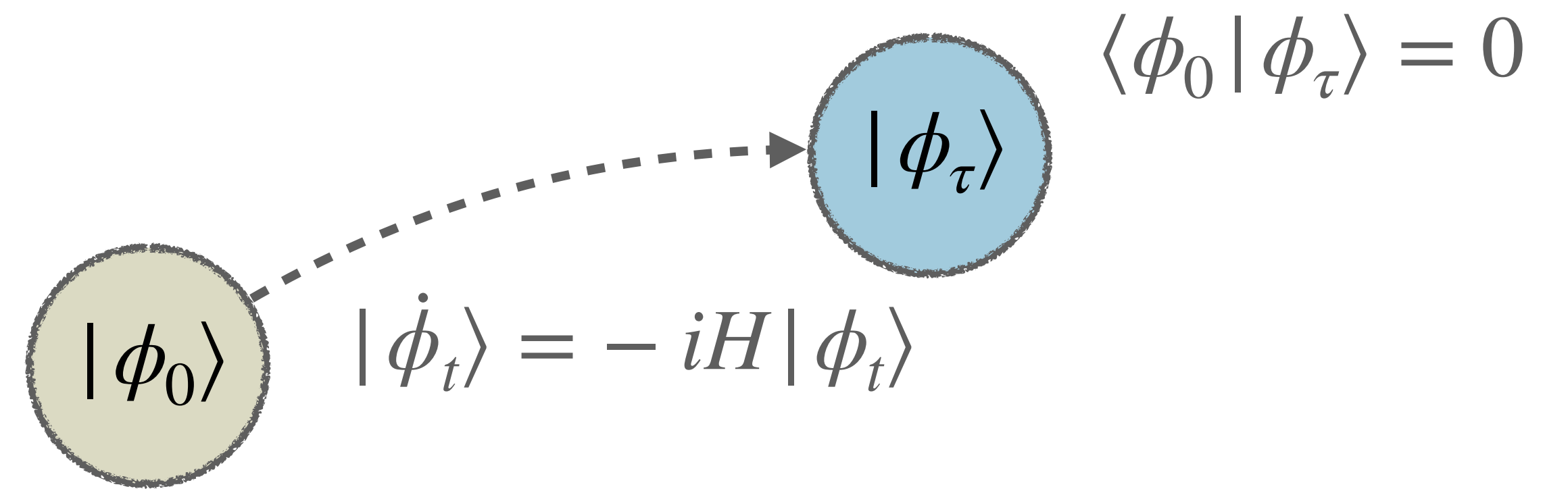
Discrete generalization of Benamou-Brenier formula

$$\mathcal{W}(p, q) = \min \sqrt{\Sigma_{\tau} \mathcal{M}_{\tau}}$$

$\Sigma_{\tau}$  : entropy production  
 $\mathcal{M}_{\tau}$  : state mobility



# Quantum speed limit



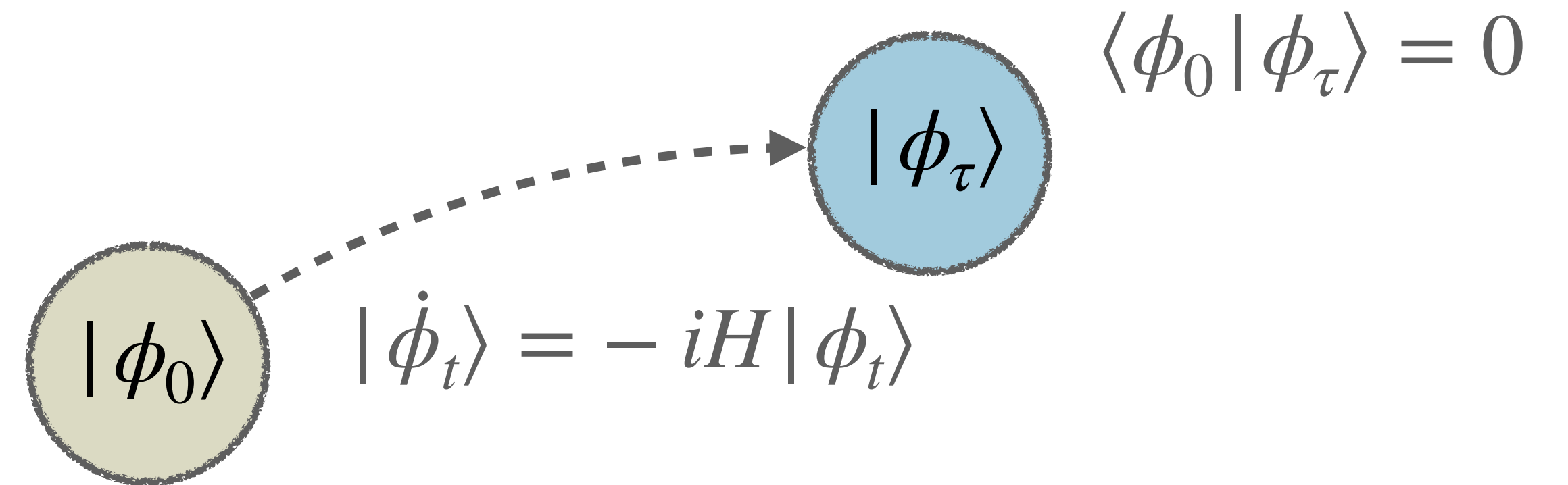
# Quantum speed limit

- Minimum time for evolution between orthogonal states

$$\tau \geq \frac{\pi}{2} \max \left\{ \frac{1}{\Delta H}, \frac{1}{\langle H \rangle} \right\}$$

Mandelstam and Tamm, J. Phys. USSR (1945)

Margolus and Levitin, Physica (1998)



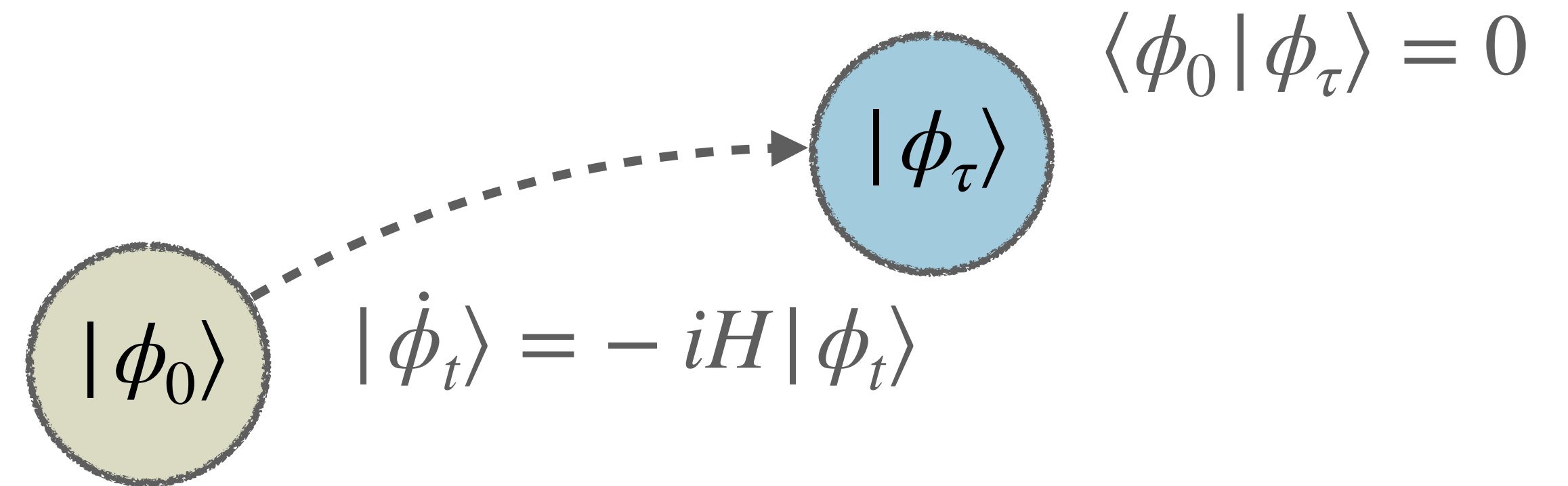
# Quantum speed limit

- Minimum time for evolution between orthogonal states

$$\tau \geq \frac{\pi}{2} \max \left\{ \frac{1}{\Delta H}, \frac{1}{\langle H \rangle} \right\}$$

Mandelstam and Tamm, J. Phys. USSR (1945)

Margolus and Levitin, Physica (1998)



- Generalizations to arbitrary states and other dynamics

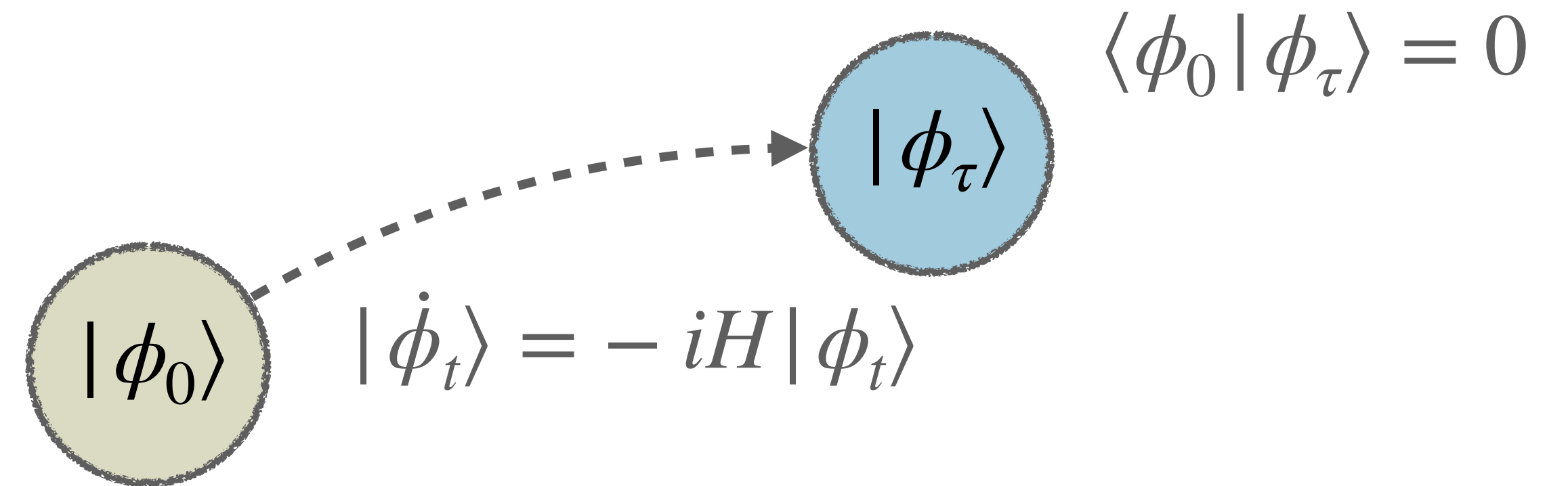
# Quantum speed limit

- Minimum time for evolution between orthogonal states

$$\tau \geq \frac{\pi}{2} \max \left\{ \frac{1}{\Delta H}, \frac{1}{\langle H \rangle} \right\}$$

Mandelstam and Tamm, J. Phys. USSR (1945)

Margolus and Levitin, Physica (1998)



- Generalizations to arbitrary states and other dynamics

$$\tau \geq \frac{\mathcal{L}(\varrho_0, \varrho_\tau)}{\bar{v}}$$

$\mathcal{L}$  : distance metric (Bures angle, Fisher information, etc.)

$\bar{v}$  : maximal average velocity

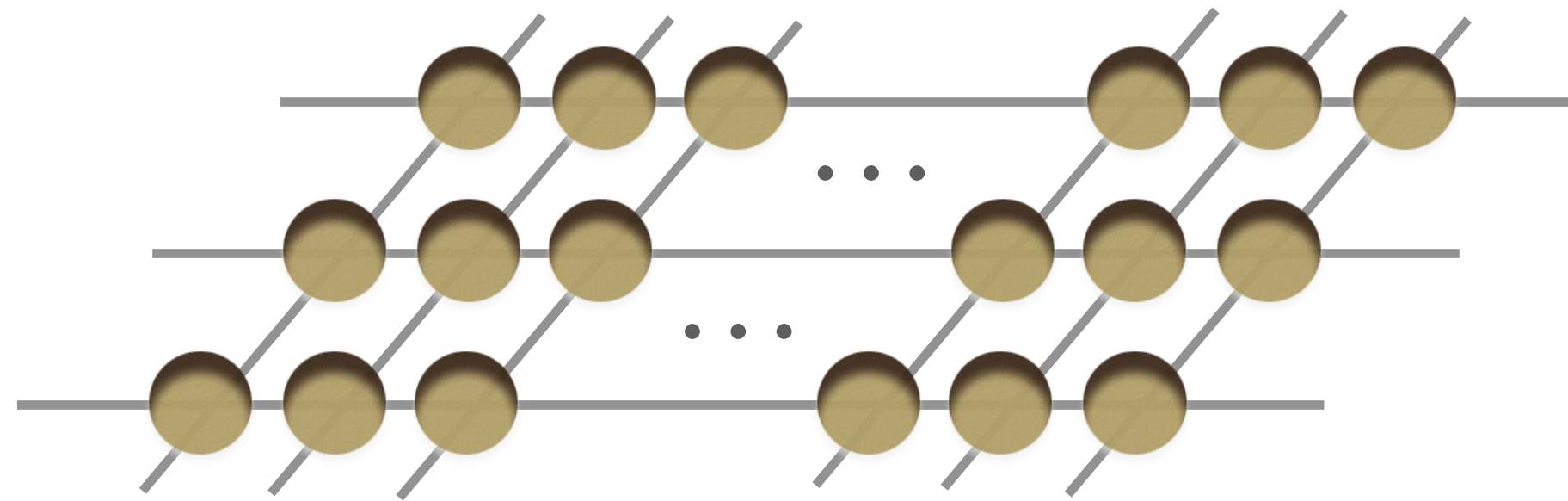
Deffner and Campbell, J. Phys. A (2017) (review)

# Strength and weakness

# Strength and weakness

$$\tau \geq \frac{\mathcal{L}(\varrho_0, \varrho_\tau)}{\bar{v}}$$

$$\mathcal{L}(\varrho_0, \varrho_\tau) = O(1)$$

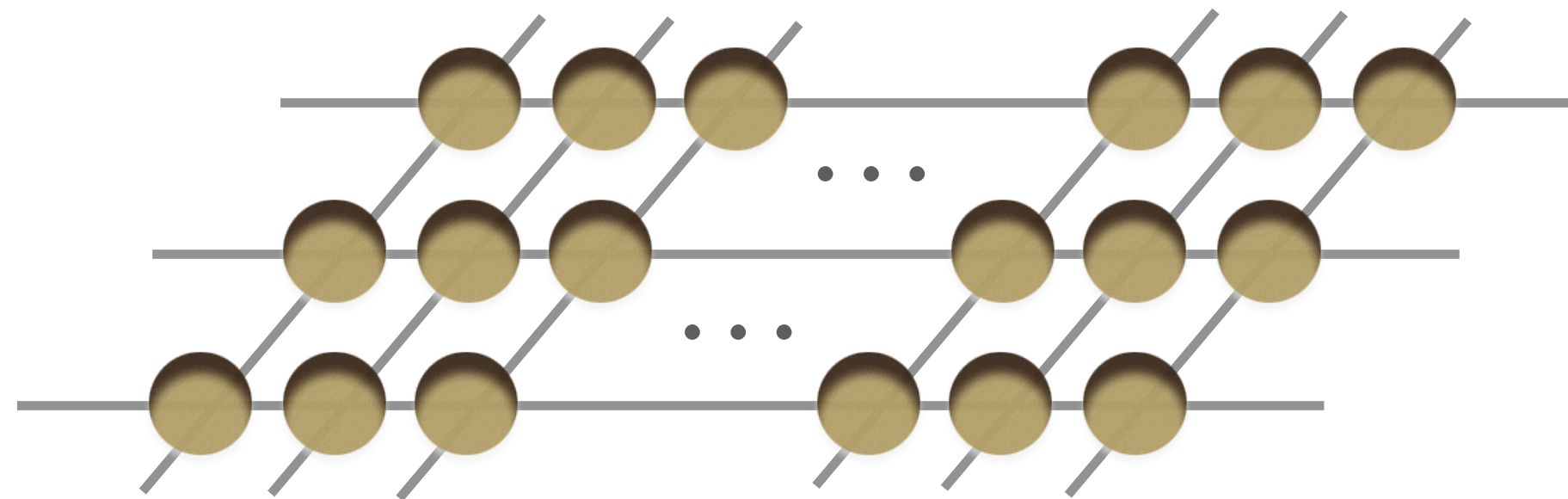


- ✓ speed limit between **specific** states
- × **not tight** for many-body systems

# Strength and weakness

$$\tau \geq \frac{\mathcal{L}(\varrho_0, \varrho_\tau)}{\bar{v}}$$

$$\mathcal{L}(\varrho_0, \varrho_\tau) = O(1)$$



- ✓ speed limit between **specific** states
- × **not tight** for many-body systems

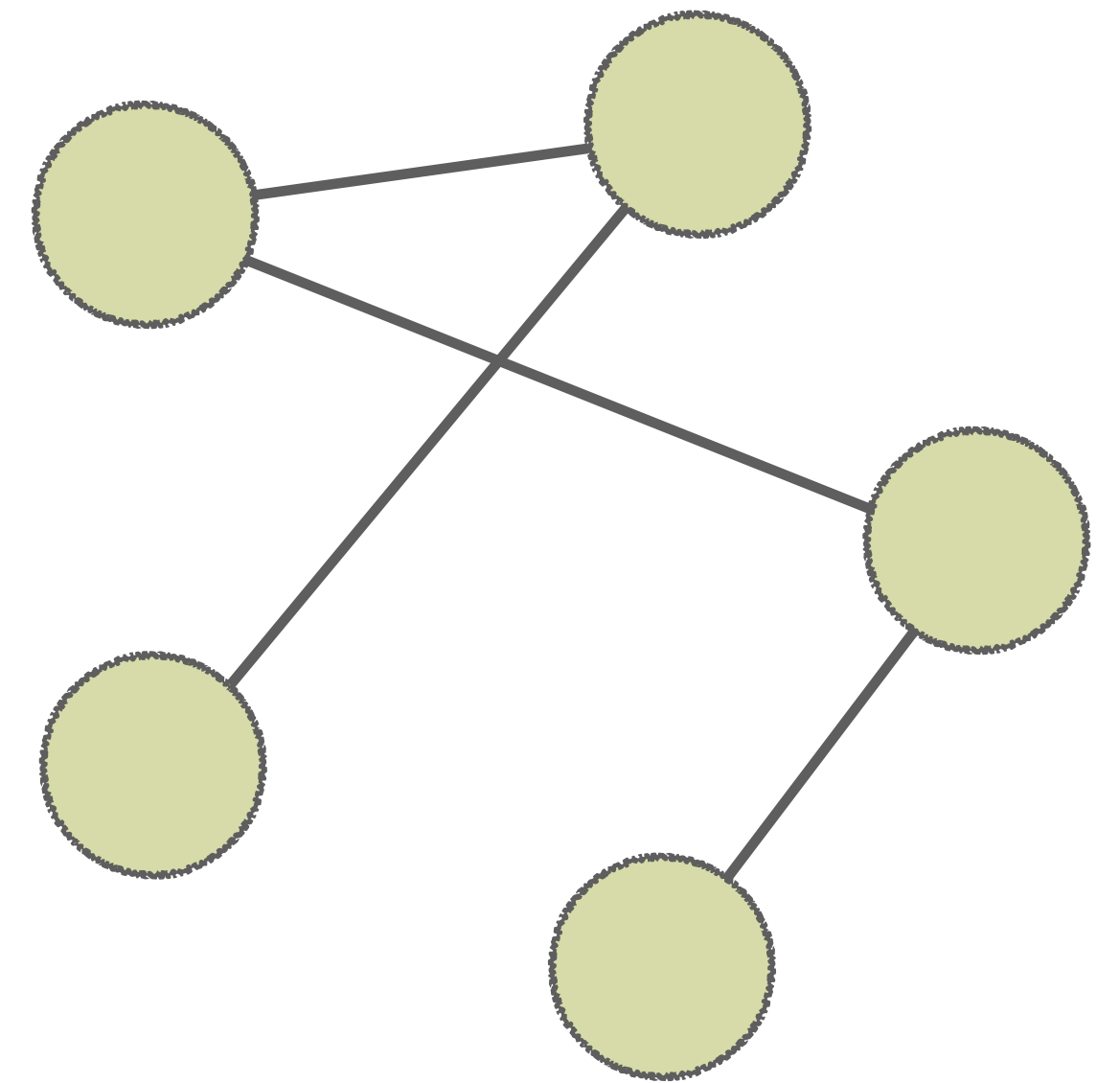
→ We need a speed limit that can incorporate the structure of many-body systems

# General speed limit

- Undirected graph  $G(V, E)$

$$V = \{1, \dots, N\}$$

$$E = \{(i, j) \mid i \text{ \& } j : \text{neighboring nodes}\}$$



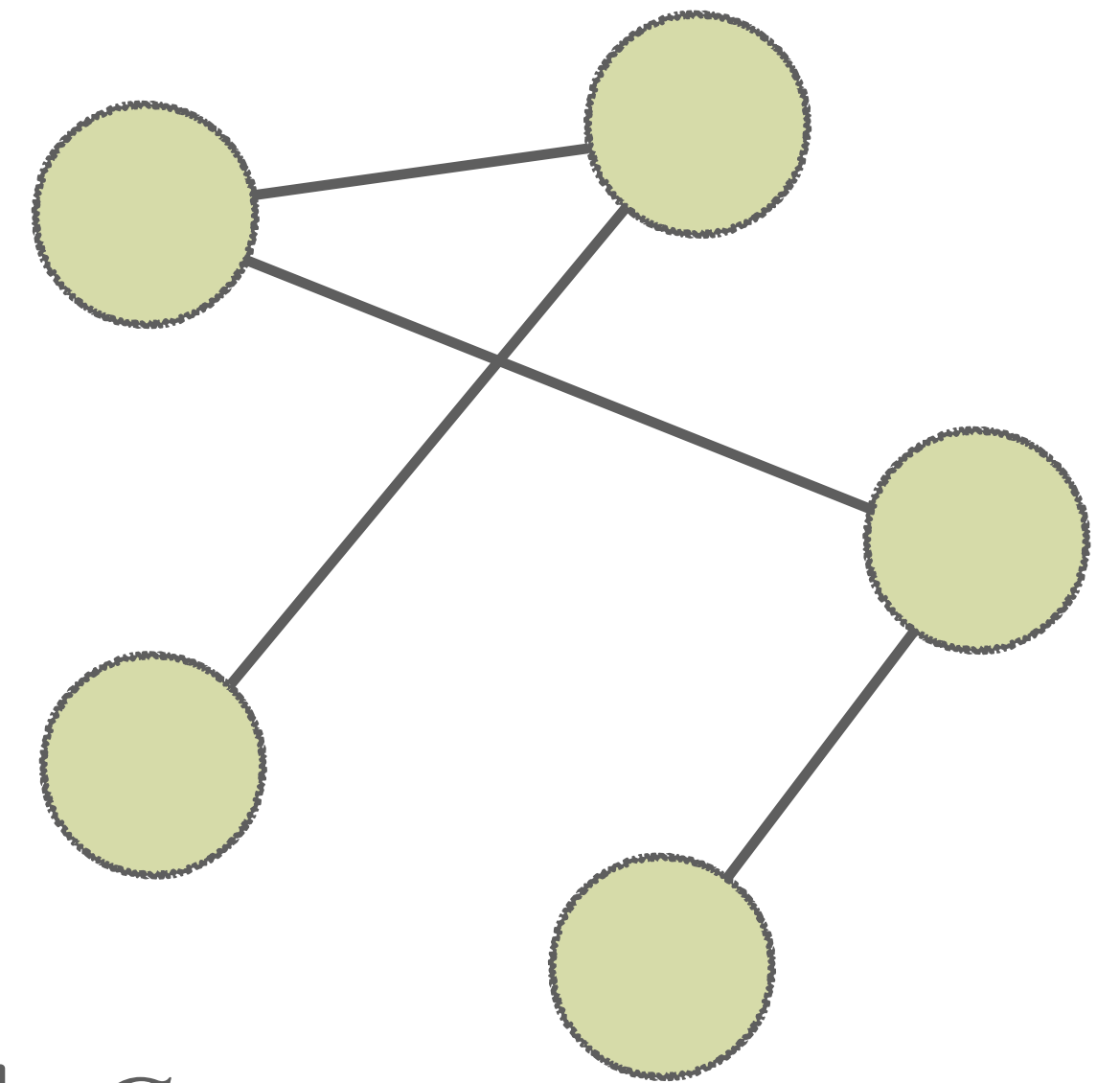
# General speed limit

- Undirected graph  $G(V, E)$

$$V = \{1, \dots, N\}$$

$$E = \{(i, j) \mid i \text{ \& } j : \text{neighboring nodes}\}$$

- Vector state  $\mathbf{x}_t = [x_1(t), \dots, x_N(t)]^\top \in \mathbb{R}_{\geq 0}^N$  evolves on graph  $G$



