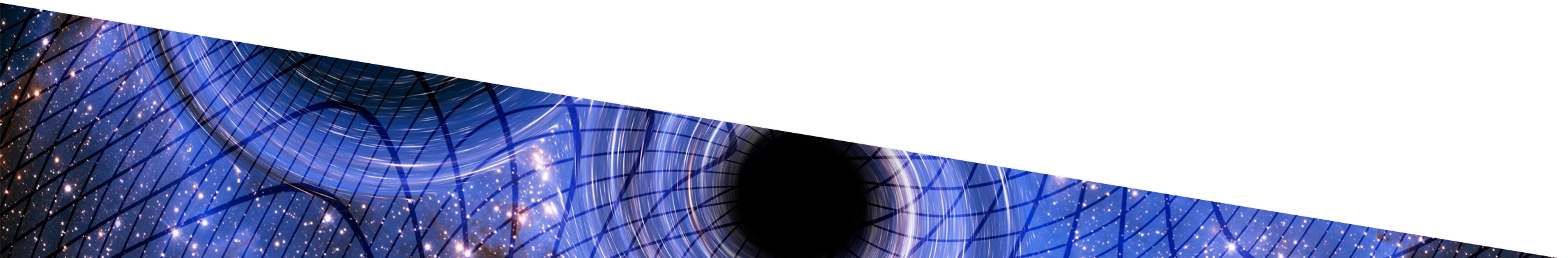
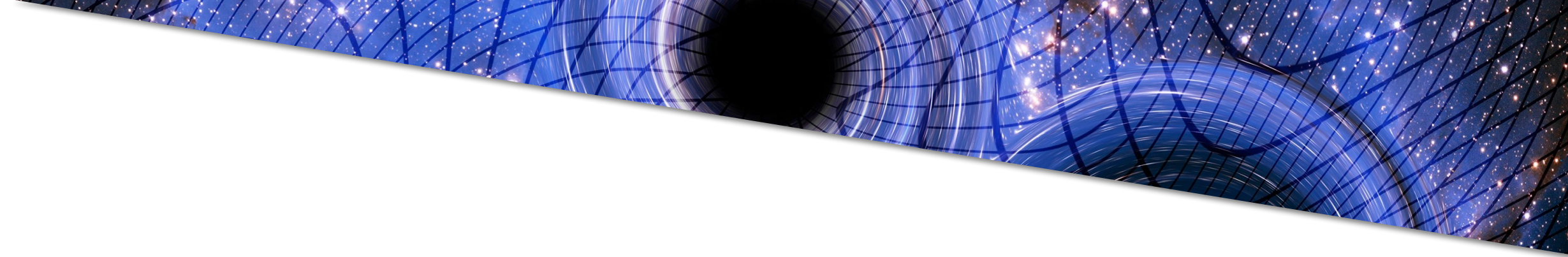


WILL WE EVER DETECT HIGH FREQUENCY GRAVITATIONAL WAVES?

Raffaele Tito D'Agnolo - CEA IPhT Saclay and ENS Paris

[Based on RTD, S. Ellis, arXiv:2412.17897, *JHEP* 04 (2025) 164]


$$\omega_g$$



ω_g

nHz

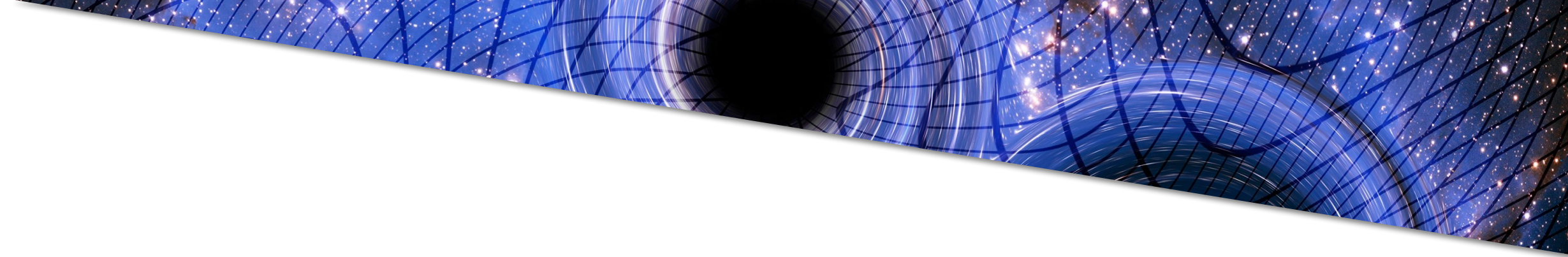
μHz

10 Hz

kHz

PULSAR TIMING

LIGO-VIRGO-KAGRA

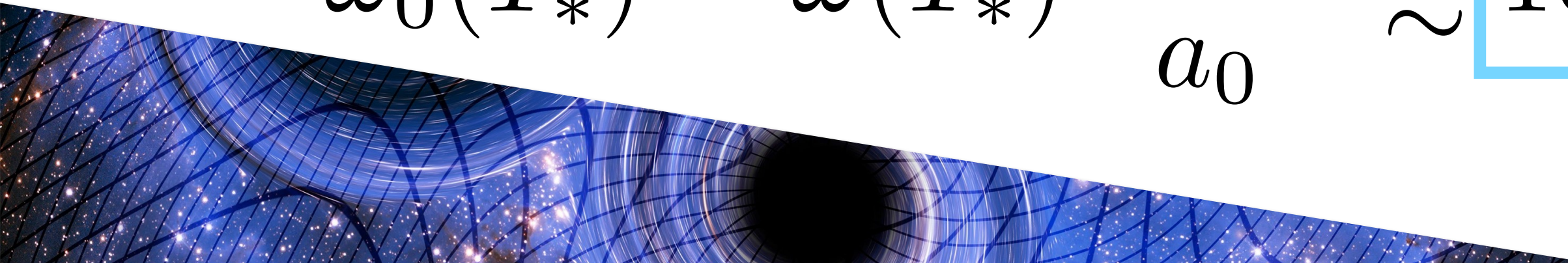

$$\omega_g$$

$$\lambda(T_*) < \frac{1}{H(T_*)} \quad \text{Causality}$$



$$\omega_g$$

$$\lambda(T_*) < \frac{1}{H(T_*)} \quad \text{Causality}$$

$$\omega_0(T_*) = \omega(T_*) \frac{a(T_*)}{a_0} \gtrsim \boxed{100 \text{ MHz}} \left(\frac{T_*}{10^{15} \text{ GeV}} \right) \left(\frac{g_*(T_*)}{100} \right)^{1/6}$$




DETECTORS

$$U_{\text{in}} \sim E_0^2 V_0$$

DETECTOR II

M



$$U_{\text{in}} \sim E_0^2 V_0$$

M



LASER



$$U_{\text{in}} \sim P_{\text{in}} \omega_L$$

$$M\ddot{\delta x} = M\ddot{h}x + F_{\text{ext}}$$

M



δx

$$M \rightarrow \infty$$

$$M\ddot{\delta x} = M\ddot{hx} + F_{\text{ext}}$$

M



δx

$$M \rightarrow \infty$$

$$M\ddot{\delta x} = M\ddot{h}x + F_{\text{ext}}$$

$$\ddot{\delta x}_{\text{sig}} \sim \ddot{h}x \sim \text{const}$$

 M  δx

$$M \rightarrow \infty$$

$$M\ddot{\delta x} = M\ddot{h}x + F_{\text{ext}}$$

$$\ddot{\delta x}_{\text{sig}} \sim \ddot{h}x \sim \text{const}$$

$$\ddot{\delta x}_{\text{noise}} \sim \frac{F_{\text{ext}}}{M} \rightarrow 0$$

 M  δx

IN THE FOLLOWING I WILL ALWAYS CONSIDER

$$M \rightarrow \infty$$

$$\ddot{\delta x}_{\text{noise}} \sim \frac{F_{\text{ext}}}{M} \rightarrow 0$$

THE BEST POSSIBLE SENSITIVITY

In this limit I can ignore noise from the test mass and focus on the signal (and noise) photons that I can detect

