

Surface Defects & Integrability

Adam Chalabi, University of Turin & INFN

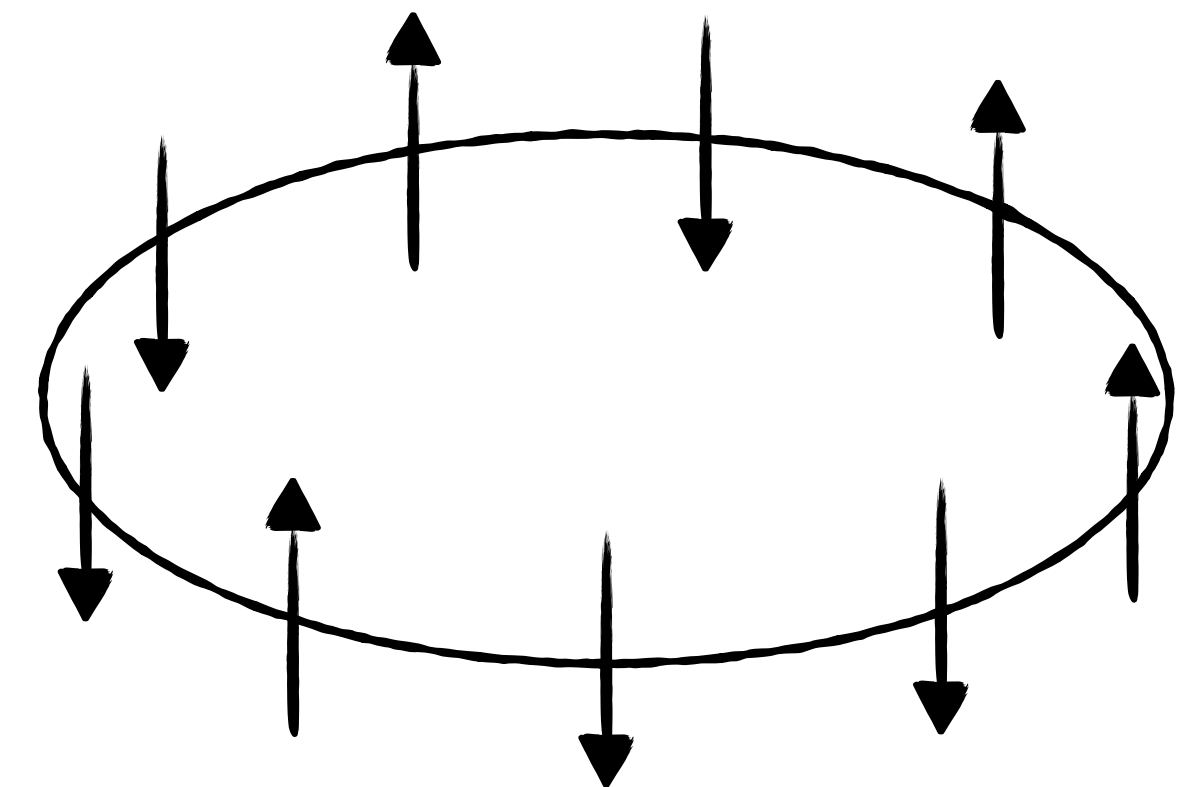
Based on 2503.22598 with C. Kristjansen and Chenliang Su

Interfaces & Symmetry, YITP, 5 March 2026



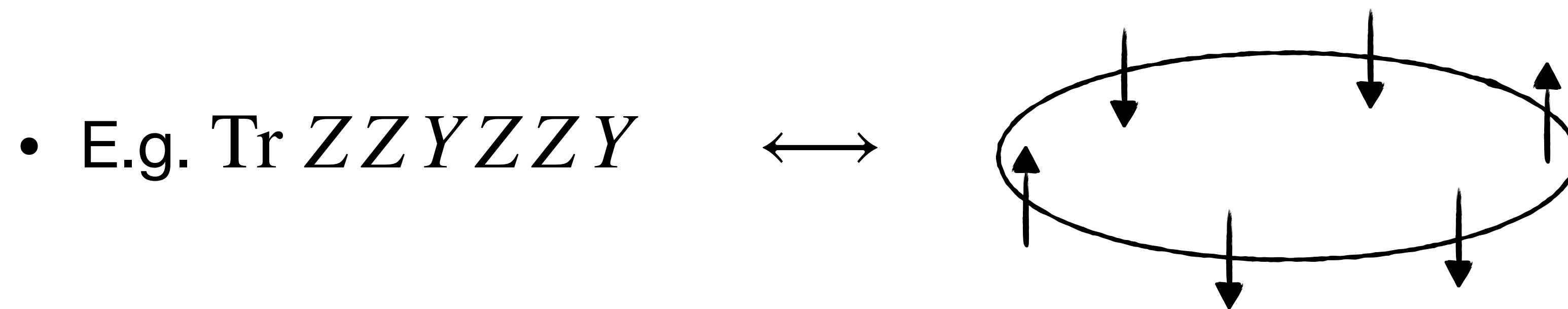
Integrability in $N=4$ SYM Theory at Large- N

- Local operators $\text{Tr } \phi^{I_1} \dots \phi^{I_L}$ of the same engineering dimension L can mix
- Diagonalise dilatation operator D to find scaling ops $\langle \mathcal{O}(x) \mathcal{O}(0) \rangle \sim |x|^{-2\Delta}$
- Minahan-Zarembo '02: $D_{1-loop} \leftrightarrow H$ Hamiltonian on a 1d spin chain
- H is integrable $\implies \exists$ tower of conserved charges Q
- H diagonalisable via Bethe ansatz



Closed Sectors and Integrability

- SU(2) sector: simplest closed sub-sector at 1-loop consists of $\{Z, Y\}$ only
- Heisenberg spin chain with identification $Z = |\downarrow\rangle$ and $Y = |\uparrow\rangle$

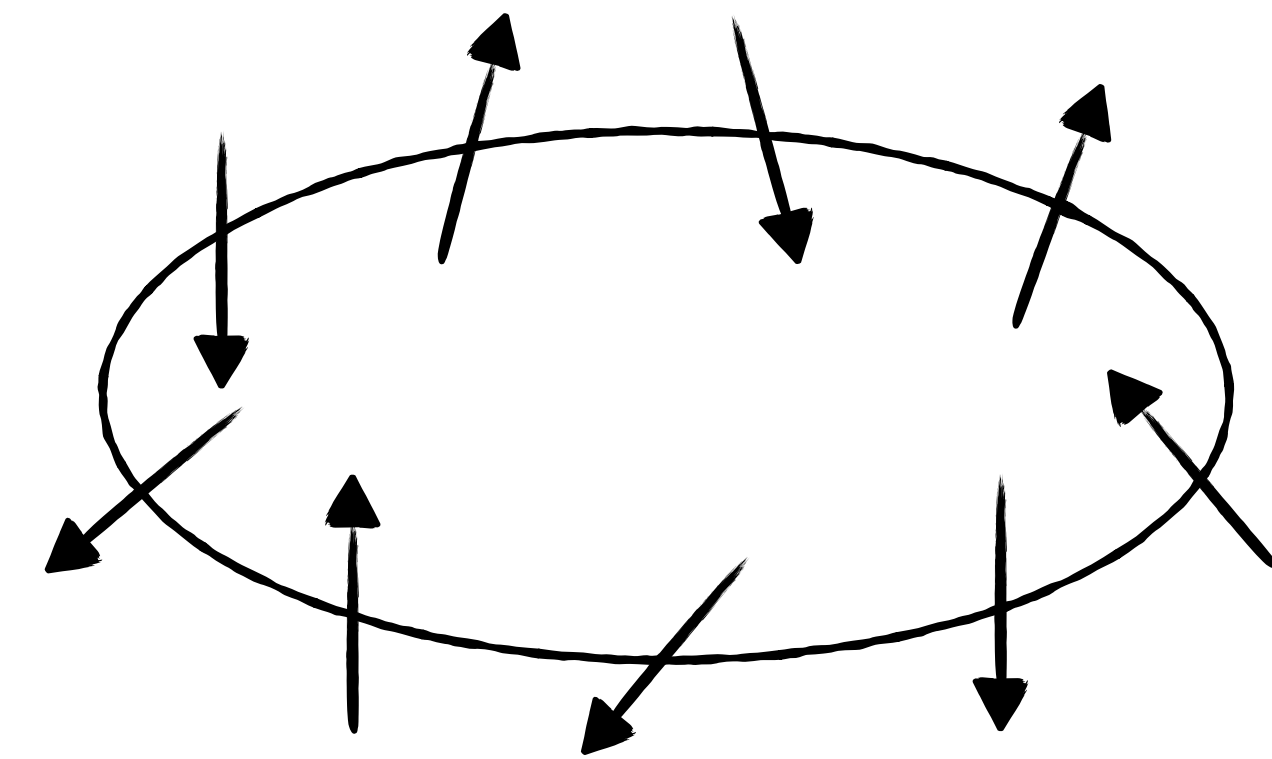


- Hamiltonian $H \propto \sum_{i=1}^L (1_{i,i+1} - P_{i,i+1})$

- Eigenstates $|\mathbf{u}\rangle$ found by solving Bethe equations $\left(\frac{u_j + i/2}{u_j - i/2}\right)^L \prod_{k \neq j}^{\mathcal{N}} \frac{u_j - u_k - i}{u_j - u_k + i} = 1$