# General speed limit

- Undirected graph  $G(V, E)$

$$V = \{1, \dots, N\}$$

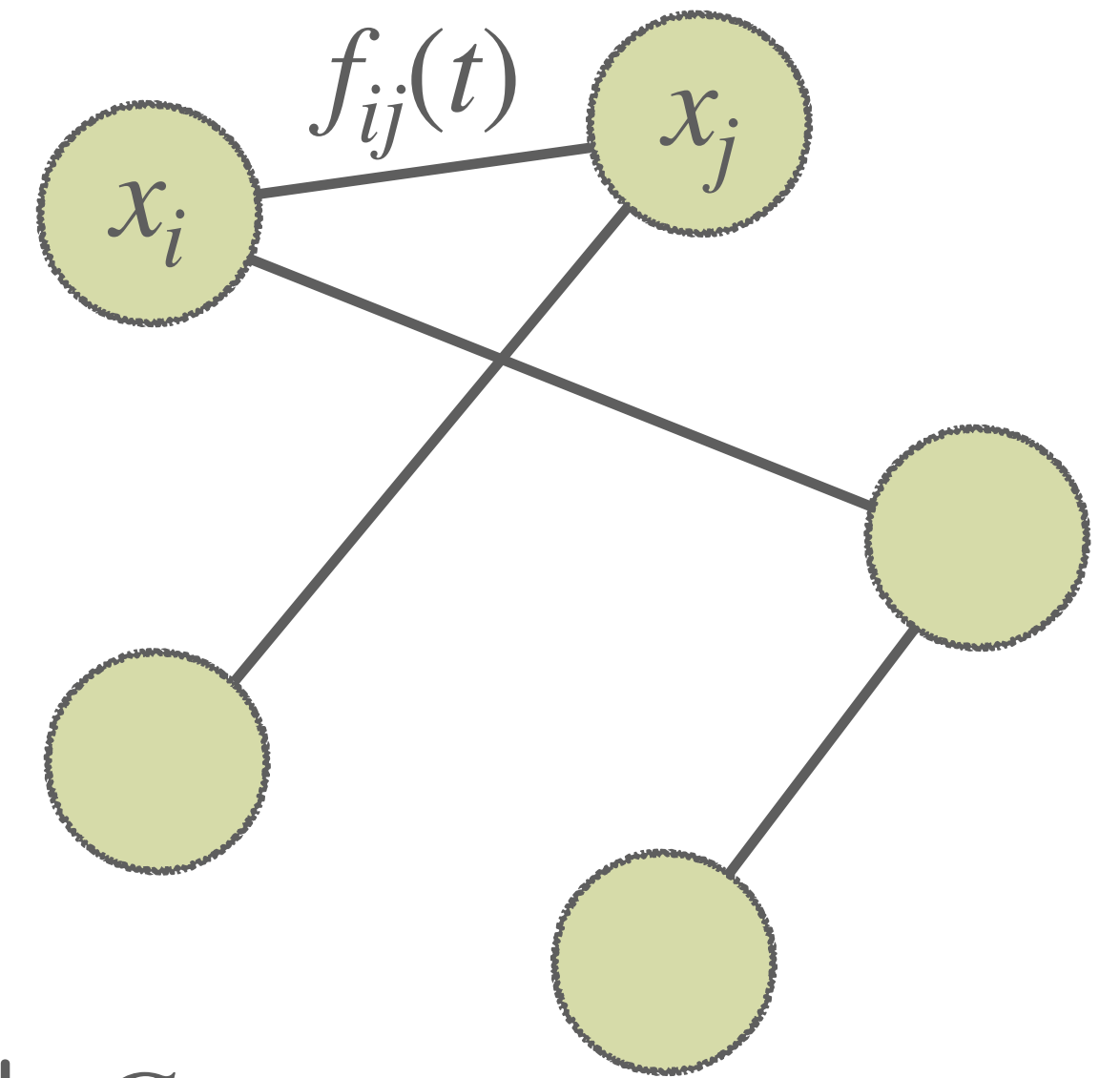
$$E = \{(i, j) \mid i \text{ \& } j : \text{neighboring nodes}\}$$

- Vector state  $\mathbf{x}_t = [x_1(t), \dots, x_N(t)]^\top \in \mathbb{R}_{\geq 0}^N$  evolves on graph  $G$

$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

$\mathcal{B}_i$  : neighboring nodes of  $i$

$f_{ij}(t) = -f_{ji}(t)$  : anti-symmetric flow



# General speed limit

- Undirected graph  $G(V, E)$

$$V = \{1, \dots, N\}$$

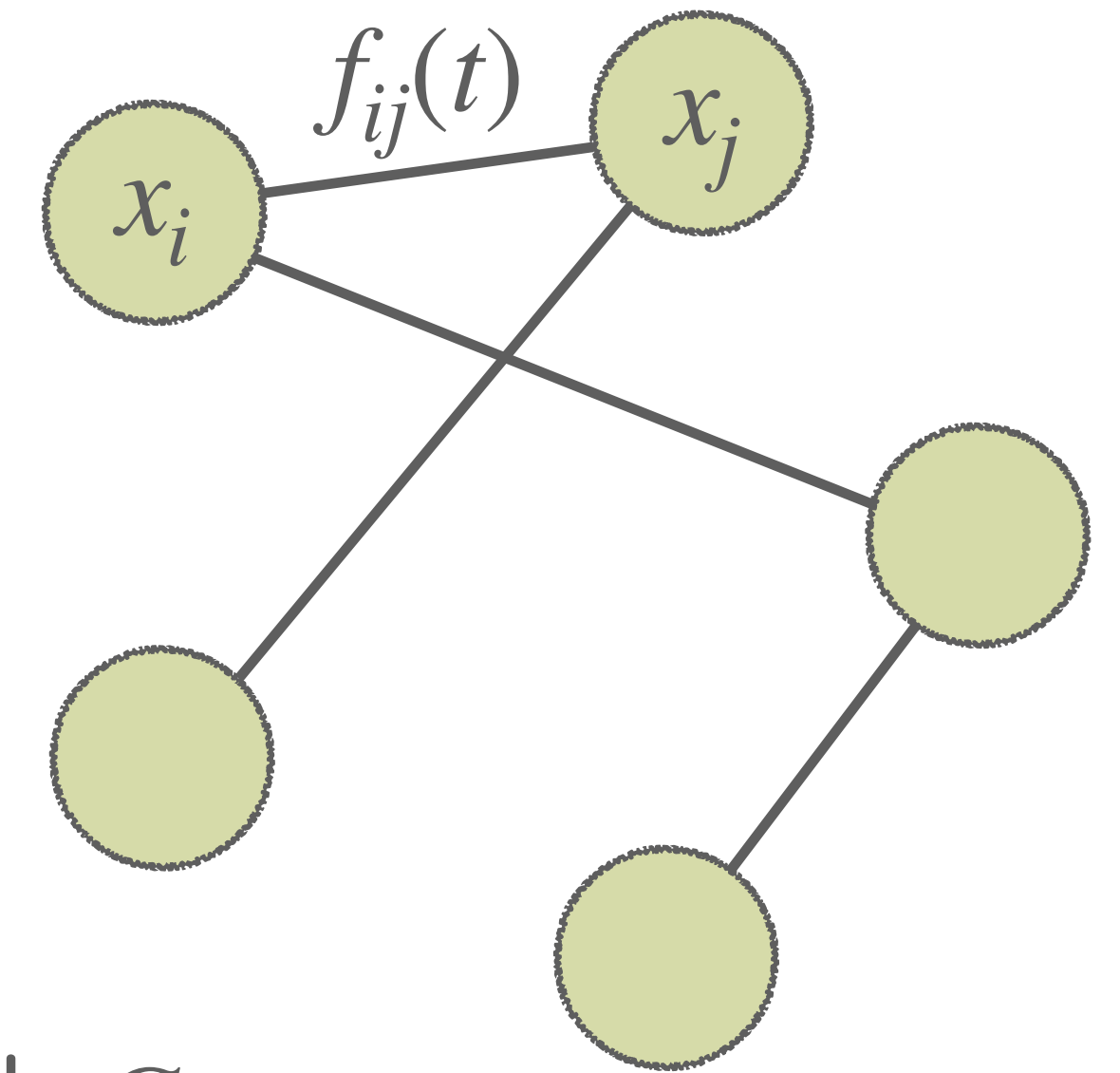
$$E = \{(i, j) \mid i \text{ \& } j : \text{neighboring nodes}\}$$

- Vector state  $\mathbf{x}_t = [x_1(t), \dots, x_N(t)]^\top \in \mathbb{R}_{\geq 0}^N$  evolves on graph  $G$

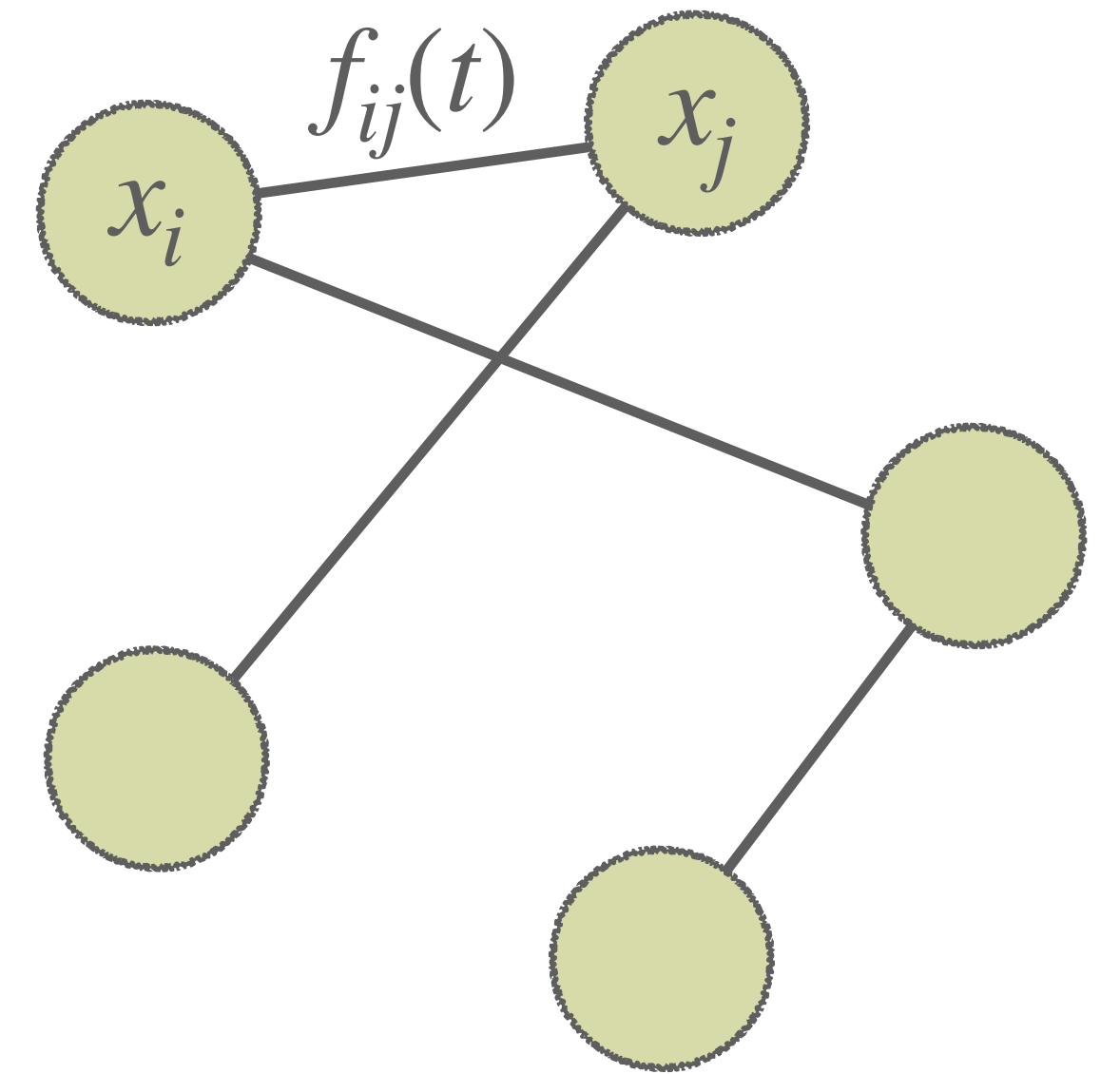
$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

$\mathcal{B}_i$  : neighboring nodes of  $i$

$$f_{ij}(t) = -f_{ji}(t) : \text{anti-symmetric flow} \rightarrow \sum_{i=1}^N x_i(t) = \text{const}$$



# General speed limit



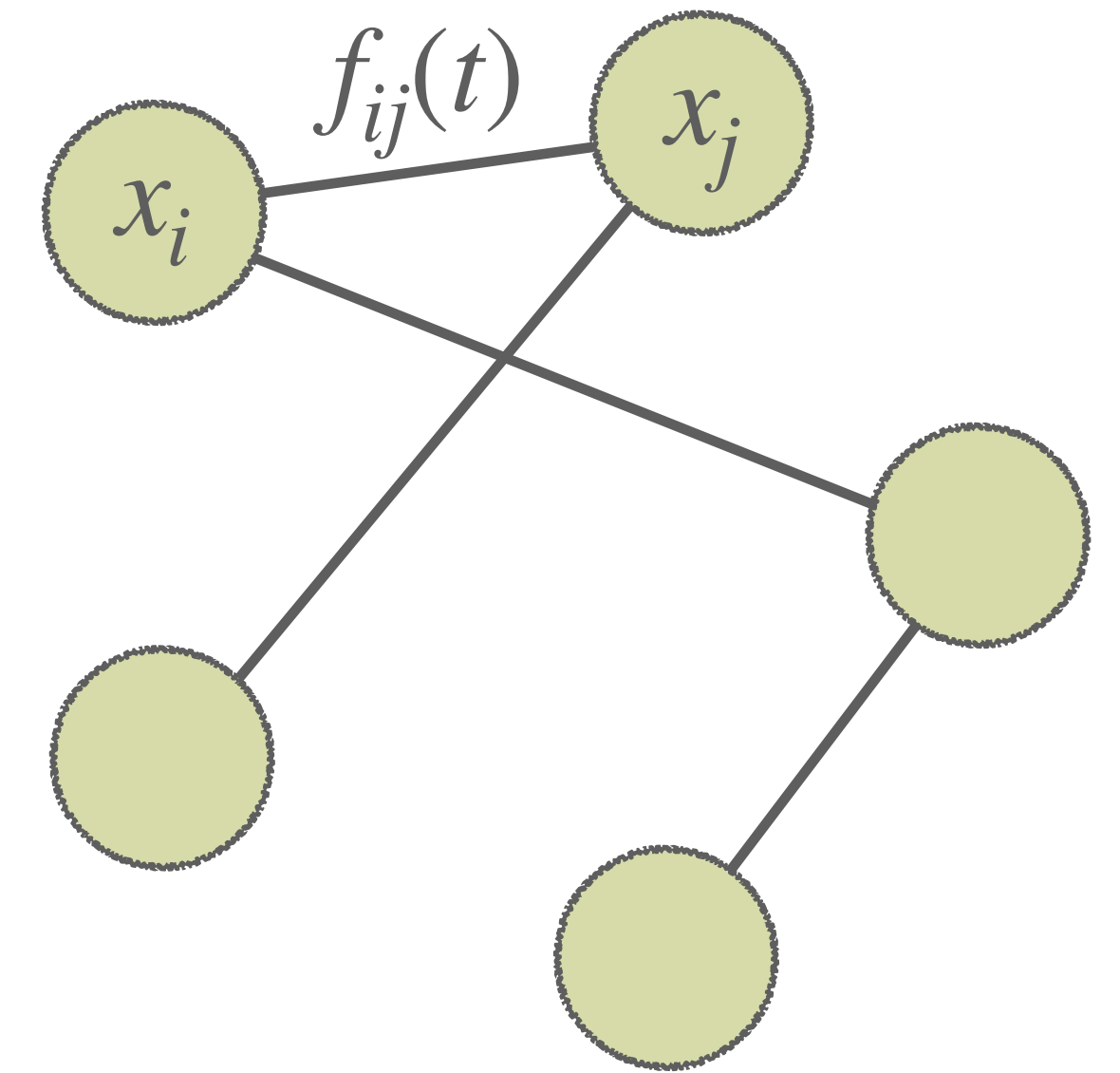
$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

# General speed limit

- Minimum time required for changing states

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

$$\bar{v} := \tau^{-1} \int_0^\tau dt \sum_{(i,j) \in E} c_{ij} |f_{ij}(t)| : \text{average velocity}$$



$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

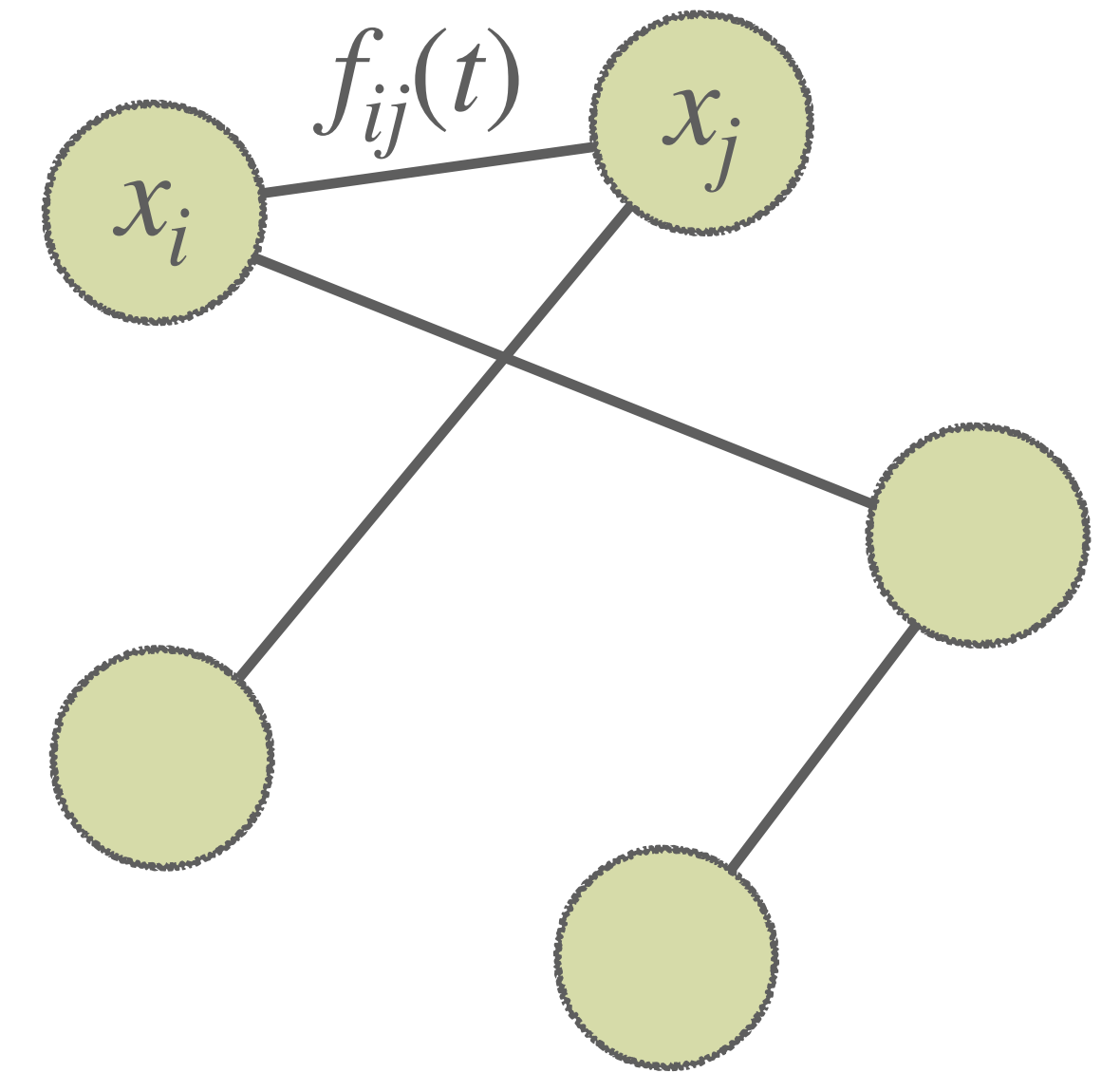
# General speed limit

- Minimum time required for changing states

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

$$\bar{v} := \tau^{-1} \int_0^\tau dt \sum_{(i,j) \in E} c_{ij} |f_{ij}(t)| : \text{average velocity}$$

