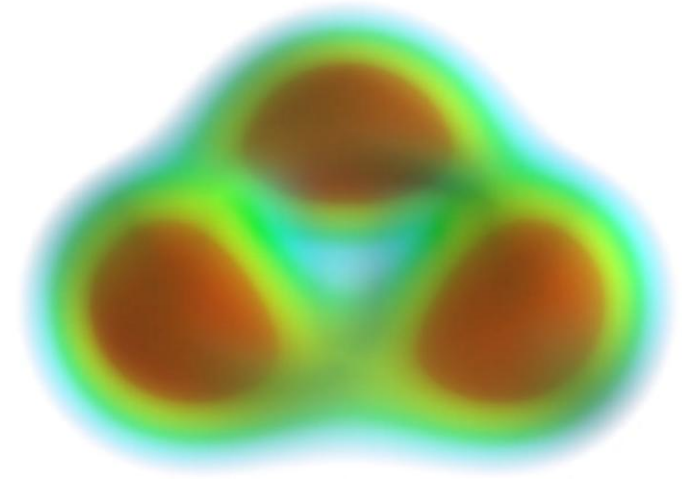
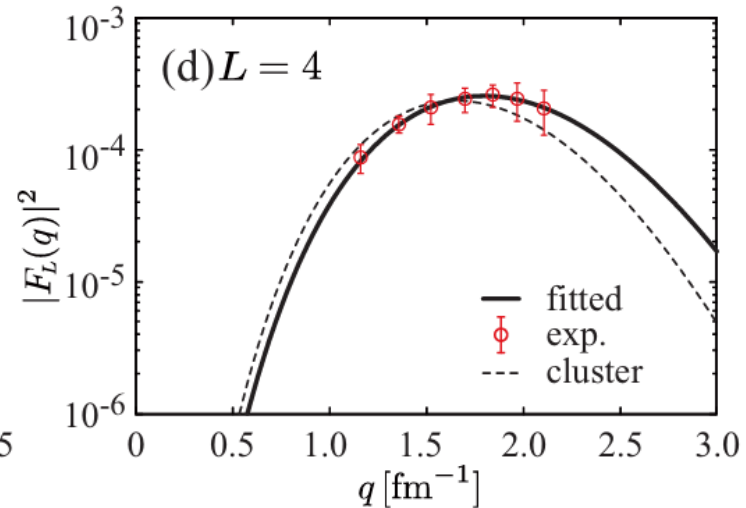
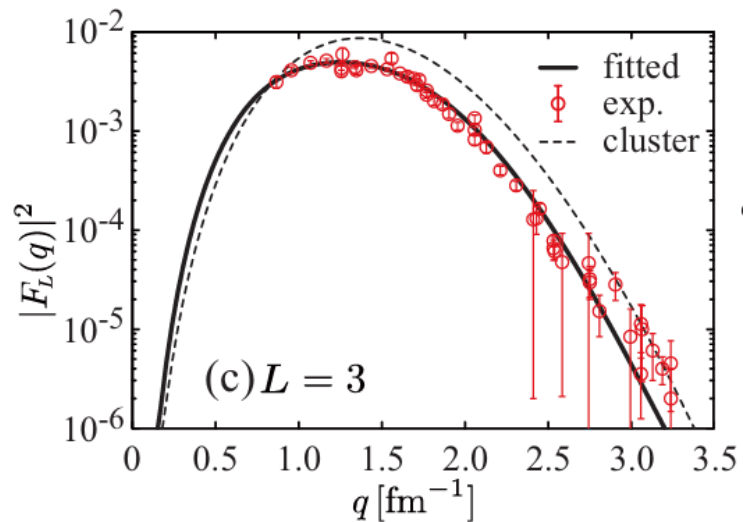
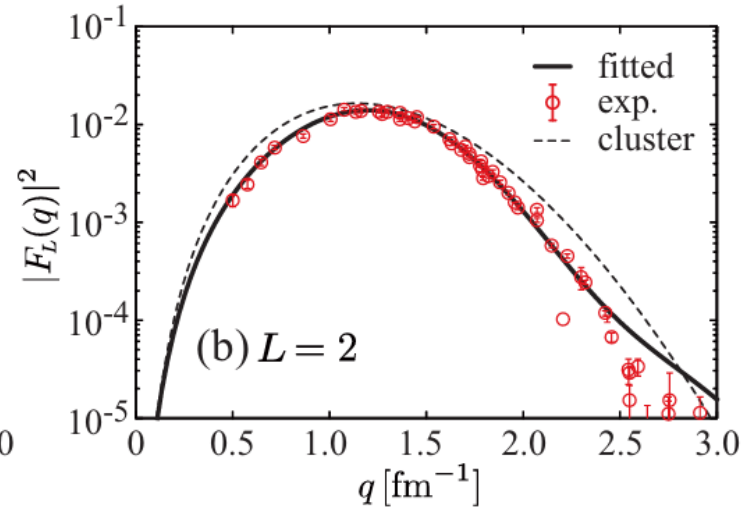
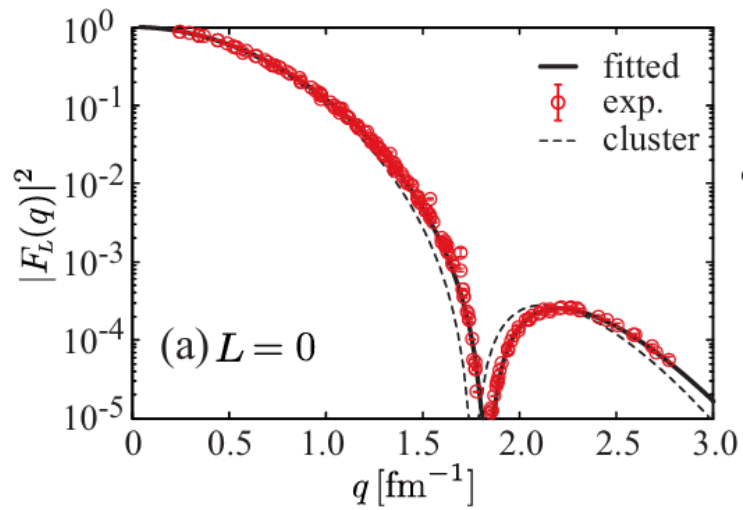


Shape of Carbon-12

M. Kimura (RIKEN)

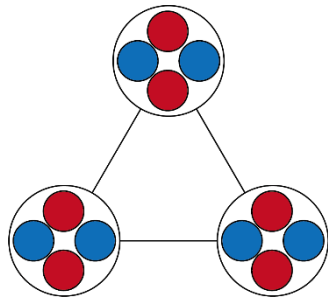


Spinning triangle of three alphas

Spinning triangle of three alphas

Three alpha particles locate in a triangle configuration (D_{3h} symmetry)

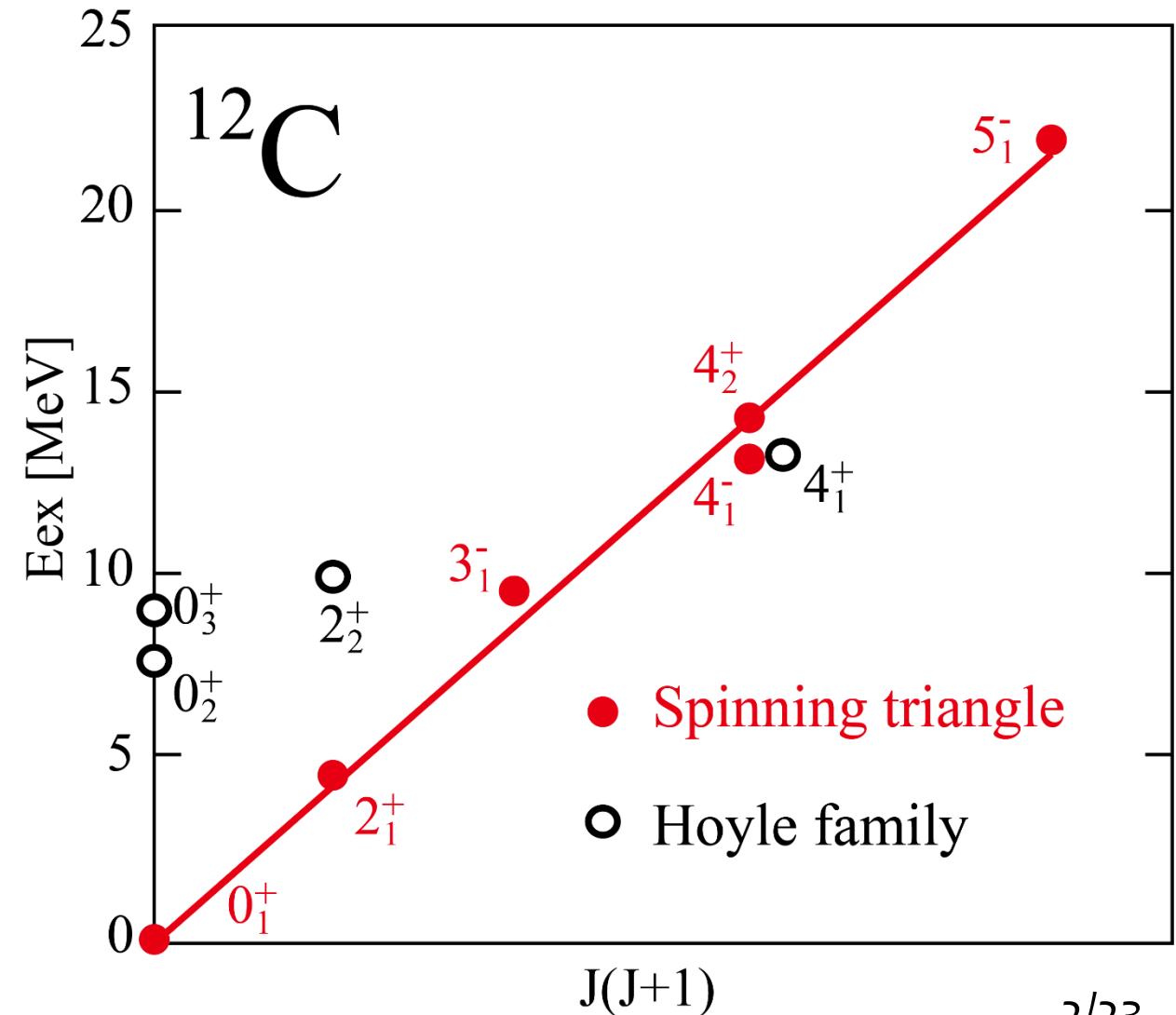
R. Bijker and F. Iachello, PRC61, 067305 (2000).