Closed Sectors and Integrability

- SO(6) sector: $\{Z, Y, X, \bar{X}, \bar{Y}, \bar{Z}\}$
 - each site in vector representation of so(6)
- SL(2) sector: $\{Z, DZ, D^2Z, \dots\}$ where $D = D_t + D_x$ and $D_\mu = \partial_\mu + i[A_\mu, \bullet]$
 - infinite-dimensional Hilbert space at each site
- Integrability extends to full planar $N=4$ SYM beyond 1-loop



Defects in CFT

- In CFTs with defects, local operators acquire 1-pt functions $\langle \mathcal{O}_\Delta \rangle = \frac{a_{\mathcal{O}}}{r^\Delta}$
- Some defects described by singularity conditions
e.g. $\phi^I = \frac{\omega^I}{r} + \text{fluctuations}$
- At leading order $\langle \text{Tr } \phi^{I_1} \dots \phi^{I_L} \rangle$ obtained by substituting classical part
- But operators mix! \longrightarrow Use integrability to find scaling ops

Integrable Defects

- In spin chain picture, encode defect as boundary state

$$\langle \mathcal{B} | = \text{tr} \left(\sum_{I=1}^d \omega^I \langle I | \right)^{\otimes L}$$

- LO 1-pt function of D -eigenstates

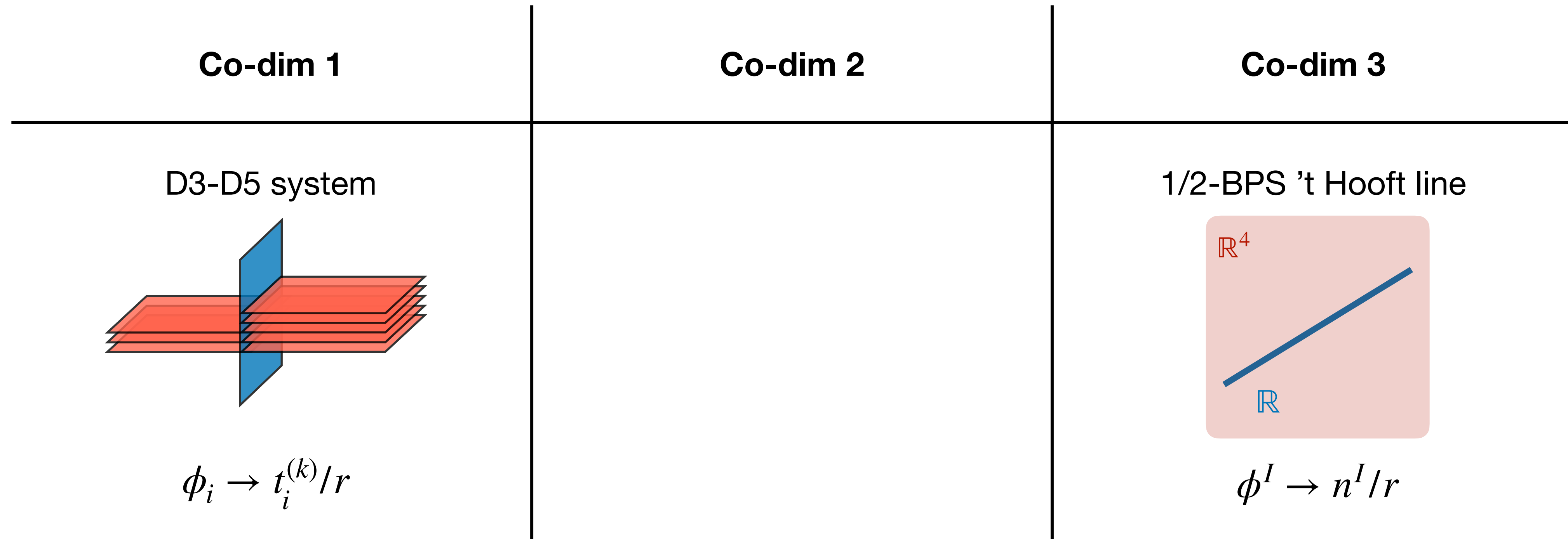
$$\langle \mathcal{O} \rangle \propto \langle \mathcal{B} | \mathbf{u} \rangle$$

Specifies defect

Specifies \mathcal{O} via Bethe roots \mathbf{u}

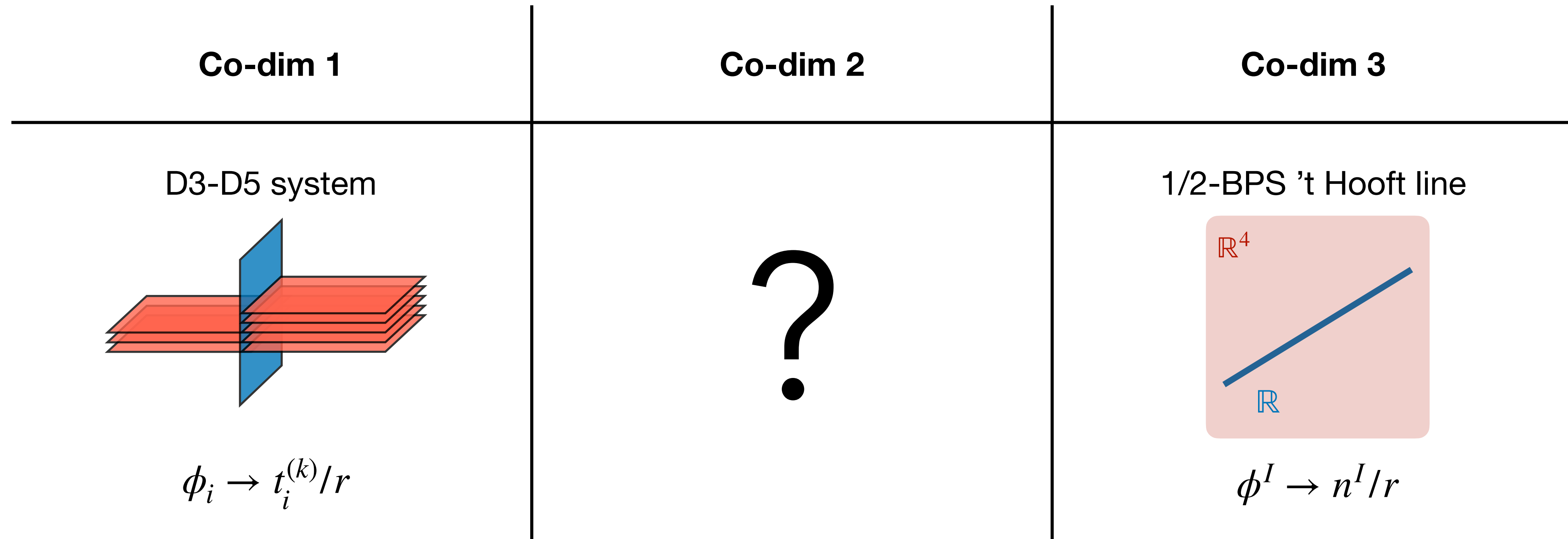
- Defect is integrable $\iff Q^{\text{odd}} | \mathcal{B} \rangle = 0$ and $\langle \mathcal{B} | \mathbf{u} \rangle$ has closed-form
- Sometimes integrability persists to all orders in λ

Integrable 1/2-BPS Defects in $N=4$ SYM



[Bajnok, Buhl-Mortensen, de Leeuw, Gombor, Ipsen, Komatsu, Kristjansen, Linardopoulos, Wang, Wilhelm, Zarembo, ... '15 — today]

Integrable 1/2-BPS Defects in $N=4$ SYM



[Bajnok, Buhl-Mortensen, de Leeuw, Gombor, Ipsen, Komatsu, Kristjansen, Linardopoulos, Wang, Wilhelm, Zarembo, ... '15 — today]