- ✓ tight and saturable



$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

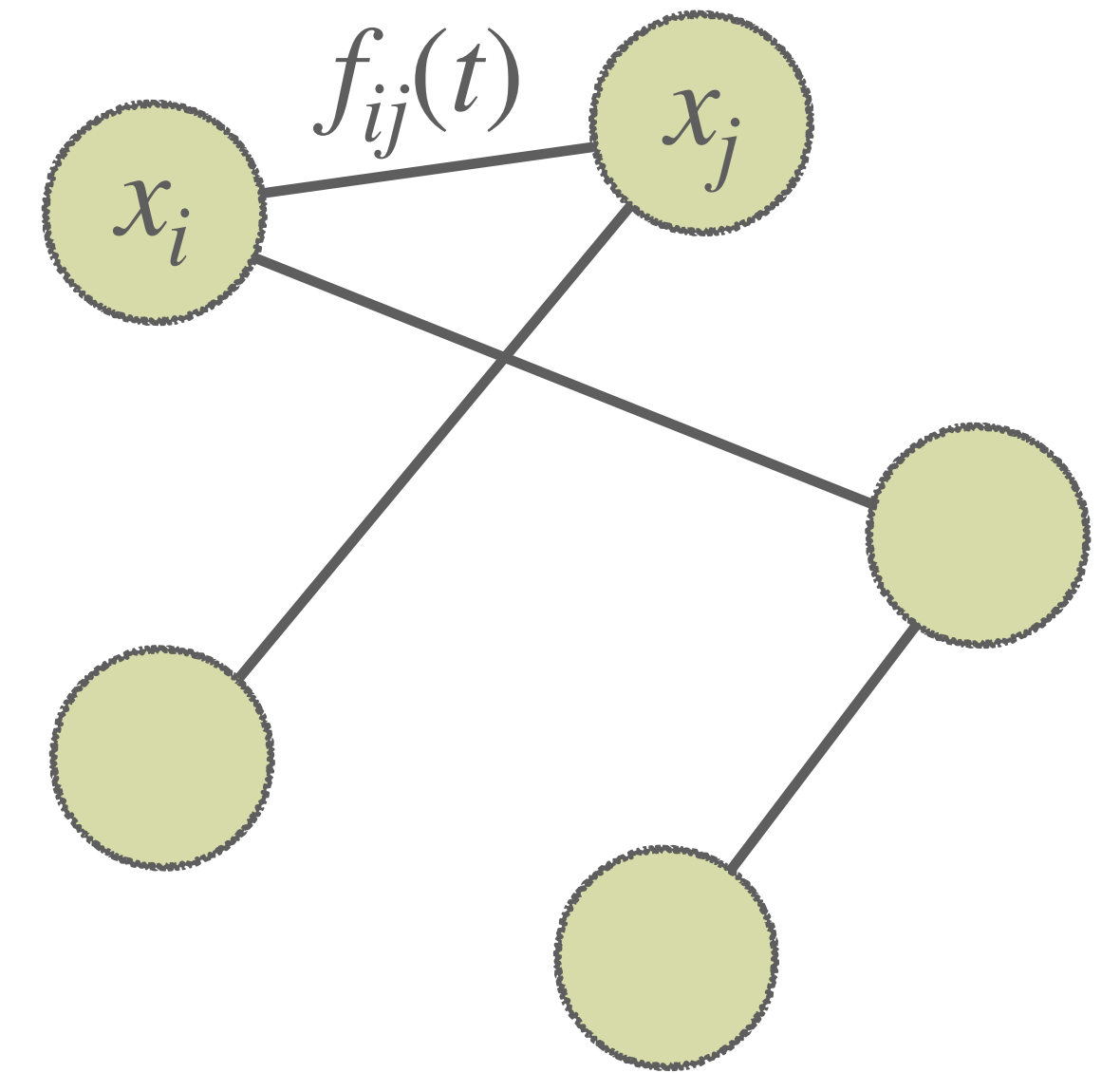
# General speed limit

- Minimum time required for changing states

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

$$\bar{v} := \tau^{-1} \int_0^\tau dt \sum_{(i,j) \in E} c_{ij} |f_{ij}(t)| : \text{average velocity}$$

- ✓ tight and saturable
- ✓ hold for a general setup and arbitrary cost matrix



$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

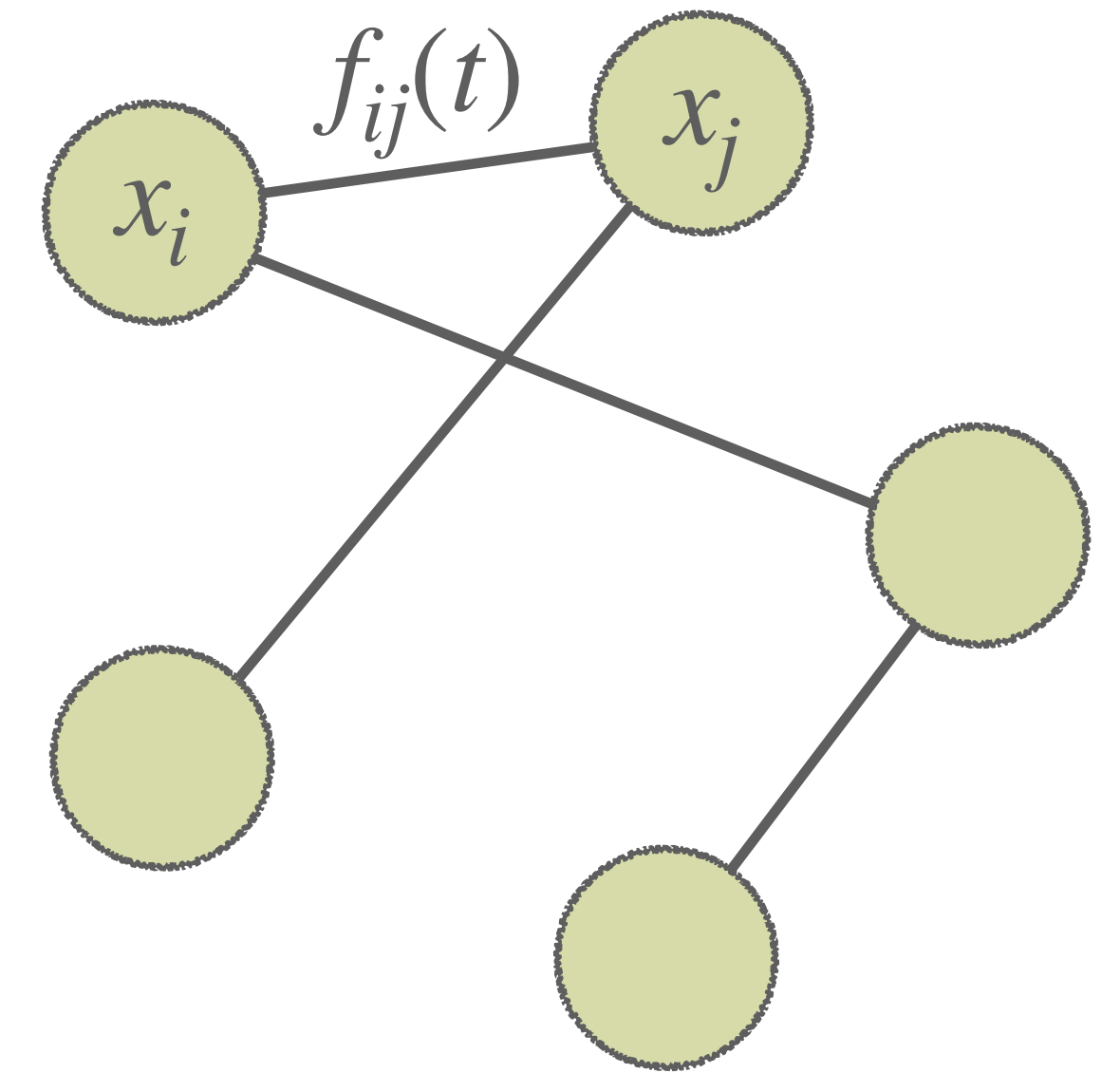
# General speed limit

- Minimum time required for changing states

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

$$\bar{v} := \tau^{-1} \int_0^\tau dt \sum_{(i,j) \in E} c_{ij} |f_{ij}(t)| : \text{average velocity}$$

- ✓ tight and saturable
- ✓ hold for a general setup and arbitrary cost matrix
- ✓ incorporate geometric structure of dynamics into  $\mathcal{W}(x_0, x_\tau)$



$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

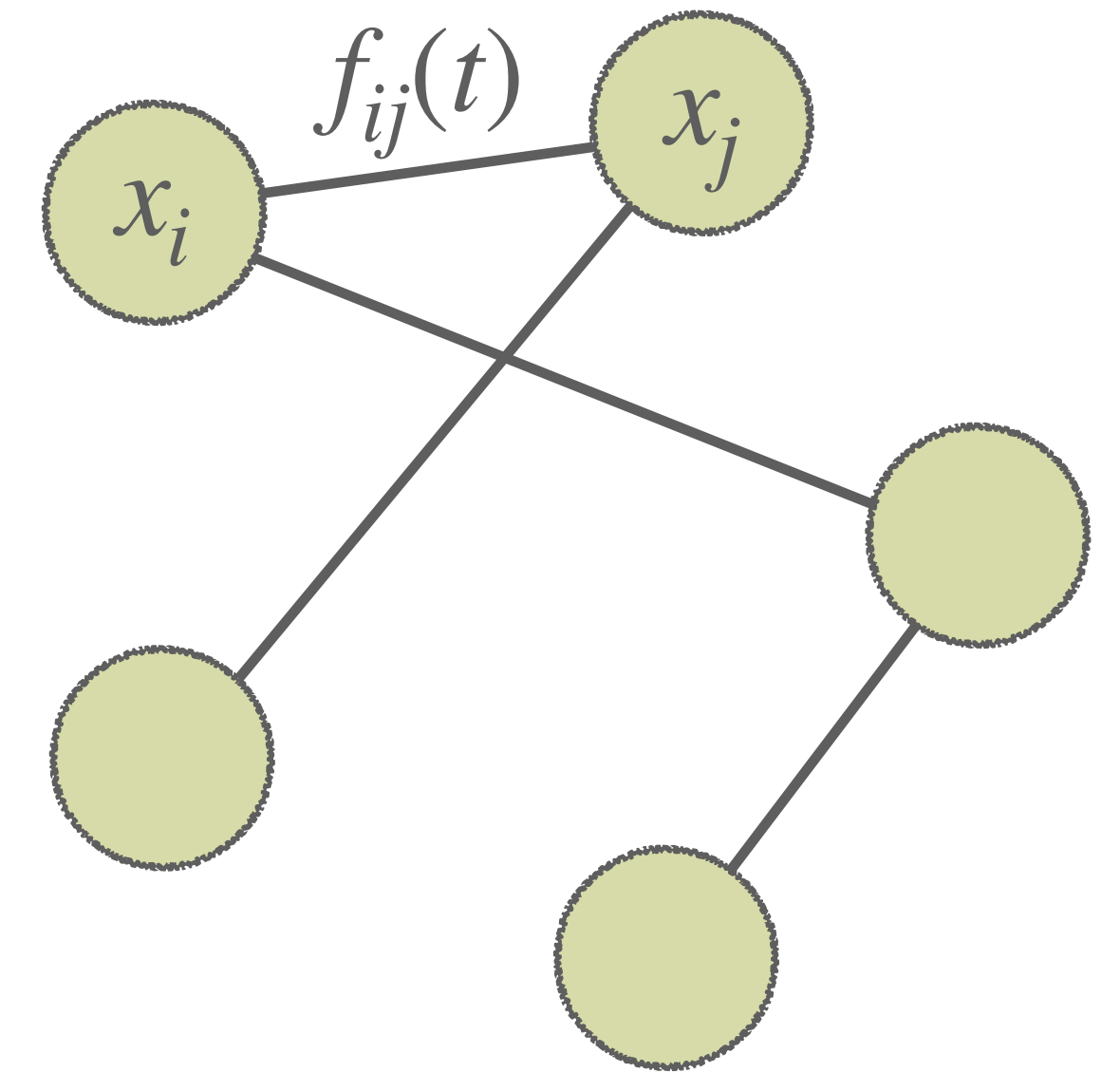
# General speed limit

- Minimum time required for changing states

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

$$\bar{v} := \tau^{-1} \int_0^\tau dt \sum_{(i,j) \in E} c_{ij} |f_{ij}(t)| : \text{average velocity}$$

- ✓ tight and saturable
- ✓ hold for a general setup and arbitrary cost matrix
- ✓ incorporate geometric structure of dynamics into  $\mathcal{W}(x_0, x_\tau)$
- ✓ suitable choices of cost matrix lead to significant implications



$$\dot{x}_i(t) = \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

# Derivation of general speed limit

# Derivation of general speed limit

- $L^1$ -Wasserstein distance

$$\mathcal{W}(x, y) := \min_{\pi} \sum_{m, n} \pi_{mn} c_{mn}$$

# Derivation of general speed limit

- Kantorovich-Rubinstein duality

$$\begin{aligned}\mathcal{W}(\mathbf{x}, \mathbf{y}) &= \max_{\boldsymbol{\phi}} \boldsymbol{\phi}^\top (\mathbf{x} - \mathbf{y}) \\ \text{s.t. } &|\phi_i - \phi_j| \leq c_{ij} \quad \forall i, j\end{aligned}$$



- $L^1$ -Wasserstein distance

$$\mathcal{W}(\mathbf{x}, \mathbf{y}) := \min_{\pi} \sum_{m,n} \pi_{mn} c_{mn}$$

# Derivation of general speed limit

- Kantorovich-Rubinstein duality

$$\begin{aligned} \mathcal{W}(\mathbf{x}, \mathbf{y}) &= \max_{\boldsymbol{\phi}} \boldsymbol{\phi}^\top (\mathbf{x} - \mathbf{y}) \\ \text{s.t. } |\phi_i - \phi_j| &\leq c_{ij} \quad \forall i, j \end{aligned}$$



- $L^1$ -Wasserstein distance

$$\mathcal{W}(\mathbf{x}, \mathbf{y}) := \min_{\pi} \sum_{m,n} \pi_{mn} c_{mn}$$

$$\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau) = \max_{\boldsymbol{\phi}} \boldsymbol{\phi}^\top (\mathbf{x}_\tau - \mathbf{x}_0)$$

$$= \max_{\boldsymbol{\phi}} \sum_i \phi_i \int_0^\tau dt \sum_{j \in \mathcal{B}_i} f_{ij}(t)$$

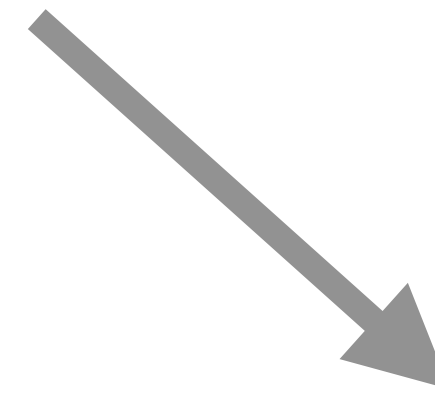
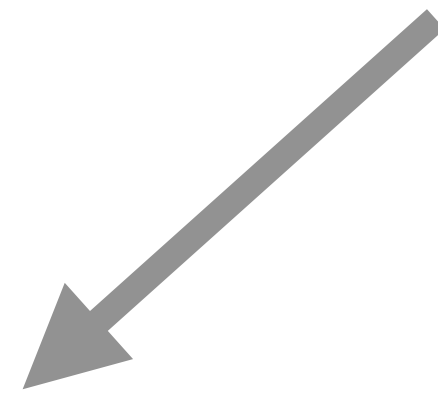
$$= \max_{\boldsymbol{\phi}} \sum_{(i,j) \in E} (\phi_i - \phi_j) \int_0^\tau dt f_{ij}(t)$$

$$\leq \max_{\boldsymbol{\phi}} \sum_{(i,j) \in E} |\phi_i - \phi_j| \int_0^\tau dt |f_{ij}(t)|$$

$$\leq \int_0^\tau dt \sum_{(i,j) \in E} c_{ij} |f_{ij}(t)|$$

# Consequences

General speed limit



Topological  
speed limit

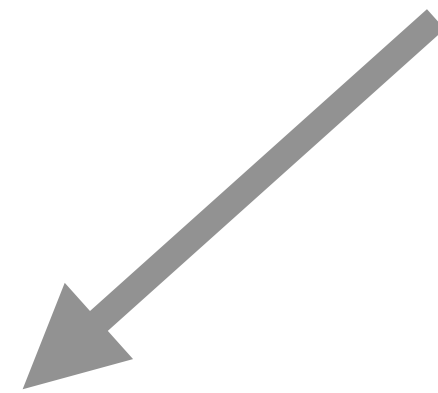
$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

Resolution of  
bosonic transport

$$\tau \gtrsim d_{XY}^\lambda$$

# Consequences

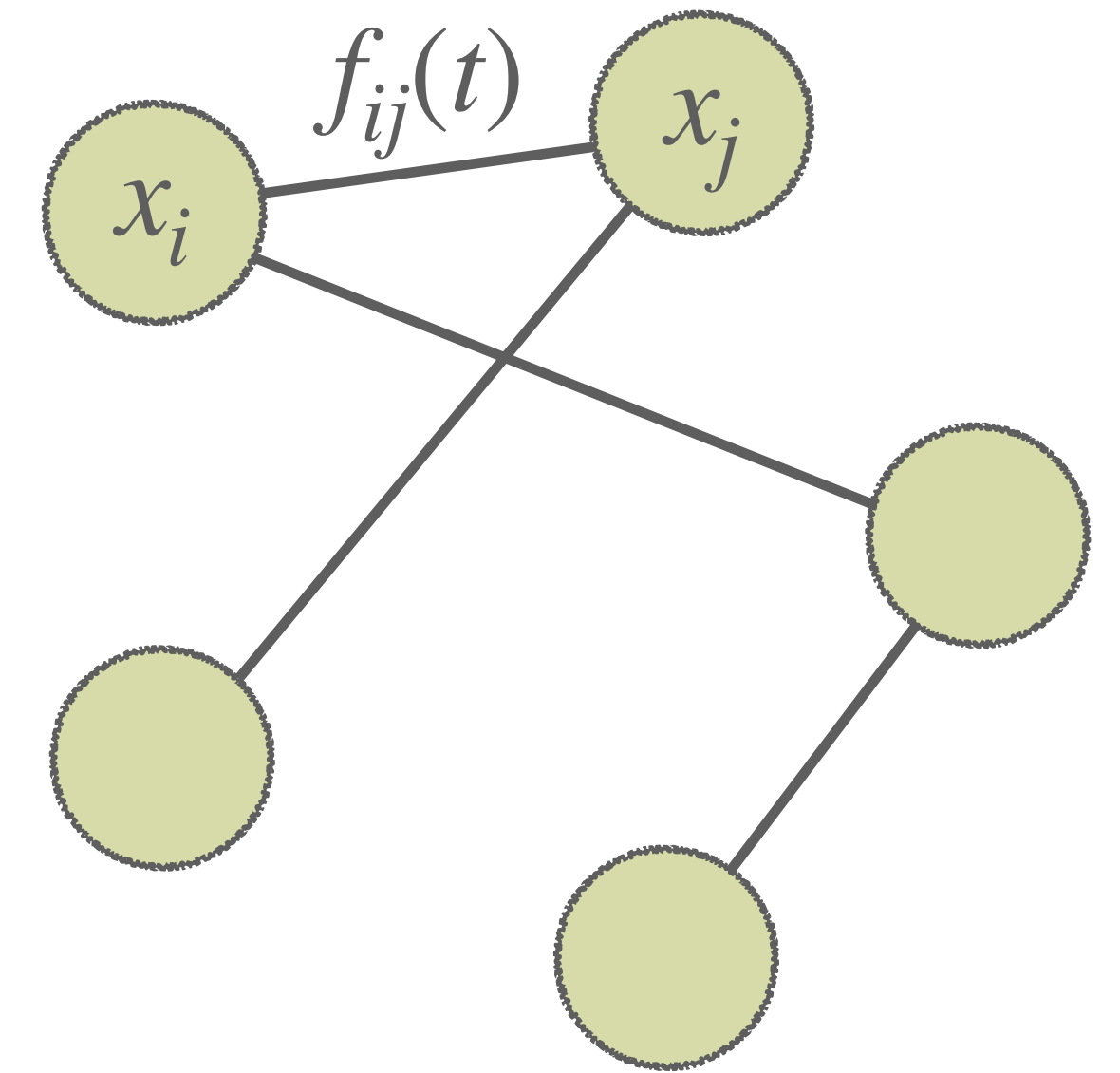
General speed limit



Topological  
speed limit

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

# Topological speed limit

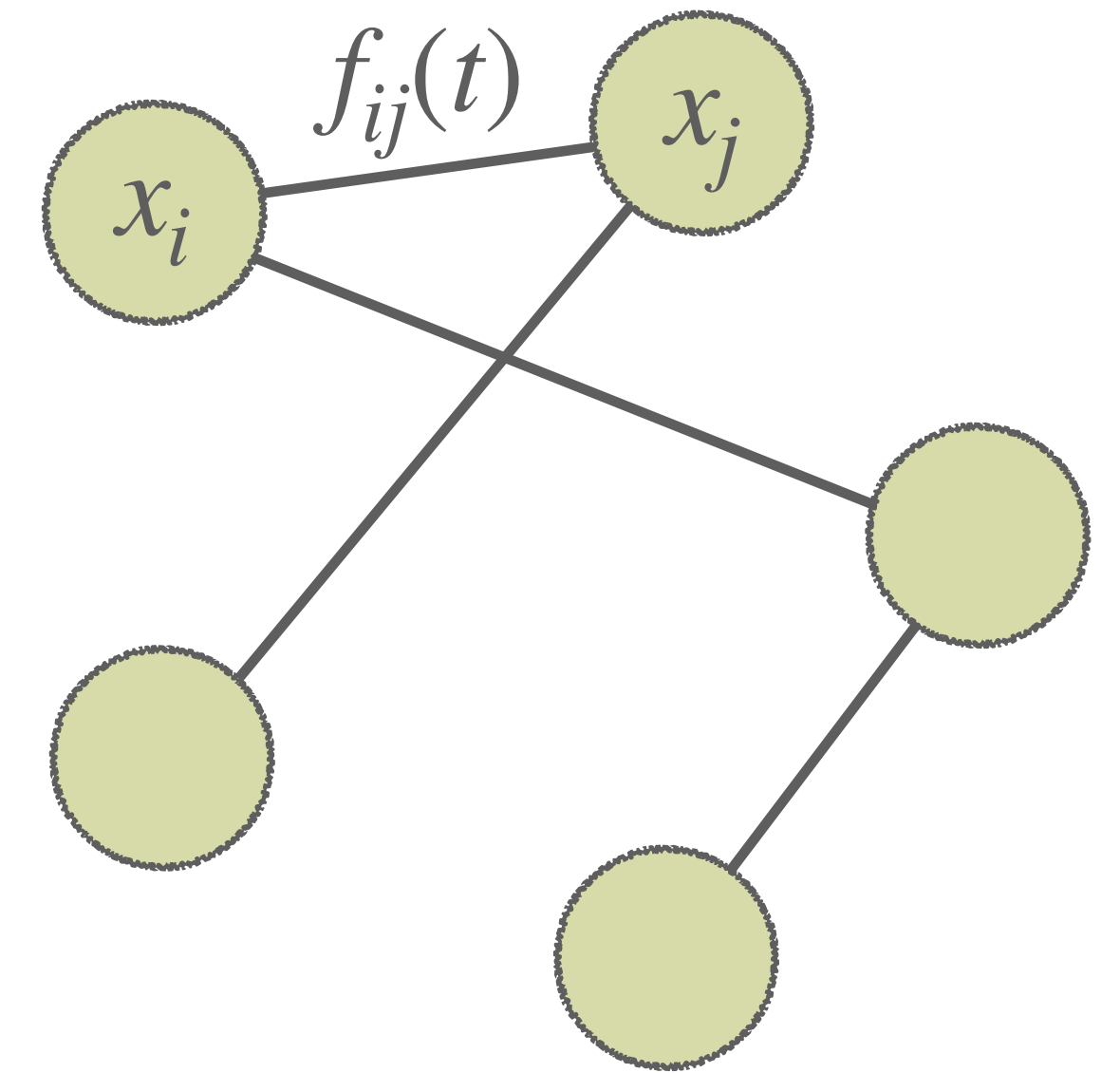


# Topological speed limit

- Shortest-path distances  $[c_{ij}]$

$$c_{ij} = 1 \quad \forall (i, j) \in E$$

$$c_{ij} = \min_{P=[i \leftrightarrow \dots \leftrightarrow j]} \text{length}(P) \quad \forall (i, j) \notin E$$