⊙ Due to the symmetry, 0^+ , 2^+ , 3^- , 4^\pm , 5^- states constitute “ground band”

⊙ 4^+ and 4^- states should degenerate due to “symmetry”

D.J. Marín-Lámbarri, et al. PRL113, 012502 (2014)

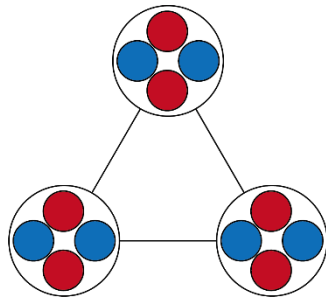


Spinning triangle of three alphas

Spinning triangle of three alphas

Three alpha particles locate
in a triangle configuration (D_{3h} symmetry)

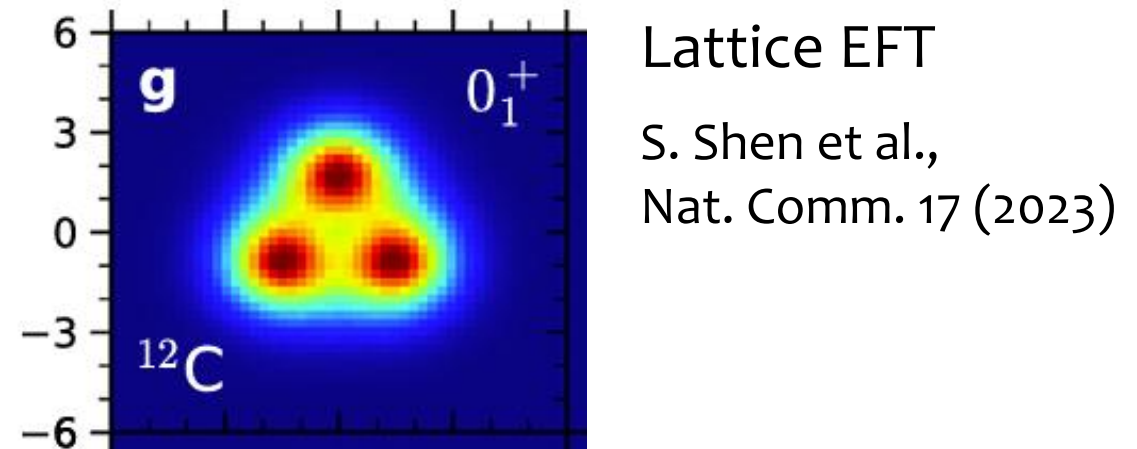
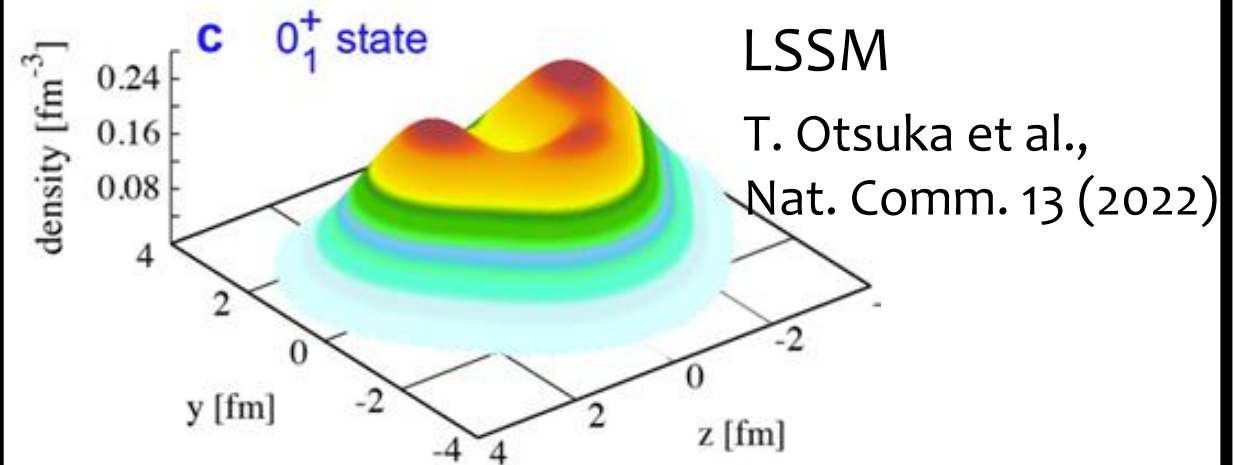
R. Bijker and F. Iachello, PRC61, 067305 (2000).



© Due to the symmetry, 0^+ , 2^+ , 3^- , 4^\pm , 5^-
states constitute “ground band”

© 4^+ and 4^- states should degenerate
due to “symmetry”

Shape of ^{12}C from theories



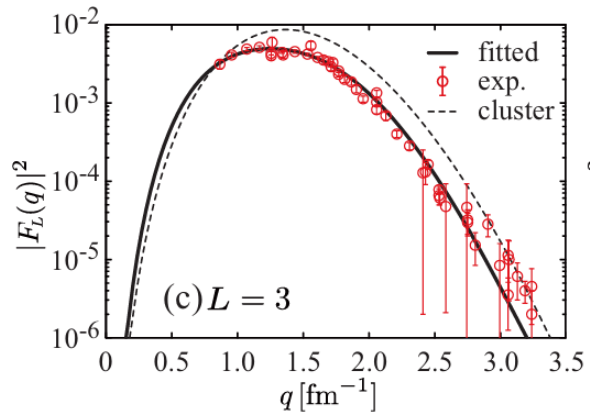
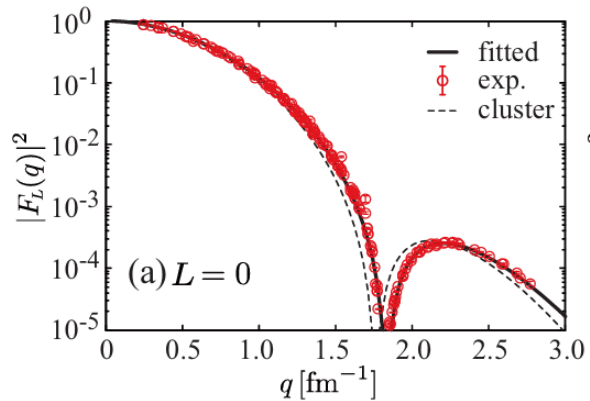
Rebuilding nuclear shape from Exp. data

M.K. and Y. Taniguchi, EPJA 60, (2024)