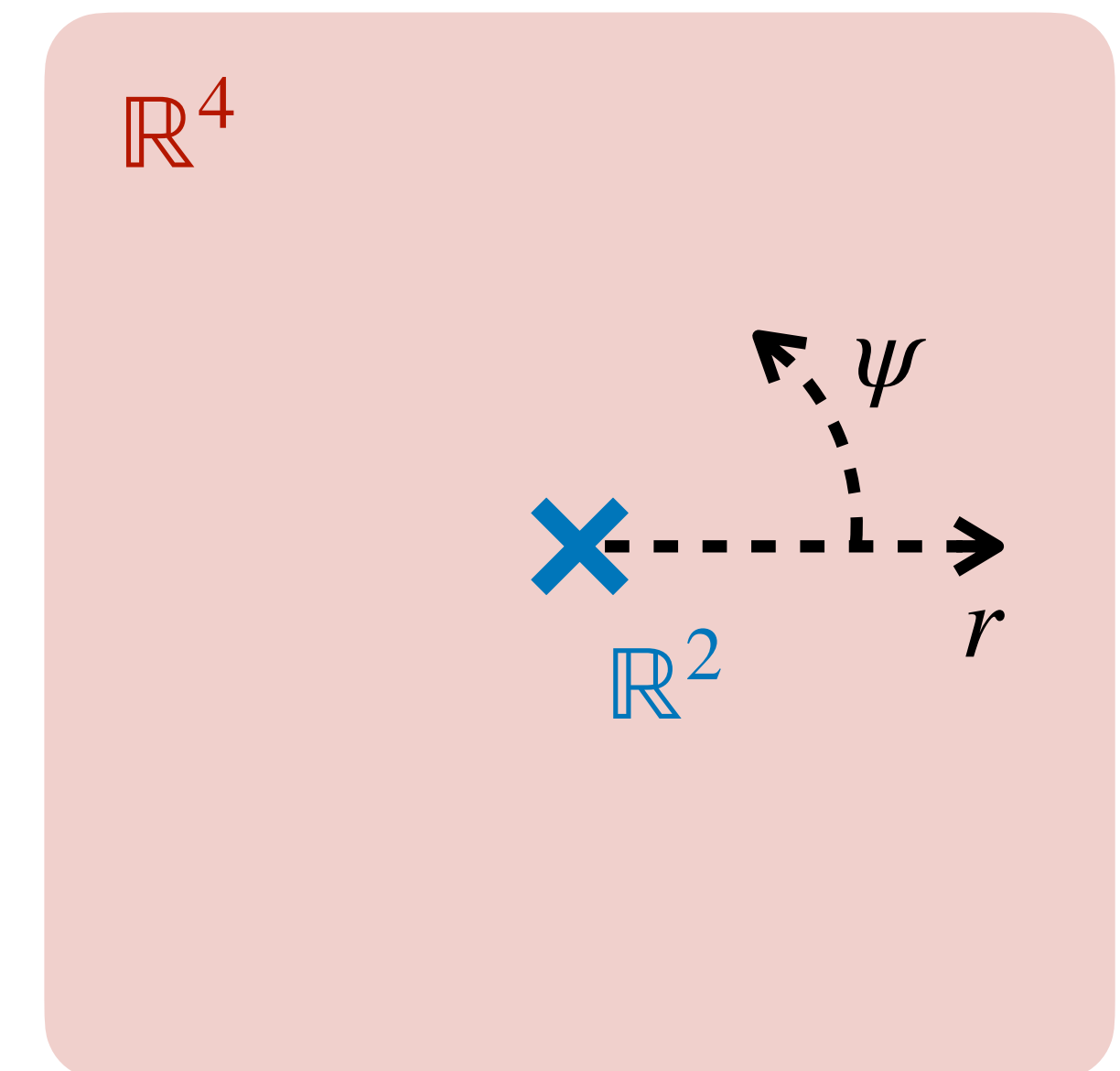
1/2-BPS Surface Defects

- 2d $N=(4,4)$ SUSY ensured by [Gukov-Witten '06]

$$A = X_3(r) d\psi \text{ and } \Phi = \frac{e^{-i\psi}}{r}(X_1(r) + iX_2(r))$$

satisfying Nahm equations

$$[X_i, X_j] + \epsilon_{ijk} \frac{\partial X_k}{\partial \log r} = 0$$



Gukov-Witten Defects

- 2d $N=(4,4)$ surface defects come in two kinds

(1) Ordinary

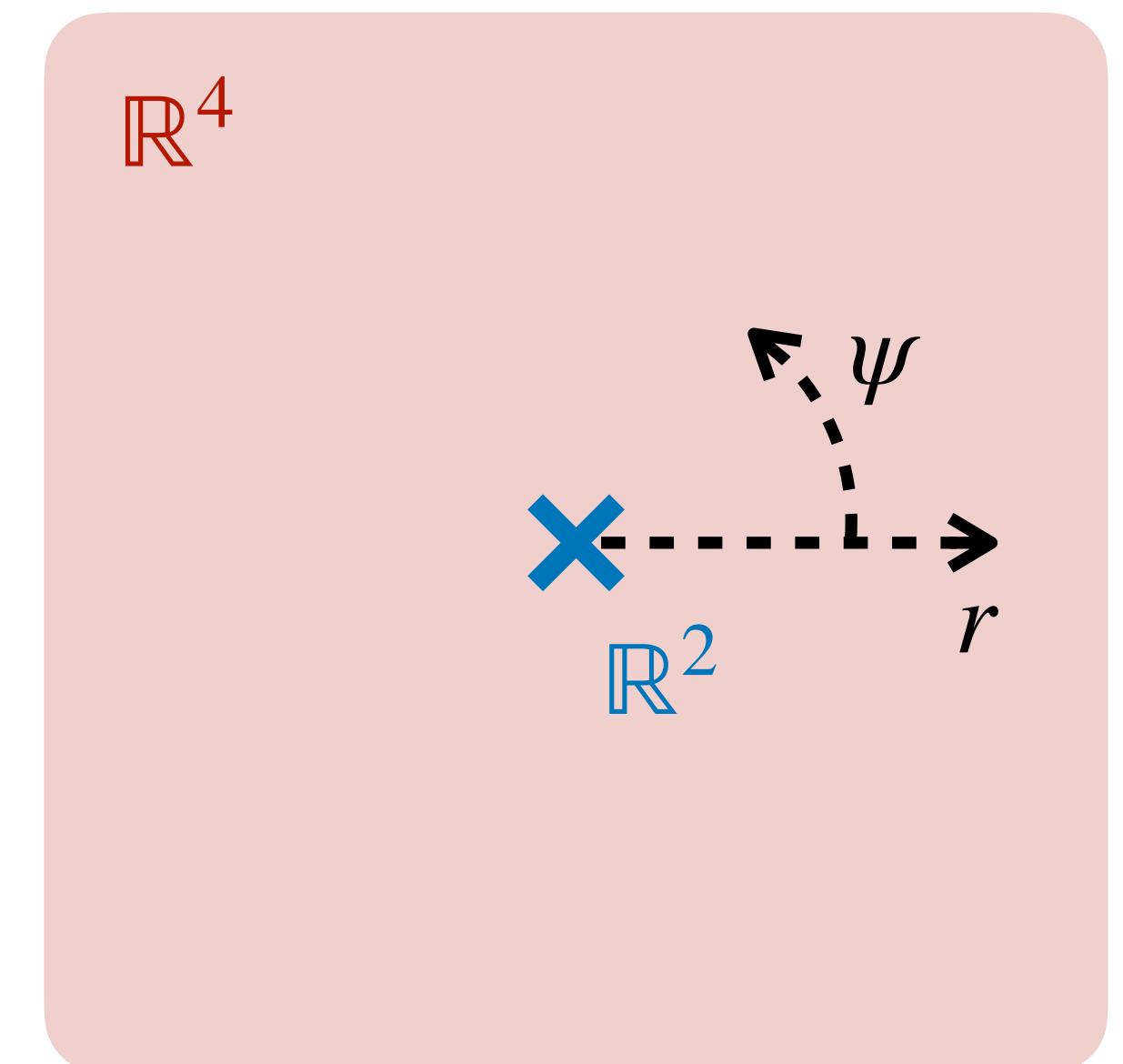
$$A = \alpha d\psi \text{ and } \Phi = \frac{e^{-i\psi}}{r}(\beta + i\gamma)$$

(2) Rigid ($\alpha, \beta, \gamma \rightarrow 0$)

$$A = \frac{t_3}{\log r} d\psi \text{ and } \Phi = \frac{e^{-i\psi}}{r \log r}(t_1 + it_2)$$

diagonal
matrices

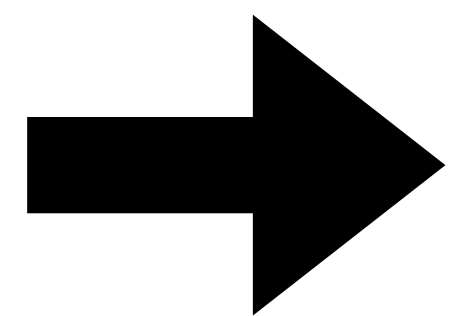
$\mathfrak{su}(2)$ representation matrices