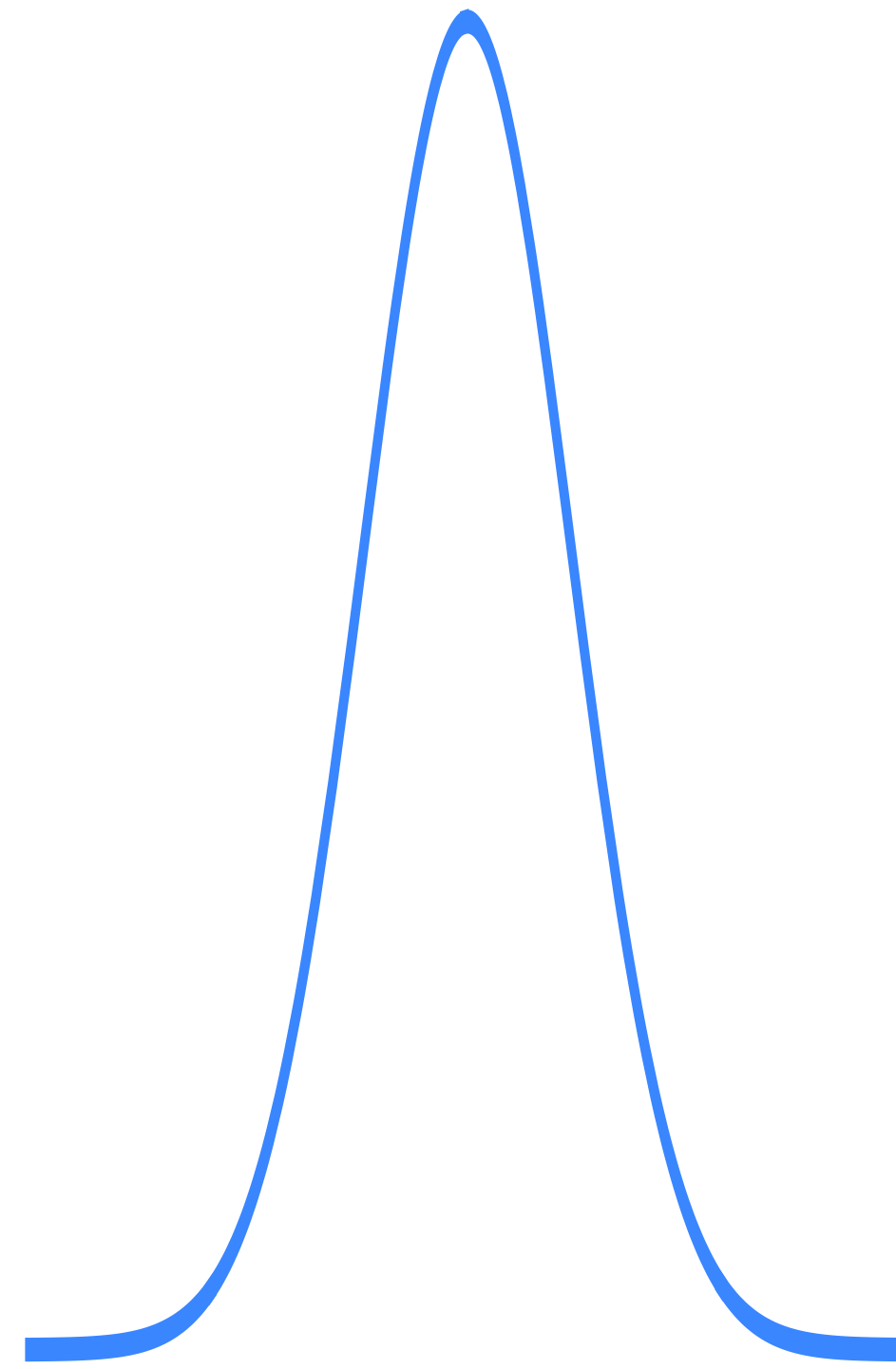
DETECTOR I



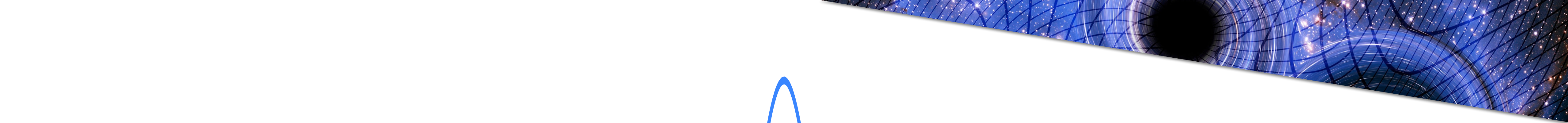
$$U_{\text{in}} \sim E_0^2 V_0$$

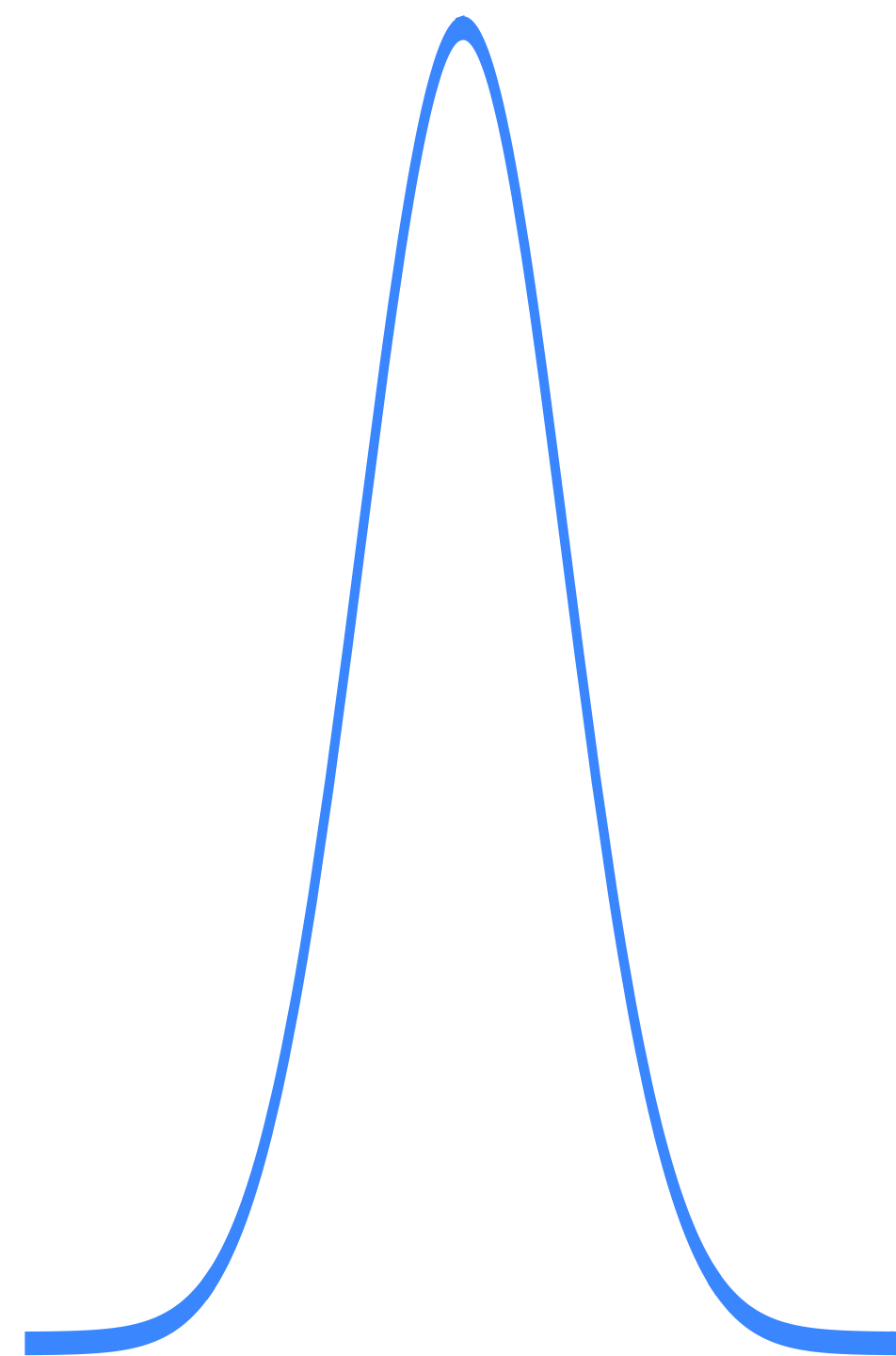
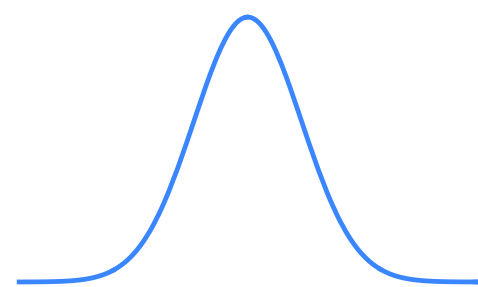