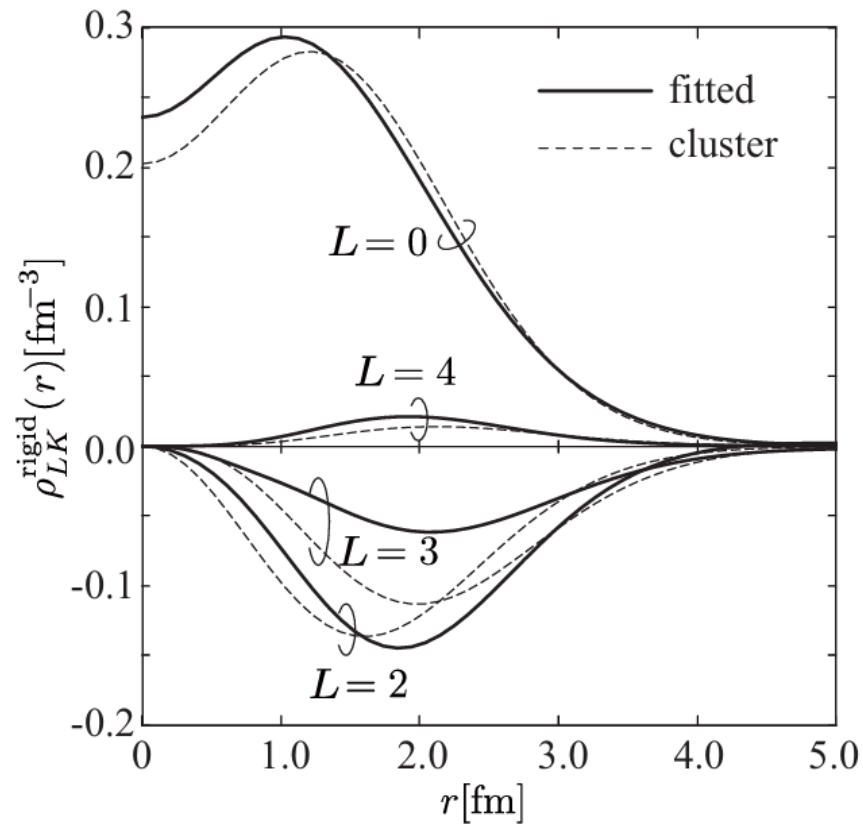


Rebuilding nuclear shape from observables

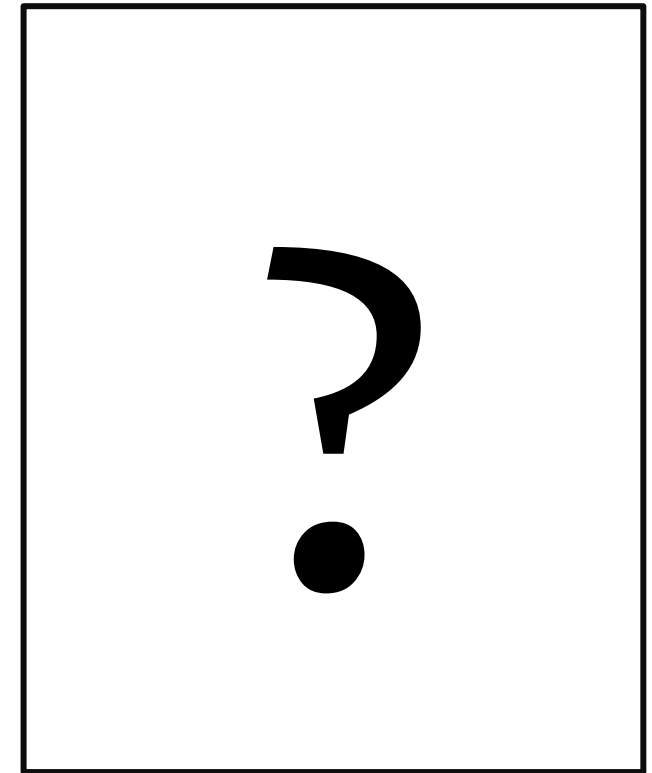
$^{12}\text{C}(e,e')$



Transition Density



Shape



Rebuilding nuclear shape from observables

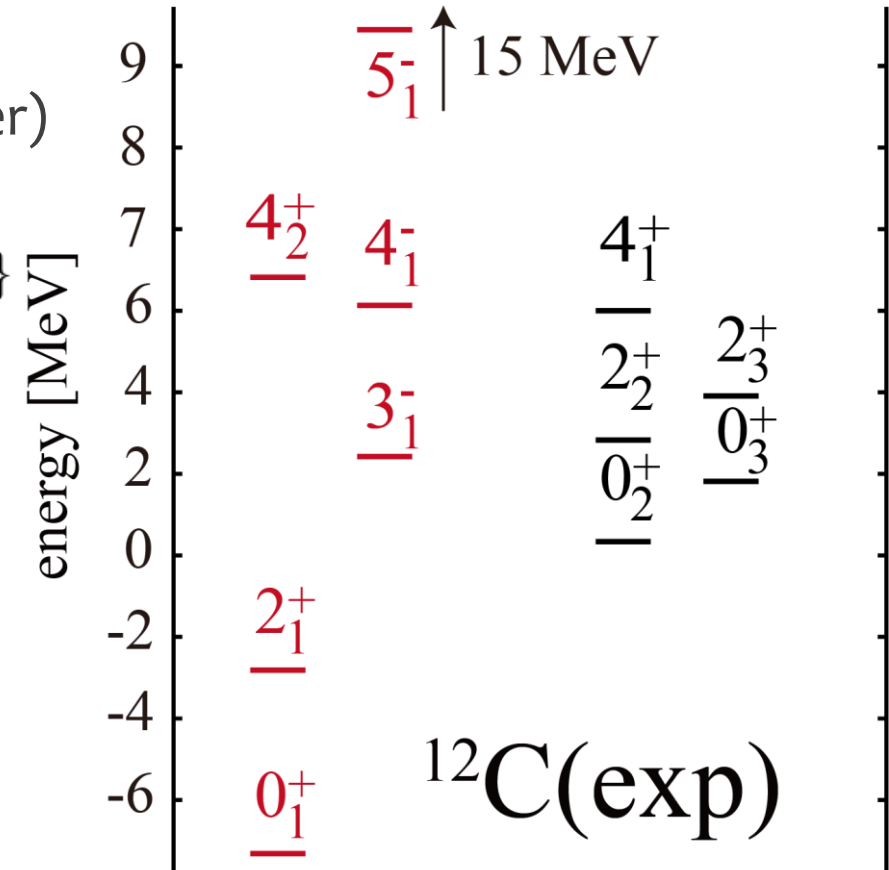
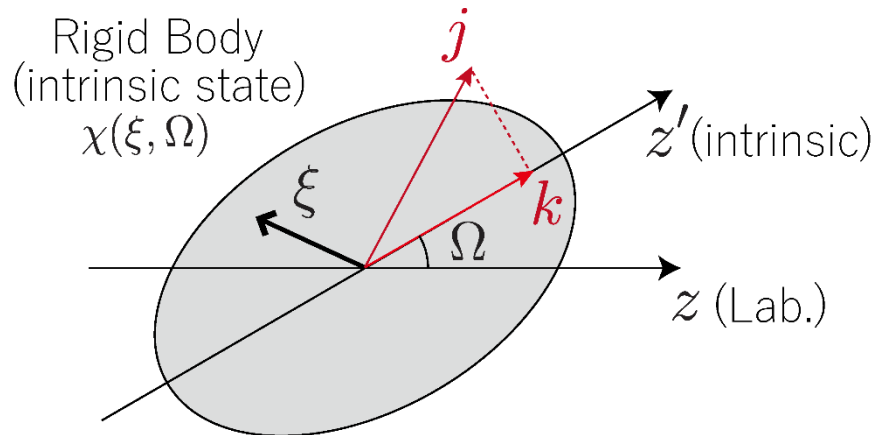
I assume that

“the 0^+ , 2^+ , 3^- and 4^+ states constitute *the ground band* sharing the same intrinsic state”

(* This is the only assumption I make in this talk.

Non-axial and parity-asymmetric rigid rotor (a textbook matter)

$$\underbrace{\Psi_{L_z K}^{L \Pi}}_{\text{laboratory (observable)}} = \sqrt{\frac{2L+1}{16\pi^2(1+\delta_{K0})}} \underbrace{\chi(\xi, \Omega)}_{\text{intrinsic (unobservable)}} \{ D_{L_z K}^L(\Omega) + \Pi(-)^{L+K} D_{L_z -K}^L(\Omega) \}$$



Rebuilding nuclear shape from observables

Electron scattering: Within PWIA, the cross section is proportional to square of Form factor

$$\frac{d\sigma_L^{\text{obs}}}{d\theta} = |F_L(q)|^2 \frac{d\sigma_L^{\text{Mott}}}{d\theta},$$

$$F_L(q) := \frac{\sqrt{4\pi(2L+1)}}{Z} \int_0^\infty r^2 dr j_L(qr) \rho^{0 \rightarrow L^\Pi}(\mathbf{r}) / Y_{L0}(\hat{r}), \quad \rho^{0^+ \rightarrow L^\Pi}(\mathbf{r}) := \langle L^\Pi, 0 | \rho(\mathbf{r}) | 0_1^+, 0 \rangle$$

Form Factor

transition density

Insert the rigid-rotor wave function into this definition, one gets

$$\rho^{0 \rightarrow L^\Pi}(\mathbf{r}) = \frac{1}{8\pi^2} \left(\frac{2L+1}{2(1+\delta_{K0})} \right)^{1/2} \int d\Omega \left\{ D_{L_z K}^{L*}(\Omega) + \Pi(-)^K D_{L_z - K}^{L*}(\Omega) \right\} \langle \chi(\xi, \Omega) | \rho(\mathbf{r}) | \chi(\xi, \Omega) \rangle.$$

transition density

rigid-body density (shape of ^{12}C)

$$= \frac{1}{\sqrt{2L+1}} \left(\frac{1+\Pi(-)^L}{1+\delta_{K0}} \right)^{1/2} \rho_{LK}^{\text{rigid}}(r) Y_{L0}(\hat{r}),$$

Multipole decomposition of the rigid-body density

(*) Multipole decomposition of the rigid-body density

$$\langle \chi(\xi, \Omega = 0) | \rho(\mathbf{r}) | \chi(\xi, \Omega = 0) \rangle = \frac{1}{2} \sum_{lm} \rho_{lm}^{\text{rigid}}(r) \{ Y_{lm}(\hat{r}) + (-)^m Y_{l-m}(\hat{r}) \}.$$

Rebuilding nuclear shape from observables

Summary of our assumption and numerical procedure

Assumption: The 0^+ , 2^+ , 3^- and 4^+ states constitute “the ground band” sharing the same intrinsic state

- ① Get the multipole decomposition of the rigid-body density from the observed transition densities

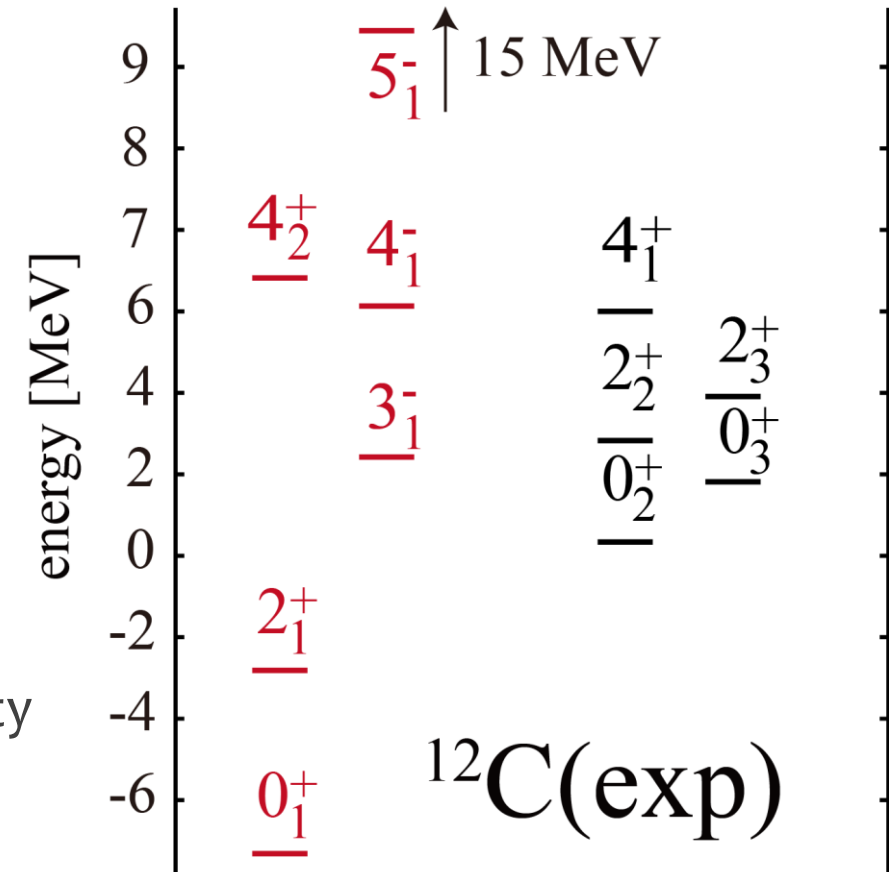
$$\underline{\rho^{0 \rightarrow L^\Pi}(\mathbf{r})} = \frac{1}{\sqrt{2L+1}} \left(\frac{1 + \Pi(-)^L}{1 + \delta_{K0}} \right)^{1/2} \underline{\rho_{LK}^{\text{rigid}}(r) Y_{L0}(\hat{r})},$$

**transition density
from $^{12}\text{C}(e,e')$ data**

**Multipole decomposition
of intrinsic density**

- ② Sum up all the multipole to reconstruct the rigid-body density

$$\rho^{\text{rigid}}(\mathbf{r}) = \frac{1}{2} \sum_{lm} \rho_{lm}^{\text{rigid}}(r) \{ Y_{lm}(\hat{r}) + (-)^m Y_{l-m}(\hat{r}) \}$$



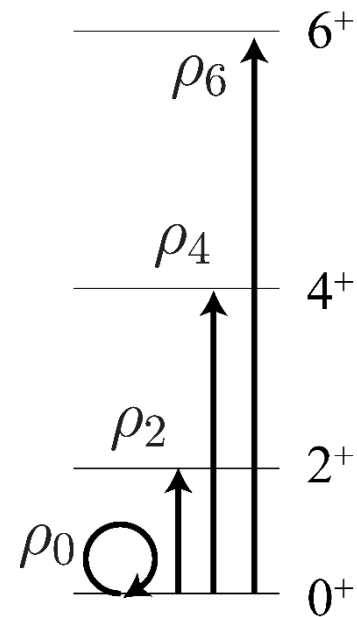
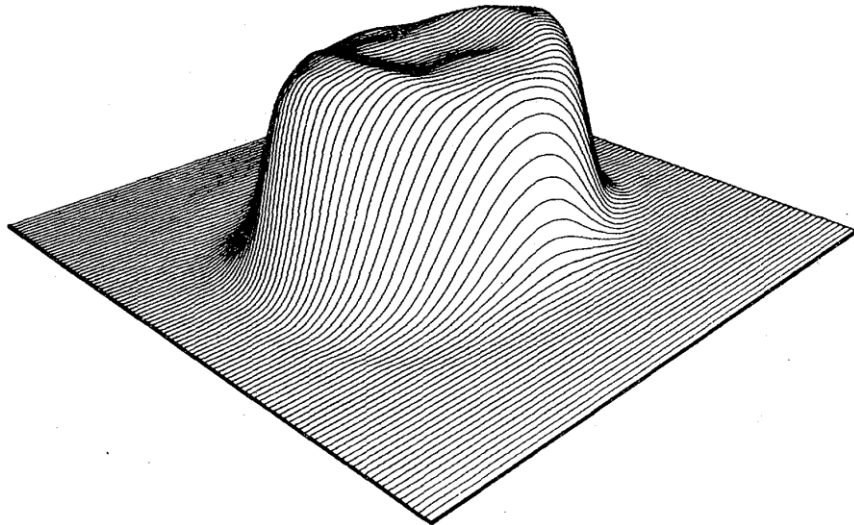
Rebuilding nuclear shape from observables

Actually, this is not a new idea.

In 1970's, it has been already applied to axial and parity-symmetric shapes.