Gukov-Witten Defects

- 2d $N=(4,4)$ surface defects come in two kinds

(1) Ordinary



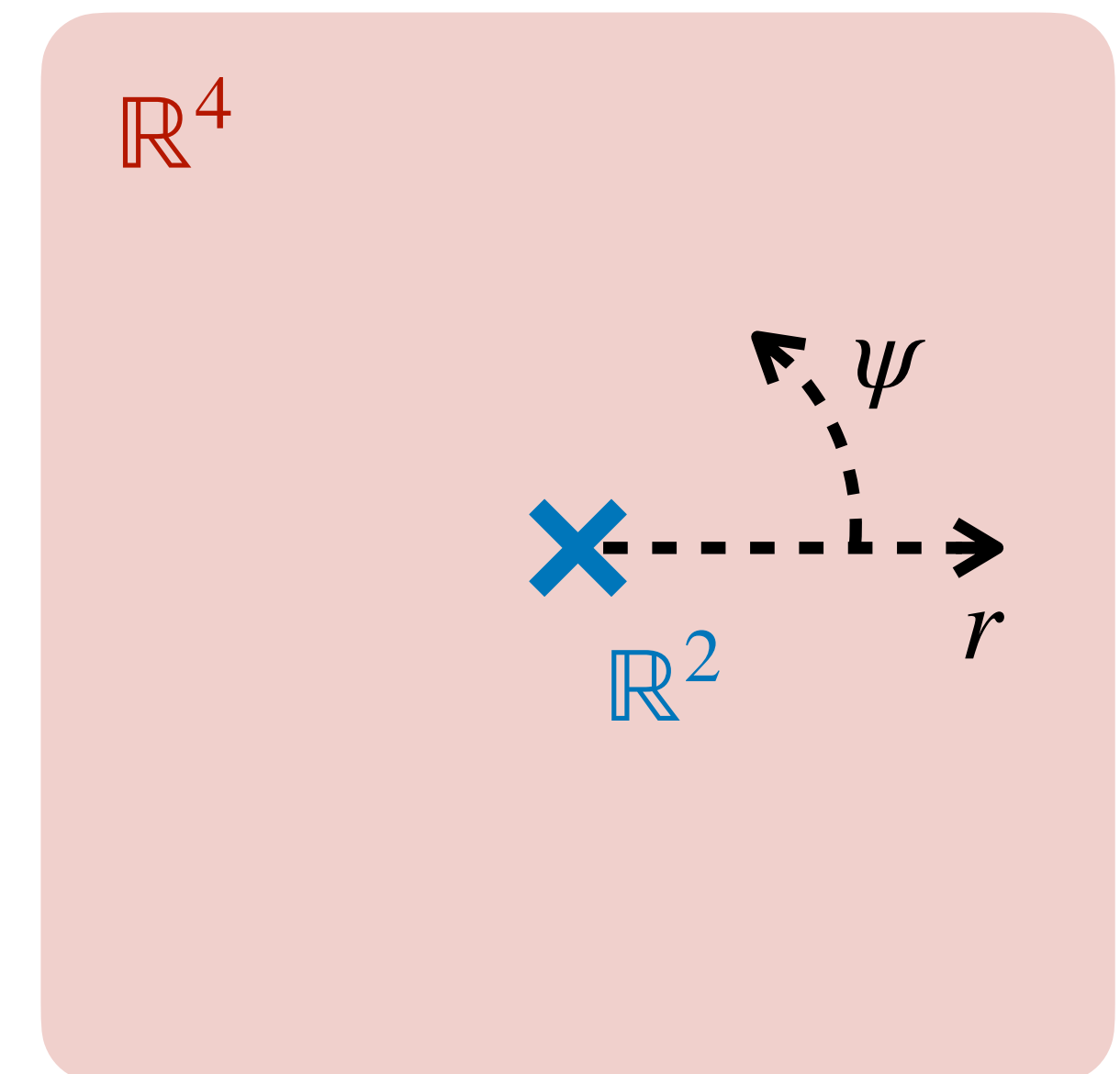
$$A = \alpha d\psi \text{ and } \Phi = \frac{e^{-i\psi}}{r}(\beta + i\gamma)$$

(2) Rigid ($\alpha, \beta, \gamma \rightarrow 0$)

$$A = \frac{t_3}{\log r} d\psi \text{ and } \Phi = \frac{e^{-i\psi}}{r \log r}(t_1 + it_2)$$

diagonal matrices

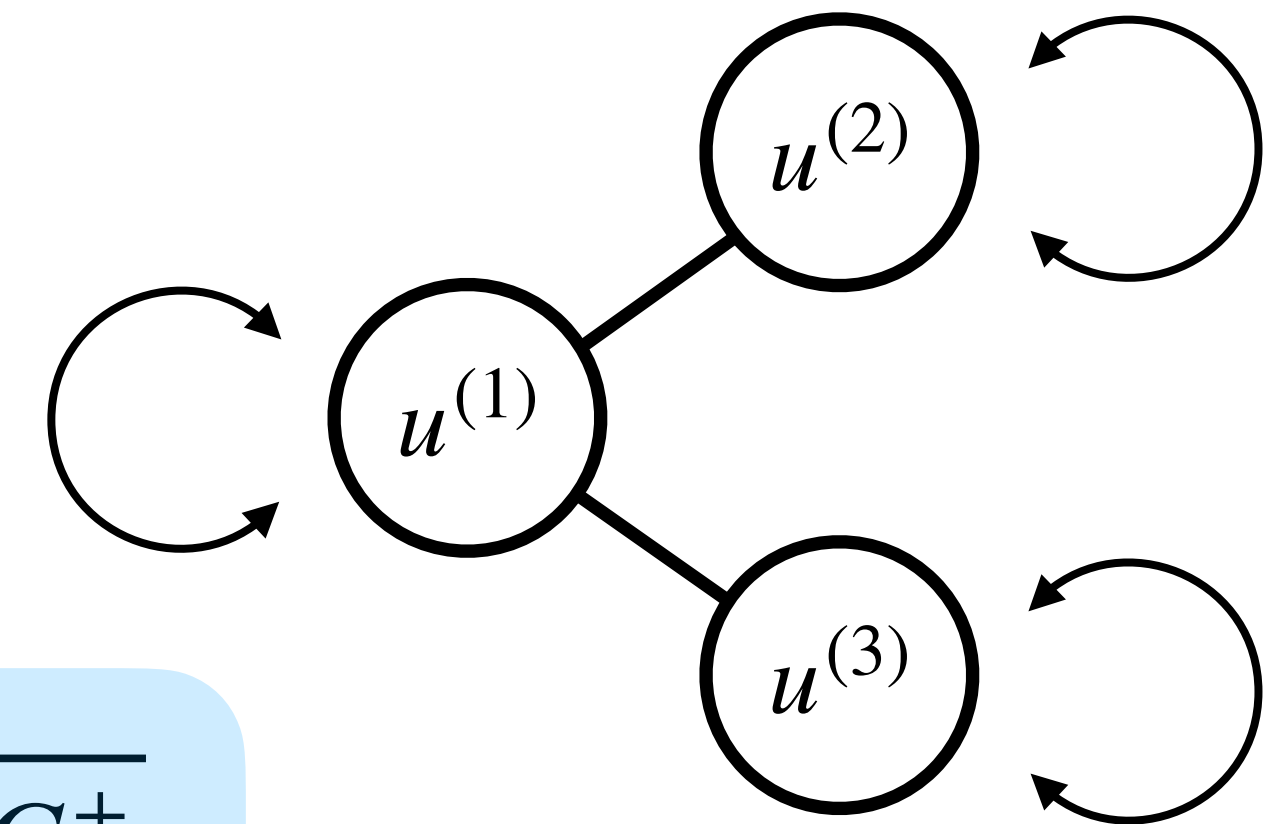
$\text{su}(2)$ representation matrices



Overlaps for Ordinary Gukov-Witten Defect

- SO(6) sector: $\langle \mathcal{B} | \mathbf{u} \rangle = 0$ unless number of Bethe roots $\mathcal{N}_1 = 2\mathcal{N}_2 = 2\mathcal{N}_3$

and roots are **chirally** paired $(u_1^{(i)}, -u_1^{(i)}, \dots, u_{\mathcal{N}_i/2}^{(i)}, -u_{\mathcal{N}_i/2}^{(i)})$