DETECTOR II

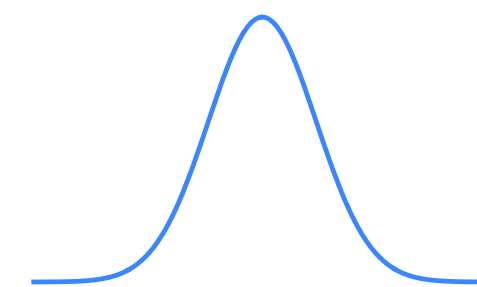


$$E_0(\omega_0)$$


$$\omega_0 - \omega_g$$

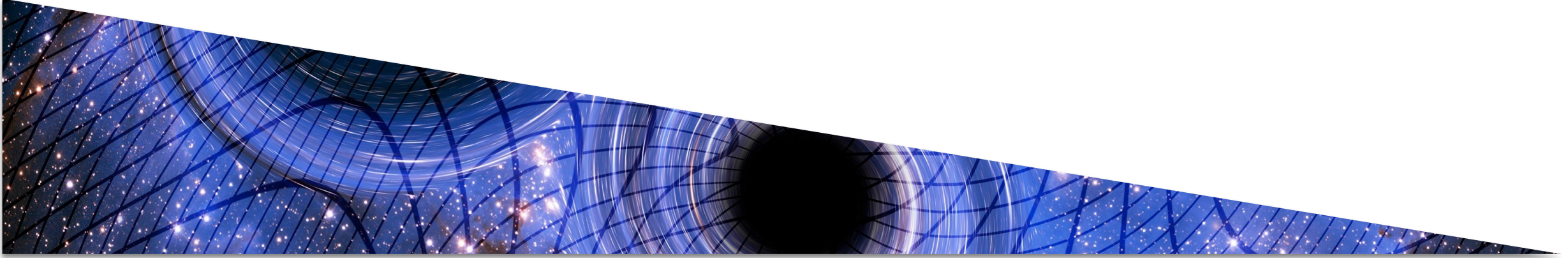


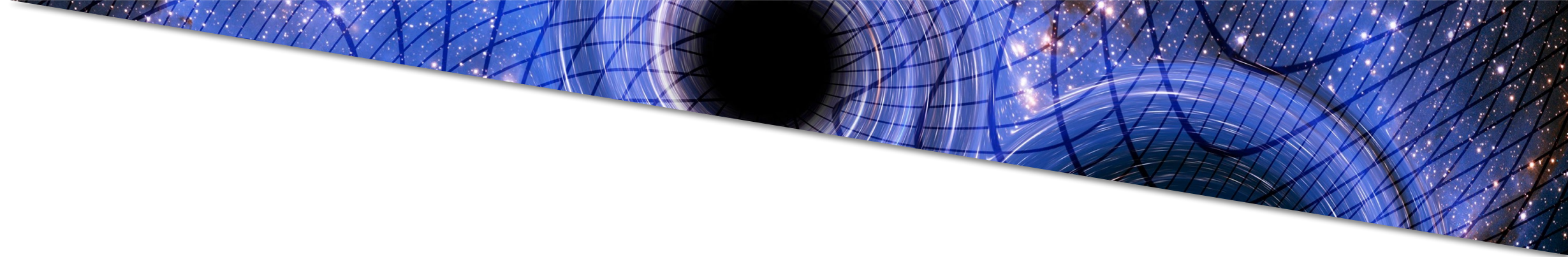
$$\omega_0 + \omega_g$$



$$E_h$$

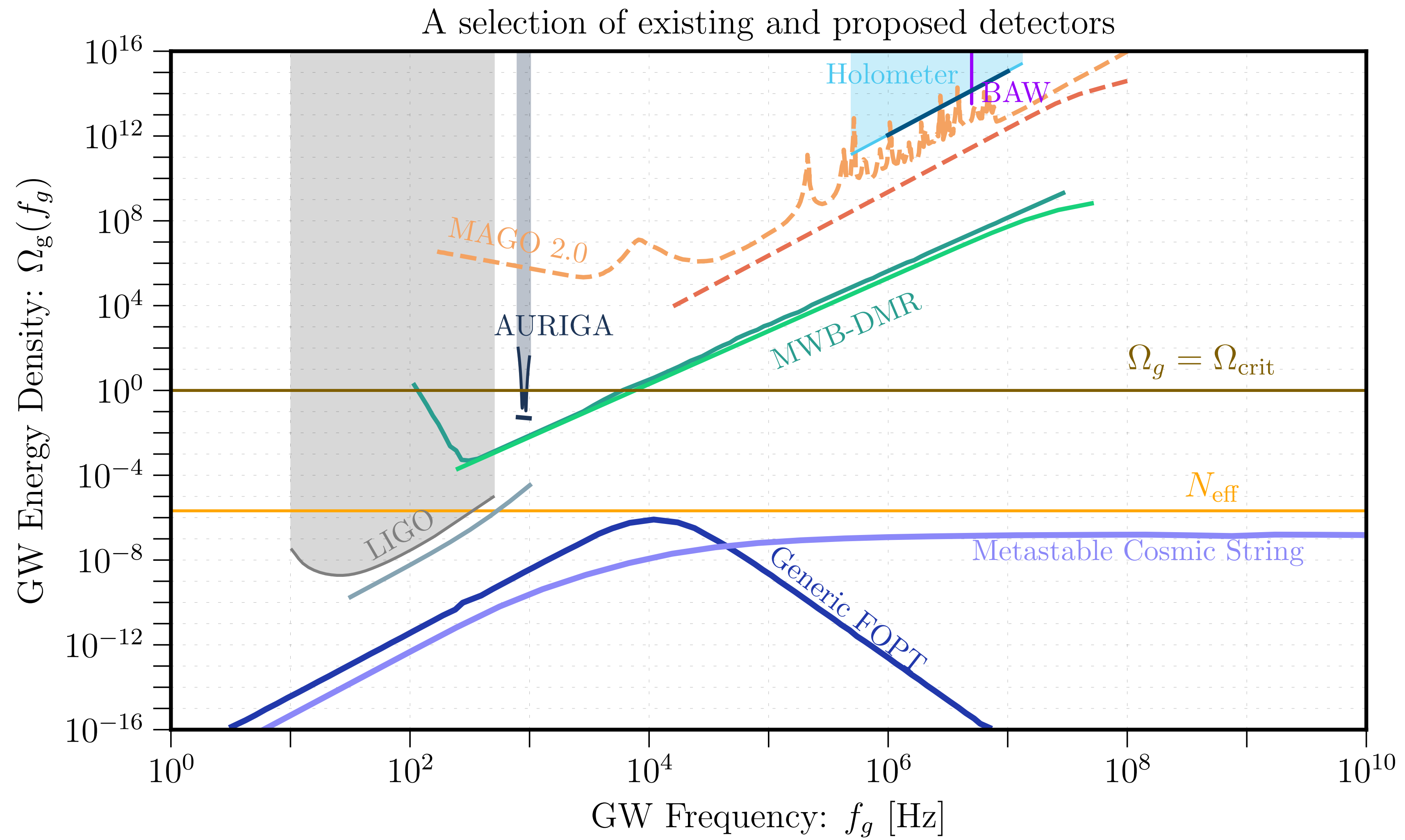
$$E_0(\omega_0)$$

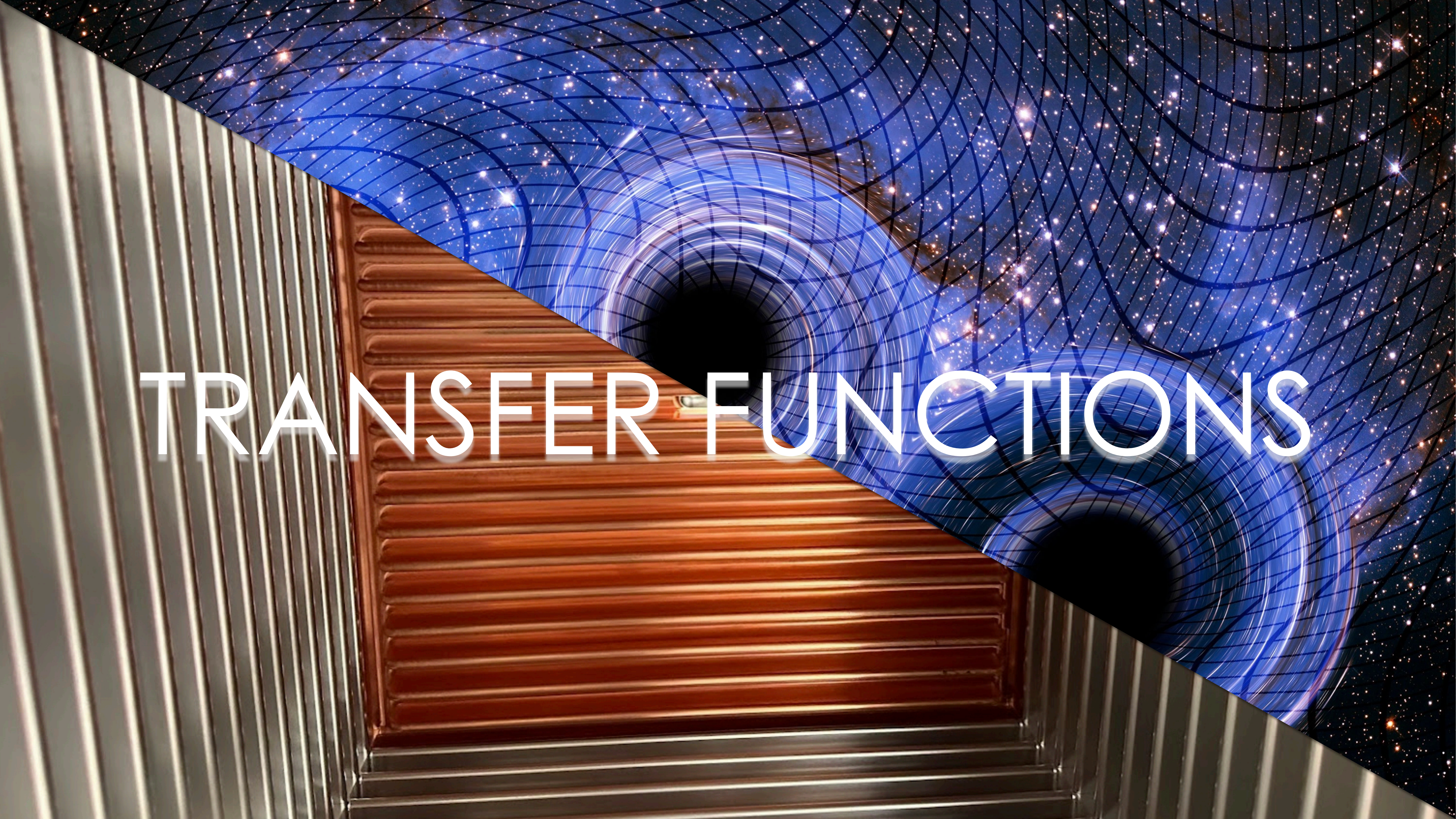
$$E_h$$



$$U_{\text{in}} \sim E_0^2 V_0$$

$$— \mathcal{T}(\omega) —$$







TRANSFER FUNCTIONS

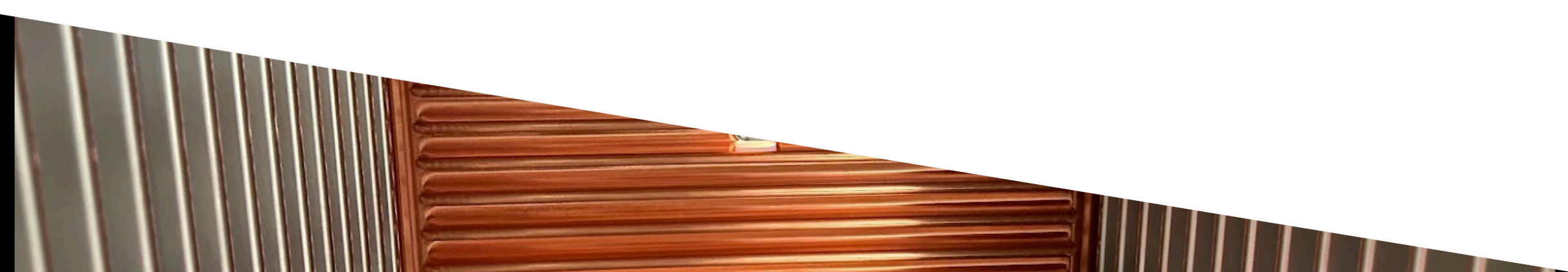


MOST DETECTORS

(LIGO, OPTOMECHANICAL, WEBER BARS...)

$$\left(\omega_m^2 - \omega^2 + i \frac{\omega \omega_m}{Q_m} \right) \tilde{u}_m(\omega) \simeq -\frac{1}{2} \omega_g^2 V^{1/3} \tilde{h}(\omega, \omega_g)$$

The wave excites a mechanical mode





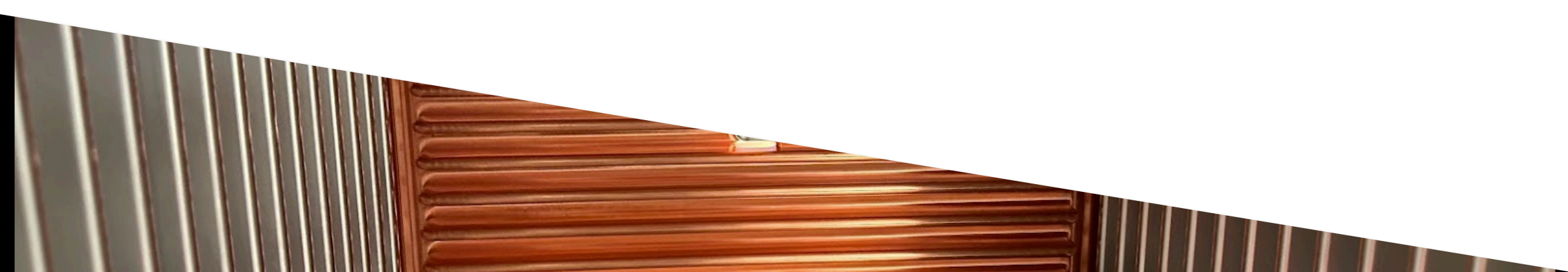
MOST DETECTORS

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$$\left(\omega_1^2 - \omega^2 + i \frac{\omega \omega_1}{Q} \right) \tilde{E}_h(\omega) \simeq -2\omega_1^2 V^{-1/3} \int d\omega' \tilde{u}_m(\omega' - \omega) \tilde{E}_0(\omega')$$

An EM resonator does the readout



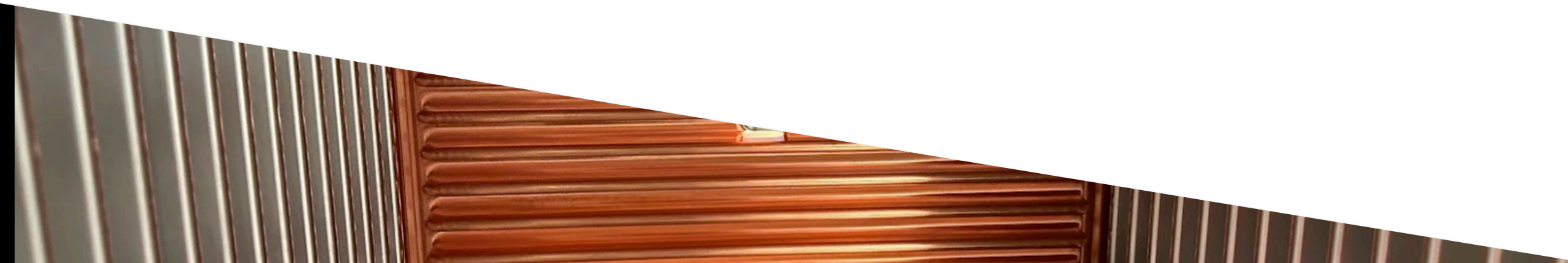
MOST DETECTORS

(LIGO, OPTOMECHANICAL, WEBER BARS...)

$$\mathcal{T}^2(\omega) = \frac{\omega_g^4 \omega_1^4}{\left((\omega_1^2 - \omega^2)^2 + \frac{\omega^2 \omega_1^2}{Q^2} \right) \left((\omega_m^2 - \omega_g^2)^2 + \frac{\omega_g^2 \omega_m^2}{Q_m^2} \right)}$$

MOST DETECTORS
(LIGO, OPTOMECHANICAL, WEBER BARS...)

$$\left(\mathcal{T}^{\text{LVK}}\right)^2(\omega_g) \simeq \frac{\omega_L^2}{\left(4\omega_g^2 + \frac{\omega_L^2}{Q^2}\right)} \simeq \frac{\omega_L^2 L_{\text{eff}}^2}{\left(4\omega_g^2 L_{\text{eff}}^2 + 1\right)}$$

$$\omega_0 \simeq \omega_1 \simeq \omega_L \gg \omega_g \gg \omega_m$$


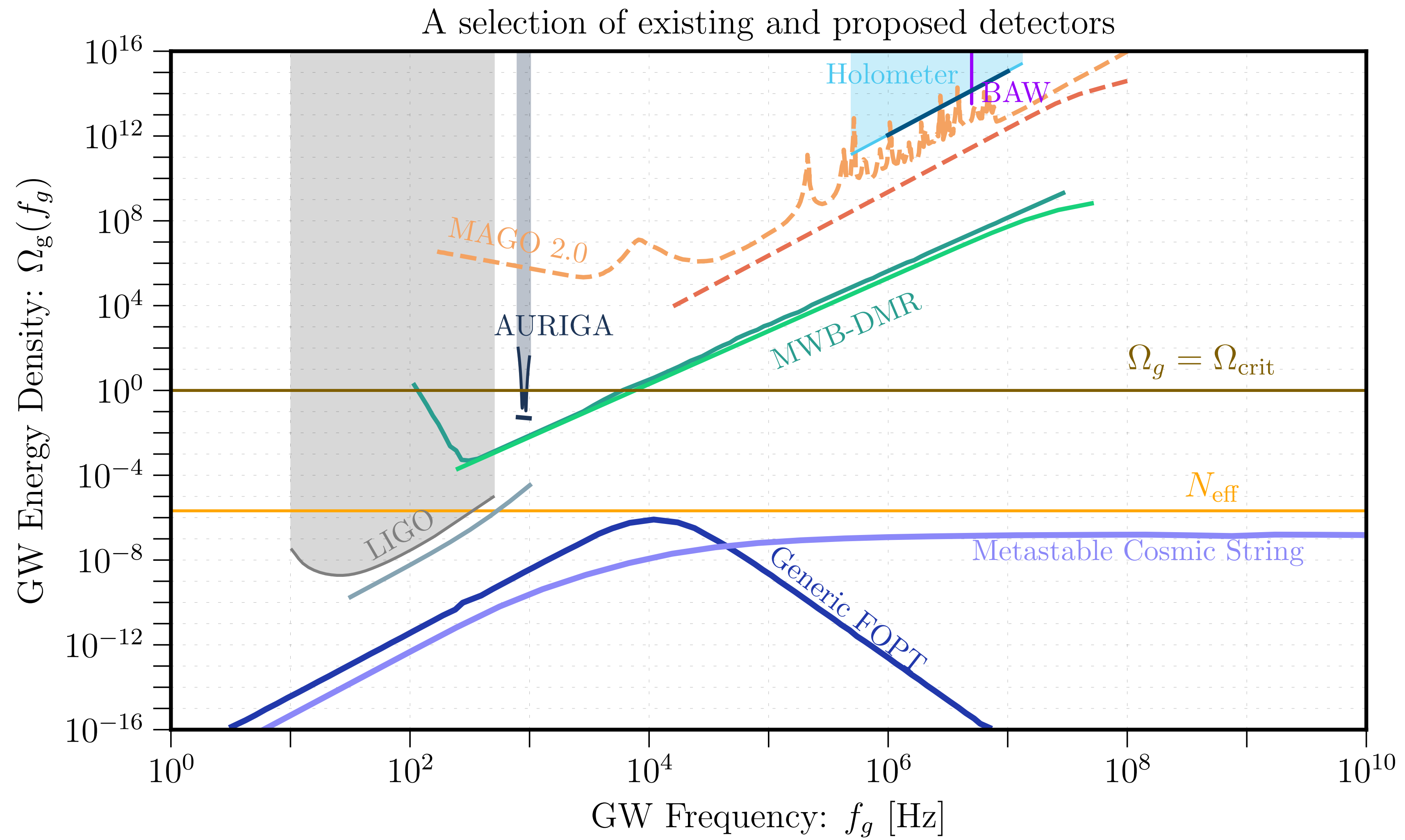
MOST DETECTORS (LIGO, OPTOMECHANICAL, WEBER BARS...)