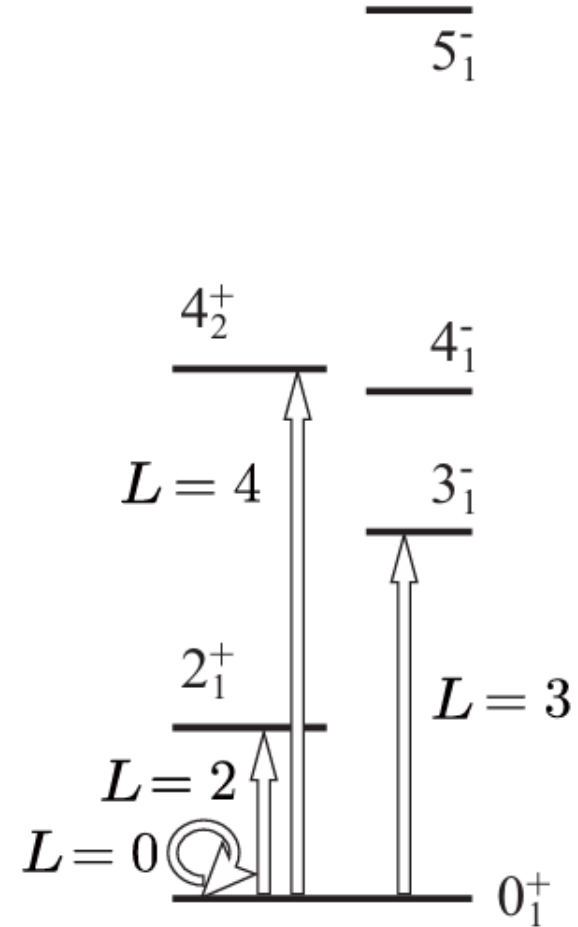
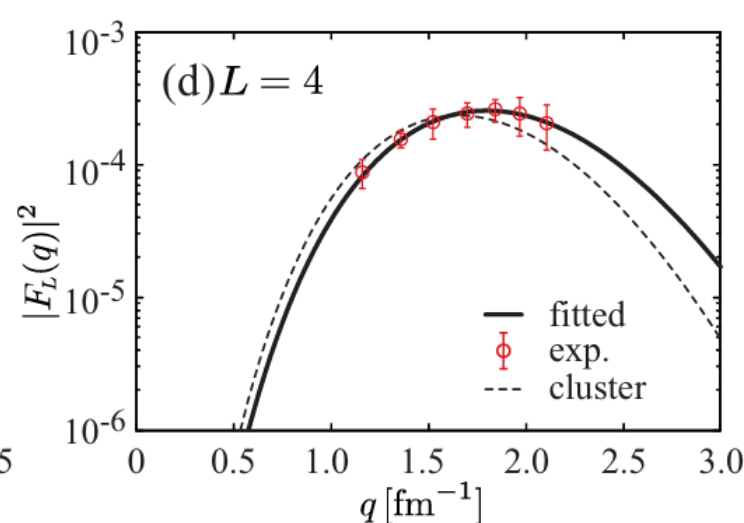
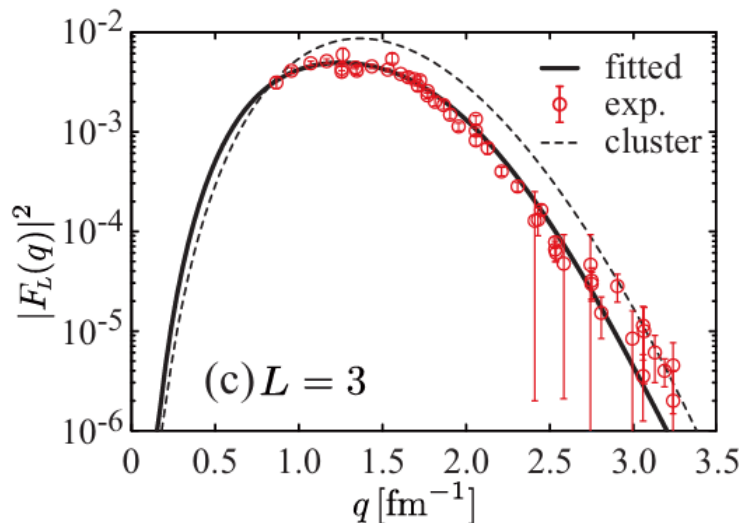
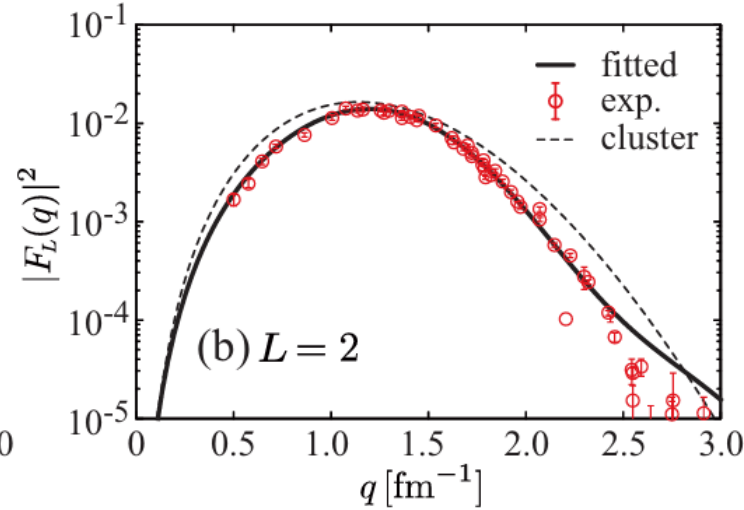
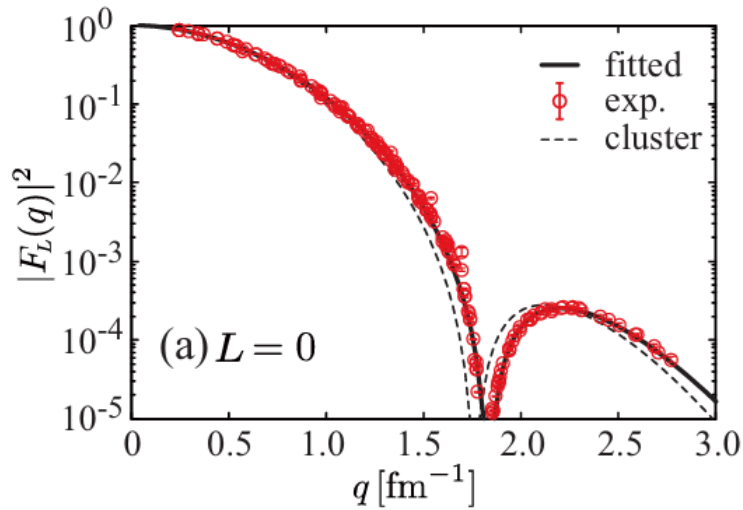
^{152}Sm ; rigid-body density

L.S. Carcman et al., PRC78, 1388 (1978).



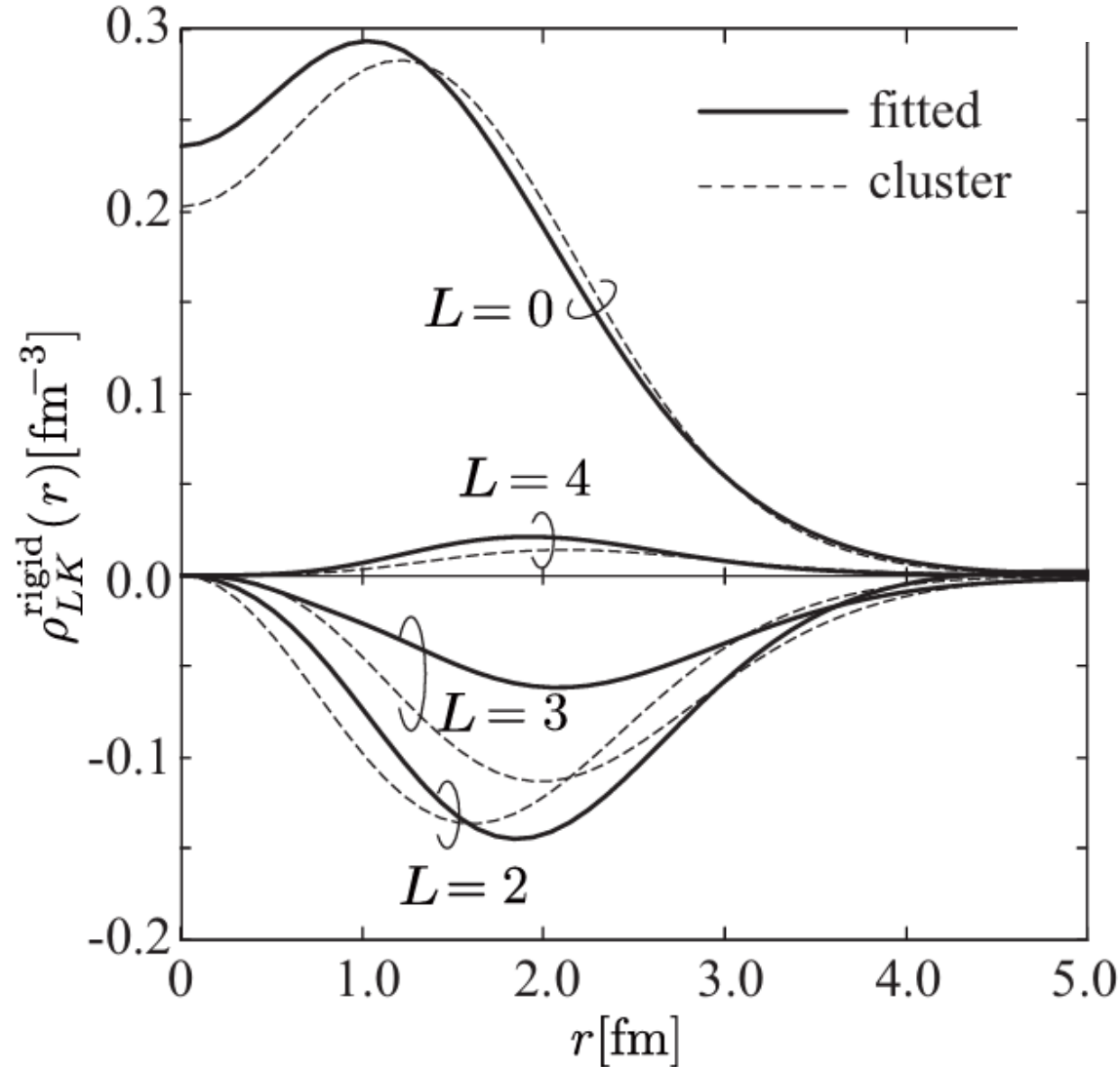
Rebuilding shape of ^{12}C

There are plenty of data for $^{12}\text{C}(e,e')$ $\frac{d\sigma_L^{\text{obs}}}{d\theta} = |F_L(q)|^2 \frac{d\sigma_L^{\text{Mott}}}{d\theta}$,