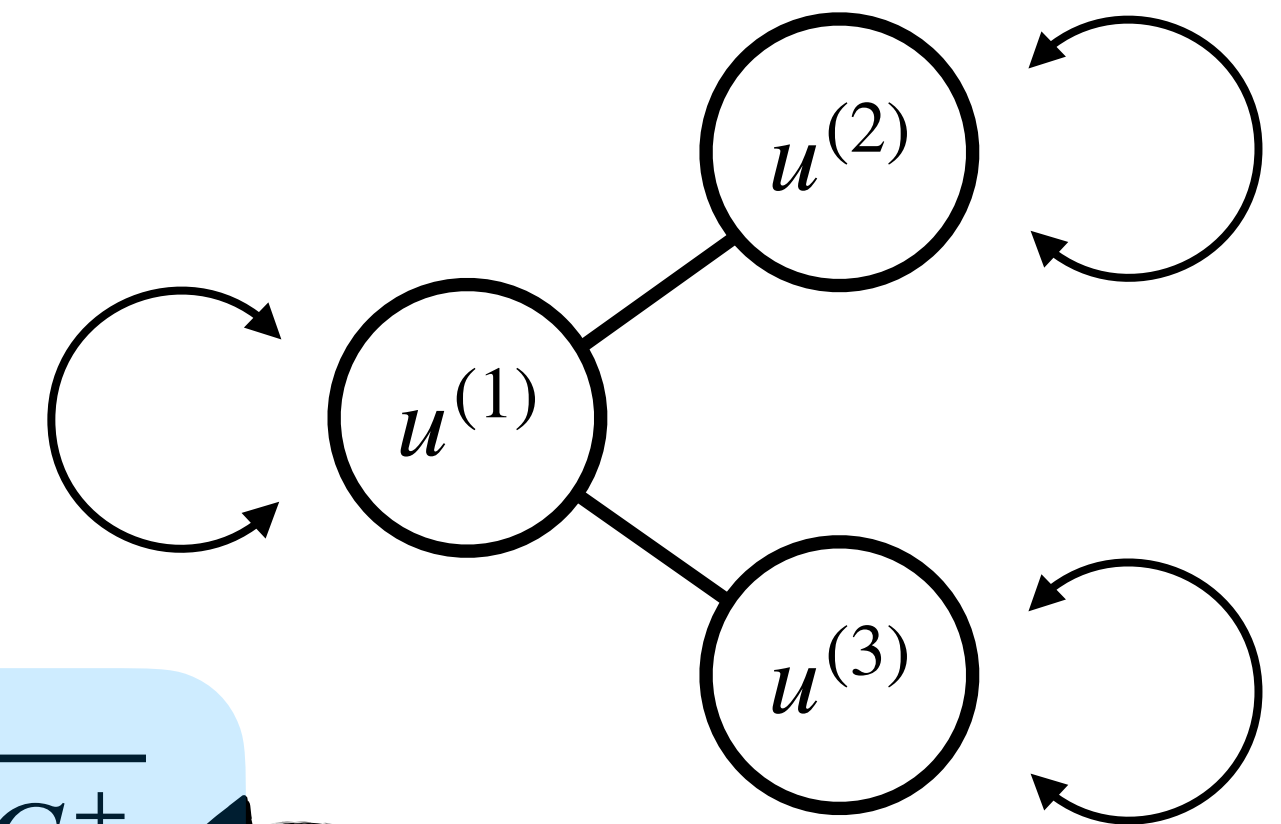


$$\langle \mathcal{B} | \mathbf{u} \rangle = \sum_{m=1}^M N_m e^{-i(L-\mathcal{N}_1)\psi_m(\beta_m^2 + \gamma_m^2)^{L/2}} \sqrt{\frac{Q_1(0)Q_1(i/2)}{Q_2(0)Q_2(i/2)Q_3(0)Q_3(i/2)} \frac{\det G^+}{\det G^-}}$$

Overlaps for Ordinary Gukov-Witten Defect

- SO(6) sector: $\langle \mathcal{B} | \mathbf{u} \rangle = 0$ unless number of Bethe roots $\mathcal{N}_1 = 2\mathcal{N}_2 = 2\mathcal{N}_3$

and roots are **chirally** paired $(u_1^{(i)}, -u_1^{(i)}, \dots, u_{\mathcal{N}_i/2}^{(i)}, -u_{\mathcal{N}_i/2}^{(i)})$



$$\langle \mathcal{B} | \mathbf{u} \rangle = \sum_{m=1}^M N_m e^{-i(L-\mathcal{N}_1)\psi_m} (\beta_m^2 + \gamma_m^2)^{L/2} \sqrt{\frac{Q_1(0)Q_1(i/2)}{Q_2(0)Q_2(i/2)Q_3(0)Q_3(i/2)} \frac{\det G^+}{\det G^-}}$$

$$\psi_m = \psi - \arg(\beta_m + i\gamma_m)$$

$$Q_i(x) = \prod_a (x - u_a^{(i)})$$

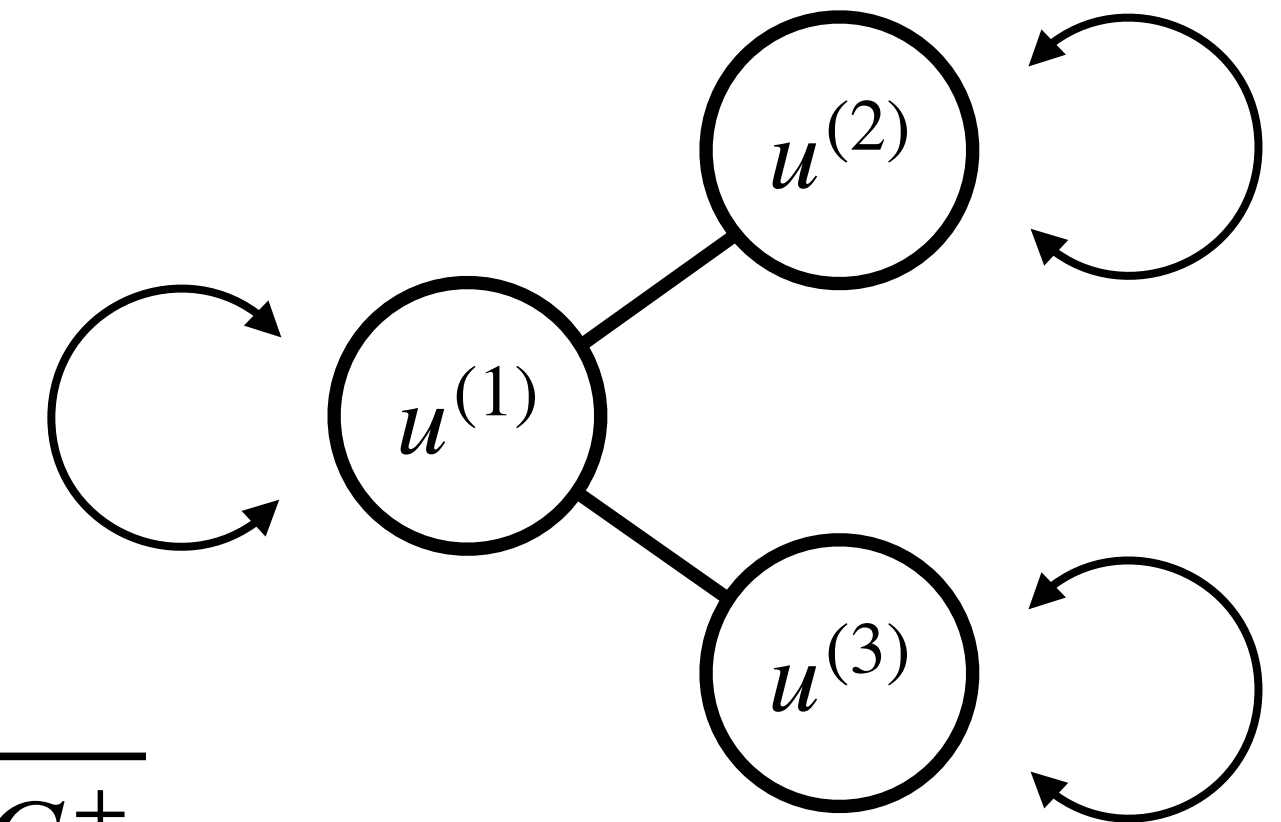
$$G^\pm = \left. \frac{\partial}{\partial u} \log (\text{Bethe equations}) \right|$$

$$\Phi = \frac{e^{-i\psi}}{\sqrt{2r}} \text{diag}((\beta_1 + i\gamma_1)1_{N_1}, \dots, (\beta_M + i\gamma_M)1_{N_M})$$

Overlaps for Ordinary Gukov-Witten Defect

- SO(6) sector: $\langle \mathcal{B} | \mathbf{u} \rangle = 0$ unless number of Bethe roots $\mathcal{N}_1 = 2\mathcal{N}_2 = 2\mathcal{N}_3$

and roots are **chirally** paired $(u_1^{(i)}, -u_1^{(i)}, \dots, u_{\mathcal{N}_1/2}^{(i)}, -u_{\mathcal{N}_1/2}^{(i)})$



$$\langle \mathcal{B} | \mathbf{u} \rangle = \sum_{m=1}^M N_m e^{-i(L-\mathcal{N}_1)\psi_m(\beta_m^2 + \gamma_m^2)^{L/2}} \sqrt{\frac{Q_1(0)Q_1(i/2)}{Q_2(0)Q_2(i/2)Q_3(0)Q_3(i/2)} \frac{\det G^+}{\det G^-}}$$

- Trivial SU(2) sector since it is reached by taking $\mathcal{N}_2 = \mathcal{N}_3 = 0 \implies \mathcal{N}_1 = 0$

Overlaps for Ordinary Gukov-Witten Defect

- $SL(2)$ sector $\{Z, DZ, D^2Z, \dots\}$: depends on how we construct $|\mathcal{B}\rangle$
- \exists choice for which $|\mathcal{B}\rangle$ not integrable! [see also Holguin-Kawai '25]

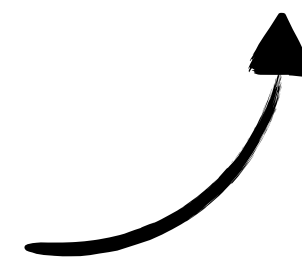
Overlaps for Ordinary Gukov-Witten Defect

- $SL(2)$ sector $\{Z, DZ, D^2Z, \dots\}$: depends on how we construct $|\mathcal{B}\rangle$
- \exists choice for which $|\mathcal{B}\rangle$ not integrable!