$$(\mathcal{T}^{\text{Weber}})^2 \lesssim Q^2 Q_m^2$$

$$\omega_m \simeq \omega_{1,0} \simeq \omega_g$$

$$\left(\omega_1^2 - \omega^2 + i\frac{\omega\omega_1}{Q}\right) \tilde{E}_1(\omega) \simeq g_e \int d\omega' \tilde{E}_0(\omega - \omega') \omega \tilde{h}^{\text{TT}}(\omega')$$

$$\mathcal{T}_{\text{EM}}^2(\omega) = \frac{\omega_g^2 \omega^2 (\omega_g L + \omega_0 L + 1)^2}{\left((\omega_1^2 - \omega^2)^2 + \frac{\omega^2 \omega_1^2}{Q^2}\right)} \min[1, \omega_g^2 L^2]$$



$$\mathcal{T}^{\text{LVK}} \lesssim 10^{10}$$

$$\mathcal{T}^{\text{res}} \simeq Q \lesssim 10^{12}$$

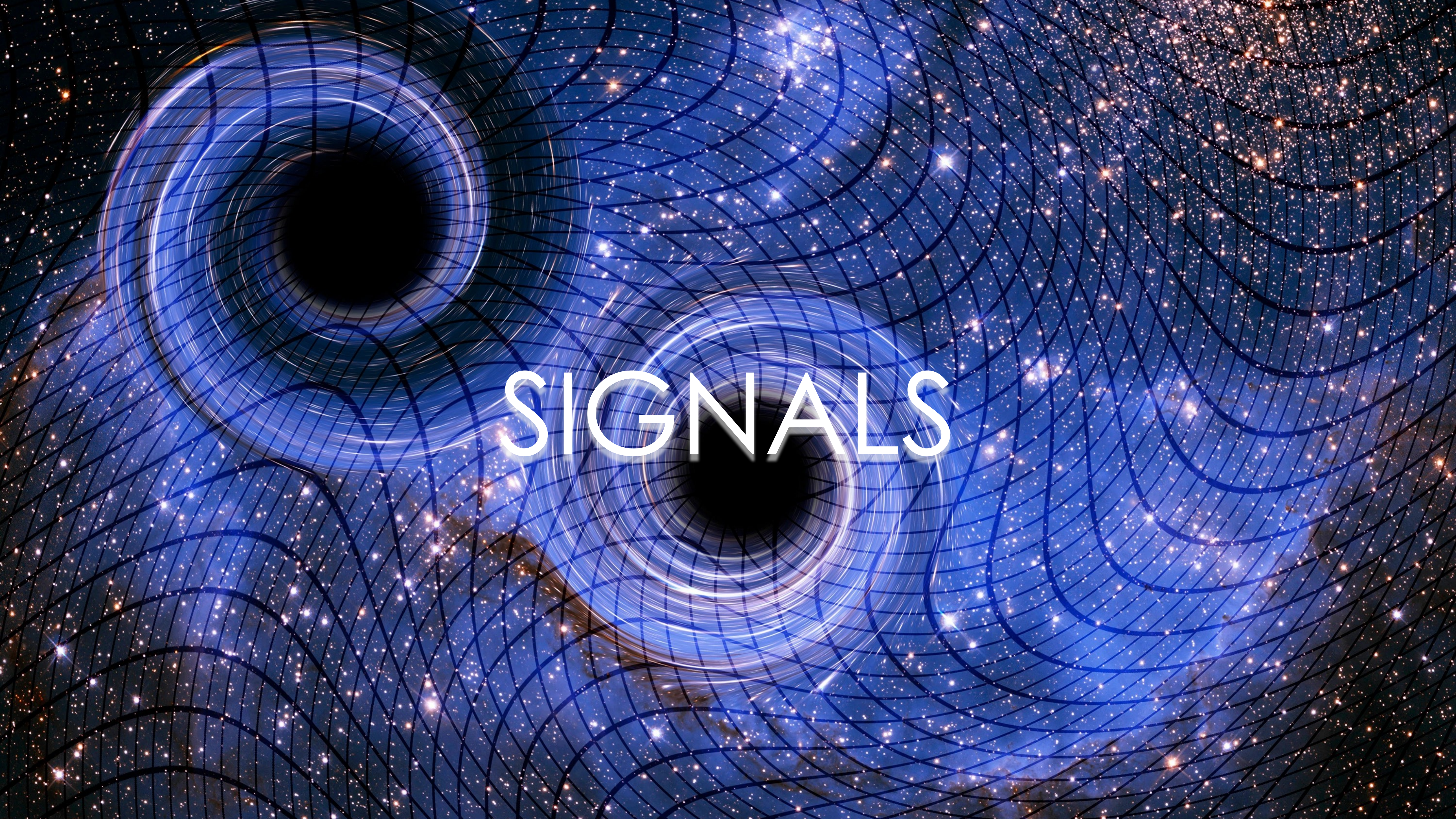
$$\Omega_g \sim \frac{h^2}{\Delta\omega}$$

$$\Omega_g \sim \frac{h^2}{\Delta\omega}$$

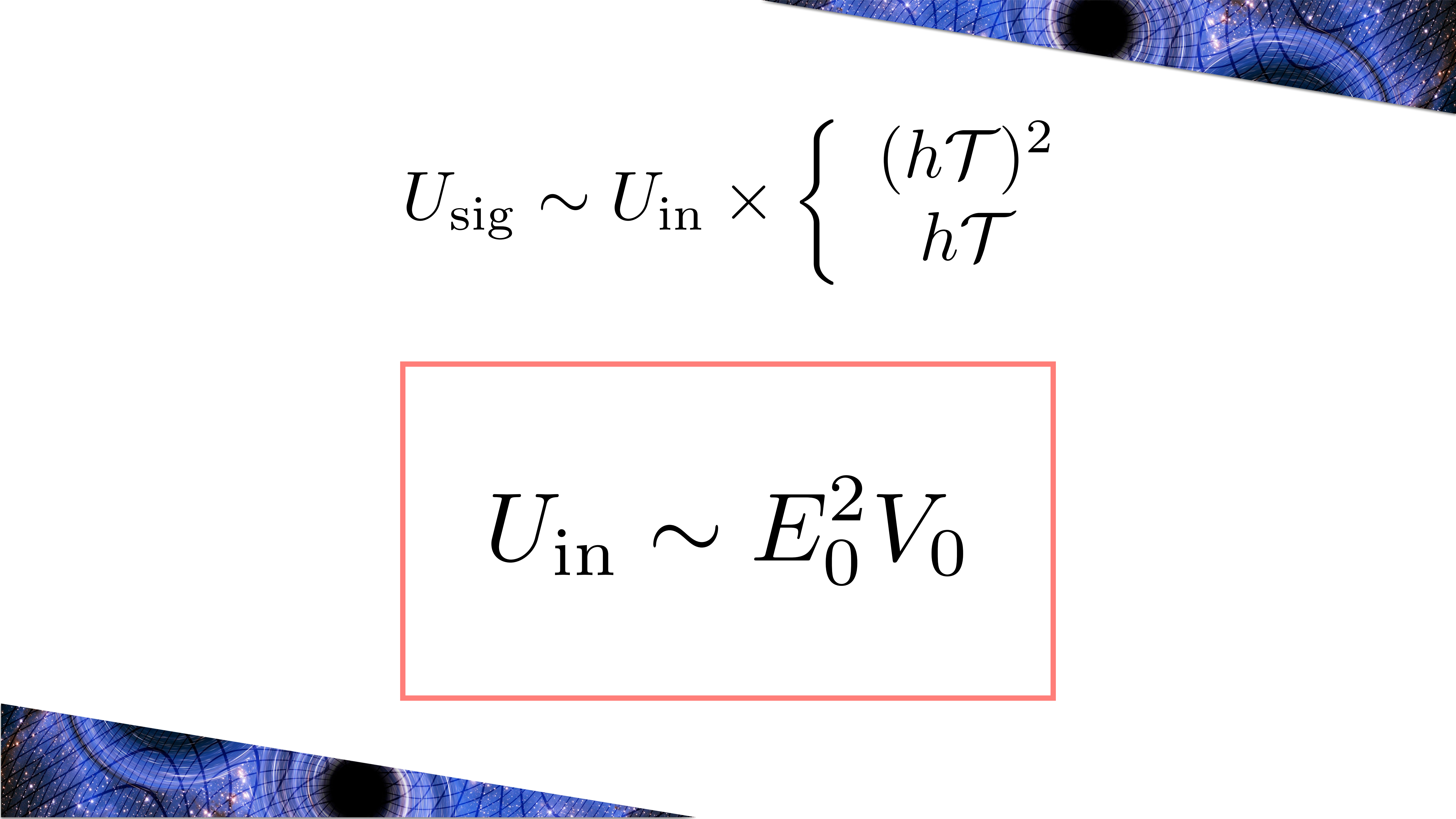
$$\Omega_g \sim \left(\mathcal{T}^{\text{res}} \sqrt{\frac{\omega}{\Delta\omega}} \cong \sqrt{Q} \right)^{-2}$$

$$\Omega_g \sim (\mathcal{T}^{\text{LVK}})^{-2}$$

$$\Omega_g \sim \left(\mathcal{T}^{\text{res}} \sqrt{\frac{\omega}{\Delta\omega}} \simeq \sqrt{Q} \right)^{-2}$$

The background is a deep blue space filled with numerous small, bright stars. Overlaid on this is a complex pattern of glowing blue lines that form a grid of concentric circles and intersecting arcs, resembling a gravitational well or a signal propagation pattern. Two prominent, larger circular structures with a dark center and a bright blue ring are visible, one in the upper left and one in the lower center. The word "SIGNALS" is centered in the middle of the image in a white, sans-serif font.