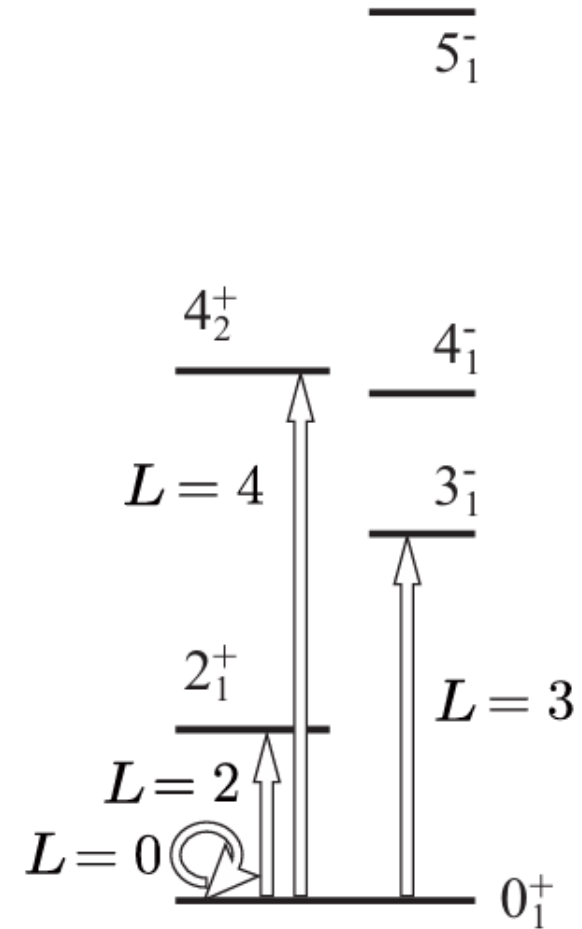


Rebuilding shape of ^{12}C

Inverse Fourier to get the densities.



$$F_L(q) := \frac{\sqrt{4\pi(2L+1)}}{Z} \int_0^\infty r^2 dr j_L(qr) \rho^{0 \rightarrow L^\Pi}(\mathbf{r}) / Y_{L0}(\hat{r}),$$



Rebuilding shape of ^{12}C

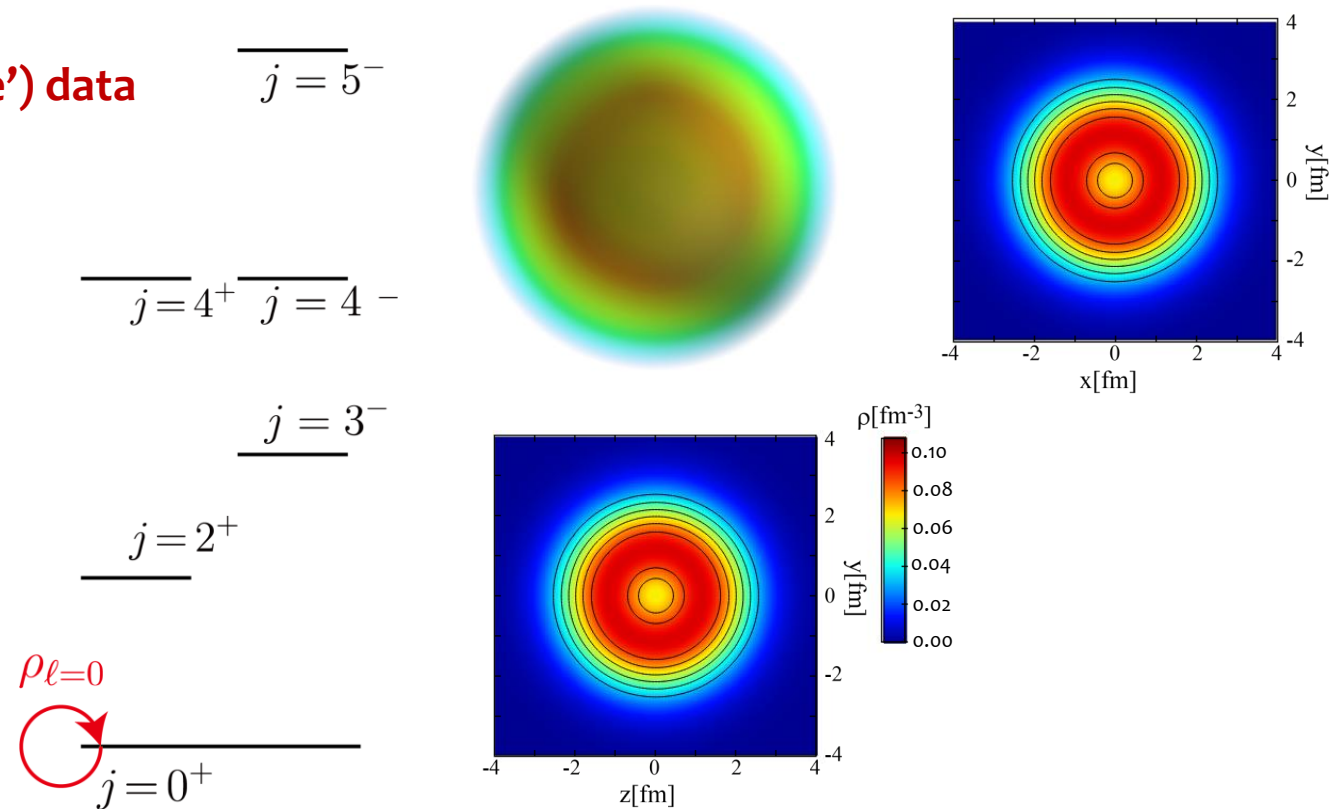
We rebuild the rigid density of ^{12}C by superposing the “observed” transition densities

$$\rho^{\text{rigid}}(\mathbf{r}) = \frac{1}{2} \sum_{lm} \rho_{lm}^{\text{rigid}}(r) \{ Y_{lm}(\hat{r}) + (-)^m Y_{l-m}(\hat{r}) \} \quad \ell = 0$$

Shape of ^{12}C

**multipole decomposition of
intrinsic density from $^{12}\text{C}(e,e')$ data**

- With only the $\ell = 0$ density, it is spherical



Rebuilding shape of ^{12}C

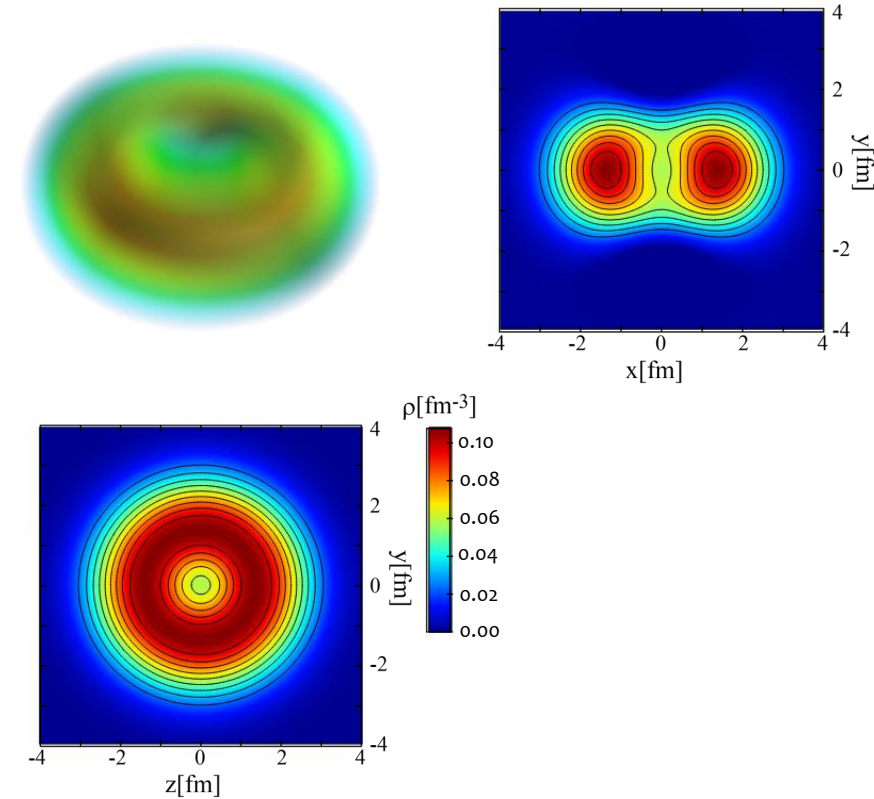
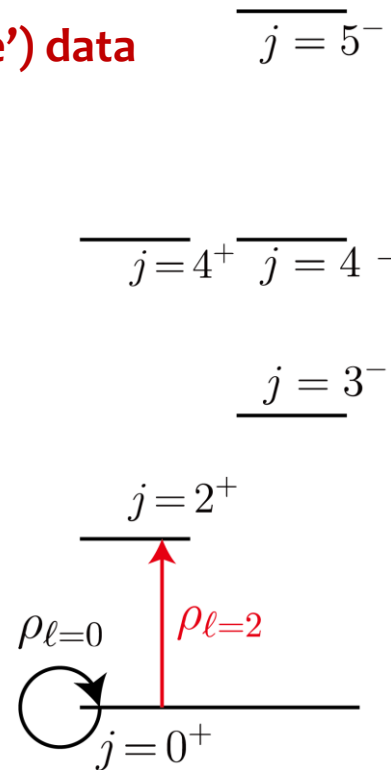
We rebuild the rigid density of ^{12}C by superposing the “observed” transition densities

$$\rho^{\text{rigid}}(\mathbf{r}) = \frac{1}{2} \sum_{lm} \rho_{lm}^{\text{rigid}}(r) \{ Y_{lm}(\hat{r}) + (-)^m Y_{l-m}(\hat{r}) \} \quad \ell = 0 + 2$$

Shape of ^{12}C

**multipole decomposition of
intrinsic density from $^{12}\text{C}(e,e')$ data**

- With only the $\ell = 0$ density, it is spherical
- The $\ell = 2$ density makes it toroidal shape



Rebuilding shape of ^{12}C

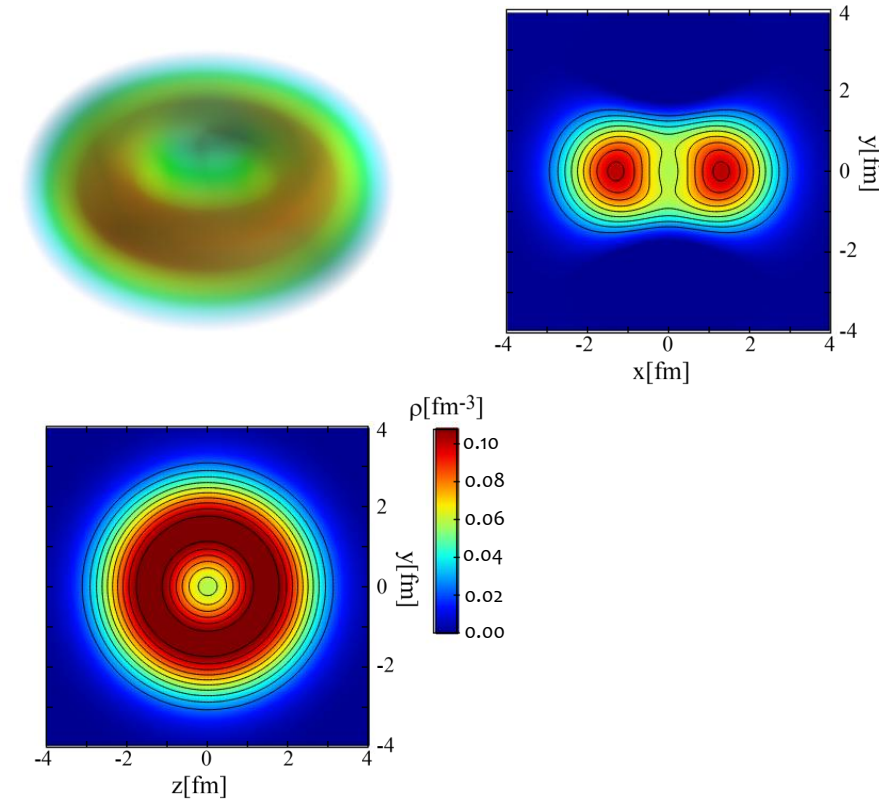
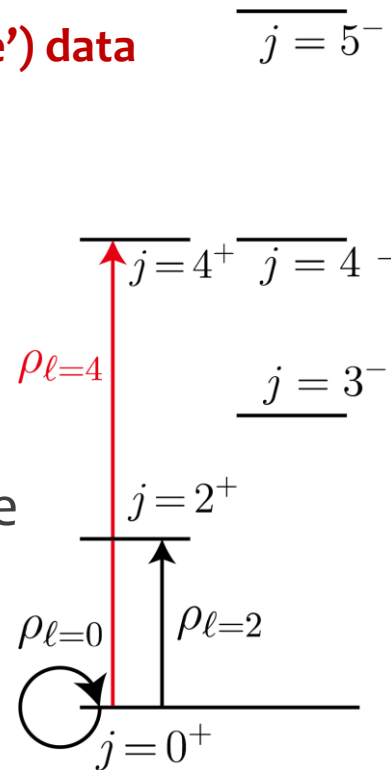
We rebuild the rigid density of ^{12}C by superposing the “observed” transition densities