[see also Holguin-Kawai '25]

SU(2)	SO(6)	SL(2)
0	✓	✗

just a coincidence at leading order



Gukov-Witten Defects

- 2d $N=(4,4)$ surface defects come in two kinds

(1) Ordinary

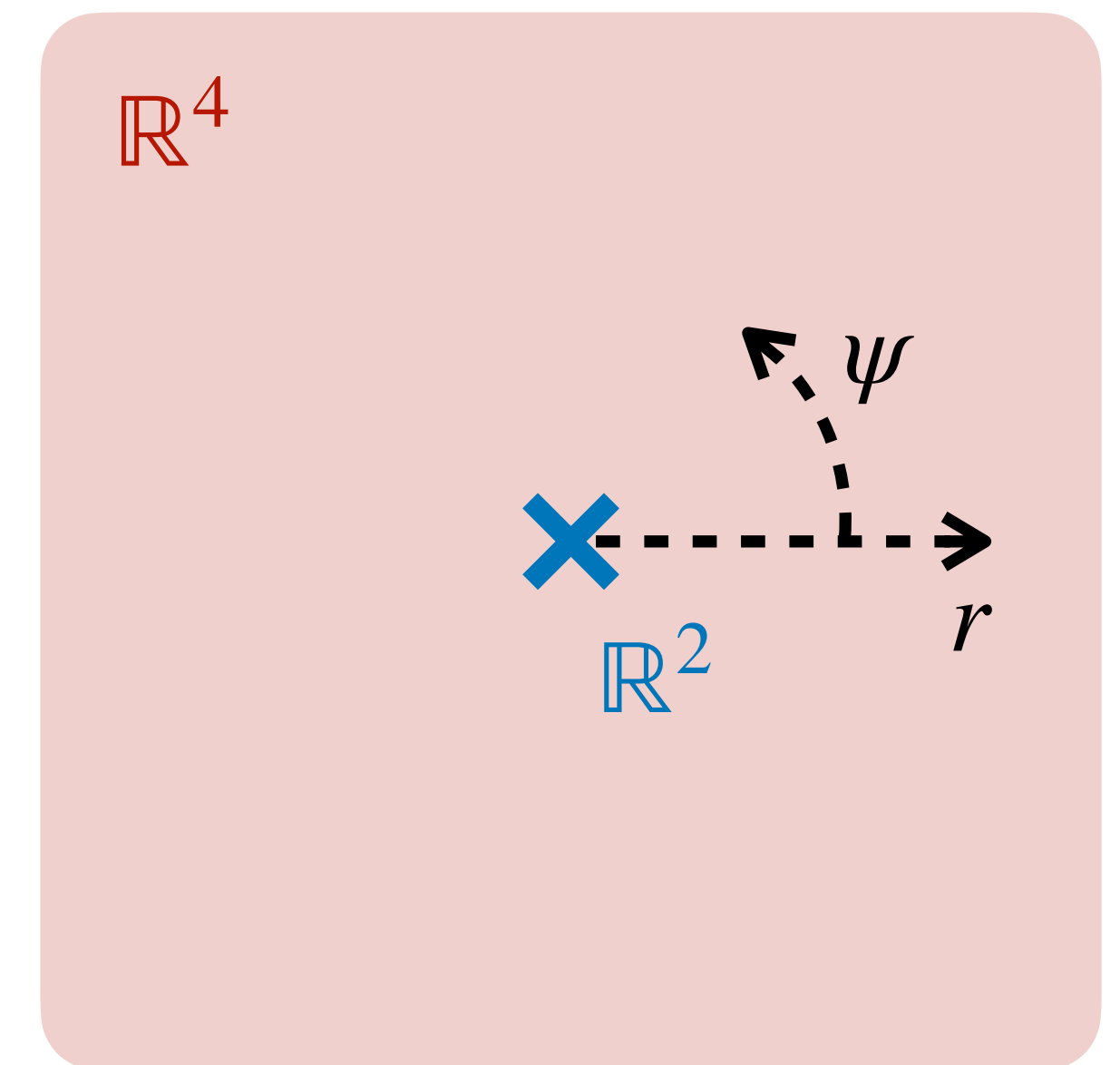
$$A = \alpha d\psi \text{ and } \Phi = \frac{e^{-i\psi}}{r}(\beta + i\gamma)$$

(2) Rigid ($\alpha, \beta, \gamma \rightarrow 0$)

➔ $A = \frac{t_3}{\log r} d\psi \text{ and } \Phi = \frac{e^{-i\psi}}{r \log r} (t_1 + it_2)$

diagonal matrices

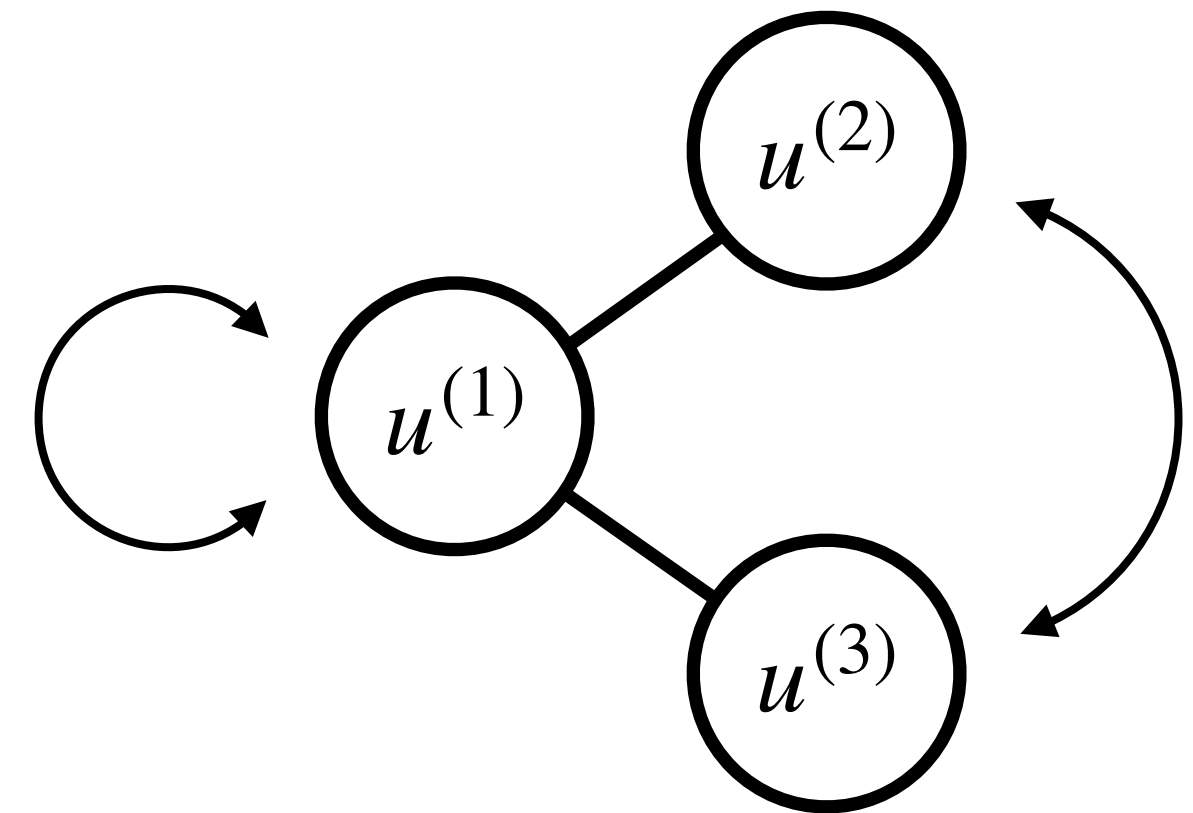
$\text{su}(2)$ representation matrices



Overlaps for Rigid Gukov-Witten Defect

- SO(6) sector: $\langle \mathcal{B} | \mathbf{u} \rangle = 0$ unless roots are **achirally** paired.