SIGNALS


$$U_{\text{sig}} \sim U_{\text{in}} \times \begin{cases} (h\mathcal{T})^2 \\ h\mathcal{T} \end{cases}$$

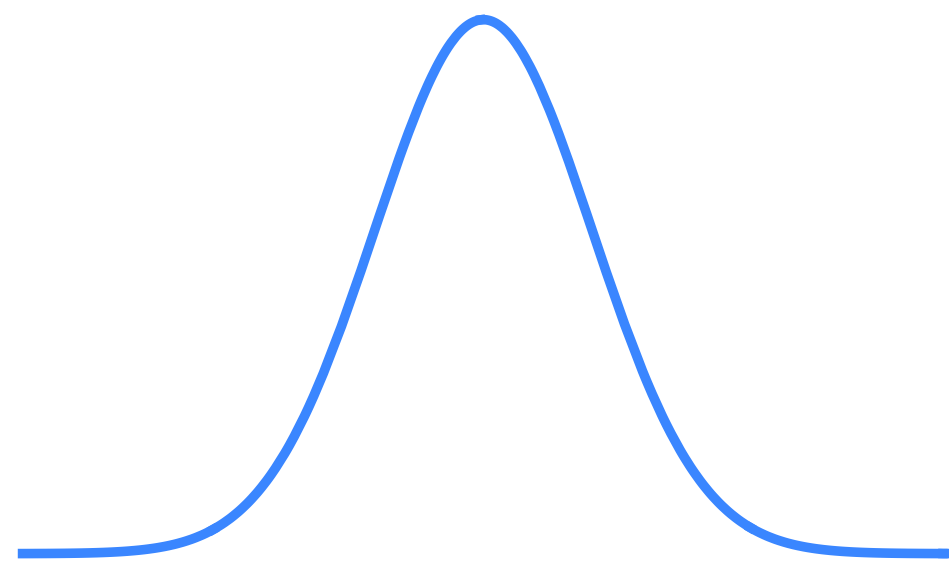
$$U_{\text{in}} \sim E_0^2 V_0$$

CASE I: QUADRATIC SIGNALS

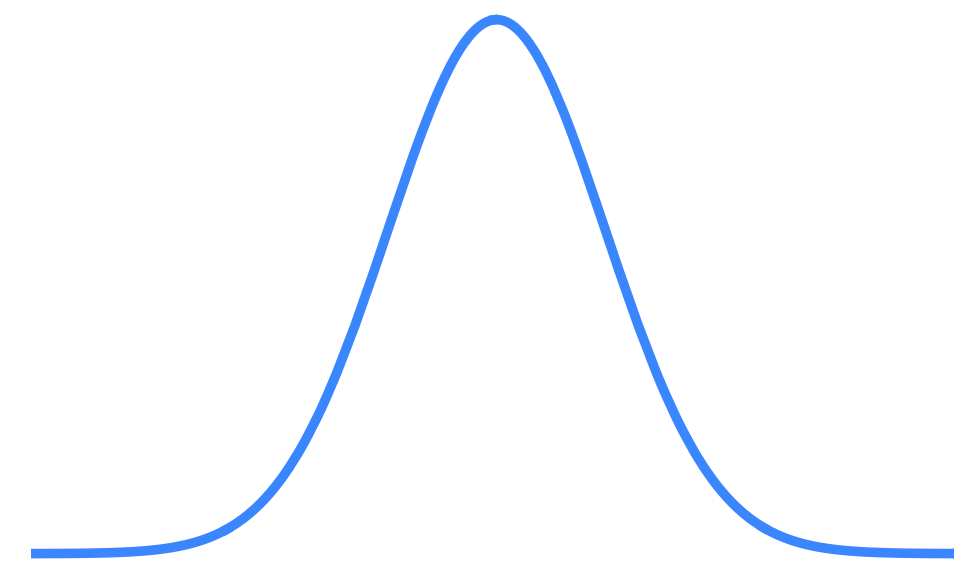
$$\langle E_h(t) E_0(t) \rangle \propto \langle \tilde{E}_h(\omega) \tilde{E}_0(\omega) \rangle = 0$$


$$\langle E_h(t) E_0(t) \rangle \propto \langle \tilde{E}_h(\omega) \tilde{E}_0(\omega) \rangle = 0$$

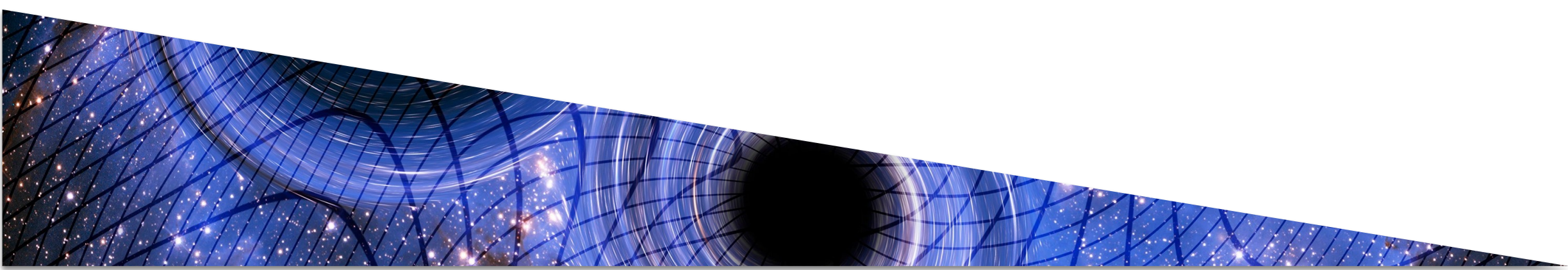
ω

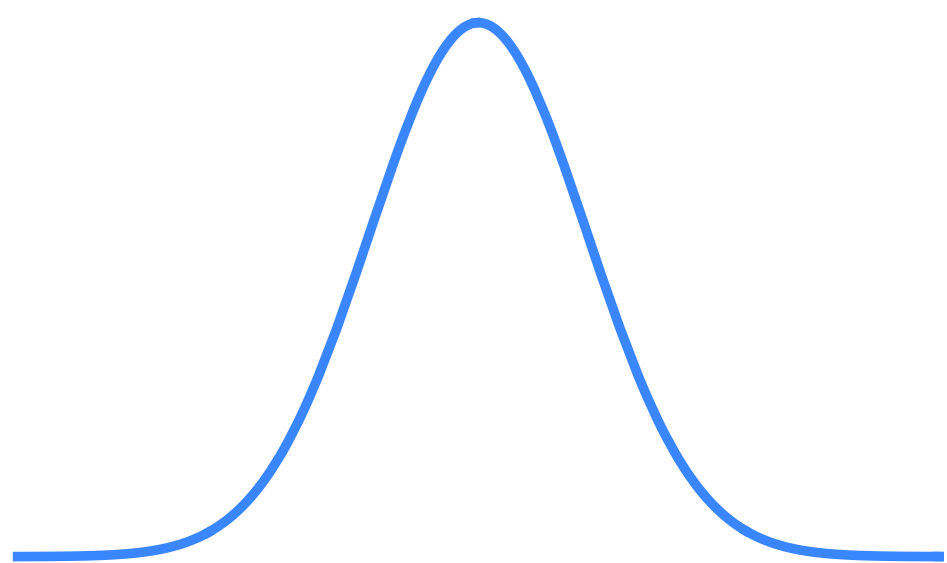


$\tilde{E}_0(\omega)$



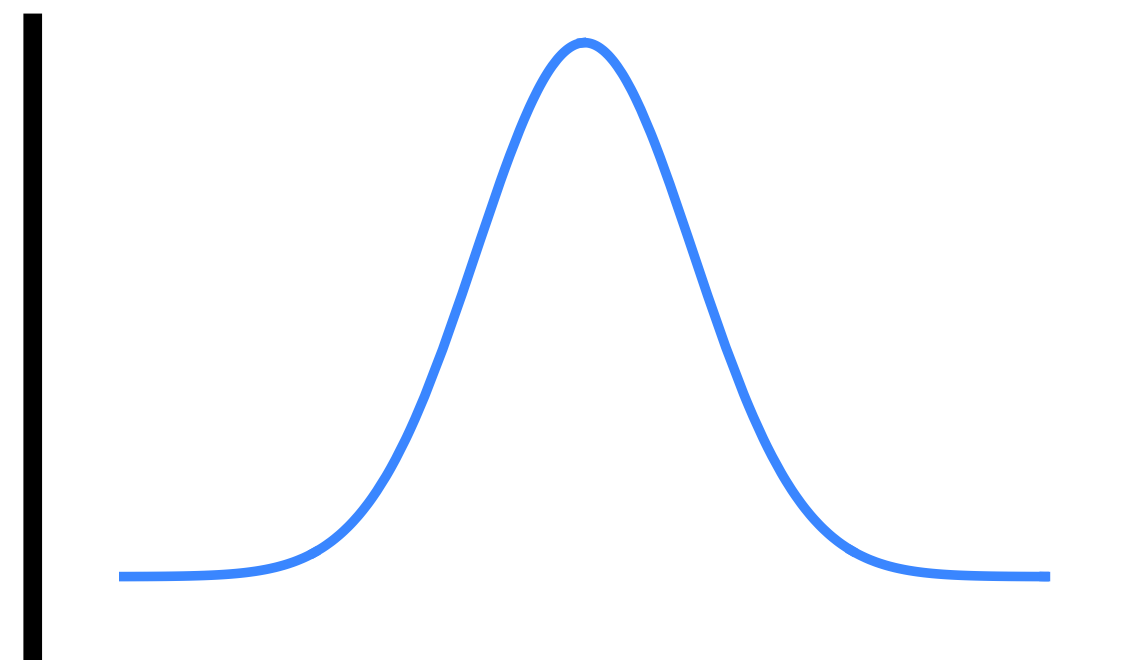
$\tilde{E}_h(\omega)$





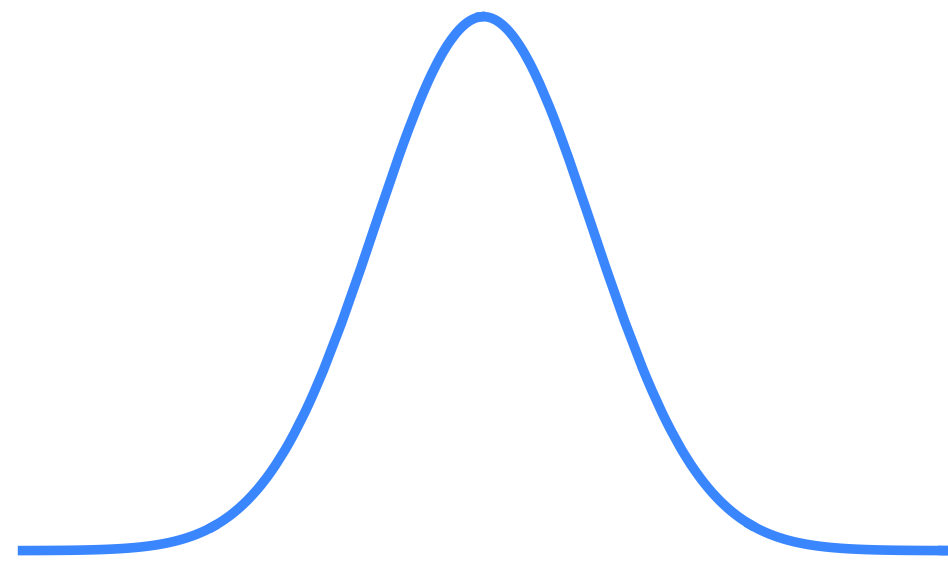
$$\tilde{E}_0(\omega)$$

$\Delta\omega_d$
Detector Bandwidth



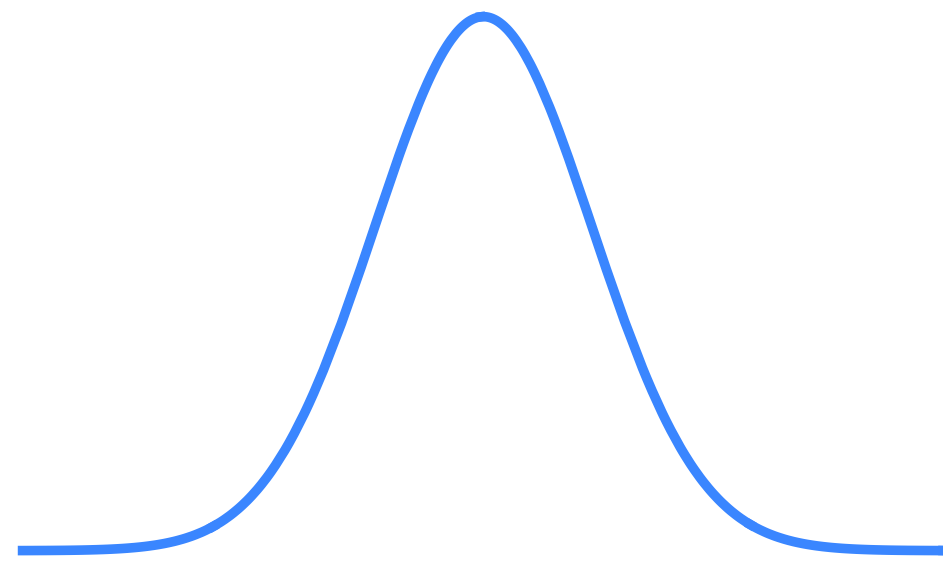
$$\tilde{E}_h(\omega)$$

In absence of a signal,
the detector is empty (classically)