# Topological speed limit

- Shortest-path distances  $[c_{ij}]$

$$c_{ij} = 1 \quad \forall (i, j) \in E$$

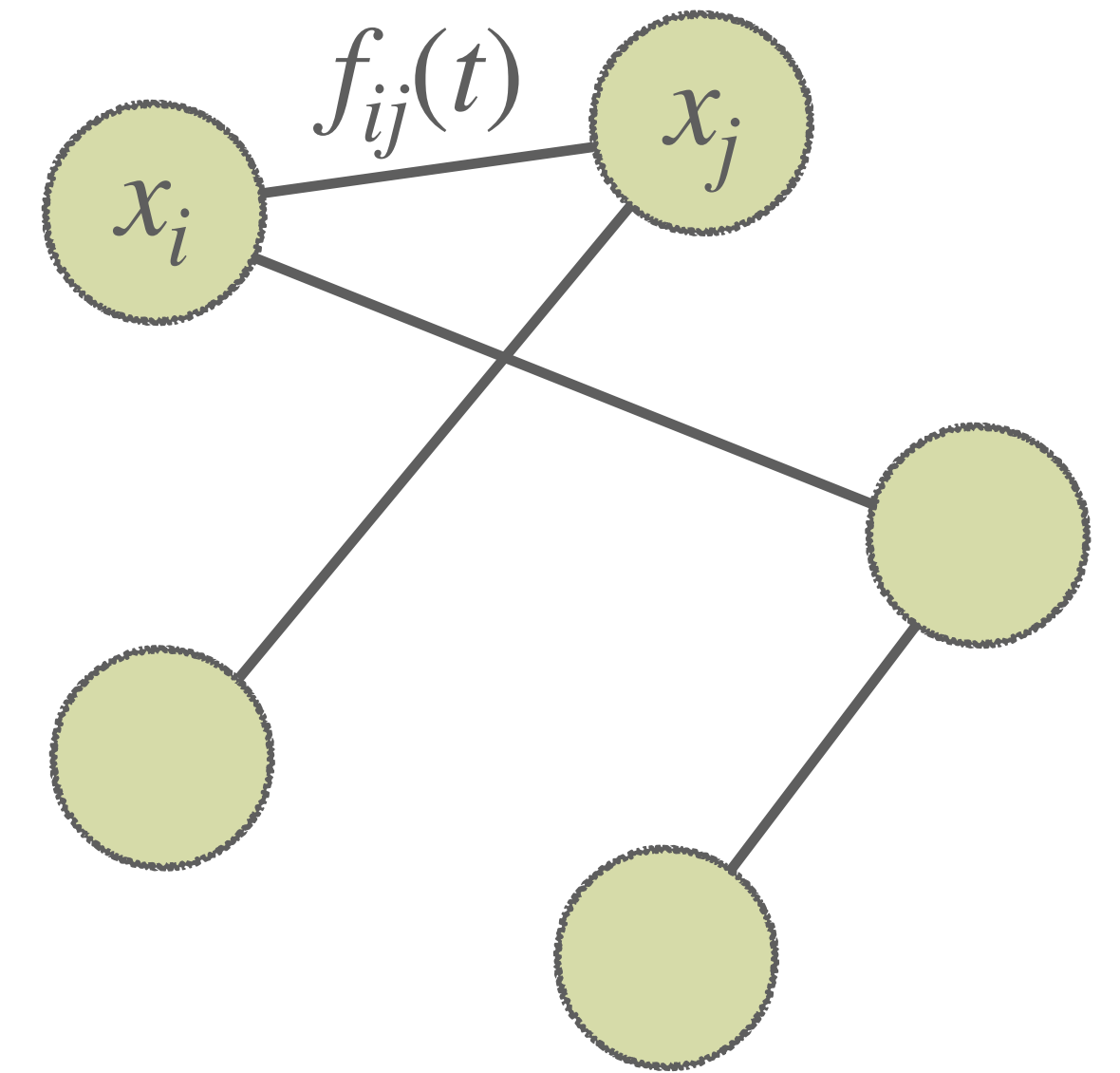
$$c_{ij} = \min_{P=[i \leftrightarrow \dots \leftrightarrow j]} \text{length}(P) \quad \forall (i, j) \notin E$$

- Topological speed limit

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

Vu and Saito, Phys. Rev. Lett. (2023)

$$\bar{v} := \tau^{-1} \int_0^\tau dt \sum_{(i,j) \in E} |f_{ij}(t)| : \text{average velocity}$$

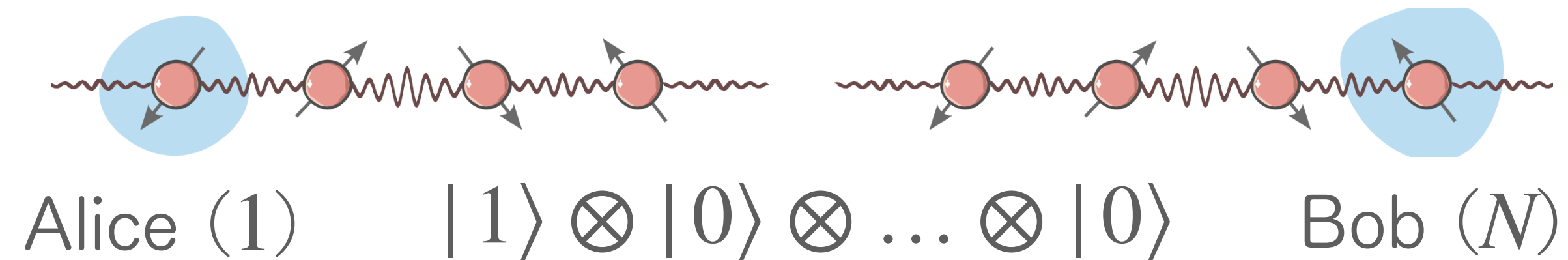


# Applications

- Quantum communication through spin chains

Bose, Phys. Rev. Lett. (2003)

$$H_t = -\frac{\gamma}{2} \sum_{n=1}^{N-1} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1} + \sum_{n=1}^N B_n(t) \sigma_n^z$$

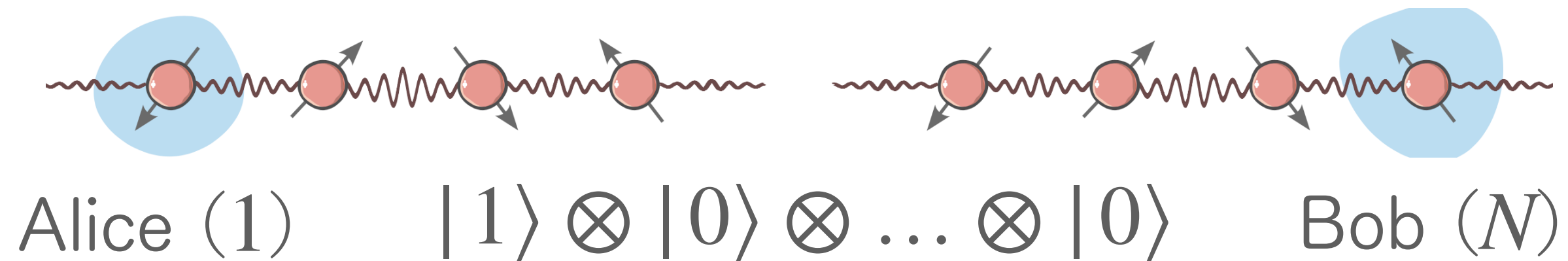


# Applications

- Quantum communication through spin chains

Bose, Phys. Rev. Lett. (2003)

$$H_t = -\frac{\gamma}{2} \sum_{n=1}^{N-1} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1} + \sum_{n=1}^N B_n(t) \sigma_n^z$$



Time needed to transfer state from spin 1 to  $N$

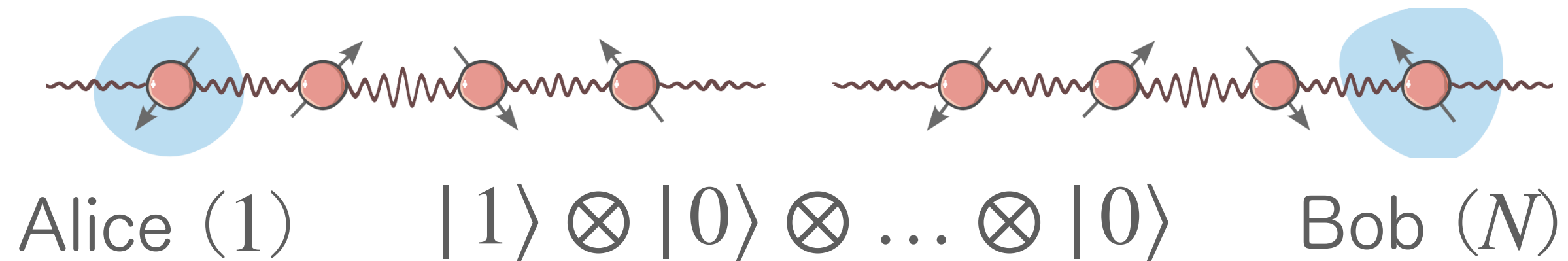
$$\tau \geq \frac{N-1}{2\gamma}$$

# Applications

- Quantum communication through spin chains

Bose, Phys. Rev. Lett. (2003)

$$H_t = -\frac{\gamma}{2} \sum_{n=1}^{N-1} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1} + \sum_{n=1}^N B_n(t) \sigma_n^z$$



Time needed to transfer state from spin 1 to  $N$

$$\tau \geq \frac{N-1}{2\gamma}$$

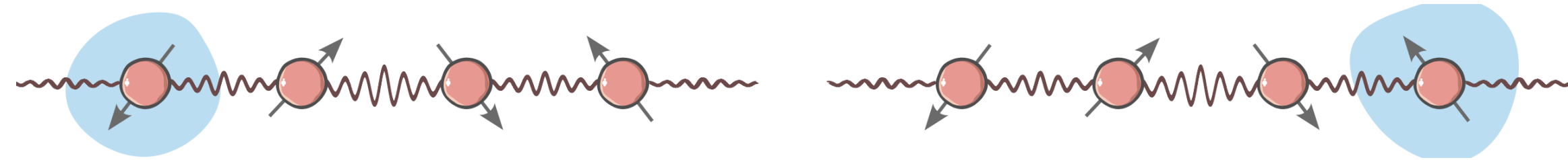
- ✓ applicable to arbitrary graph of spins

# Applications

- Quantum communication through spin chains

Bose, Phys. Rev. Lett. (2003)

$$H_t = -\frac{\gamma}{2} \sum_{n=1}^{N-1} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1} + \sum_{n=1}^N B_n(t) \sigma_n^z$$

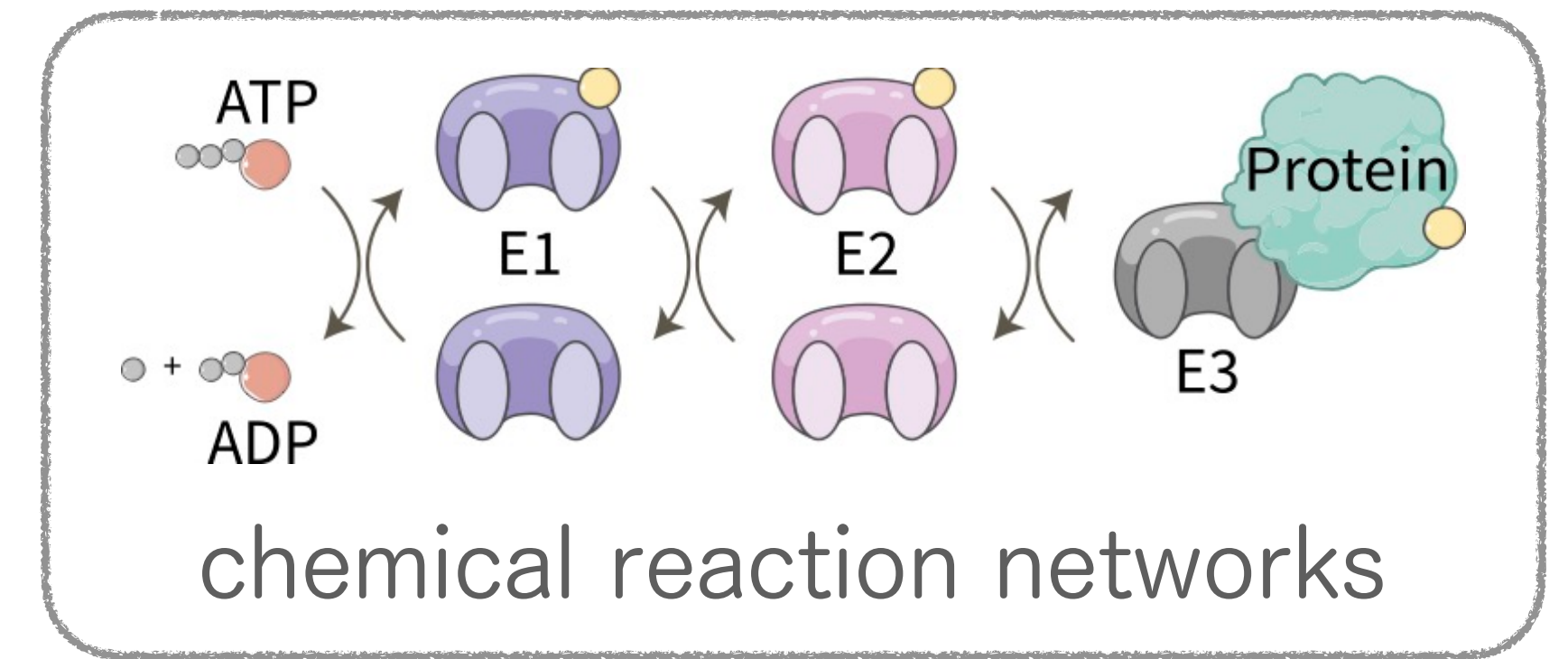


Alice (1)  $|1\rangle \otimes |0\rangle \otimes \dots \otimes |0\rangle$  Bob (N)

Time needed to transfer state from spin 1 to N

$$\tau \geq \frac{N-1}{2\gamma}$$

✓ applicable to arbitrary graph of spins

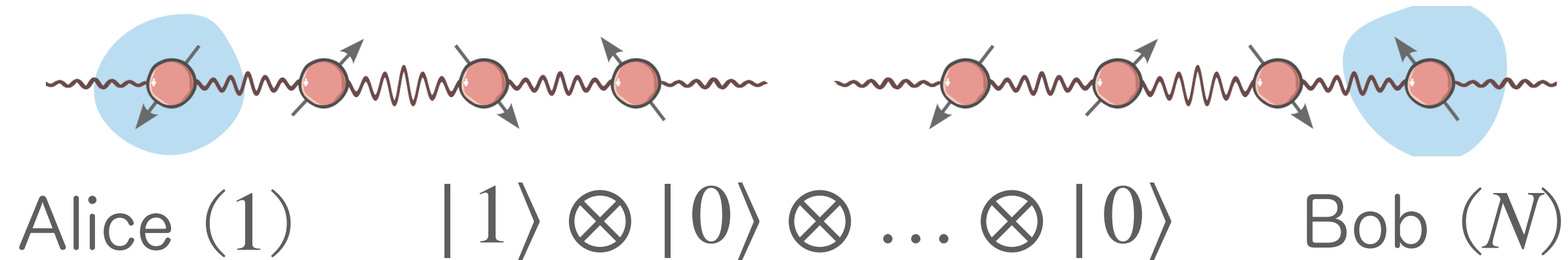


# Applications

- Quantum communication through spin chains

Bose, Phys. Rev. Lett. (2003)

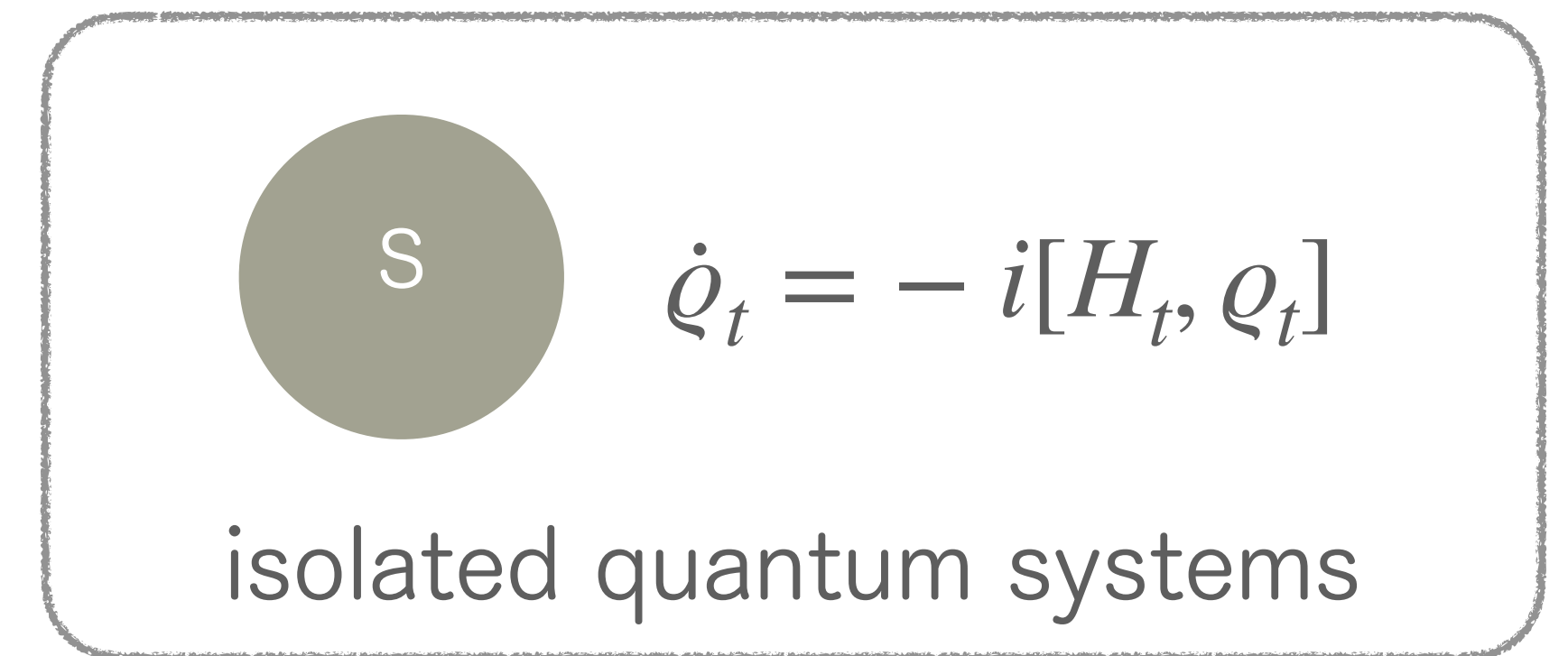
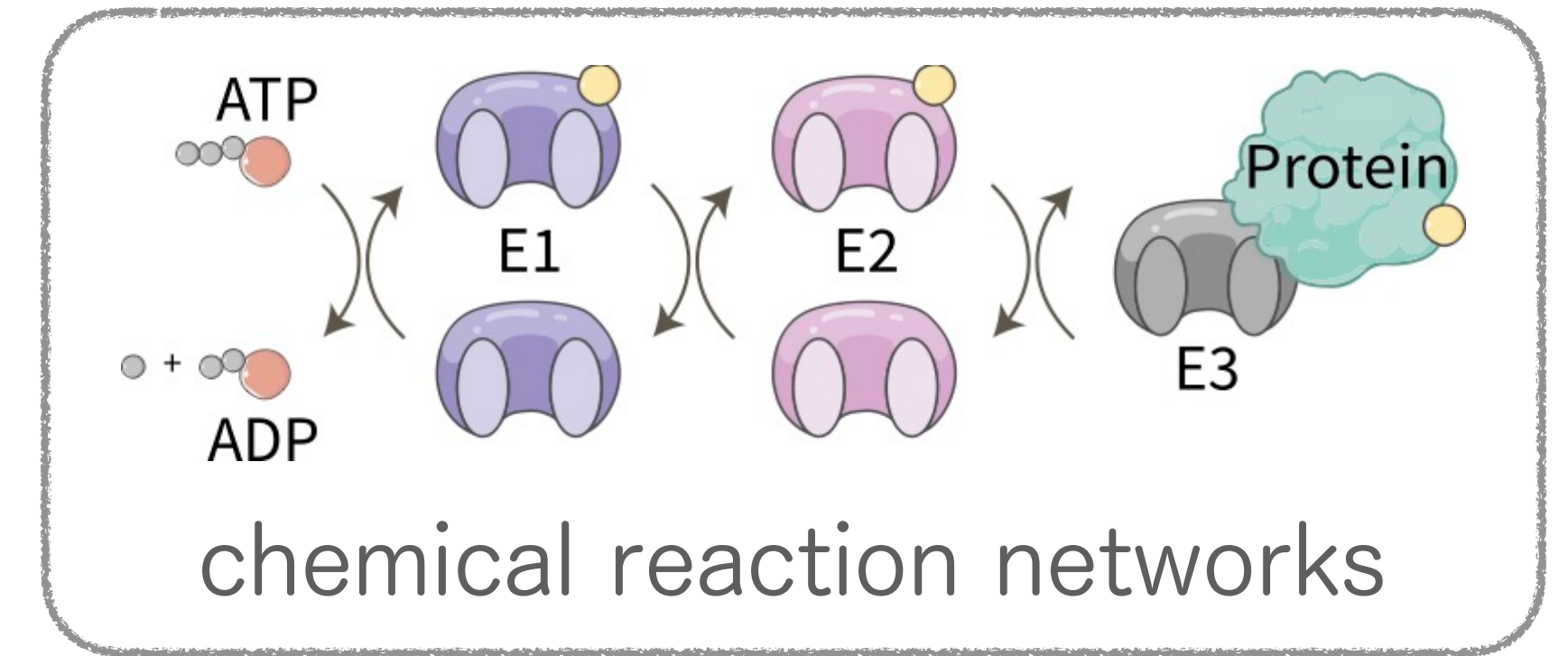
$$H_t = -\frac{\gamma}{2} \sum_{n=1}^{N-1} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1} + \sum_{n=1}^N B_n(t) \sigma_n^z$$



Time needed to transfer state from spin 1 to  $N$

$$\tau \geq \frac{N-1}{2\gamma}$$

- ✓ applicable to arbitrary graph of spins

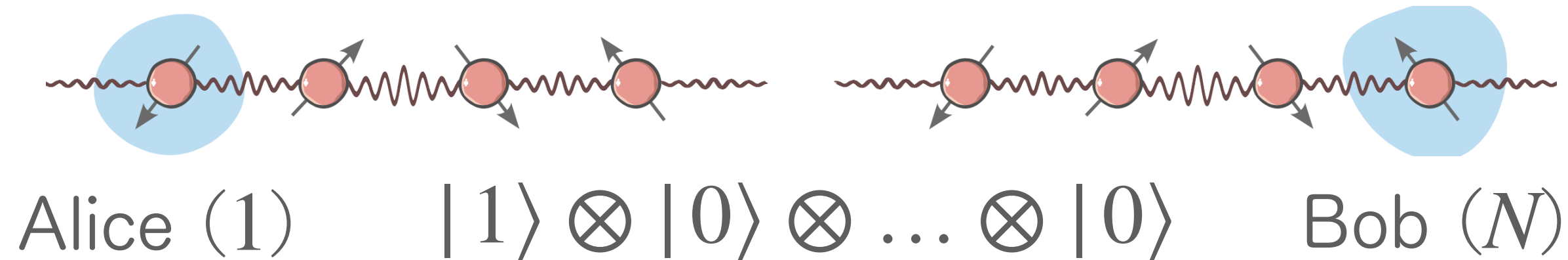


# Applications

- Quantum communication through spin chains

Bose, Phys. Rev. Lett. (2003)