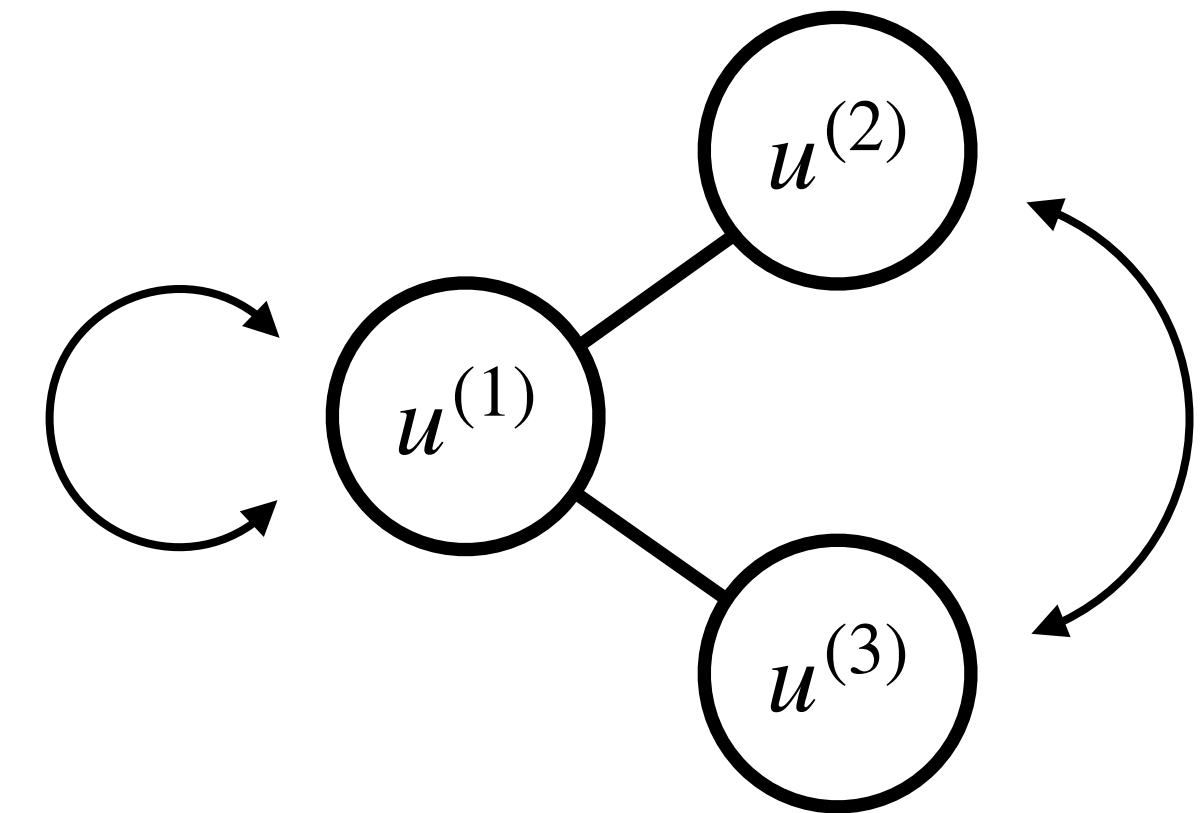
When $k = 2$ only, we find $\langle \mathcal{B} | \mathbf{u} \rangle = \frac{1}{2^{L-1}} \sqrt{\frac{Q_1(i/2) \det G^+}{Q_1(0) \det G^-}}$



Overlaps for Rigid Gukov-Witten Defect

- SO(6) sector: $\langle \mathcal{B} | \mathbf{u} \rangle = 0$ unless roots are **achirally** paired.

When $k = 2$ only, we find $\langle \mathcal{B} | \mathbf{u} \rangle = \frac{1}{2^{L-1}} \sqrt{\frac{Q_1(i/2)}{Q_1(0)} \frac{\det G^+}{\det G^-}}$



- SU(2) sector actually integrable for all k

$$|\langle \mathcal{B} | \mathbf{u} \rangle| = S_k Q\left(\frac{ik}{2}\right) \sqrt{Q(i/2)Q(0) \frac{\det G^+}{\det G^-}}, \text{ where } S_k = \sum_{q=-\frac{k-1}{2}}^{\frac{k-1}{2}} \frac{q^L}{Q\left(\frac{2q+1}{2}i\right) Q\left(\frac{2q-1}{2}i\right)}$$

Overlaps for Rigid Gukov-Witten Defect

- For SL(2) sector $\{Z, DZ, D^2Z, \dots\}$ restrict to leading singularity $\sim 1/\log^L r$
- $\langle \mathcal{B} | \mathbf{u} \rangle = 0$ unless all \mathcal{N} roots are paired,

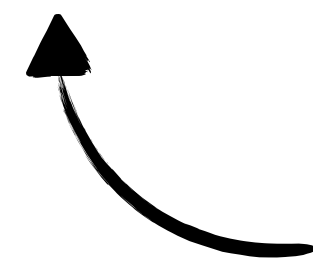
$$\langle \mathcal{B} | \mathbf{u} \rangle = \frac{\sin^{\mathcal{N}} \psi}{r^{\mathcal{N}}} \frac{1}{2^{L-1}} \sqrt{\frac{Q(i/2) \det G^+}{Q(0) \det G^-}} \quad \text{for } k = 2 \text{ only}$$

Overlaps for Rigid Gukov-Witten Defect

- For SL(2) sector $\{Z, DZ, D^2Z, \dots\}$ restrict to leading singularity $\sim 1/\log^L r$
- $\langle \mathcal{B} | \mathbf{u} \rangle = 0$ unless all \mathcal{N} roots are paired,

$$\langle \mathcal{B} | \mathbf{u} \rangle = \frac{\sin^{\mathcal{N}} \psi}{r^{\mathcal{N}}} \frac{1}{2^{L-1}} \sqrt{\frac{Q(i/2) \det G^+}{Q(0) \det G^-}} \quad \text{for } k = 2 \text{ only}$$

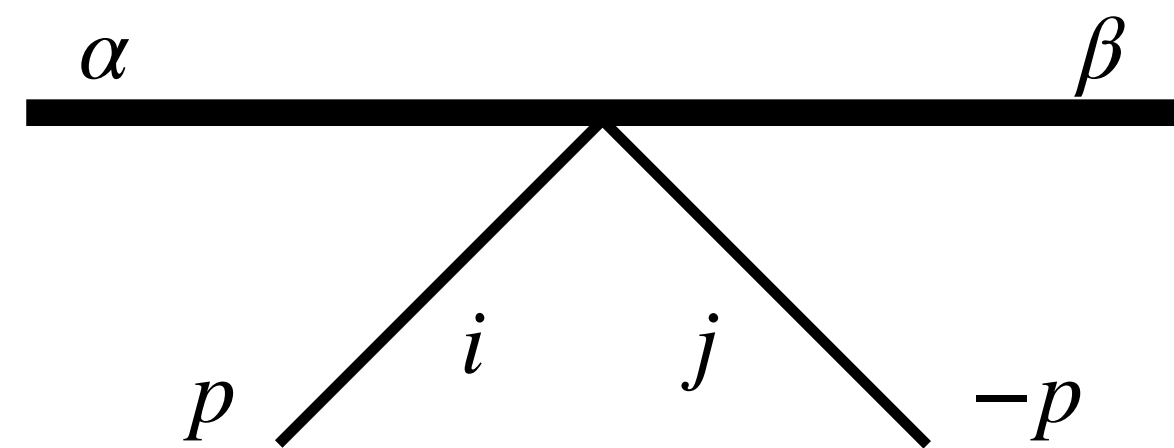
SU(2)	SO(6)	SL(2)
✓ all k	✓ k=2 only	✓ k=2 only



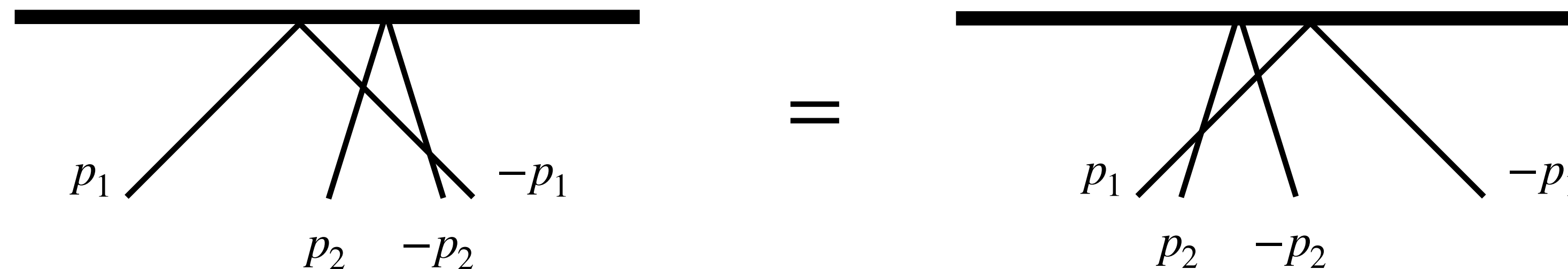
$k > 2$ just a coincidence at leading order

Sketch of Derivation

- Key ingredient is the K-matrix $(K_{i,j}(u))^{\alpha,\beta}$
- Amplitude of two excitations with opposite momenta annihilated by $|\mathcal{B}\rangle$



- Integrable scattering off a boundary needs to satisfy the K-Yang-Baxter eq



Sketch of Derivation

- Boundary state, K-matrix and monodromy matrix T satisfy KT-relation

- $K_{i,\ell}(u)\omega_I \mathcal{L}_{\ell,j;I,J}(u) = \omega_I K_{\ell,j}(u) \mathcal{L}_{i,\ell;I,J}(-u)$

$\langle \mathcal{B} | = \text{tr} \left(\sum_{I=1}^d \omega^I \langle I | \right)^{\otimes L}$