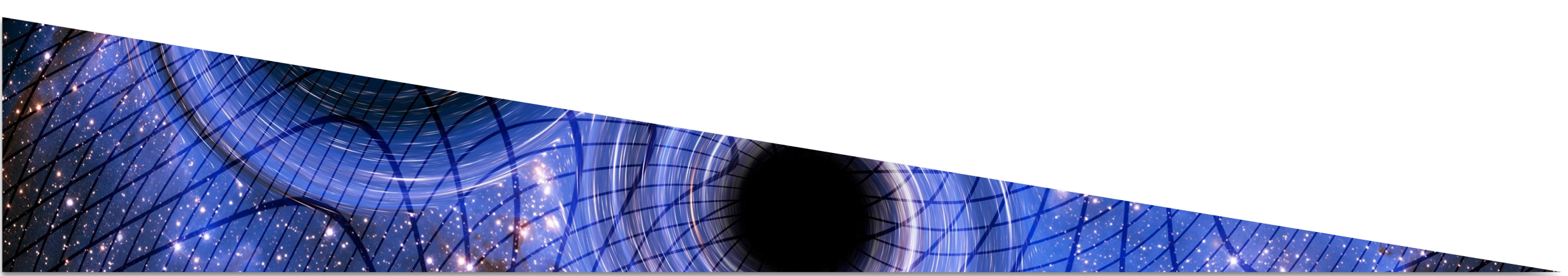


In absence of a signal,
the detector is empty (classically)



$$P_{\min} \simeq \frac{2\pi\omega}{t_{\text{int}}}$$

But we have one quantum of noise
from vacuum fluctuations



$$P_{\text{sig}} \lesssim \underline{h^2 U_{\text{in}} \omega_s \mathcal{T}^2(\omega_s)}$$

Signal Energy

$$P_{\text{sig}} \lesssim h^2 U_{\text{in}} \underline{\omega_s} \mathcal{T}^2(\omega_s)$$

Maximum power
from Poynting's theorem

$$\frac{P_{\text{sig}}}{P_{\text{noise}}} \simeq 1 \quad \rightarrow \quad h_{\text{min}} \gtrsim \sqrt{\frac{2\pi}{U_{\text{in}} t_{\text{int}}}} \frac{1}{\mathcal{T}}$$

$$\Omega_g^{\text{min}}(\omega) \simeq \frac{\omega^3 h_{\text{min}}^2 t_{\text{int}}}{3H_0^2}$$

$$\Omega_g^{\text{min}}(\omega) \simeq \frac{\omega^3 h_{\text{min}}^2 t_{\text{int}}}{3H_0^2}$$

$$U_{\text{in}}^{\text{BBN}} = 10^{12} J \left(\frac{\omega}{2\pi \times 80\text{kHz}} \right) \left(\frac{10^{-6}}{\Omega_g} \right) \left(\frac{10^{12}}{\mathcal{T}^2 \frac{\omega}{\Delta\omega}} \right)$$

$$\Omega_g^{\text{min}}(\omega) \simeq \frac{\omega^3 h_{\text{min}}^2 t_{\text{int}}}{3H_0^2}$$

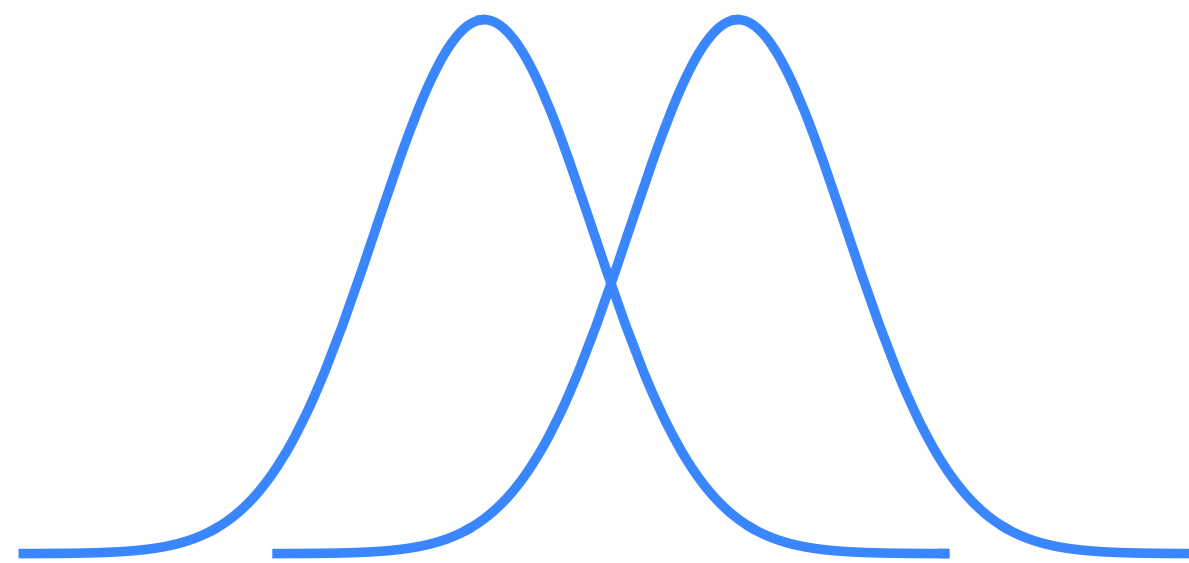
$$U_{\text{in}}^{\text{BBN}} \simeq U_{\text{ITER}}$$

$$U_{\text{in}}^{\text{BBN}} = 10^{12} J \left(\frac{\omega}{2\pi \times 80\text{kHz}} \right) \left(\frac{10^{-6}}{\Omega_g} \right) \left(\frac{10^{12}}{\mathcal{T}^2 \frac{\omega}{\Delta\omega}} \right)$$

CASE II: LINEAR SIGNALS

$$\langle E_h(t) E_0(t) \rangle \propto \langle \tilde{E}_h(\omega) \tilde{E}_0(\omega) \rangle \neq 0$$

ω

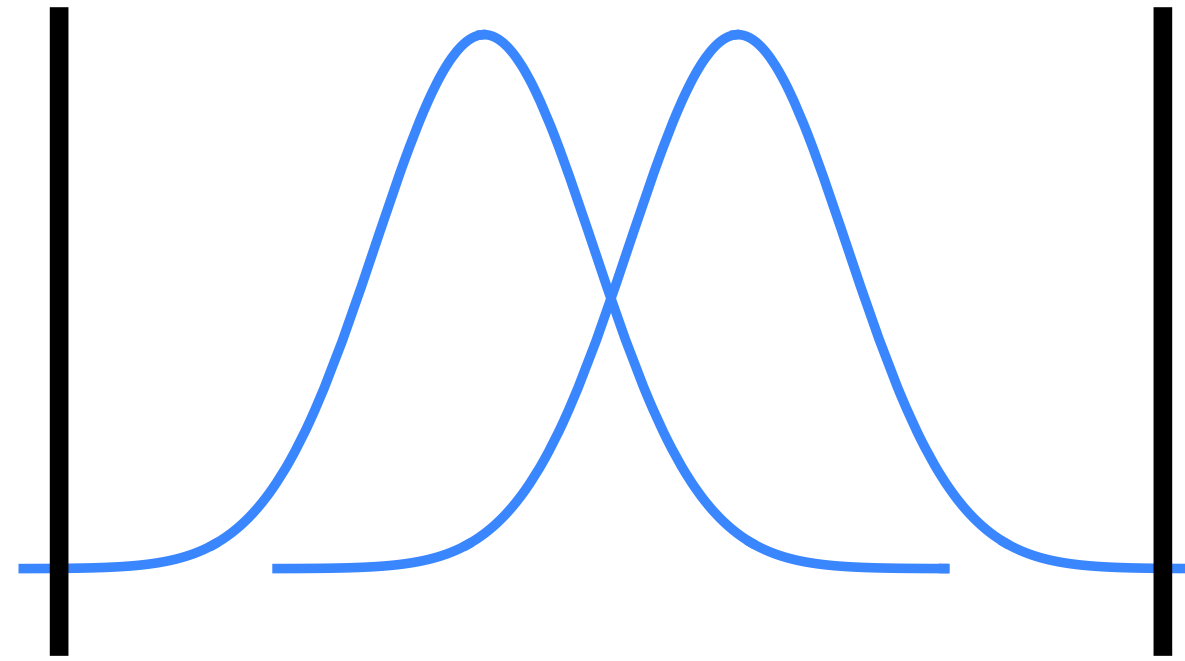


$\tilde{E}_0(\omega)$

$\tilde{E}_h(\omega)$

CASE II: LINEAR SIGNALS

$\Delta\omega_d$
Detector Bandwidth

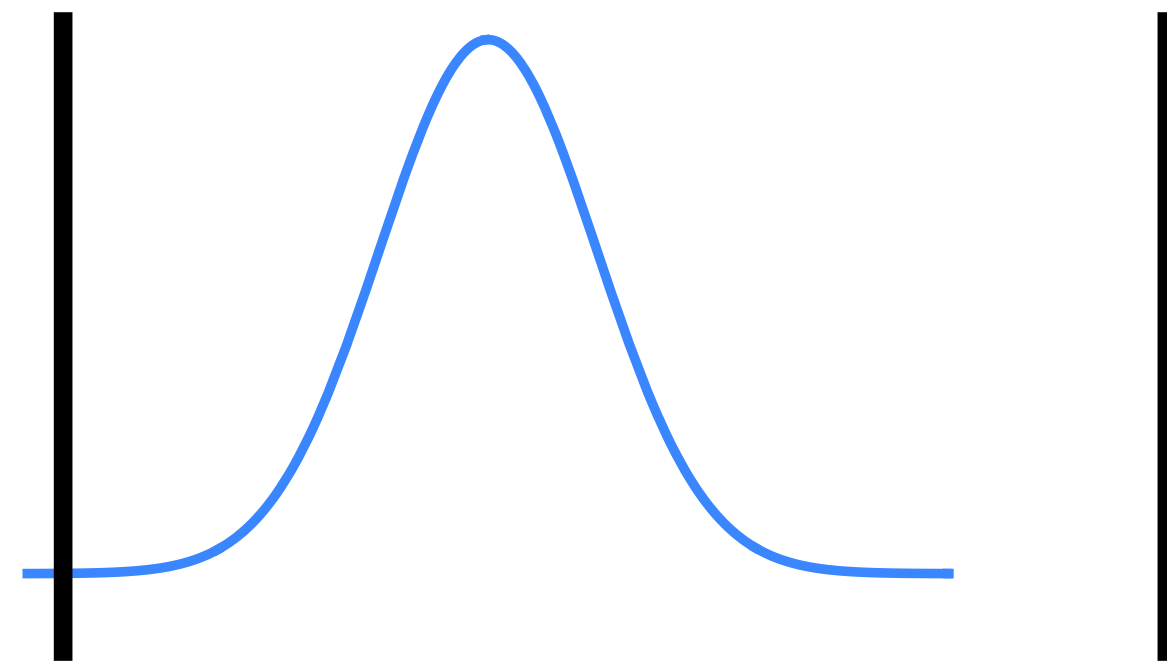


$$\tilde{E}_0(\omega)$$

$$\tilde{E}_h(\omega)$$

CASE II: LINEAR SIGNALS

In absence of a signal,
the detector is not empty



$$P_{\text{noise}}^{\text{min}} \simeq \frac{2\pi\omega}{t_{\text{int}}} \left(1 + \sqrt{\frac{P_{\text{in}} t_{\text{int}}}{2\pi\omega}} \right)$$

$$\frac{P_{\text{sig}}}{P_{\text{noise}}} \simeq 1 \quad \rightarrow \quad h_{\text{min}} \gtrsim \sqrt{\frac{2\pi}{U_{\text{in}} t_{\text{int}}}} \frac{1}{\mathcal{T}}$$

Same as quadratic!