$$\rho^{\text{rigid}}(\mathbf{r}) = \frac{1}{2} \sum_{lm} \rho_{lm}^{\text{rigid}}(r) \{ Y_{lm}(\hat{r}) + (-)^m Y_{l-m}(\hat{r}) \} \quad \ell = 0^+ + 2^+ + 4^+$$

Shape of ^{12}C

multipole decomposition of intrinsic density from $^{12}\text{C}(e,e')$ data

- With only the $\ell = 0$ density, it is spherical
- The $\ell = 2$ density makes it toroidal shape
- The $\ell = 4$ density emphasizes toroidal shape



Rebuilding shape of ^{12}C

We rebuild the rigid density of ^{12}C by superposing the “observed” transition densities

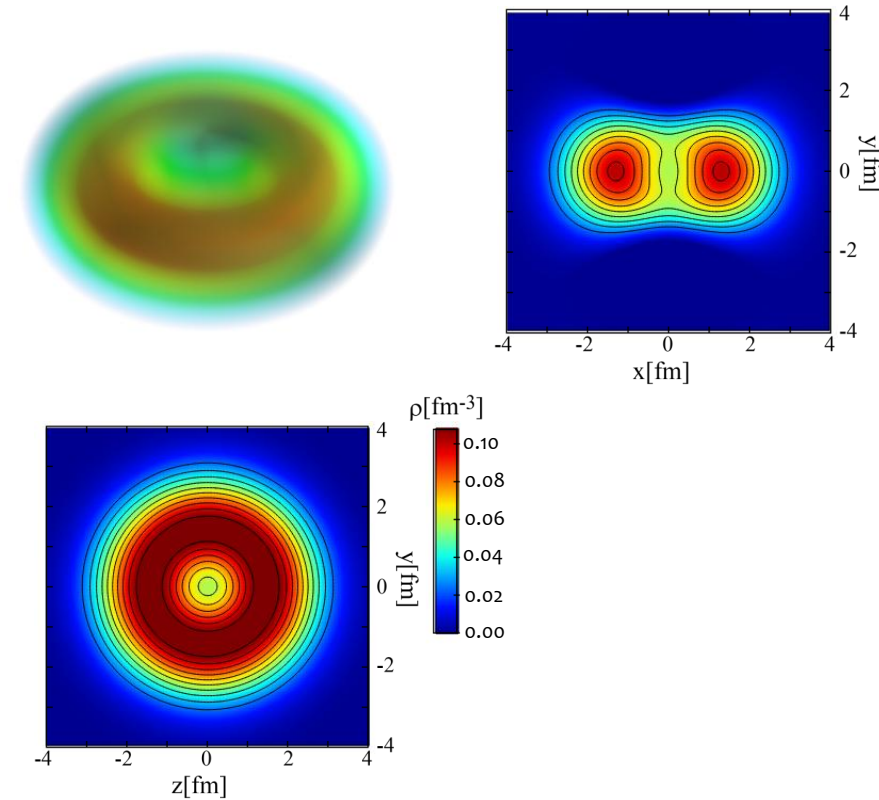
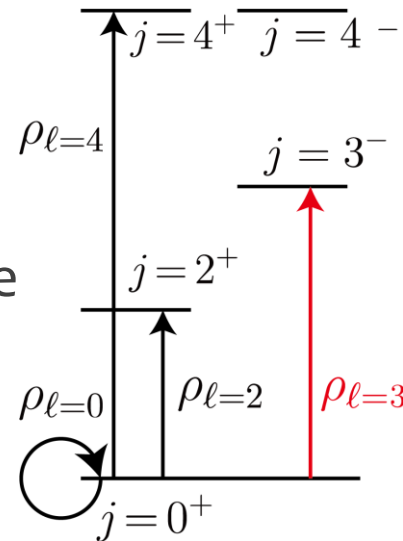
$$\rho^{\text{rigid}}(\mathbf{r}) = \frac{1}{2} \sum_{lm} \rho_{lm}^{\text{rigid}}(r) \{ Y_{lm}(\hat{r}) + (-)^m Y_{l-m}(\hat{r}) \} \quad \ell = 0^+ + 2^+ + 4^+$$

Shape of ^{12}C

**multipole decomposition of
intrinsic density from $^{12}\text{C}(e,e')$ data**

$j = 5^-$

- With only the $\ell = 0$ density, it is spherical
- The $\ell = 2$ density makes it toroidal shape
- The $\ell = 4$ density emphasizes toroidal shape



Rebuilding shape of ^{12}C

We rebuild the rigid density of ^{12}C by superposing the “observed” transition densities

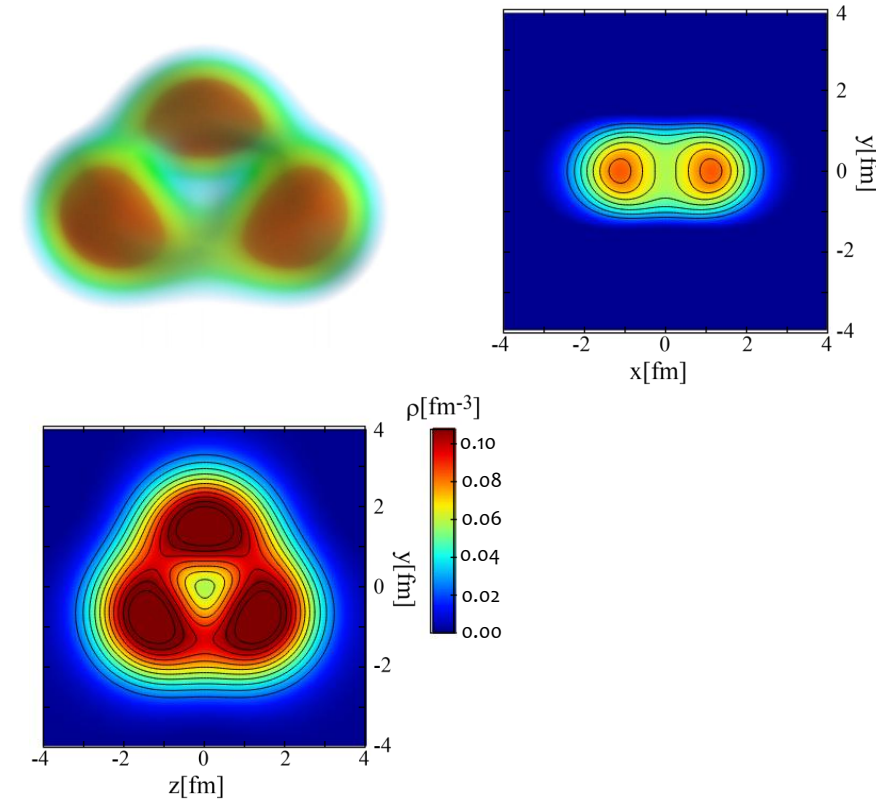
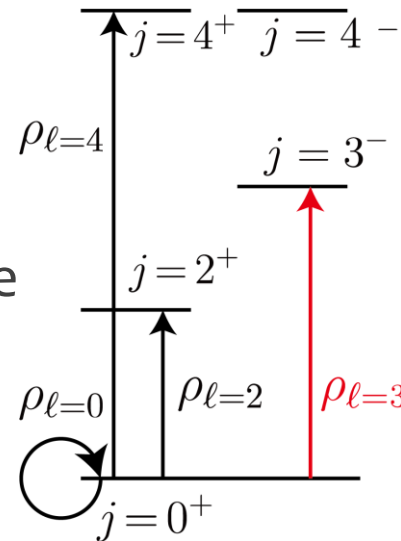
$$\rho^{\text{rigid}}(\mathbf{r}) = \frac{1}{2} \sum_{lm} \rho_{lm}^{\text{rigid}}(r) \{ Y_{lm}(\hat{r}) + (-)^m Y_{l-m}(\hat{r}) \} \quad \ell = 0^+ + 2^+ + 4^+ + 3^-$$

Shape of ^{12}C

multipole decomposition of intrinsic density from $^{12}\text{C}(e,e')$ data

$j = 5^-$

- With only the $\ell = 0$ density, it is spherical
- The $\ell = 2$ density makes it toroidal shape
- The $\ell = 4$ density emphasizes toroidal shape
- The $\ell = 3$ density makes it very exotic.
 - Emphasized triangular shape
 - Three prominent peaks of density



Summary: shape of ^{12}C

Assumption: The 0^+ , 2^+ , 3^- and 4^+ states constitute “the ground band” sharing the same intrinsic state