$T_0 = \mathcal{L}_{0,L} \cdots \mathcal{L}_{0,1}$

- Combines K-Yang-Baxter equation and $Q^{\text{odd}} | \mathcal{B} \rangle = 0$ condition

Sketch of Derivation

- Gombor '24: general method to find $\langle \mathcal{B} | \mathbf{u} \rangle$ from solution of the KT-relation

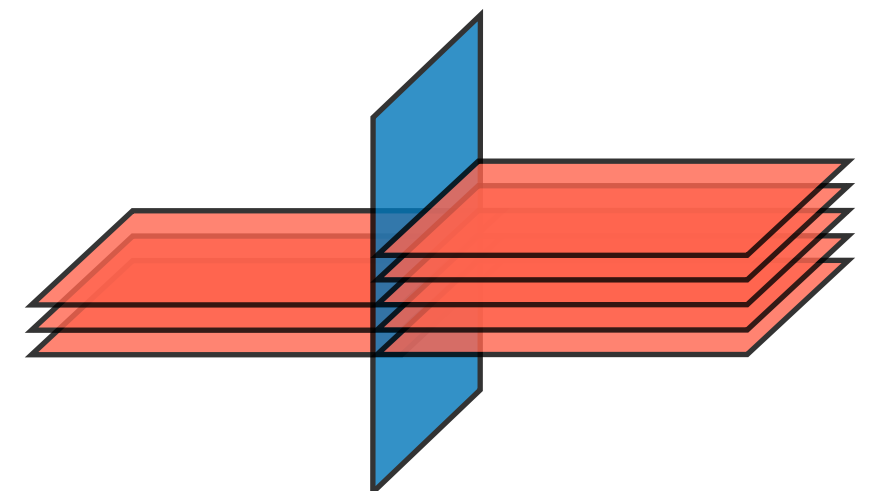
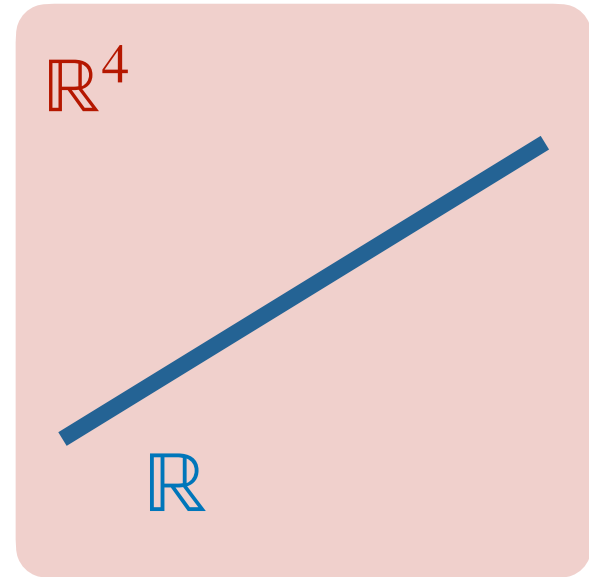
size of matrix ω^I # distinct types of roots # distinct roots of a given type

$$\langle \mathcal{B} | \mathbf{u} \rangle = \sum_{\gamma=1}^k \beta_{\gamma} \left(\prod_{a=1}^{n_+} \prod_{j=1}^{n_a} \tilde{\mathcal{F}}_{\gamma}^{(a)}(u_{a,j}) \right) \sqrt{\frac{\det G^+}{\det G^-}}$$

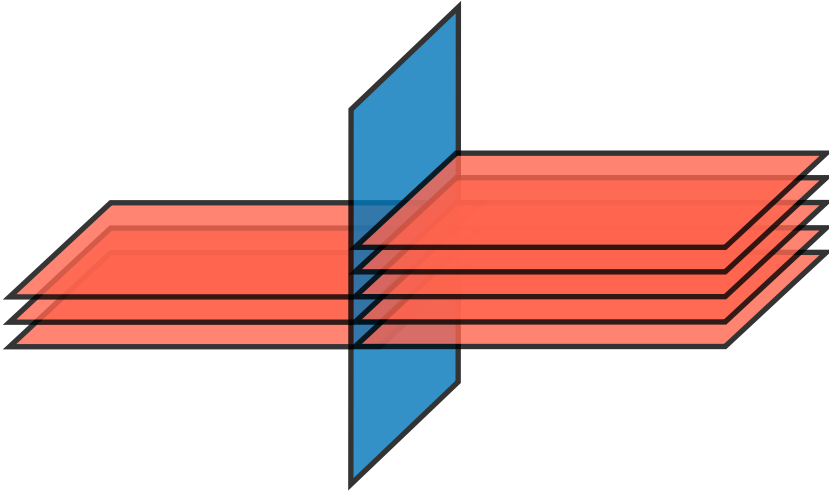
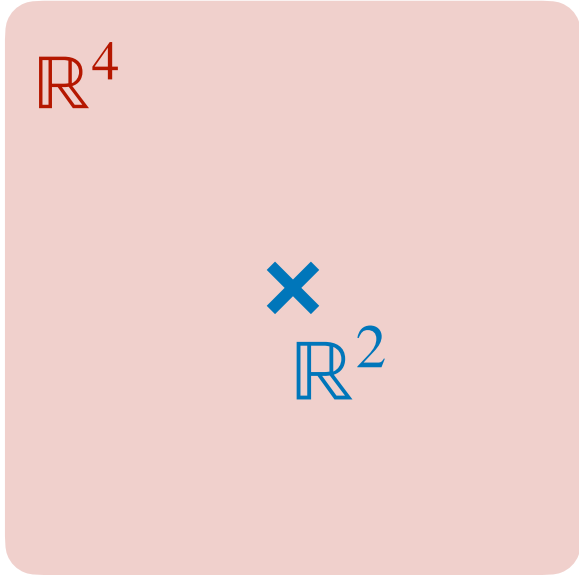
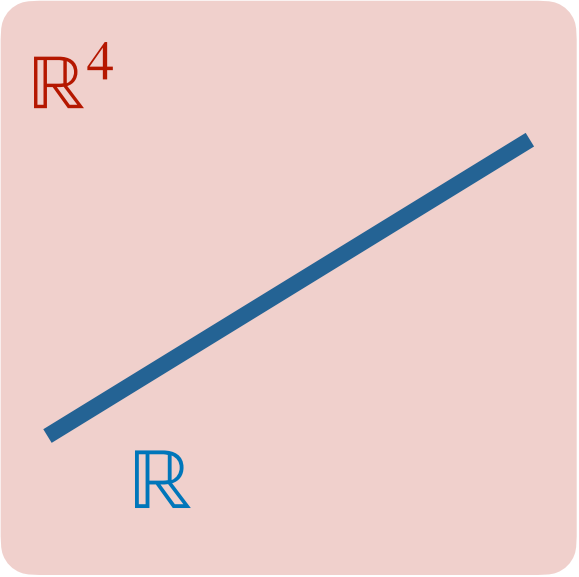
eigenvalues of $\langle \mathcal{B} | 0 \rangle$ constructed out of K universal part

- Pair structure determined by reflection algebra $Y(\mathfrak{g}, \mathfrak{h})$

Integrable 1/2-BPS Defects in $N=4$ SYM

Co-dim 1	Co-dim 2	Co-dim 3
<p data-bbox="616 909 982 966">D3-D5 system</p>  <p data-bbox="633 1397 932 1491">$\phi_i \rightarrow t_i^{(k)}/r$</p>	<p data-bbox="1582 1153 1749 1416">?</p>	<p data-bbox="2282 909 2815 966">1/2-BPS 't Hooft line</p>  <p data-bbox="2399 1397 2698 1491">$\phi^I \rightarrow n^I/r$</p>

Integrable 1/2-BPS Defects in $N=4$ SYM

Co-dim 1	Co-dim 2	Co-dim 3
<p data-bbox="616 915 972 971">D3-D5 system</p>  <p data-bbox="636 1399 926 1487">$\phi_i \rightarrow t_i^{(k)}/r$</p>	<p data-bbox="1319 915 2009 971">Rigid Gukov-Witten defect?</p>  <p data-bbox="1392 1399 1725 1581">$A_\psi \rightarrow \frac{t_3^{(2)}}{\log r},$</p> <p data-bbox="1392 1643 1975 1761">$\Phi \rightarrow \frac{e^{-i\psi}}{r \log r} (t_1^{(2)} + it_2^{(2)})$</p>	<p data-bbox="2285 915 2795 971">1/2-BPS 't Hooft line</p>  <p data-bbox="2409 1399 2692 1487">$\phi^I \rightarrow n^I/r$</p>

Outlook

- Loop corrections?
- Rigid defect has same symmetries as $\langle \mathcal{D}\mathcal{D}\mathcal{O} \rangle$ [Jiang-Komatsu-Vescovi '19]
→ bootstrap all-loop asymptotic formula?
- Rigid Gukov-Witten defects in holography, SUSY localisation, VOA, AGT?
[WIP with R. Izquierdo Garcia, C. Kristjansen, A. O'Bannon, James Ratcliffe, R. Rodgers]
- Space of integrable defects? 2d $N=(8,0)$ SUSY?