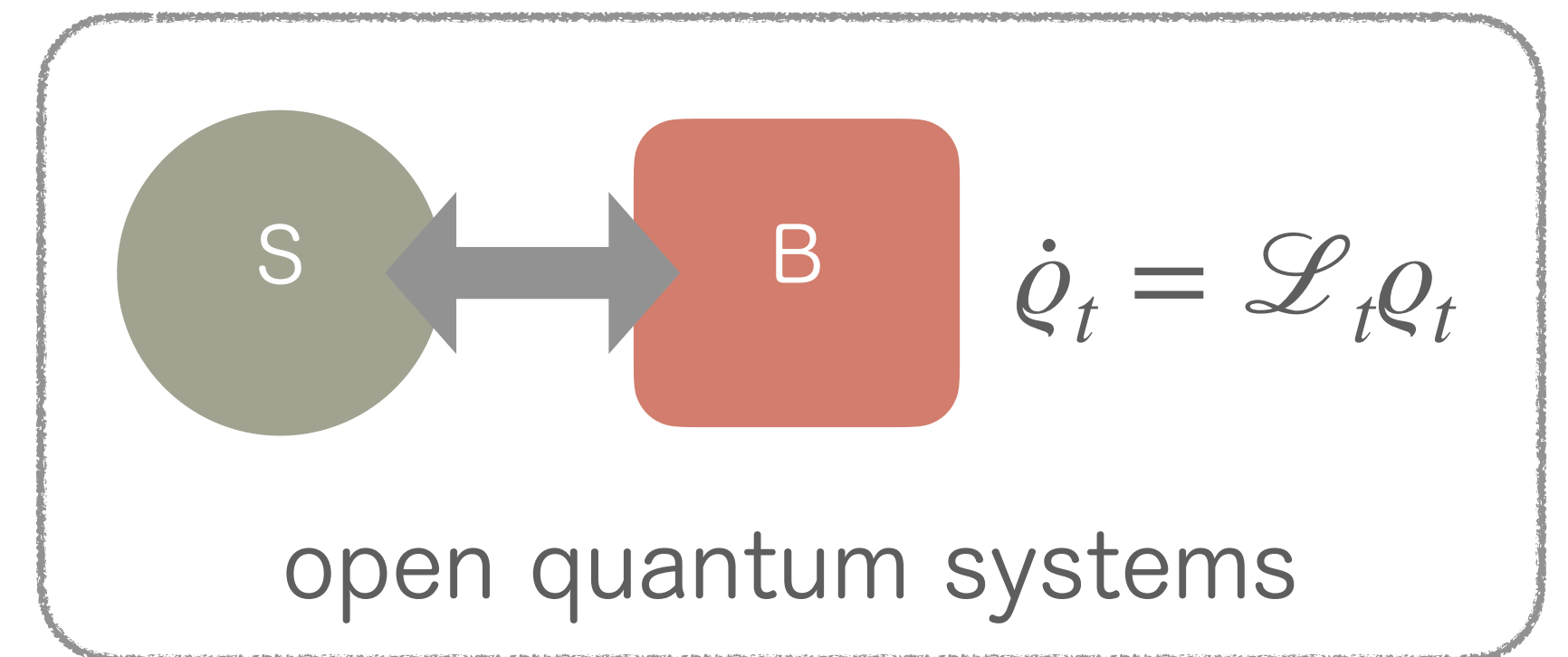
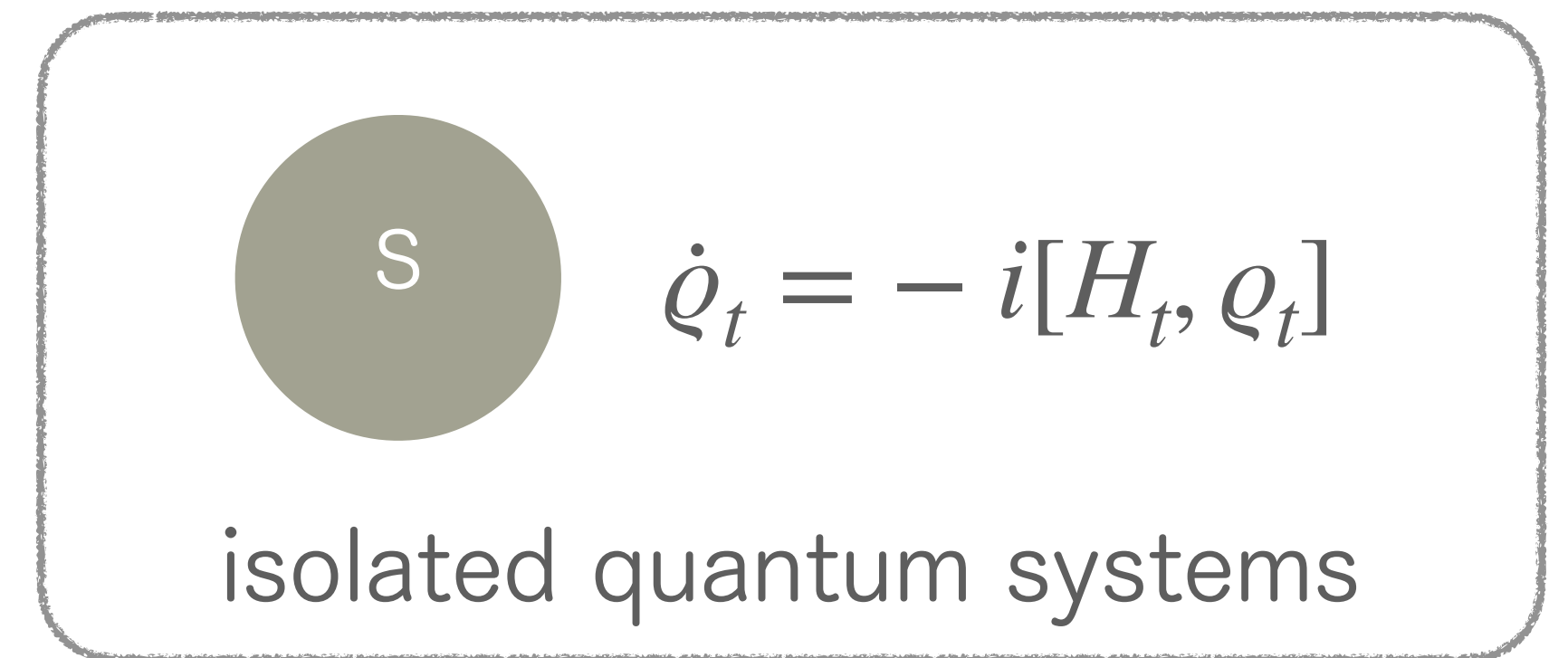
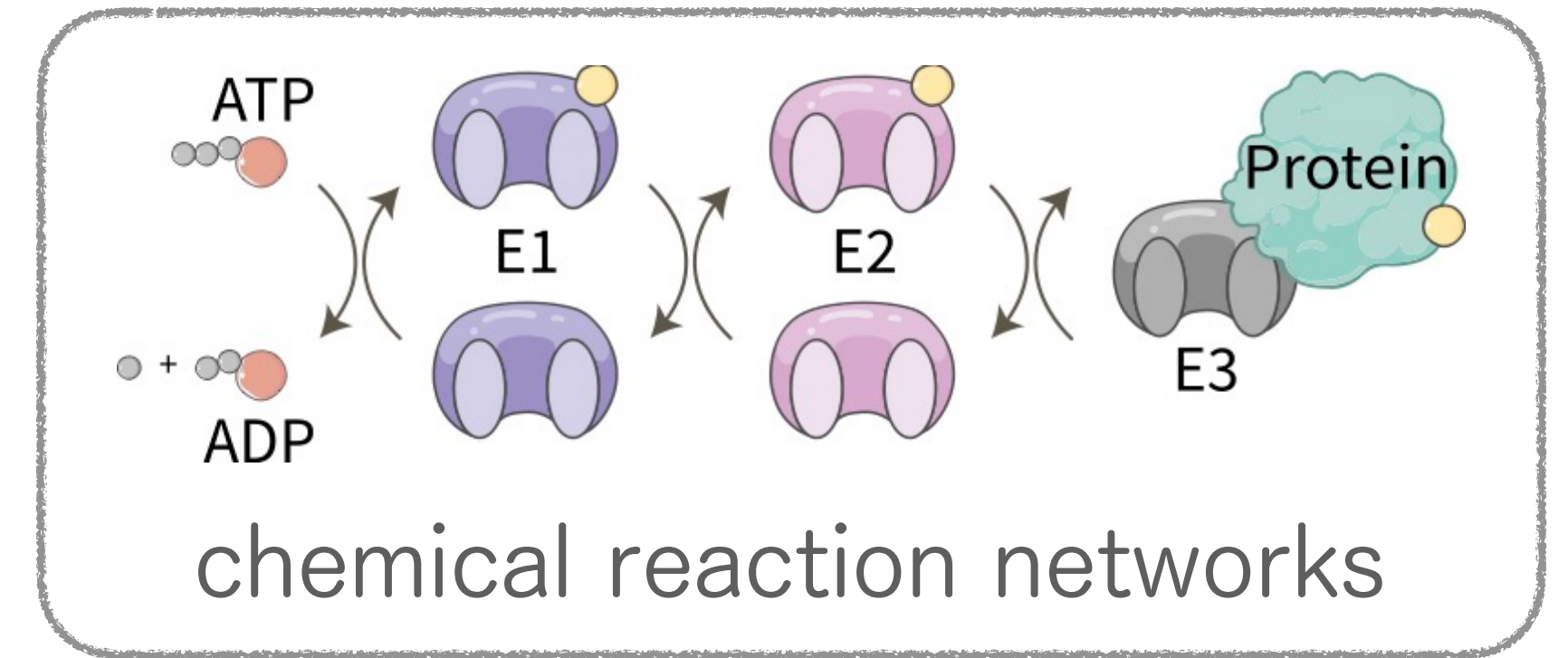
$$H_t = -\frac{\gamma}{2} \sum_{n=1}^{N-1} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1} + \sum_{n=1}^N B_n(t) \sigma_n^z$$



Time needed to transfer state from spin 1 to  $N$

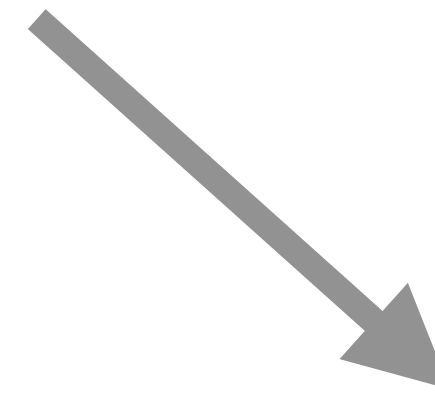
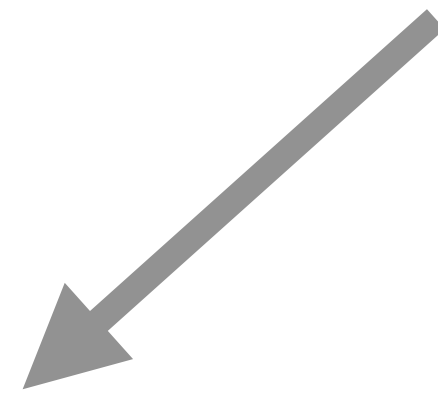
$$\tau \geq \frac{N-1}{2\gamma}$$

- ✓ applicable to arbitrary graph of spins



# Consequences

General speed limit



Topological  
speed limit

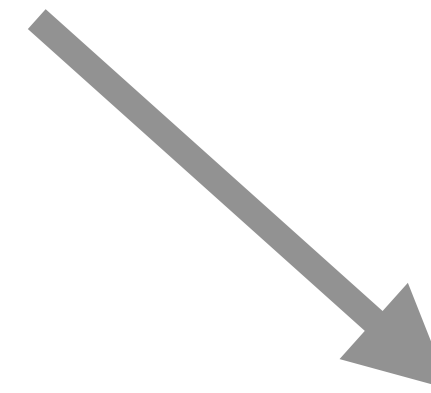
$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}}$$

Resolution of  
bosonic transport

$$\tau \gtrsim d_{XY}^\lambda$$

# Consequences

General speed limit



Resolution of  
bosonic transport

$$\tau \gtrsim d_{XY}^\lambda$$

# Particle transport in closed systems

- Boson system with long-range hopping and long-range interactions

$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_{Z \subset \Lambda} h_Z(t)$$

$|J_{ij}(t)| \leq J/\|i - j\|^\alpha$ : power-law decay  $\alpha > D$  (spatial dimension)

$h_Z(t)$ : arbitrary function of  $\{\hat{n}_i\}_{i \in Z}$

# Particle transport in closed systems

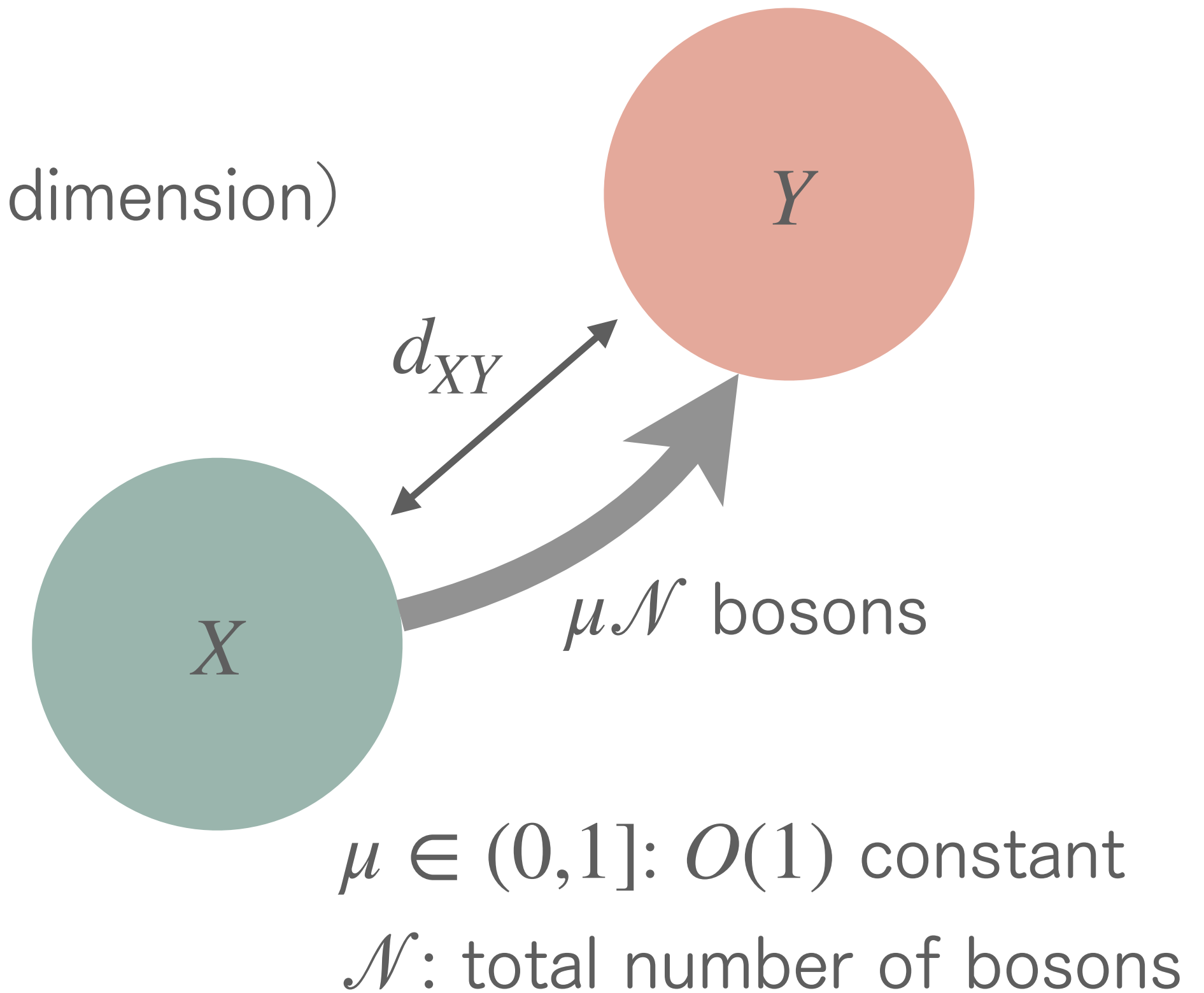
- Boson system with long-range hopping and long-range interactions

$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_{Z \subset \Lambda} h_Z(t)$$

$|J_{ij}(t)| \leq J/\|i - j\|^\alpha$ : power-law decay  $\alpha > D$  (spatial dimension)

$h_Z(t)$ : arbitrary function of  $\{\hat{n}_i\}_{i \in Z}$

- Speed of **macroscopic** particle transport



# Particle transport in closed systems

- Boson system with long-range hopping and long-range interactions

$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_{Z \subset \Lambda} h_Z(t)$$

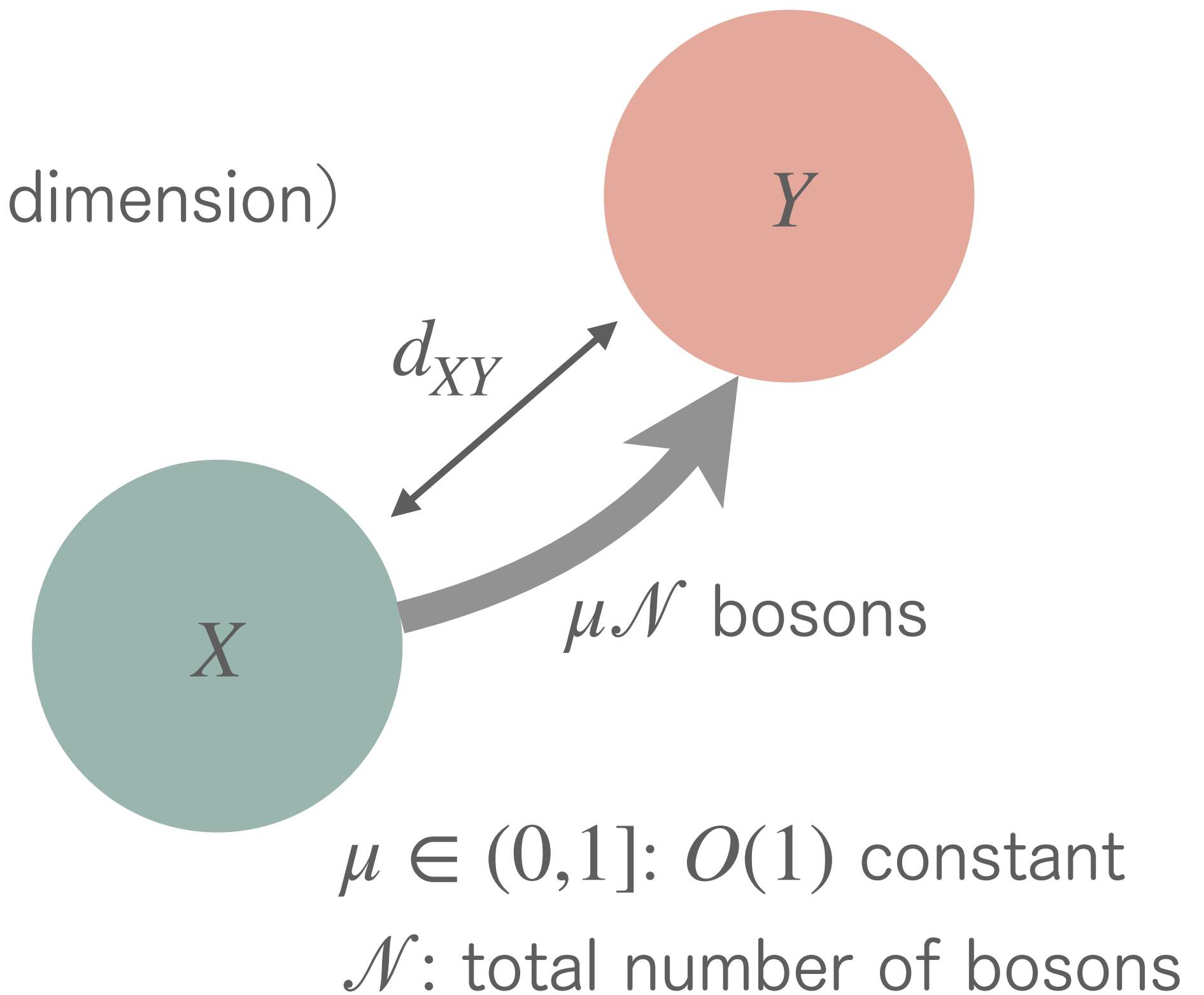
$|J_{ij}(t)| \leq J/\|i - j\|^\alpha$ : power-law decay  $\alpha > D$  (spatial dimension)

$h_Z(t)$ : arbitrary function of  $\{\hat{n}_i\}_{i \in Z}$

- Speed of **macroscopic** particle transport

$$\alpha > D + 2 \rightarrow \tau \gtrsim d_{XY}$$

Faupin, Lemm, and Sigal,  
Phys. Rev. Lett. (2022)



# Particle transport in closed systems

- Boson system with long-range hopping and long-range interactions

$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_{Z \subset \Lambda} h_Z(t)$$

$|J_{ij}(t)| \leq J/\|i - j\|^\alpha$ : power-law decay  $\alpha > D$  (spatial dimension)

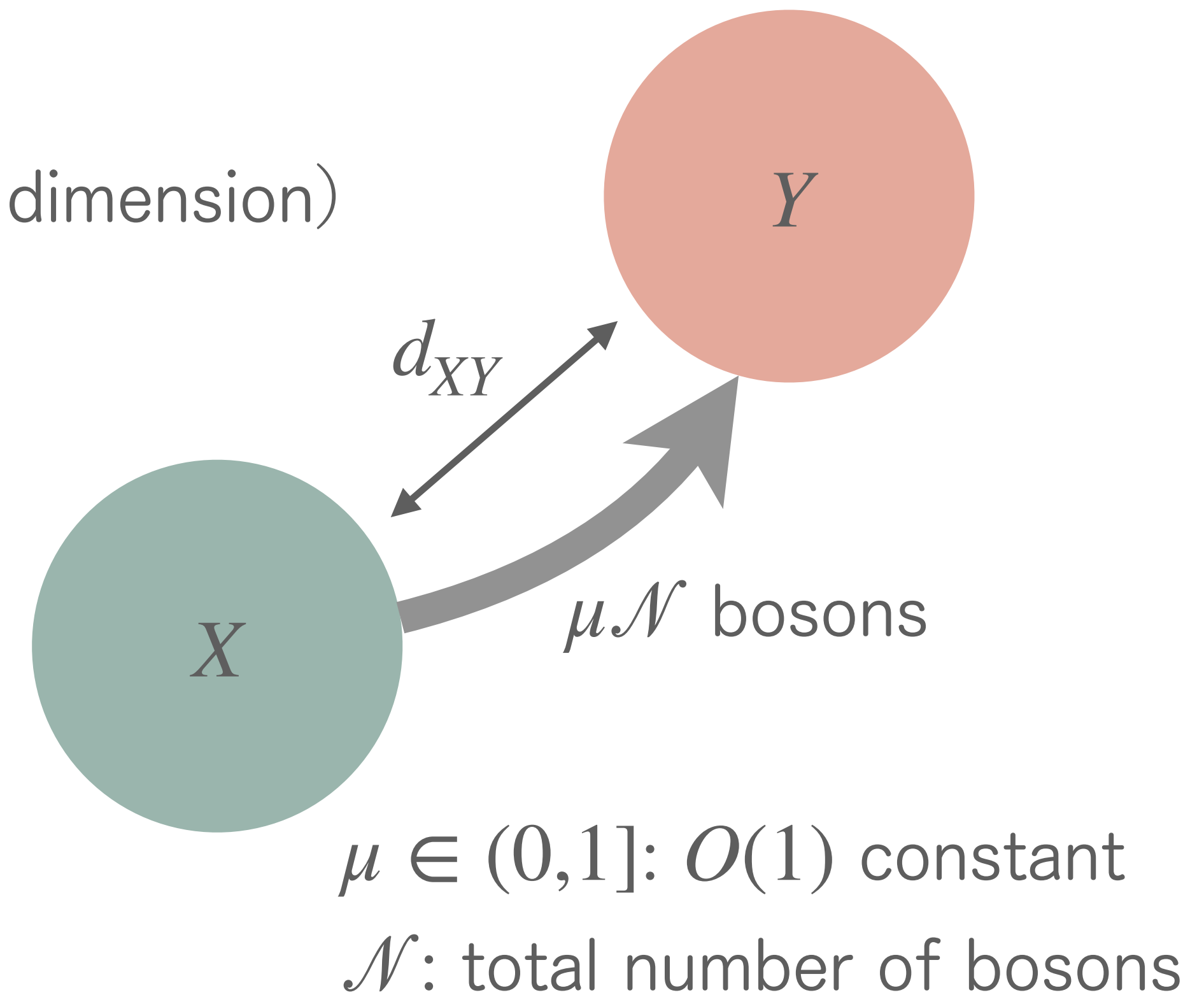
$h_Z(t)$ : arbitrary function of  $\{\hat{n}_i\}_{i \in Z}$