An exotic triangular shape of ^{12}C emerges

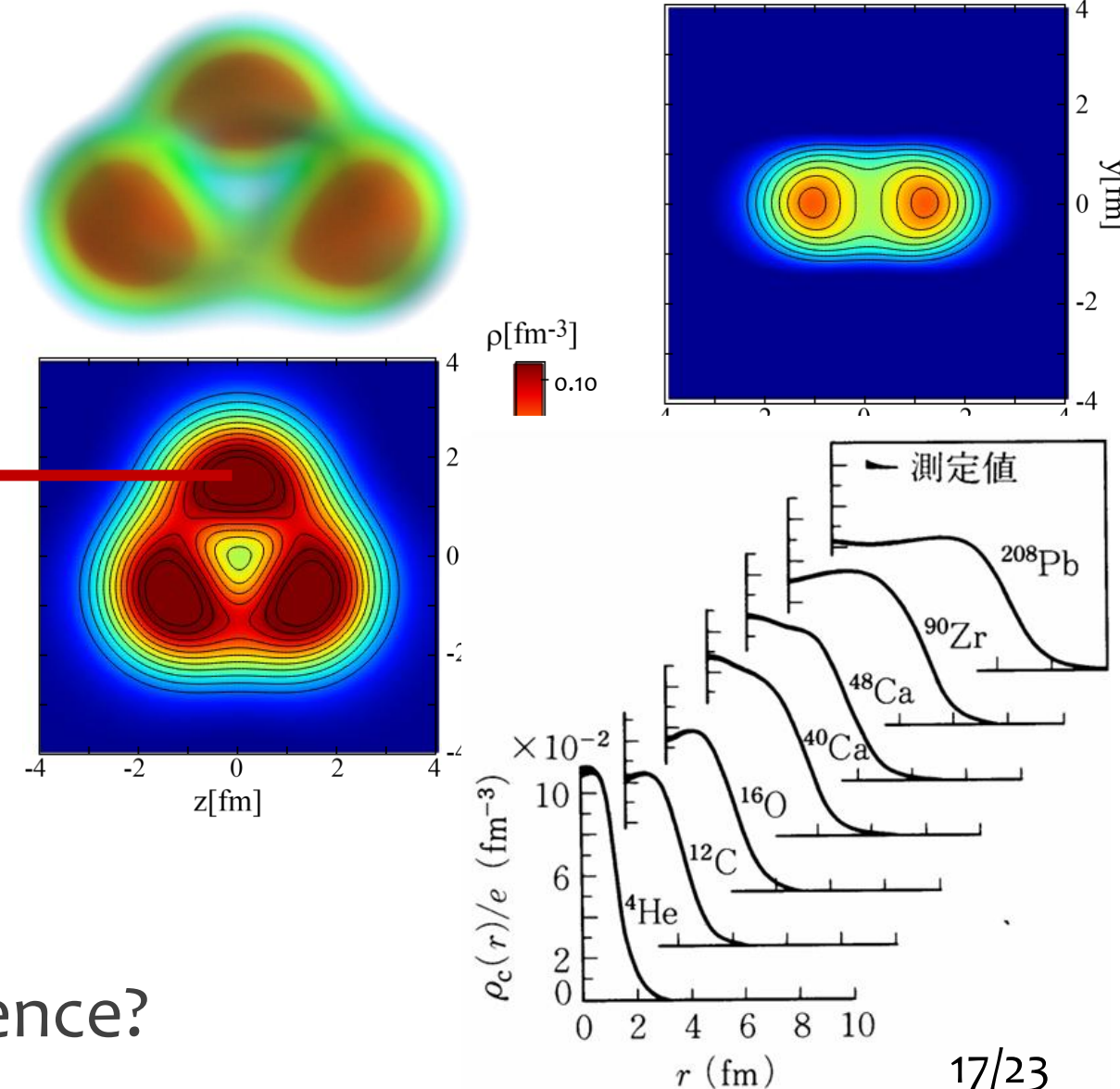
Is this real?

Density at the vertex is as high as 0.11 fm^{-3} which is much larger than typical density ($\sim 0.08 \text{ fm}^{-3}$).

Among all nuclei, there is only one nucleus which has such high central density.

The central density of alpha particle is 0.11 fm^{-3} !!!

Question: Is this an accidental coincidence?



How can we confirm the shape of ^{12}C ?

A novel experimental probe: α knockout reaction

REPORT

J. Tanaka, Z. Yang et al., Science 2021

NUCLEAR PHYSICS

Formation of α clusters in dilute neutron-rich matter

Junki Tanaka^{1,2,3*}, Zaihong Yang^{3,4*}, Stefan Typel^{1,2}, Satoshi Adachi⁴, Shiwei Bai⁵, Patrik van Beek¹,

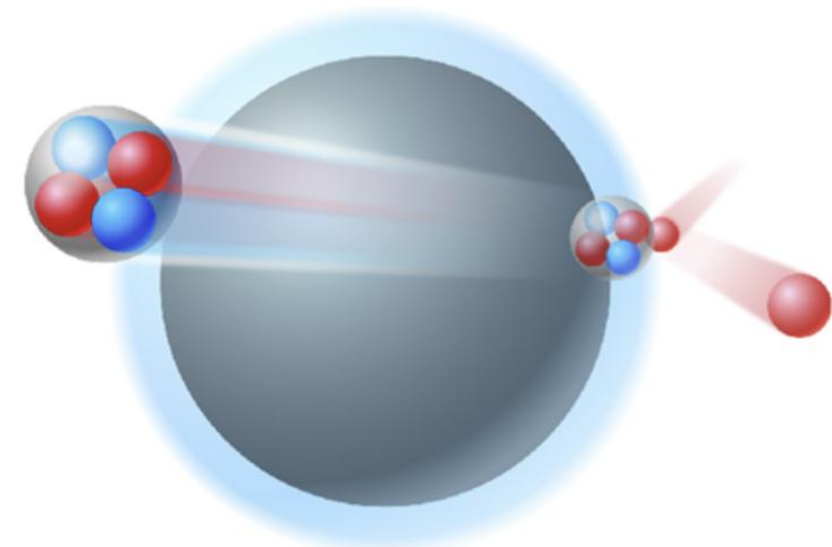
Theoretical works

$^{20}\text{Ne}(p,p\alpha)$: K. Yoshida et al., PRC 100, 044601 (2019)

$^{48}\text{Ti}(p,p\alpha)$: Y. Taniguchi et al., PRC 103, L031305 (2021).

$^A\text{Be}(p,p\alpha)$, $^A\text{B}(p,p\alpha)$, : H.Motoki et al., PTEP2022, 113D01 (2022)

Q. Zhao et al., PRC 106, 054313 (2022)



(p,p α) reaction as a probe for α particle preformed in nuclei

High energy (p,p α) reaction is a clean probe for the alpha-particle preformation on the surface of nuclei

DWIA for (p,p α) reaction

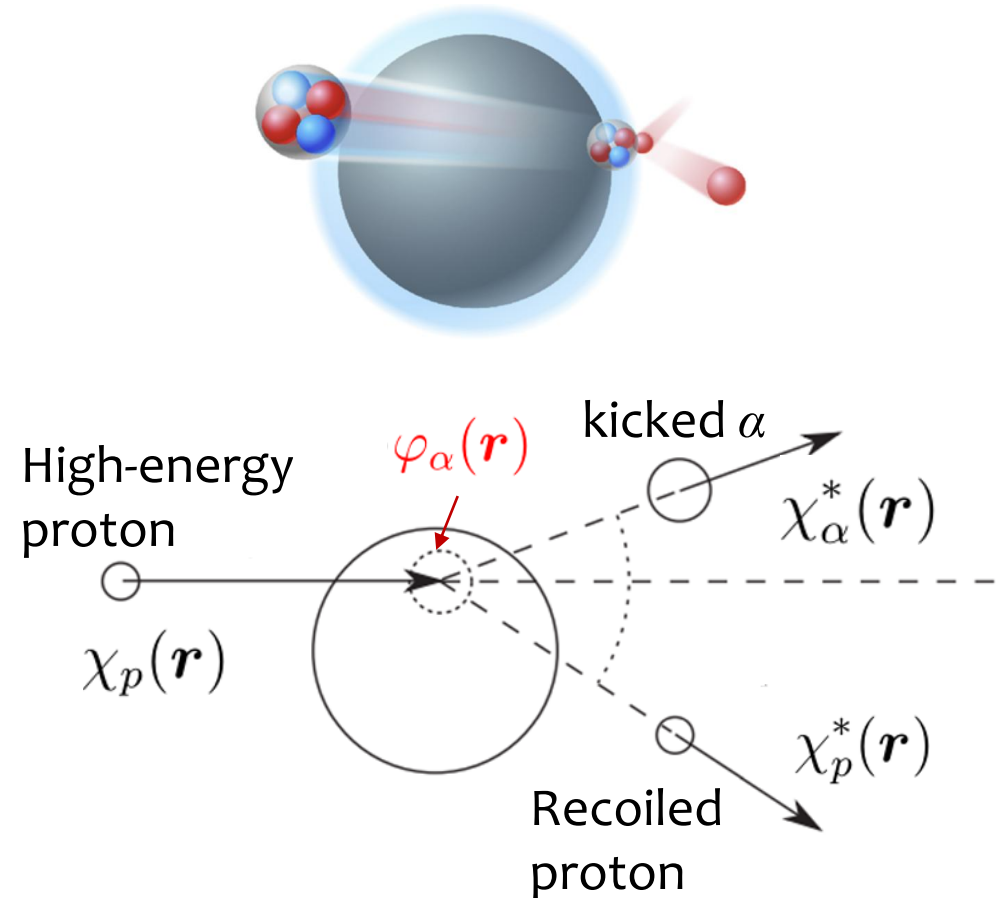
K. Yoshida and K. Ogata PRC 94, 044604 (2016)

$$\frac{d^3\sigma}{dt_p d\Omega_p d\Omega_\alpha} = F_{\text{kin}} \frac{d\sigma_{p\alpha}}{d\Omega_{p\alpha}} |T|^2$$

Transition matrix

$$T = \int \chi_p^*(\mathbf{r}) \chi_\alpha^*(\mathbf{r}) \chi_p(\mathbf{r}) \varphi_\alpha(\mathbf{r}) e^{-ik_0 \cdot \mathbf{r}}$$

α particle wave function inside nucleus



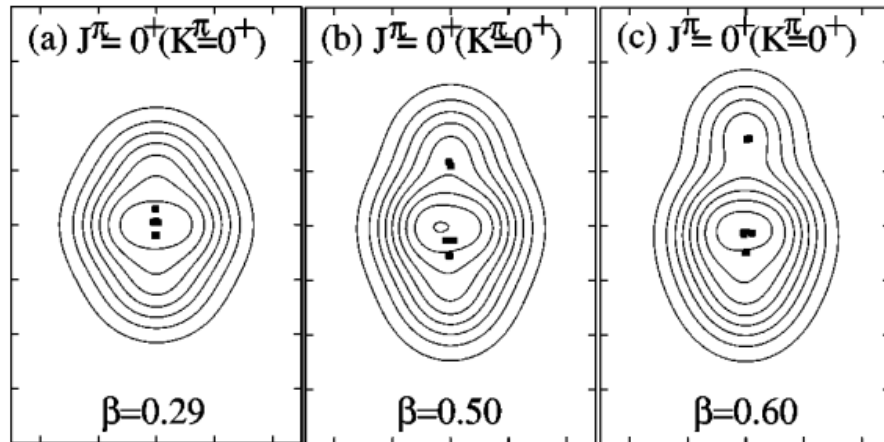
A benchmark calculation for $^{20}\text{Ne} (\alpha + ^{16}\text{O})$

Alpha particle wave function preformed in the target nucleus is defined as the overlap of target nucleus and decay channel = probability amplitude to observe alpha particle at distance r

$$\varphi(\mathbf{r}) = \sqrt{AC_4} \langle \delta^3(\mathbf{r}' - \mathbf{r}) \Phi_\alpha \Phi_{\text{residue}} | \Phi_{\text{target}} \rangle = \left\langle \begin{array}{c} \text{---} \left(\text{---} \left(\text{---} \right) \text{---} \right) \text{---} \\ \alpha \quad \text{residue} \quad \text{target} \end{array} \right\rangle$$