STATISTICS INTERLUDE

$$\frac{v \frac{dr}{dr} = -\Omega_k^2 r + \Omega_c^2}{2}$$



$$\frac{h''}{\rho} \left(2rp + p \frac{\partial r}{\partial H} \right)$$

$$\frac{\partial_r(\rho c_s^2)}{\rho} = c_s^2 \frac{\partial_r \rho}{\rho} = c_s^2 \frac{\partial_r \rho}{\rho}$$

$$\frac{V^2 - c_s^2}{V}$$

$$V = W V_0 \rightarrow \dots$$

$$\sqrt{V_0} (V_0 \partial_r W + W \partial_r V_0)$$

Stationary Gaussian Process with Zero Mean

$$\langle h(t) \rangle = 0$$

$$\langle h(t)h(t') \rangle = H(t' - t)$$

$$P_{\text{sig,noise}} = \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} S_{\text{sig,noise}}(\omega)$$

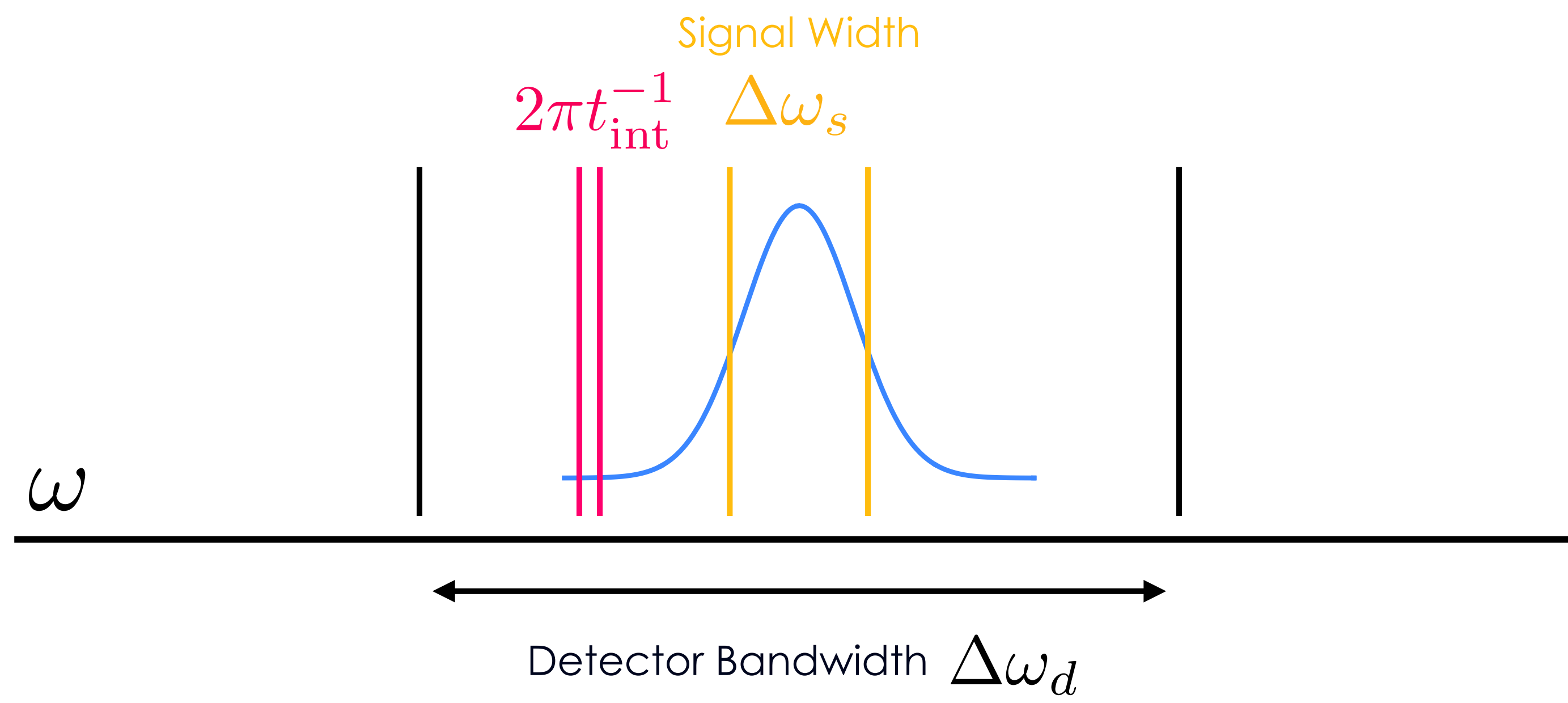
In practice we are comparing two Gaussians with zero mean and different variance

$$\text{SNR} = \left(t_{\text{int}} \int \frac{d\omega}{2\pi} \frac{S_{\text{sig}}^2(\omega)}{S_{\text{noise}}^2(\omega)} \right)^{1/2} \simeq 1$$

From maximum likelihood

$$h_{\min} \gtrsim \sqrt{\frac{1}{U_{\text{in}}}} \left(\frac{2\pi \Delta\omega}{t_{\text{int}}} \right)^{1/4} \frac{1}{\mathcal{T}(\omega)}$$

$$\Delta\omega \equiv \max[\min[\Delta\omega_d, \Delta\omega_s], (2\pi) t_{\text{int}}^{-1}]$$

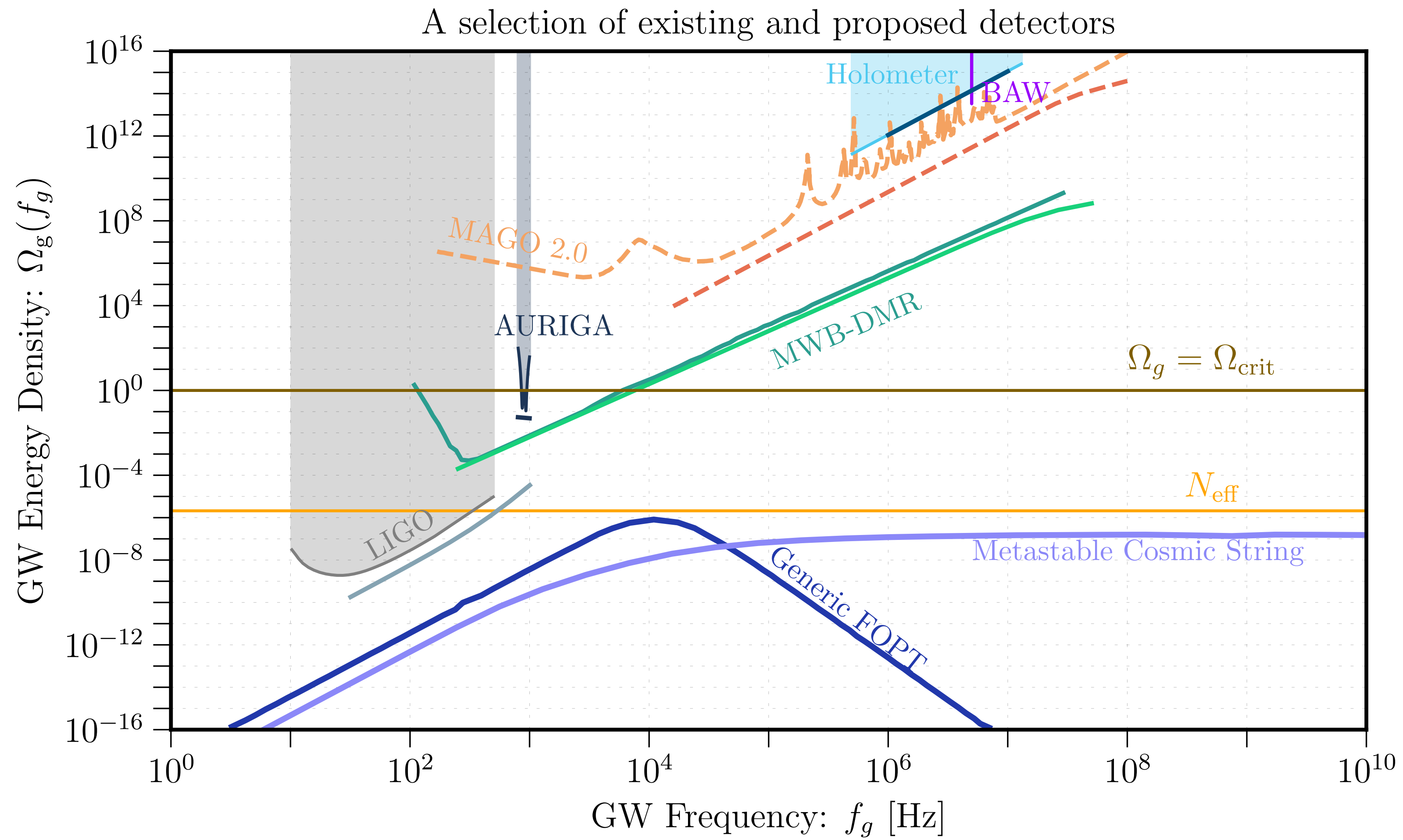


$$\Delta\omega \equiv \max[\min[\Delta\omega_d, \Delta\omega_s], (2\pi) t_{\text{int}}^{-1}]$$

Exception

$$\text{SNR} = \left(t_{\text{int}} \int \frac{d\omega}{2\pi} \frac{S_{\text{sig}}(\omega)}{S_{\text{noise}}(\omega)} \right)^{1/2} \simeq 1$$

Photon counting [2211.04016, 2404.07524]



WHAT DID WE LEARN?

$$\frac{v dr}{dr} = \frac{-\Omega_k^2 r + \Omega_k^2}{2}$$

$$\frac{h''}{\rho} \left(2rp + p \frac{\partial r}{\partial H} \right)$$

$$\frac{\partial_r(\rho c_s^2)}{\rho} = c_s^2 \frac{\partial_r \rho}{\rho} = c_s^2 \frac{\partial_r \rho}{\rho}$$