- Speed of **macroscopic** particle transport

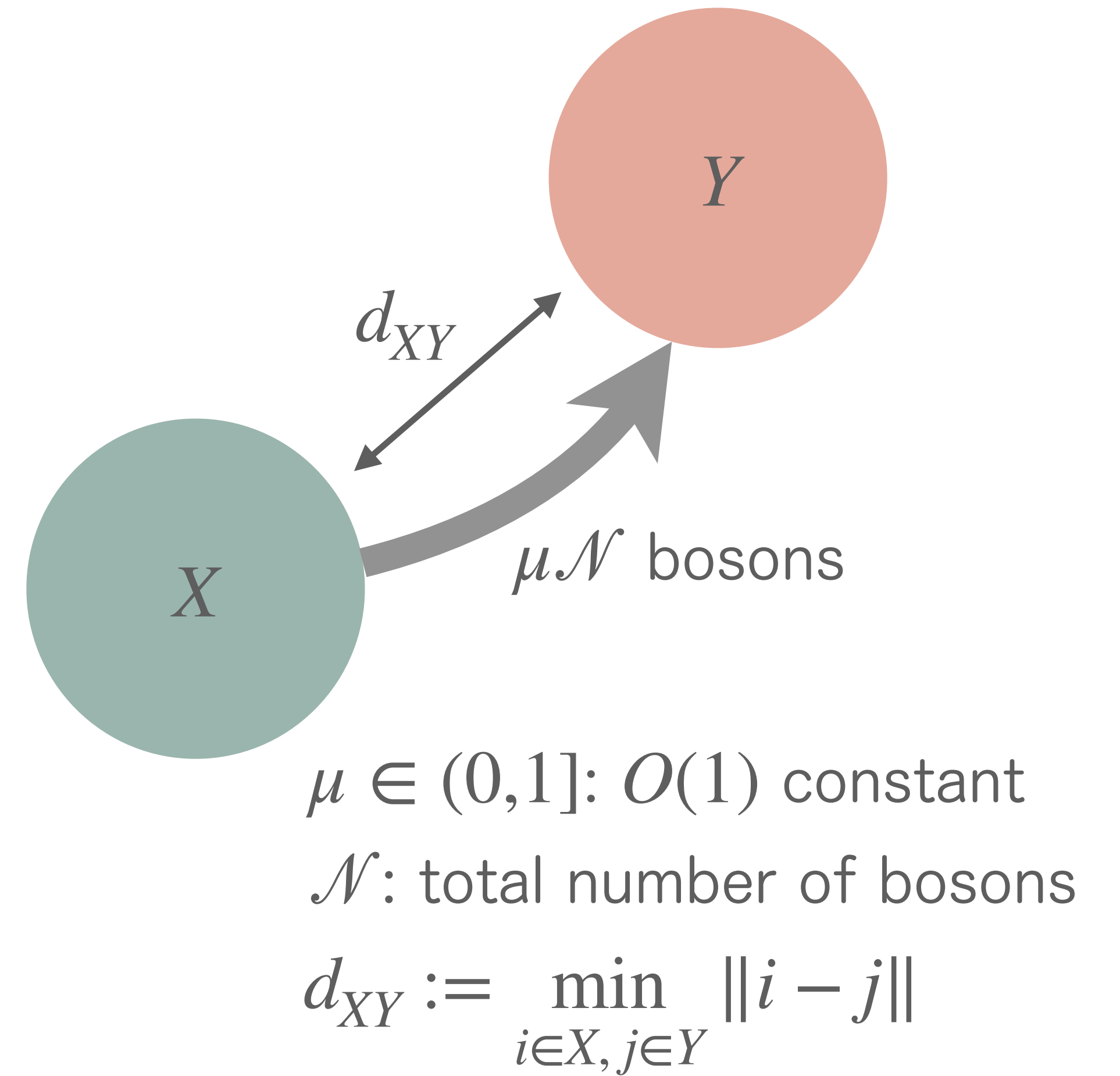
$$\alpha > D + 2 \rightarrow \tau \gtrsim d_{XY}$$

Faupin, Lemm, and Sigal,  
Phys. Rev. Lett. (2022)

x open problem for full range  $\alpha > D$



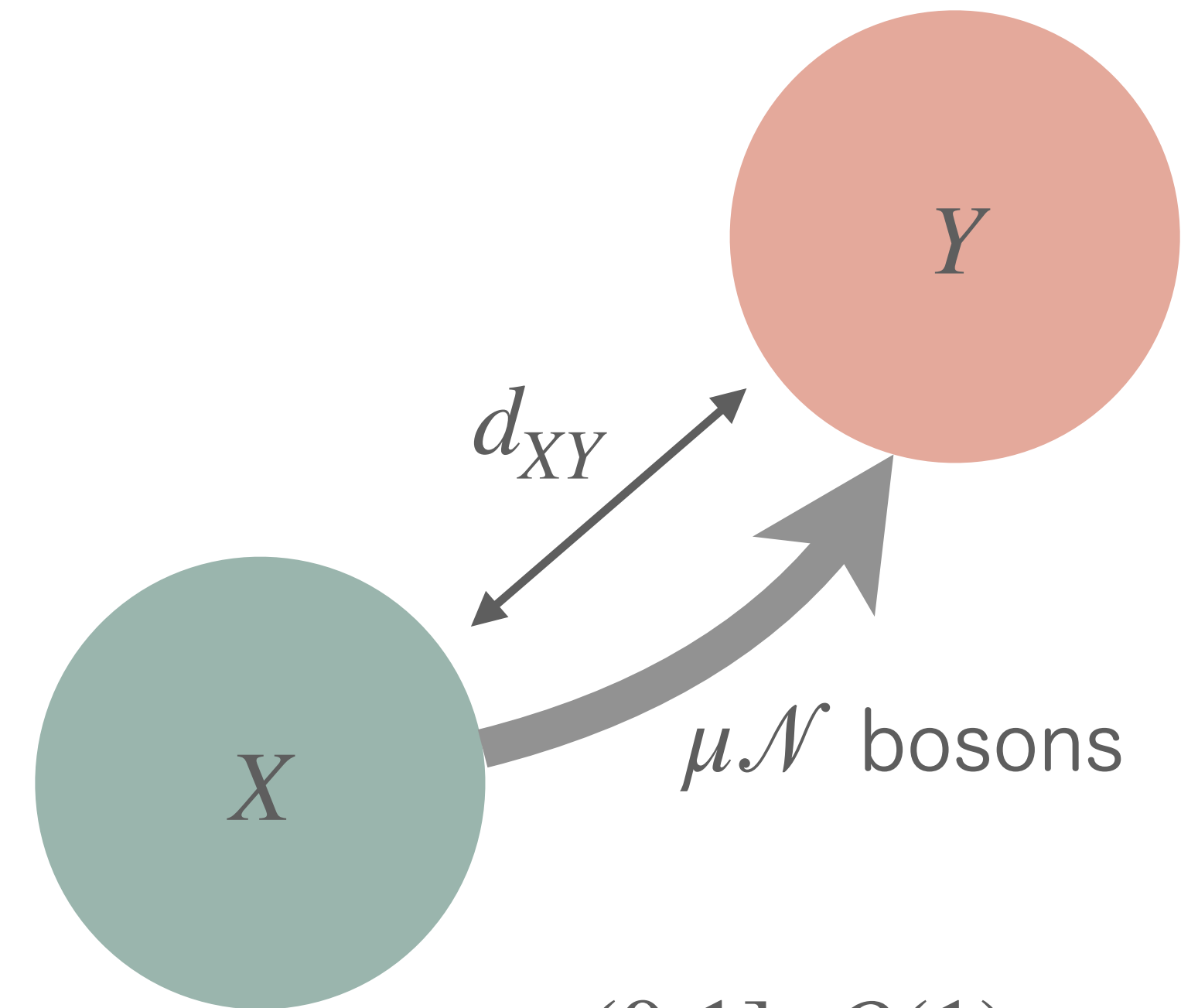
# Resolution of particle transport



# Resolution of particle transport

- Boson concentrations at each site

$$x_i(t) := \mathcal{N}^{-1} \text{tr} \{ \hat{n}_i \varrho_t \}$$



$\mu \in (0,1]$ :  $O(1)$  constant

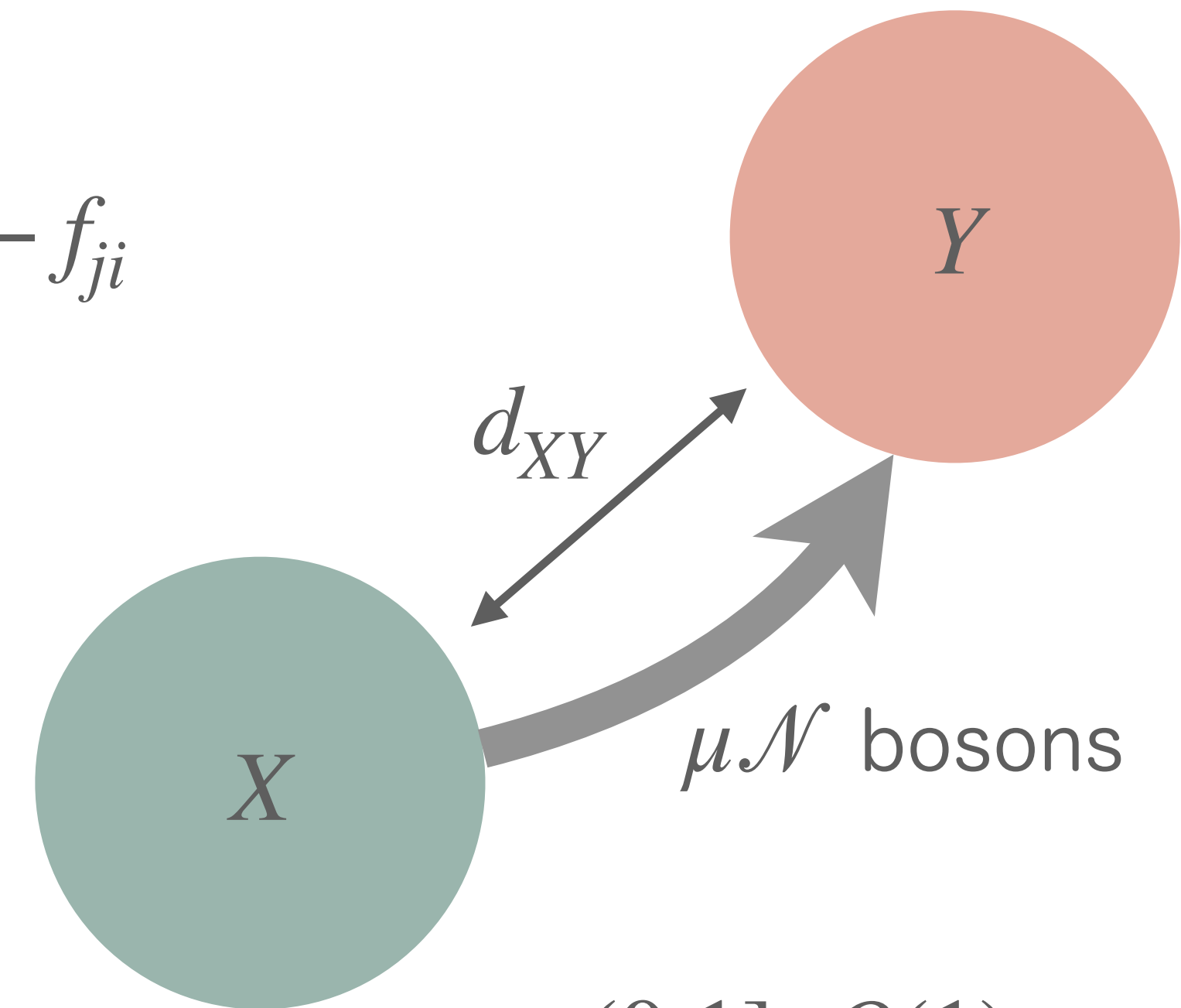
$\mathcal{N}$ : total number of bosons

$$d_{XY} := \min_{i \in X, j \in Y} \|i - j\|$$

# Resolution of particle transport

- Boson concentrations at each site

$$x_i(t) := \mathcal{N}^{-1} \text{tr}\{\hat{n}_i \varrho_t\} \rightarrow \dot{x}_i(t) = \sum_{j(\neq i)} f_{ij}(t), \quad f_{ij} = -f_{ji}$$



$\mu \in (0,1]$ :  $O(1)$  constant

$\mathcal{N}$ : total number of bosons

$$d_{XY} := \min_{i \in X, j \in Y} \|i - j\|$$

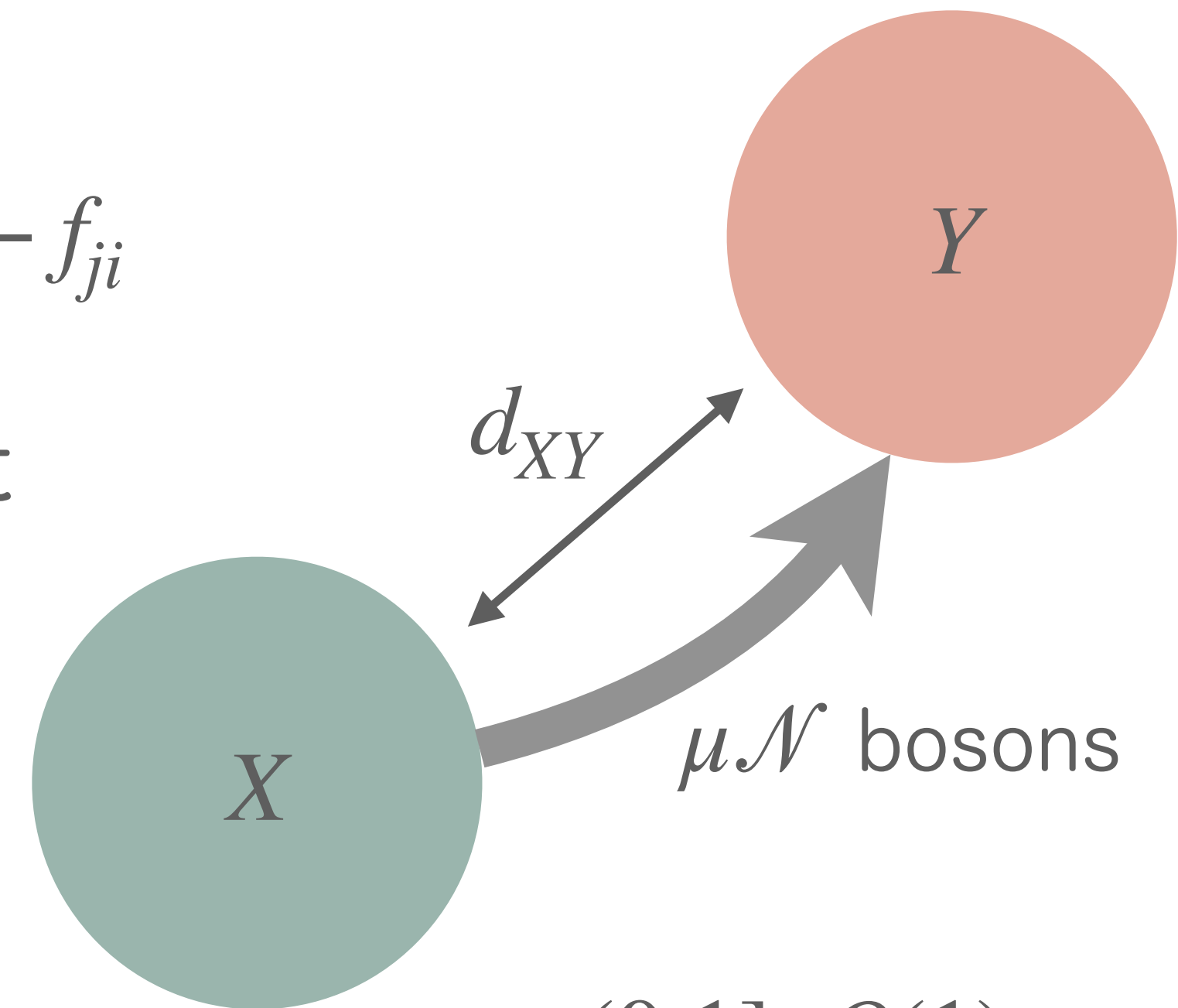
# Resolution of particle transport

- Boson concentrations at each site

$$x_i(t) := \mathcal{N}^{-1} \text{tr}\{\hat{n}_i \mathcal{Q}_t\} \rightarrow \dot{x}_i(t) = \sum_{j(\neq i)} f_{ij}(t), \quad f_{ij} = -f_{ji}$$

- Finite speed of macroscopic particle transport

$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}}$$



$\mu \in (0,1]$ :  $O(1)$  constant  
 $\mathcal{N}$ : total number of bosons

$$d_{XY} := \min_{i \in X, j \in Y} \|i - j\|$$

# Resolution of particle transport

- Boson concentrations at each site

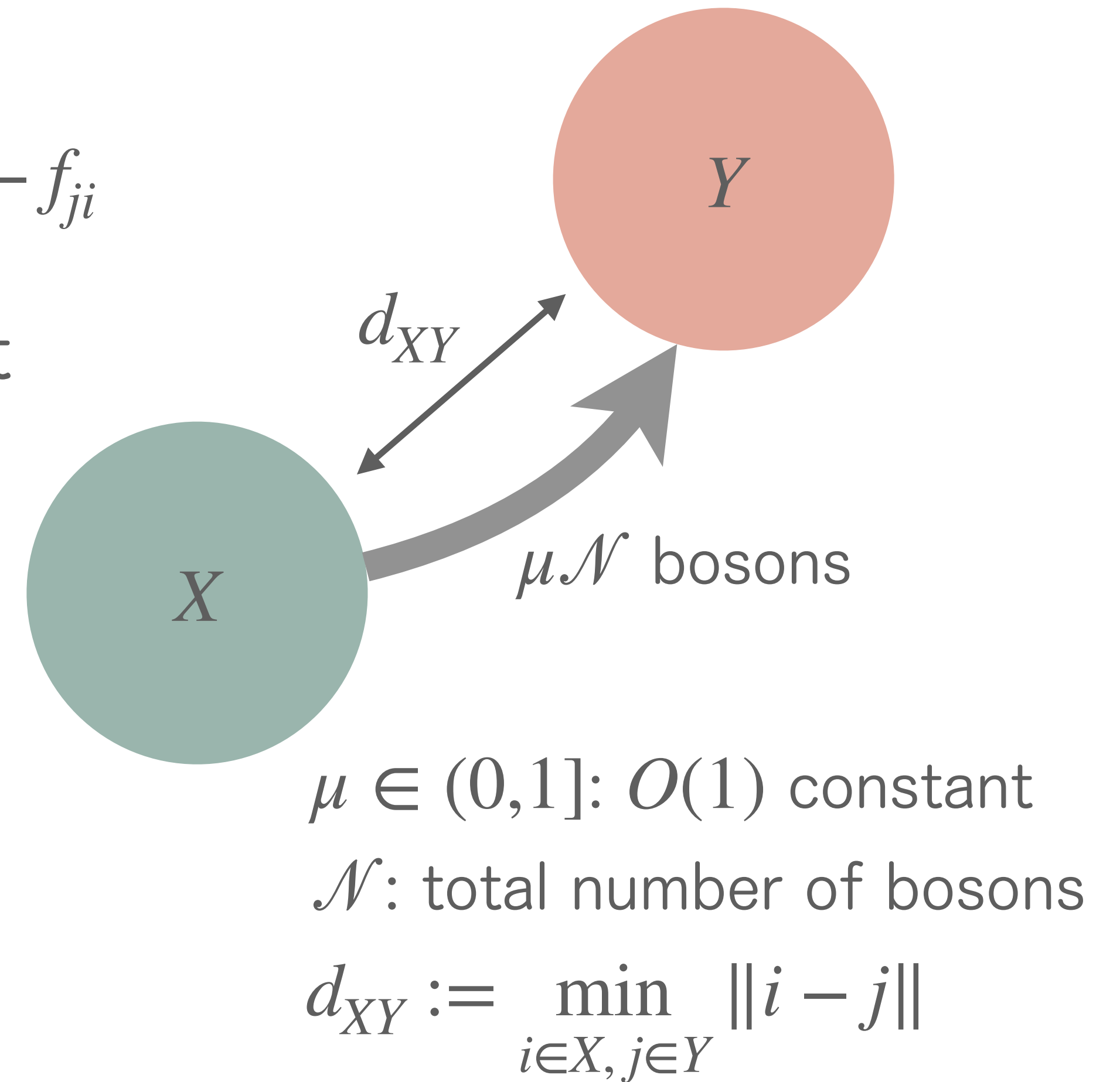
$$x_i(t) := \mathcal{N}^{-1} \text{tr} \{ \hat{n}_i \varrho_t \} \rightarrow \dot{x}_i(t) = \sum_{j(\neq i)} f_{ij}(t), \quad f_{ij} = -f_{ji}$$

- Finite speed of macroscopic particle transport

$$\text{cost matrix } c_{ij} = \|i - j\|^{\min(1, \alpha - D - \epsilon)}$$

$$\tau \geq \frac{\mathcal{W}(x_0, x_\tau)}{\bar{v}} \longrightarrow$$

Vu, Kuwahara, and Saito, Quantum (2024)



# Resolution of particle transport

- Boson concentrations at each site

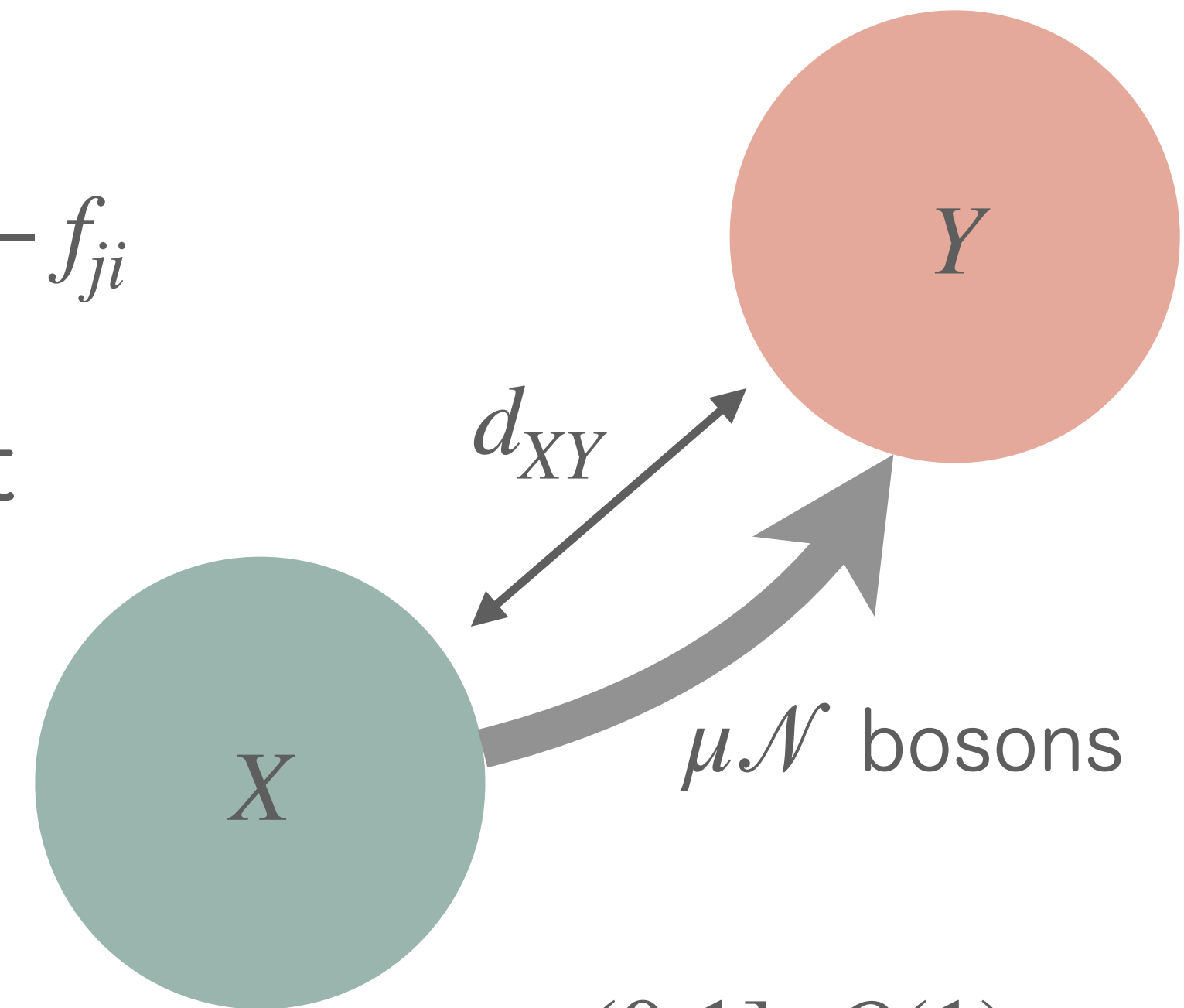
$$x_i(t) := \mathcal{N}^{-1} \text{tr}\{\hat{n}_i \varrho_t\} \rightarrow \dot{x}_i(t) = \sum_{j(\neq i)} f_{ij}(t), \quad f_{ij} = -f_{ji}$$

- Finite speed of macroscopic particle transport

$$\text{cost matrix } c_{ij} = \|i - j\|^{\min(1, \alpha - D - \epsilon)}$$

$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}} \longrightarrow \tau \gtrsim d_{XY}^{\min(1, \alpha - D)}$$

Vu, Kuwahara, and Saito, Quantum (2024)



$\mu \in (0, 1]$ :  $O(1)$  constant

$\mathcal{N}$ : total number of bosons

$$d_{XY} := \min_{i \in X, j \in Y} \|i - j\|$$

$x \gtrsim y$  means  $x \geq O(1) \times y$

# Resolution of particle transport

- Boson concentrations at each site

$$x_i(t) := \mathcal{N}^{-1} \text{tr}\{\hat{n}_i \mathcal{Q}_t\} \rightarrow \dot{x}_i(t) = \sum_{j(\neq i)} f_{ij}(t), \quad f_{ij} = -f_{ji}$$

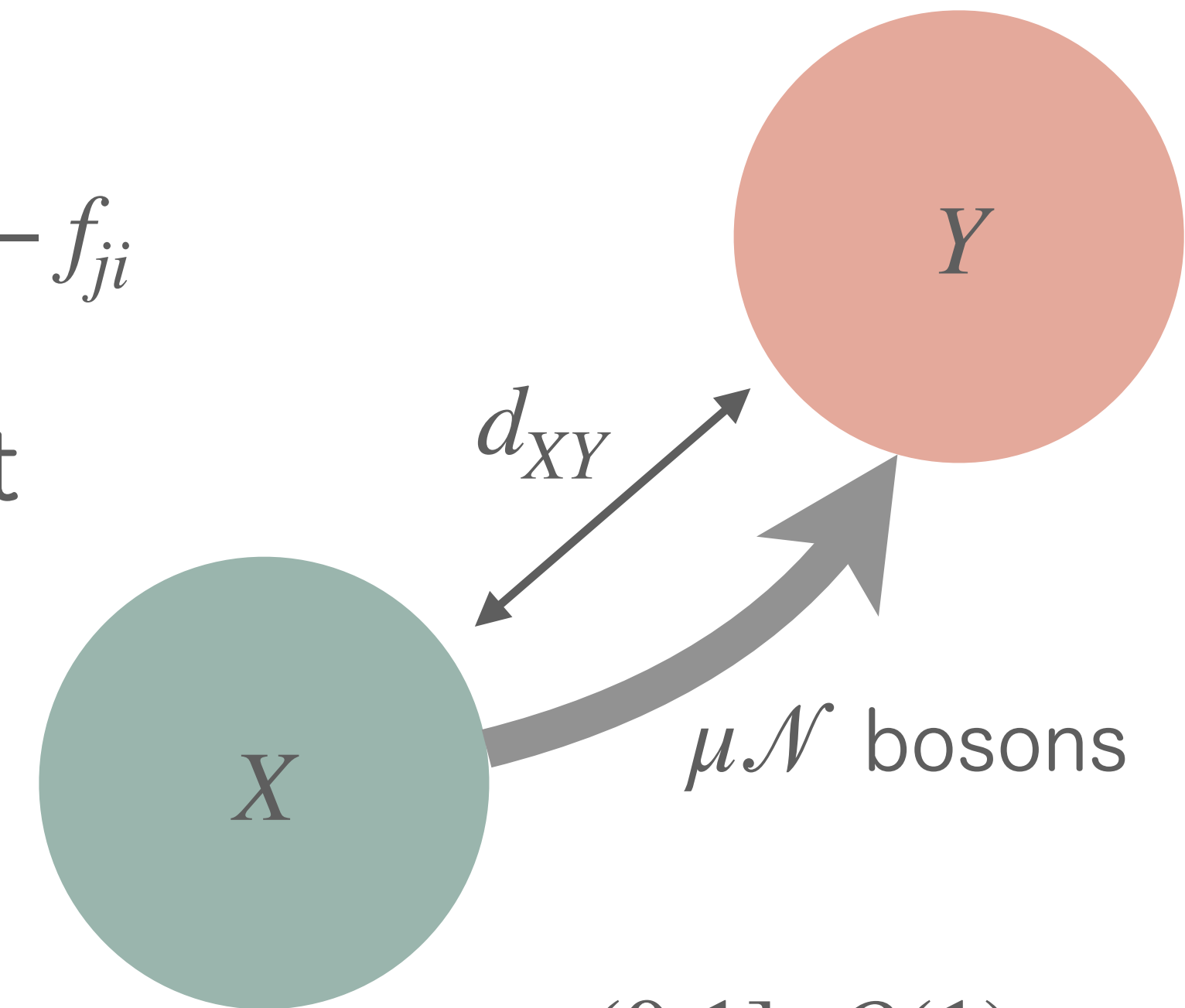
- Finite speed of macroscopic particle transport

$$\text{cost matrix } c_{ij} = \|i - j\|^{\min(1, \alpha - D - \epsilon)}$$

$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}} \longrightarrow \tau \gtrsim d_{XY}^{\min(1, \alpha - D)}$$

Vu, Kuwahara, and Saito, Quantum (2024)

$$\tau \gtrsim \begin{cases} d_{XY} & \text{if } \alpha > D + 1 \\ d_{XY}^{\alpha - D} & \text{if } D + 1 \geq \alpha > D \end{cases}$$



$\mu \in (0, 1]$ :  $O(1)$  constant

$\mathcal{N}$ : total number of bosons

$$d_{XY} := \min_{i \in X, j \in Y} \|i - j\|$$

$x \gtrsim y$  means  $x \geq O(1) \times y$

# Resolution of particle transport

- Boson concentrations at each site

$$x_i(t) := \mathcal{N}^{-1} \text{tr}\{\hat{n}_i \varrho_t\} \rightarrow \dot{x}_i(t) = \sum_{j(\neq i)} f_{ij}(t), \quad f_{ij} = -f_{ji}$$

- Finite speed of macroscopic particle transport

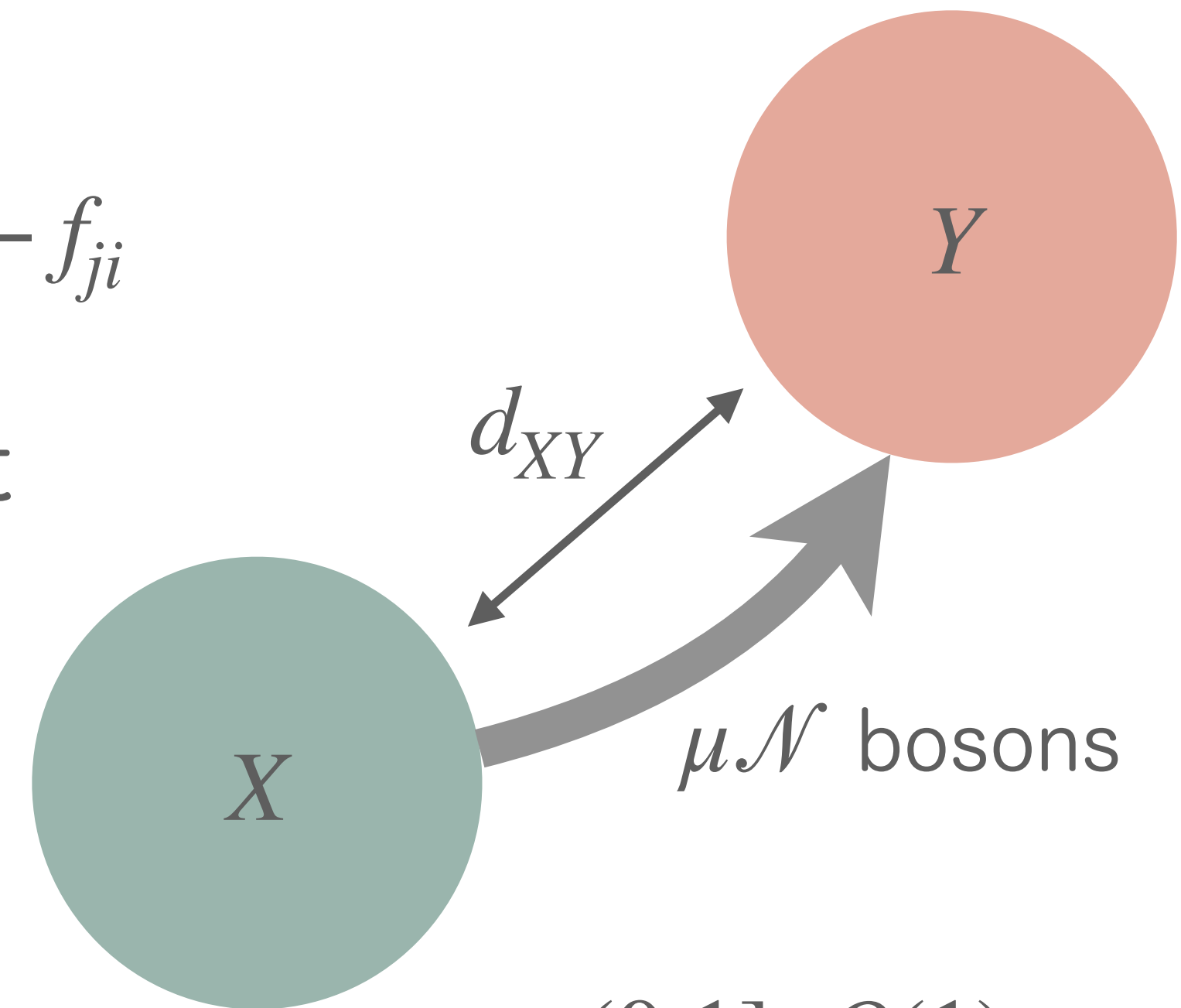
$$\text{cost matrix } c_{ij} = \|i - j\|^{\min(1, \alpha - D - \epsilon)}$$

$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}} \longrightarrow \tau \gtrsim d_{XY}^{\min(1, \alpha - D)}$$

Vu, Kuwahara, and Saito, Quantum (2024)

✓ bound is optimal

$$\tau \gtrsim \begin{cases} d_{XY} & \text{if } \alpha > D + 1 \\ d_{XY}^{\alpha - D} & \text{if } D + 1 \geq \alpha > D \end{cases}$$



$\mu \in (0, 1]$ :  $O(1)$  constant

$\mathcal{N}$ : total number of bosons

$$d_{XY} := \min_{i \in X, j \in Y} \|i - j\|$$

$x \gtrsim y$  means  $x \geq O(1) \times y$

# Resolution of particle transport

- Boson concentrations at each site

$$x_i(t) := \mathcal{N}^{-1} \text{tr} \{ \hat{n}_i \varrho_t \} \rightarrow \dot{x}_i(t) = \sum_{j(\neq i)} f_{ij}(t), \quad f_{ij} = -f_{ji}$$

- Finite speed of macroscopic particle transport

$$\text{cost matrix } c_{ij} = \|i - j\|^{\min(1, \alpha - D - \epsilon)}$$

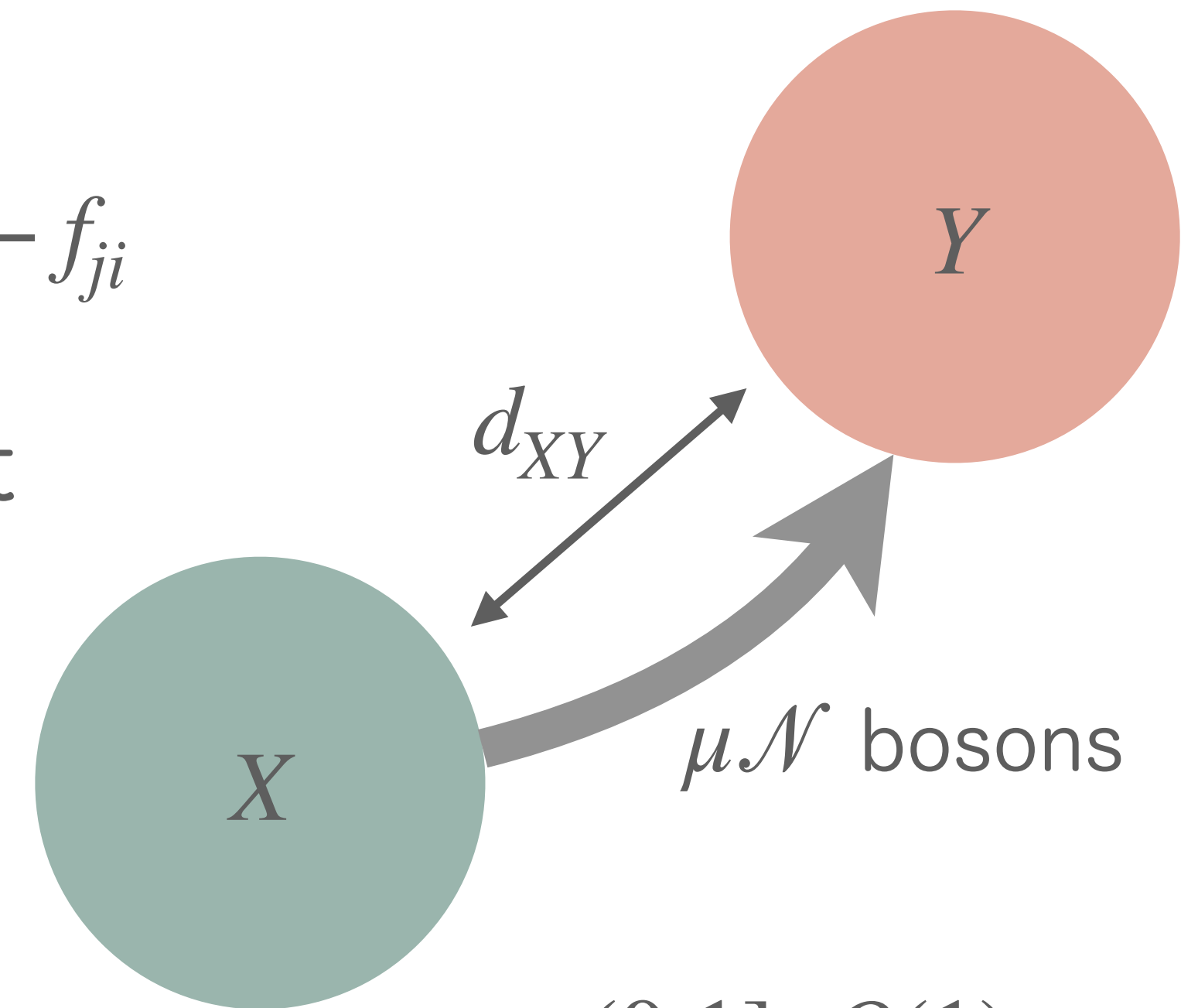
$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}} \longrightarrow \tau \gtrsim d_{XY}^{\min(1, \alpha - D)}$$

Vu, Kuwahara, and Saito, Quantum (2024)

$$\tau \gtrsim \begin{cases} d_{XY} & \text{if } \alpha > D + 1 \\ d_{XY}^{\alpha - D} & \text{if } D + 1 \geq \alpha > D \end{cases}$$

✓ bound is optimal

✓ resolve open problem of macroscopic particle transport



$\mu \in (0, 1]$ :  $O(1)$  constant

$\mathcal{N}$ : total number of bosons

$$d_{XY} := \min_{i \in X, j \in Y} \|i - j\|$$

$x \gtrsim y$  means  $x \geq O(1) \times y$

# Sketch proof

cost matrix  $c_{ij} = \|i - j\|^{\alpha_\epsilon}$

$$\alpha_\epsilon := \min(1, \alpha - D - \epsilon)$$

$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}} \longrightarrow \tau \gtrsim d_{XY}^{\min(1, \alpha - D)}$$

# Sketch proof

cost matrix  $c_{ij} = \|i - j\|^{\alpha_\epsilon}$

$$\alpha_\epsilon := \min(1, \alpha - D - \epsilon)$$

$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}} \longrightarrow \tau \gtrsim d_{XY}^{\min(1, \alpha - D)}$$

- Observation #1: Wasserstein distance is bounded by transport distance

$$n_Y(\tau) \geq n_{X^c}(0) + \mu \mathcal{N} \Rightarrow \mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau) \geq \mu d_{XY}^{\alpha_\epsilon}$$

# Sketch proof

cost matrix  $c_{ij} = \|i - j\|^{\alpha_\epsilon}$

$$\alpha_\epsilon := \min(1, \alpha - D - \epsilon)$$

$$\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}} \longrightarrow \tau \gtrsim d_{XY}^{\min(1, \alpha - D)}$$

- Observation #1: Wasserstein distance is bounded by transport distance

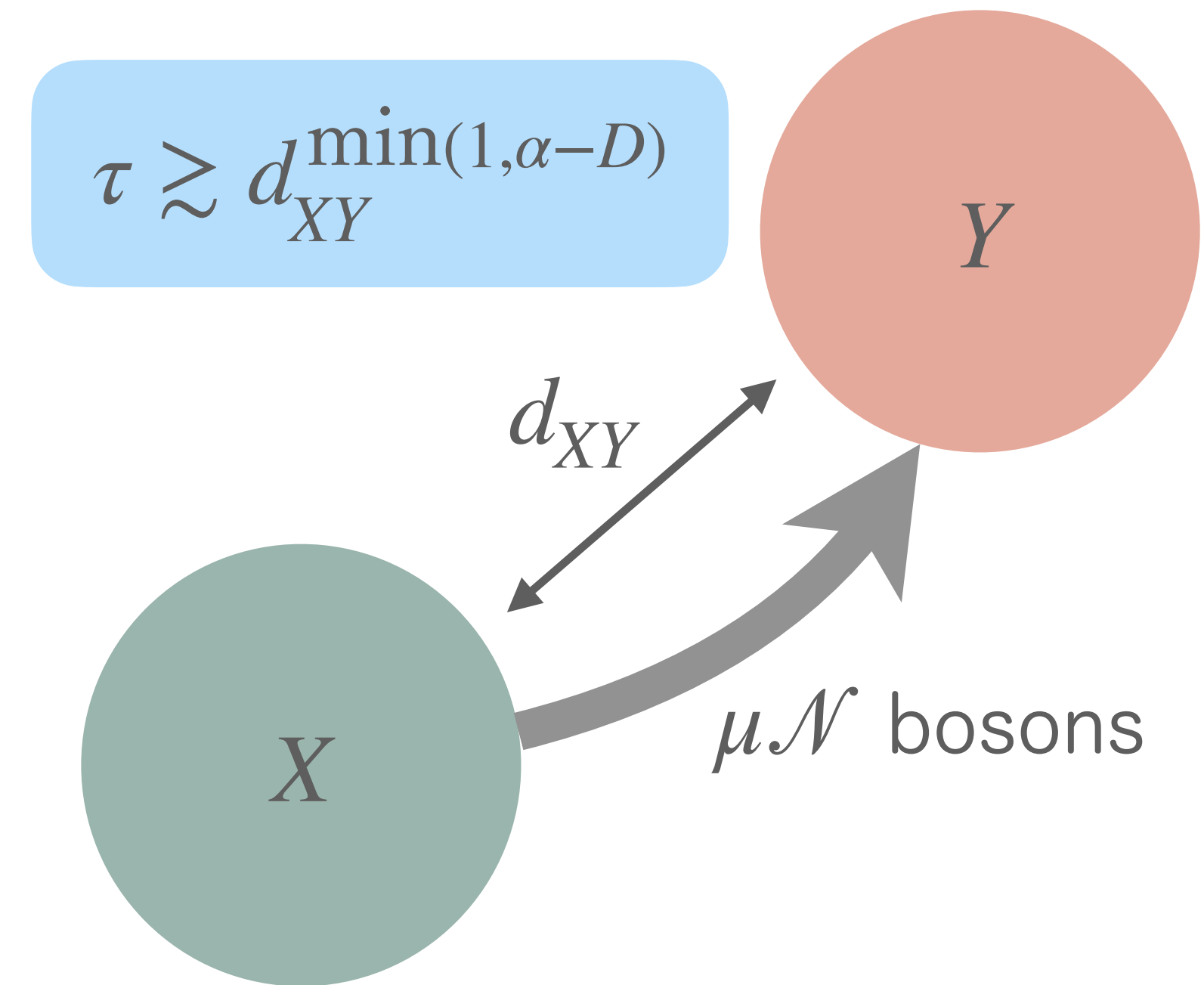
$$n_Y(\tau) \geq n_{X^c}(0) + \mu \mathcal{N} \Rightarrow \mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau) \geq \mu d_{XY}^{\alpha_\epsilon}$$