$$\frac{v^2 - c_s^2}{v}$$

$$v = w v_0 \rightarrow \dots$$

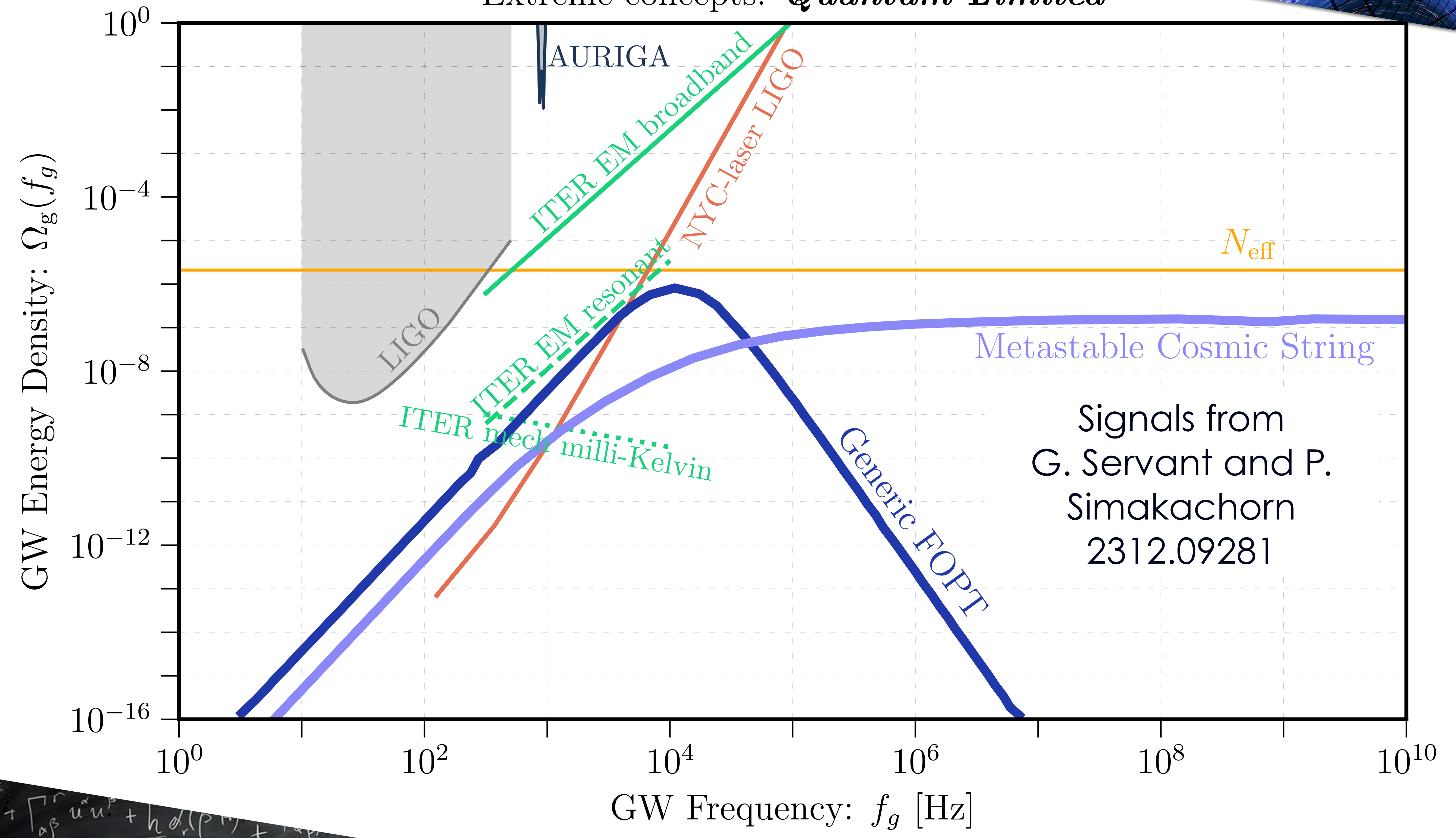
$$\sqrt{v_0} (v_0 \partial_r w + w \partial_r v_0)$$

$$h_{\min} \gtrsim \sqrt{\frac{2\pi}{U_{\text{in}} t_{\text{int}}}} \frac{1}{\mathcal{T}}$$

$$\Omega_g^{\text{min}}(\omega) \simeq \frac{\omega^3 h_{\text{min}}^2 t_{\text{int}}}{3H_0^2}$$

$$U_{\text{in}}^{\text{BBN}} = 10^{12} J \left(\frac{\omega}{2\pi \times 80\text{kHz}} \right) \left(\frac{10^{-6}}{\Omega_g} \right) \left(\frac{10^{12}}{\mathcal{T}^2 \frac{\omega}{\Delta\omega}} \right)$$

Extreme concepts: *Quantum-Limited*



Signals from
G. Servant and P.
Simakachorn
2312.09281

QUADRATIC vs LINEAR

QUADRATIC

$$P_{\text{noise}}^{\text{min}} \simeq \frac{2\pi\omega}{t_{\text{int}}}$$

LINEAR

$$P_{\text{noise}}^{\text{min}} \simeq \frac{2\pi\omega}{t_{\text{int}}} \left(1 + \sqrt{\frac{P_{\text{in}} t_{\text{int}}}{2\pi\omega}} \right)$$

We can gain a lot
from
quantum techniques

BEYOND THE QUANTUM LIMIT

$$v \frac{dr}{dr} = -\Omega_k^2 r + \dots$$



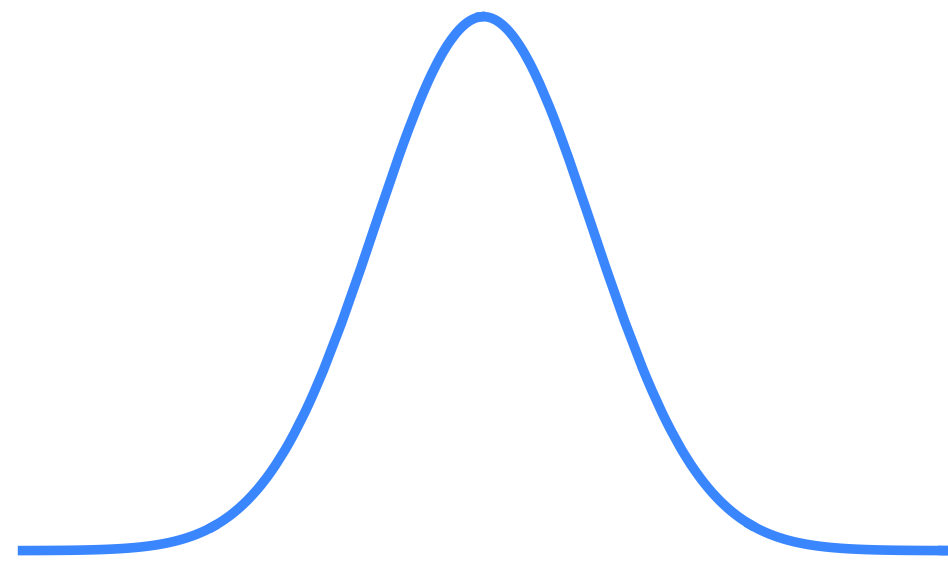
$$p \pm p \frac{\partial}{\partial r} (p H) +$$

$$\frac{h}{p} \left(2rp + p \frac{\partial}{\partial r} \left(\frac{h}{H} \right) \right) \frac{V^2 - c_s^2}{V}$$
$$\frac{\partial_r(p c_s^2)}{p} = c_s^2 \frac{\partial_r p}{p}$$
$$= c_s^2 \frac{\partial_r \rho}{\rho}$$

$$v = w V_0 \rightarrow \dots$$

$$\sqrt{V_0} (v_0 \partial_r w + w \partial_r V_0)$$

NO QUALITATIVE CHANGE FOR QUADRATIC SIGNALS



$$P_{\text{SQL}} \simeq \frac{2\pi\omega}{t_{\text{int}}}$$

$$P_{\text{min}} \simeq \frac{2\pi\omega}{t_{\text{int}}}$$

LINEAR

$$P_{\text{noise}}^{\text{min}} \simeq \frac{2\pi\omega}{t_{\text{int}}} \left(1 + \sqrt{\frac{P_{\text{in}} t_{\text{int}}}{2\pi\omega}} \right)$$
$$\simeq \frac{2\pi\omega}{t_{\text{int}}} \left(1 + \sqrt{N_{\gamma}} \right)$$

In principle

$$P_{\text{noise}}^{\text{min}} \rightarrow \frac{2\pi\omega}{t_{\text{int}}}$$

Heisenberg Limit

In practice

$$P_{\text{noise}}^{\text{min}} \rightarrow \frac{2\pi\omega}{t_{\text{int}}} \quad \longrightarrow \quad (\sqrt{N_{\gamma}} + 1) \rightarrow 0$$

You need to control an Avogadro's number of photons
at the single photon level

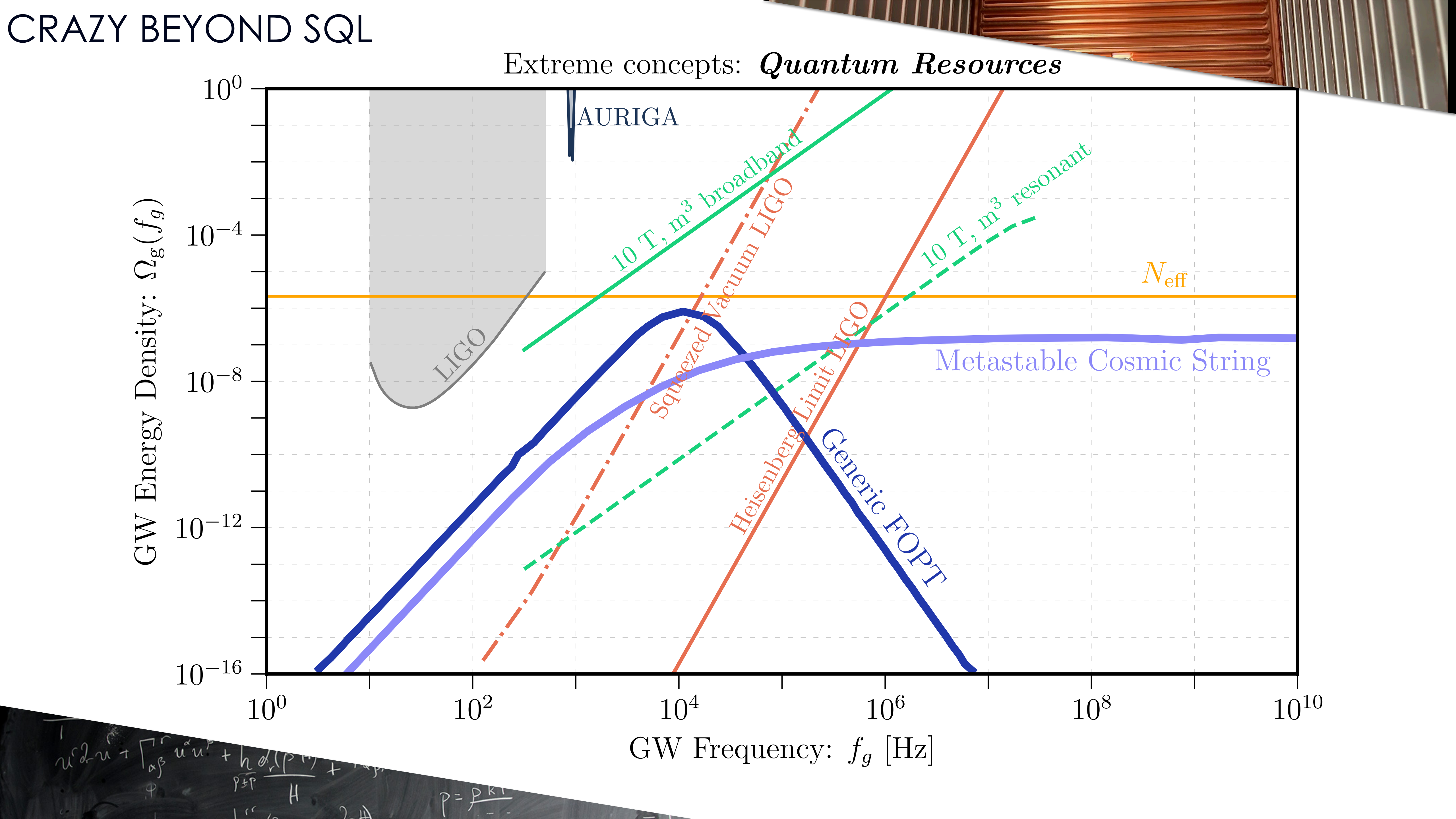
Heisenberg limited sensors
+
Classical computer

$$t_{\text{int}} = \mathcal{O} \left(\frac{1}{U_{\text{in}} h_{\text{min}}} \left\lceil \frac{|\Delta\omega|}{U_{\text{in}} h_{\text{min}}} \right\rceil \right)$$

Heisenberg limited sensors
+
Grover algorithm

$$t_{\text{int}} = \mathcal{O} \left(\frac{1}{U_{\text{in}} h_{\text{min}}} \sqrt{\left\lceil \frac{|\Delta\omega|}{U_{\text{in}} h_{\text{min}}} \right\rceil} \right)$$

[2501.07625]



$$\Omega_g(\omega_g) \sim \omega_g^3 h^2$$