- Observation #2: velocity  $\bar{v}$  is finite

$$\bar{v} \leq J\gamma\zeta(\alpha - \alpha_\epsilon - D + 1)$$

$\zeta$  : zeta function

# Beyond Bose-Hubbard-type models



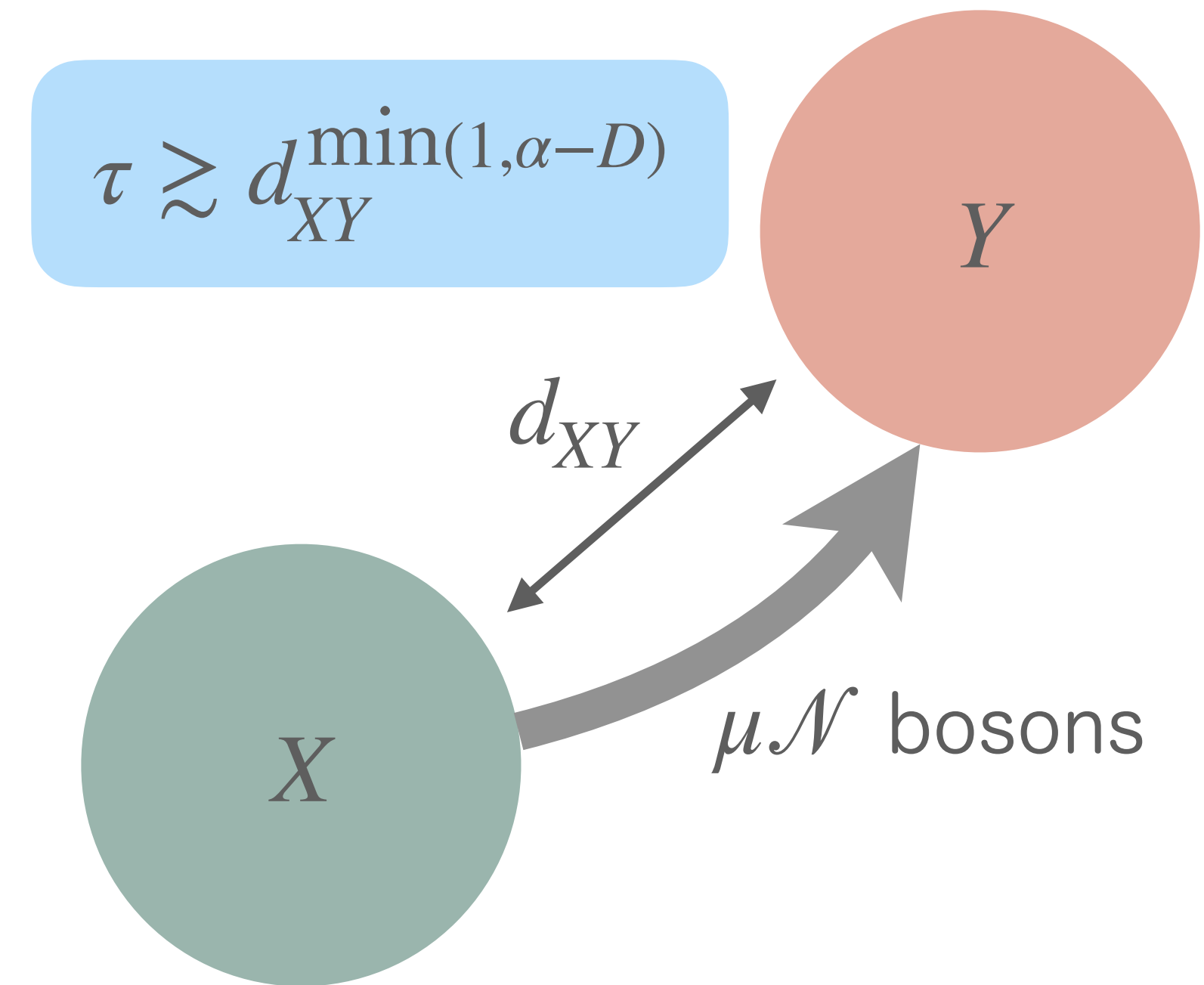
$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_Z h_Z(t)$$

# Beyond Bose-Hubbard-type models

- Interaction-induced tunneling terms

Sowinski+, Phys. Rev. Lett. (2012)

$$H_t = \sum_{|i-j|=1} J_{ij}(t) \hat{b}_i^\dagger \hat{b}_j + \sum_{|i-j|=1} T_{ij}(t) \hat{b}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{b}_j \\ + \frac{1}{2} \sum_i U_i(t) \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i(t) \hat{n}_i$$



$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_Z h_Z(t)$$

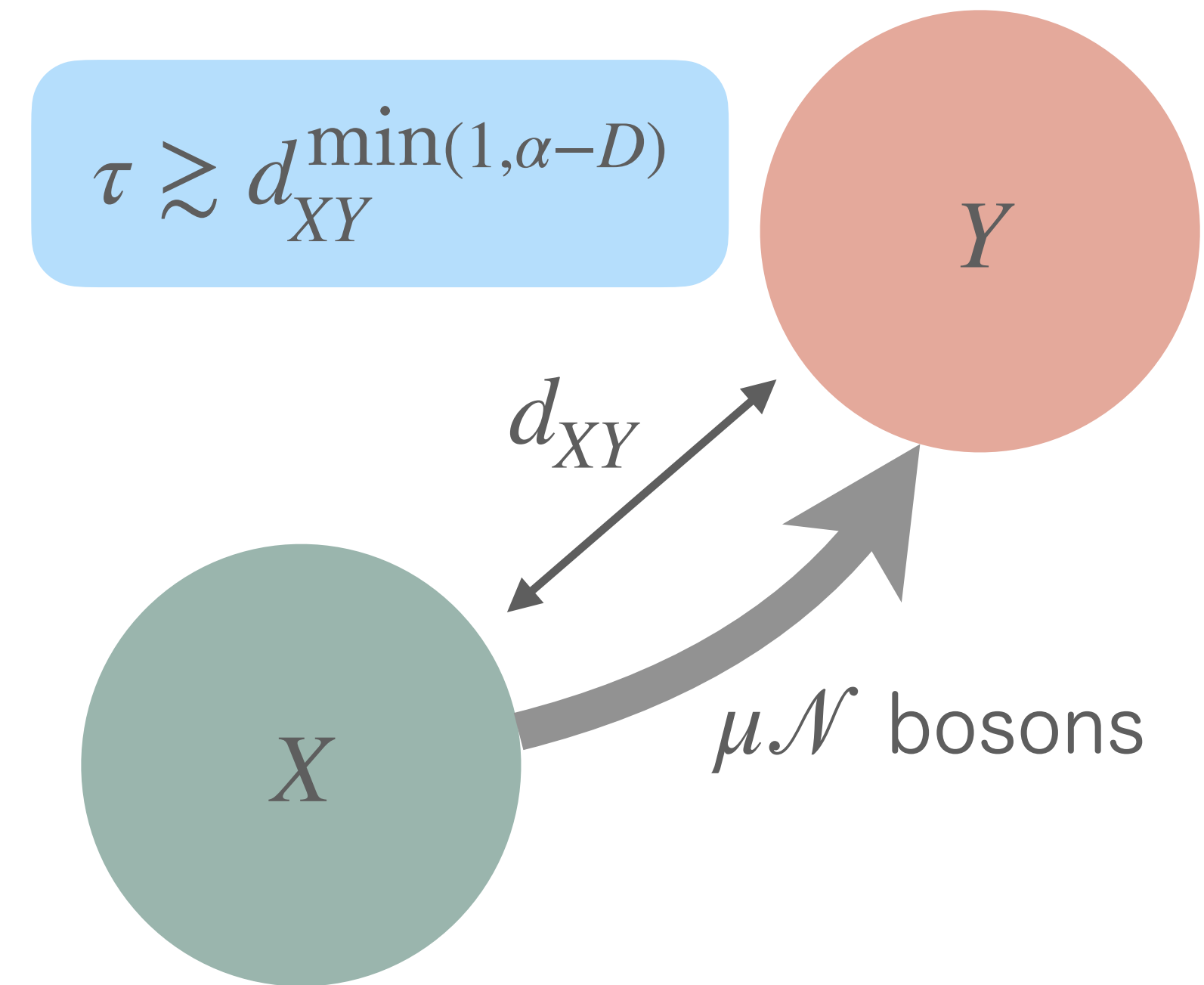
# Beyond Bose-Hubbard-type models

- Interaction-induced tunneling terms

Sowinski+, Phys. Rev. Lett. (2012)

$$H_t = \sum_{|i-j|=1} J_{ij}(t) \hat{b}_i^\dagger \hat{b}_j + \sum_{|i-j|=1} T_{ij}(t) \hat{b}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{b}_j \\ + \frac{1}{2} \sum_i U_i(t) \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i(t) \hat{n}_i$$

- Speed of macroscopic particle transport?



$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_Z h_Z(t)$$

# Beyond Bose-Hubbard-type models

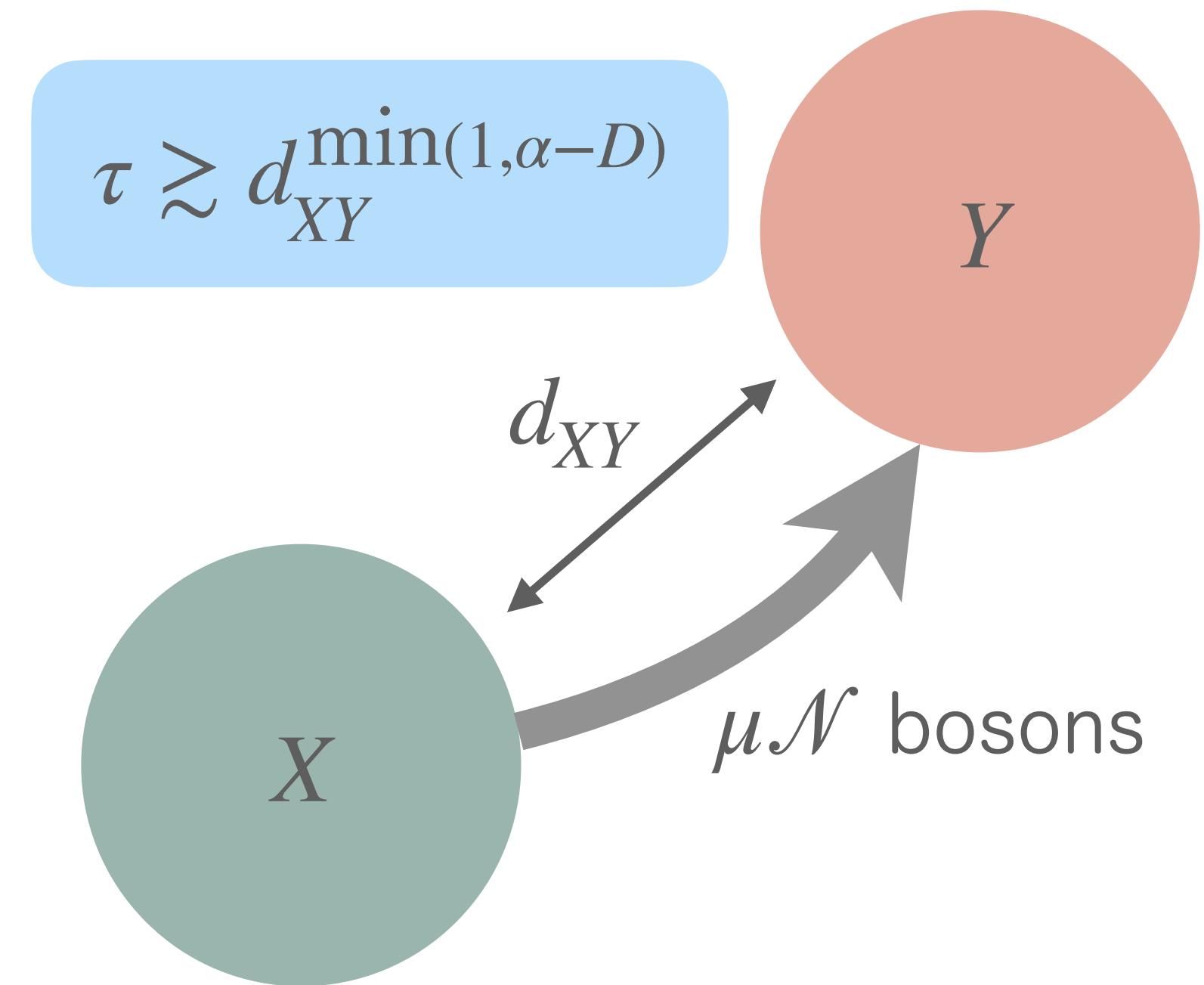
- Interaction-induced tunneling terms

Sowinski+, Phys. Rev. Lett. (2012)

$$H_t = \sum_{|i-j|=1} J_{ij}(t) \hat{b}_i^\dagger \hat{b}_j + \sum_{|i-j|=1} T_{ij}(t) \hat{b}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{b}_j \\ + \frac{1}{2} \sum_i U_i(t) \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i(t) \hat{n}_i$$

- Speed of macroscopic particle transport?

✓ speed can be infinite



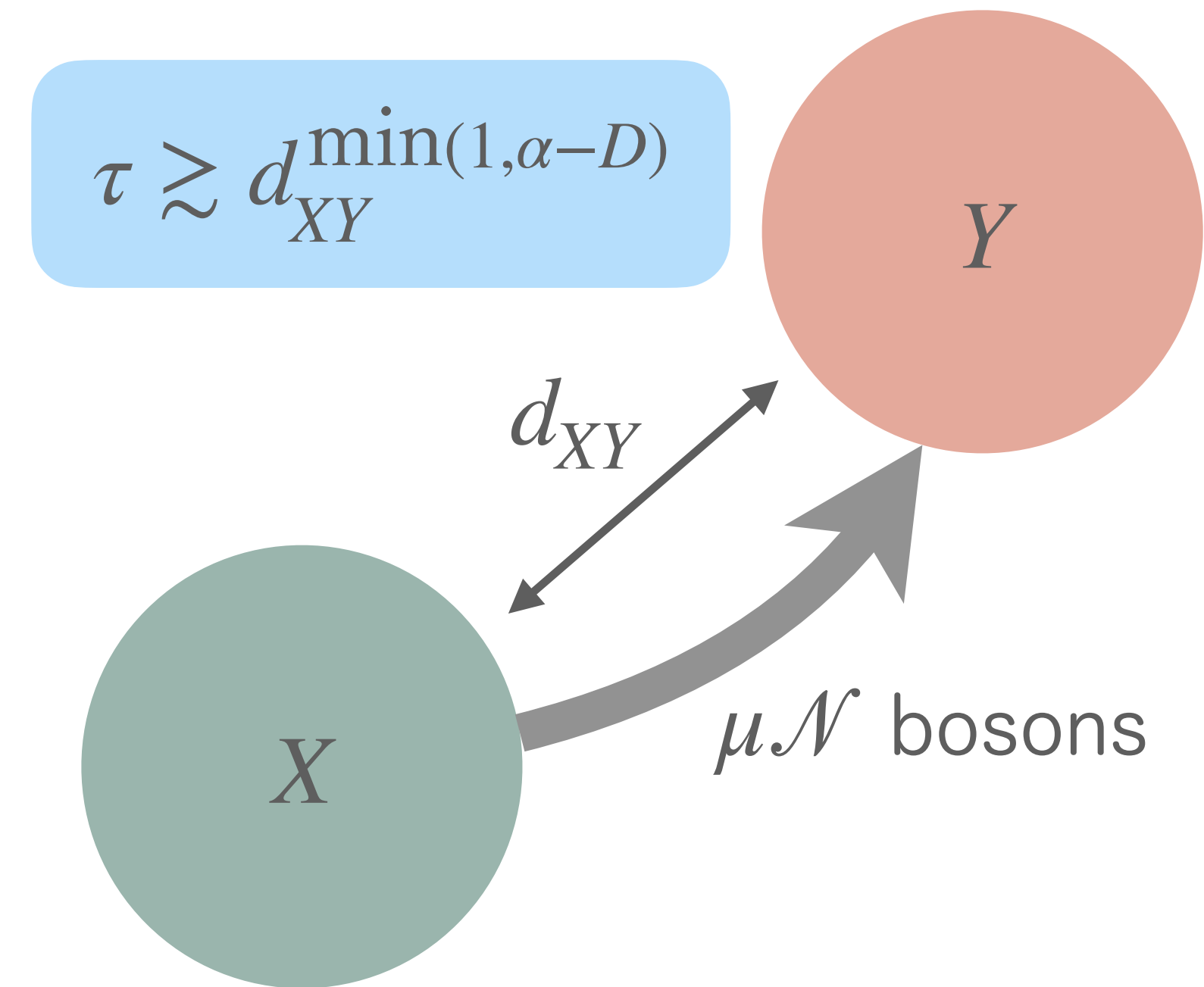
$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_Z h_Z(t)$$

# Beyond Bose-Hubbard-type models

- Interaction-induced tunneling terms

Sowinski+, Phys. Rev. Lett. (2012)

$$H_t = \sum_{|i-j|=1} J_{ij}(t) \hat{b}_i^\dagger \hat{b}_j + \sum_{|i-j|=1} T_{ij}(t) \hat{b}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{b}_j \\ + \frac{1}{2} \sum_i U_i(t) \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i(t) \hat{n}_i$$



- Speed of macroscopic particle transport?

✓ speed can be infinite

→ all bosons can be transported within a constant of time

$$H_t = \sum_{i \neq j} J_{ij}(t) \hat{b}_i \hat{b}_j^\dagger + \sum_Z h_Z(t)$$

# Beyond Bose-Hubbard-type models

- Initial Mott state

$$|L\rangle \otimes |0\rangle \otimes \dots \otimes |0\rangle \rightarrow |0\rangle \otimes |0\rangle \otimes \dots \otimes |L\rangle$$

$$H_t = \sum_{|i-j|=1} J_{ij}(t) \hat{b}_i^\dagger \hat{b}_j + \sum_{|i-j|=1} T_{ij}(t) \hat{b}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{b}_j \\ + \frac{1}{2} \sum_i U_i(t) \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i(t) \hat{n}_i$$

# Beyond Bose-Hubbard-type models

- Initial Mott state

$$|L\rangle \otimes |0\rangle \otimes \dots \otimes |0\rangle \rightarrow |0\rangle \otimes |0\rangle \otimes \dots \otimes |L\rangle$$

- Protocol for transporting bosons

$$\text{step } k = 1, \dots, L-1 : |L\rangle_k \otimes |0\rangle_{k+1} \rightarrow |0\rangle_k \otimes |L\rangle_{k+1}$$

$$H_t = \sum_{|i-j|=1} J_{ij}(t) \hat{b}_i^\dagger \hat{b}_j + \sum_{|i-j|=1} T_{ij}(t) \hat{b}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{b}_j \\ + \frac{1}{2} \sum_i U_i(t) \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i(t) \hat{n}_i$$

# Beyond Bose-Hubbard-type models

- Initial Mott state

$$|L\rangle \otimes |0\rangle \otimes \dots \otimes |0\rangle \rightarrow |0\rangle \otimes |0\rangle \otimes \dots \otimes |L\rangle$$

- Protocol for transporting bosons

$$\text{step } k = 1, \dots, L-1 : |L\rangle_k \otimes |0\rangle_{k+1} \rightarrow |0\rangle_k \otimes |L\rangle_{k+1}$$

3 stages

$$|L\rangle_k \otimes |0\rangle_{k+1} \rightarrow |L-1\rangle_k \otimes |1\rangle_{k+1} \rightarrow |1\rangle_k \otimes |L-1\rangle_{k+1} \rightarrow |0\rangle_k \otimes |L\rangle_{k+1}$$

$$H_t = \sum_{|i-j|=1} J_{ij}(t) \hat{b}_i^\dagger \hat{b}_j + \sum_{|i-j|=1} T_{ij}(t) \hat{b}_i^\dagger (\hat{n}_i + \hat{n}_j) \hat{b}_j \\ + \frac{1}{2} \sum_i U_i(t) \hat{n}_i (\hat{n}_i - 1) + \sum_i \mu_i(t) \hat{n}_i$$

# Beyond Bose-Hubbard-type models

- Protocol for each stage

# Beyond Bose-Hubbard-type models

- Protocol for each stage

$|L-1\rangle_1 \otimes |1\rangle_2 \rightarrow |1\rangle_1 \otimes |L-1\rangle_2$  within time  $t = \pi/(2JL)$

$$H = J(\hat{b}_1^\dagger \hat{b}_2 + \hat{b}_2^\dagger \hat{b}_1) + J[\hat{b}_2^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_1 + \hat{b}_1^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_2]$$

# Beyond Bose-Hubbard-type models

- Protocol for each stage

$$|L-1\rangle_1 \otimes |1\rangle_2 \rightarrow |1\rangle_1 \otimes |L-1\rangle_2 \text{ within time } t = \pi/(2JL)$$

$$H = J(\hat{b}_1^\dagger \hat{b}_2 + \hat{b}_2^\dagger \hat{b}_1) + J[\hat{b}_2^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_1 + \hat{b}_1^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_2]$$

$$|1\rangle_1 \otimes |L-1\rangle_2 \leftrightarrow |0\rangle_1 \otimes |L\rangle_2 \text{ within time } t = \pi/(2JL\sqrt{L})$$

$$H = J(\hat{b}_1^\dagger \hat{b}_2 + \hat{b}_2^\dagger \hat{b}_1) + J[\hat{b}_2^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_1 + \hat{b}_1^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_2] \\ - U\hat{n}_2^2 + U(2L-1)\hat{n}_2 \quad (U \rightarrow +\infty)$$

# Beyond Bose-Hubbard-type models

- Protocol for each stage

$$|L-1\rangle_1 \otimes |1\rangle_2 \rightarrow |1\rangle_1 \otimes |L-1\rangle_2 \text{ within time } t = \pi/(2JL)$$

$$H = J(\hat{b}_1^\dagger \hat{b}_2 + \hat{b}_2^\dagger \hat{b}_1) + J[\hat{b}_2^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_1 + \hat{b}_1^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_2]$$

$$|1\rangle_1 \otimes |L-1\rangle_2 \leftrightarrow |0\rangle_1 \otimes |L\rangle_2 \text{ within time } t = \pi/(2JL\sqrt{L})$$

$$H = J(\hat{b}_1^\dagger \hat{b}_2 + \hat{b}_2^\dagger \hat{b}_1) + J[\hat{b}_2^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_1 + \hat{b}_1^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_2] \\ -U\hat{n}_2^2 + U(2L-1)\hat{n}_2 \quad (U \rightarrow +\infty)$$

$$\text{total time } \tau = \frac{\pi}{2J} \sum_{k=1}^{L-1} \left( \frac{1}{L} + \frac{2}{L\sqrt{L}} \right) < \frac{\pi}{J}$$

# Beyond Bose-Hubbard-type models

- Protocol for each stage

$$|L-1\rangle_1 \otimes |1\rangle_2 \rightarrow |1\rangle_1 \otimes |L-1\rangle_2 \text{ within time } t = \pi/(2JL)$$

$$H = J(\hat{b}_1^\dagger \hat{b}_2 + \hat{b}_2^\dagger \hat{b}_1) + J[\hat{b}_2^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_1 + \hat{b}_1^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_2]$$

$$|1\rangle_1 \otimes |L-1\rangle_2 \leftrightarrow |0\rangle_1 \otimes |L\rangle_2 \text{ within time } t = \pi/(2JL\sqrt{L})$$

$$H = J(\hat{b}_1^\dagger \hat{b}_2 + \hat{b}_2^\dagger \hat{b}_1) + J[\hat{b}_2^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_1 + \hat{b}_1^\dagger(\hat{n}_2 + \hat{n}_1)\hat{b}_2] \\ - U\hat{n}_2^2 + U(2L-1)\hat{n}_2 \quad (U \rightarrow +\infty)$$

$$\text{total time } \tau = \frac{\pi}{2J} \sum_{k=1}^{L-1} \left( \frac{1}{L} + \frac{2}{L\sqrt{L}} \right) < \frac{\pi}{J}$$

$|L\rangle \otimes |0\rangle \otimes \dots \otimes |0\rangle \rightarrow |0\rangle \otimes |0\rangle \otimes \dots \otimes |L\rangle \Rightarrow$  transport to distant place within a constant of time

# Summary

- Reveal maximal speed of macroscopic particle transport using optimal transport theory for both closed systems

General speed limit  $\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}}$  

Resolution of  
bosonic transport

$$\tau \gtrsim d_{XY}^\lambda$$

# Summary

- Reveal maximal speed of macroscopic particle transport using optimal transport theory for both closed systems

General speed limit  $\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}}$  

Resolution of  
bosonic transport

$$\tau \gtrsim d_{XY}^\lambda$$

- Future works: exploit full quantum states to uncover new constraints

# Summary

- Reveal maximal speed of macroscopic particle transport using optimal transport theory for both closed systems

General speed limit  $\tau \geq \frac{\mathcal{W}(\mathbf{x}_0, \mathbf{x}_\tau)}{\bar{v}}$  

Resolution of  
bosonic transport

$$\tau \gtrsim d_{XY}^\lambda$$

- Future works: exploit full quantum states to uncover new constraints

Hongchao's talk on generalization to open systems (next Monday)  
Li, Shang, Kuwahara, and Vu, arXiv:2503.13731

Thank you for your attention!