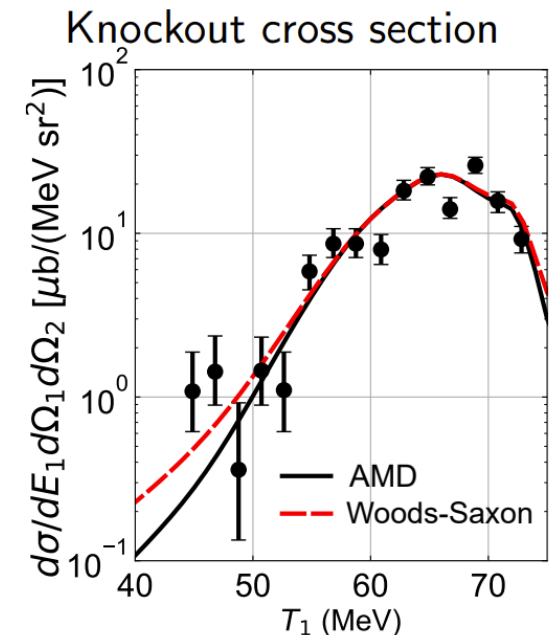
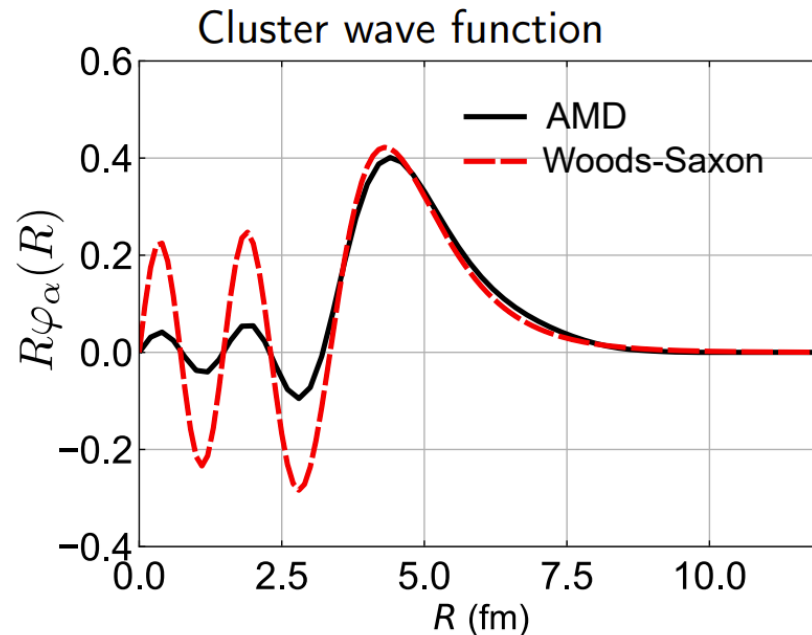
We have a reliable wave function for ^{20}Ne

- Spectrum of ^{20}Ne
- α decay width of the excited states
- $^{16}\text{O}(^6\text{Li},d)^{20}\text{Ne}$ reaction



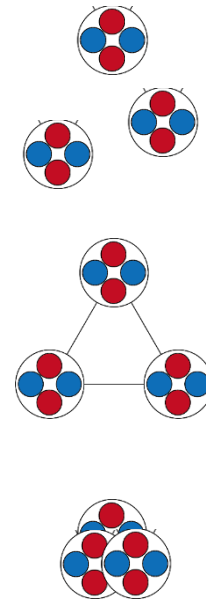
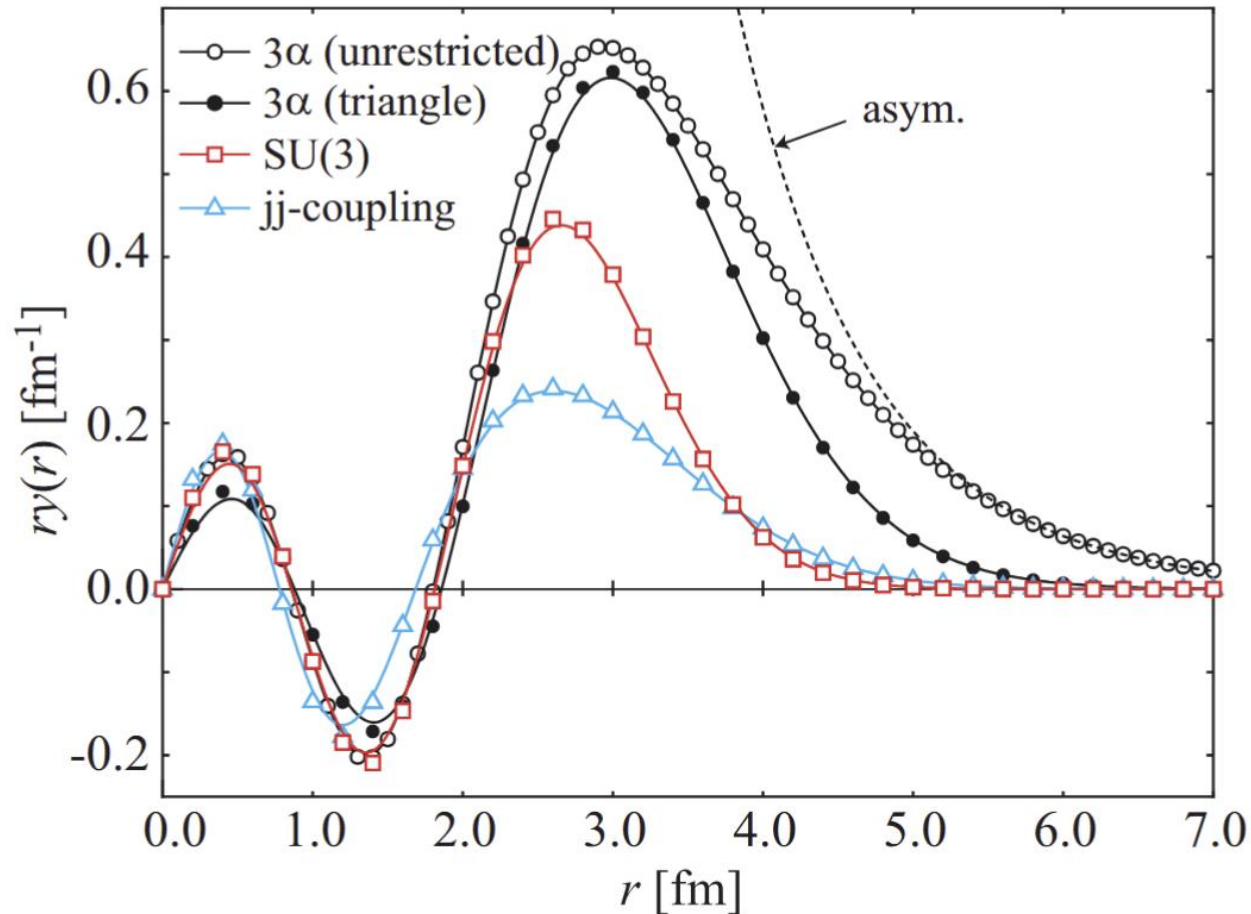
M. Kimura, PRC69, 044319 (2004)

$^{20}\text{Ne}(p,p\alpha)$: K. Yoshida et al., PRC 100, 044601 (2019)



Numerical experiments for ^{12}C

Numerical experiments.



3 α cluster model (RGM, GCM, THSR)
 3 α system solved
 without any symmetry assumption

Fixed 3 α triangle (\approx Algebraic cluster model)

A simple shell model (SU3 limit)
 (os)⁴(op)⁸ Harmonic oscillator
 jj-coupling

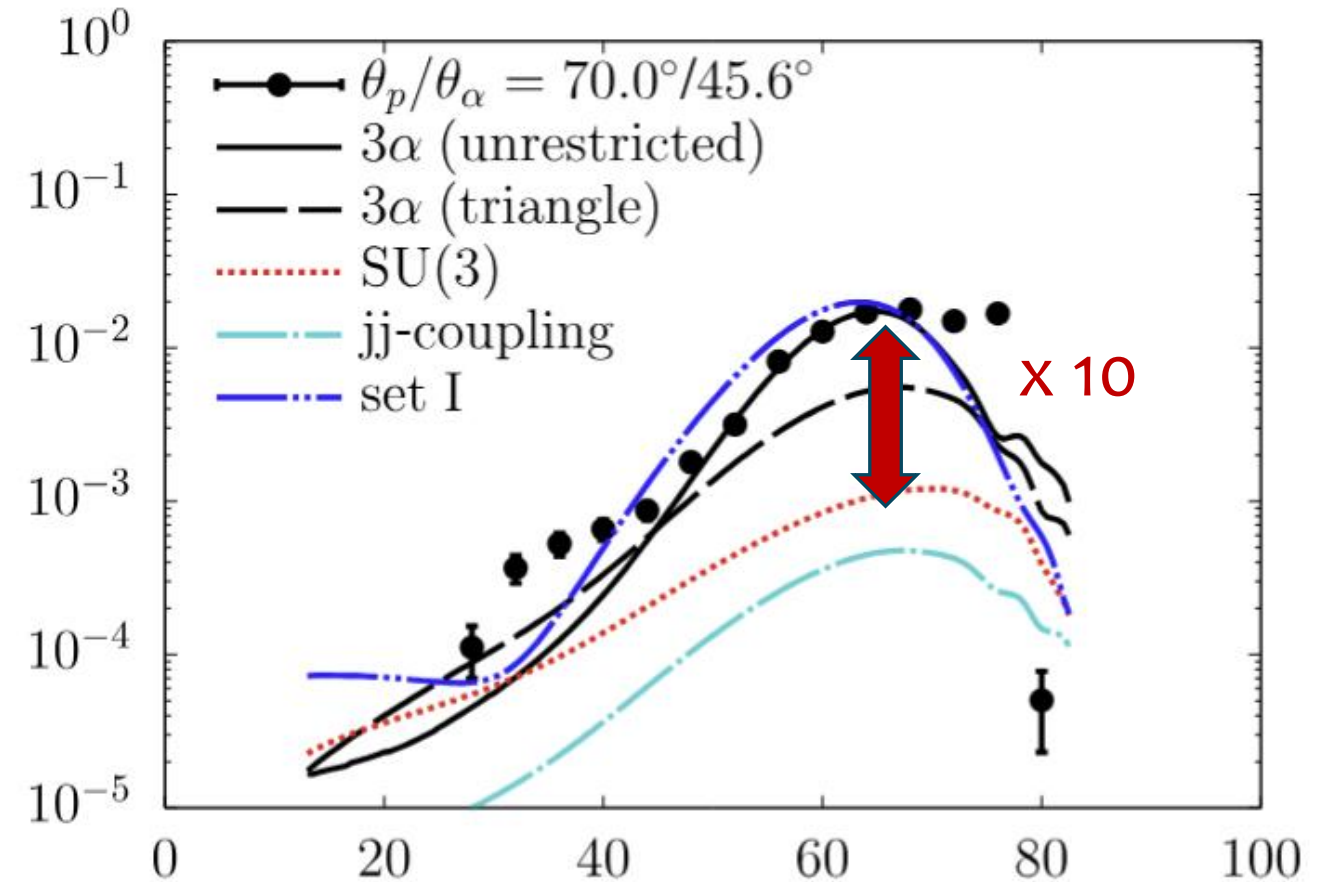
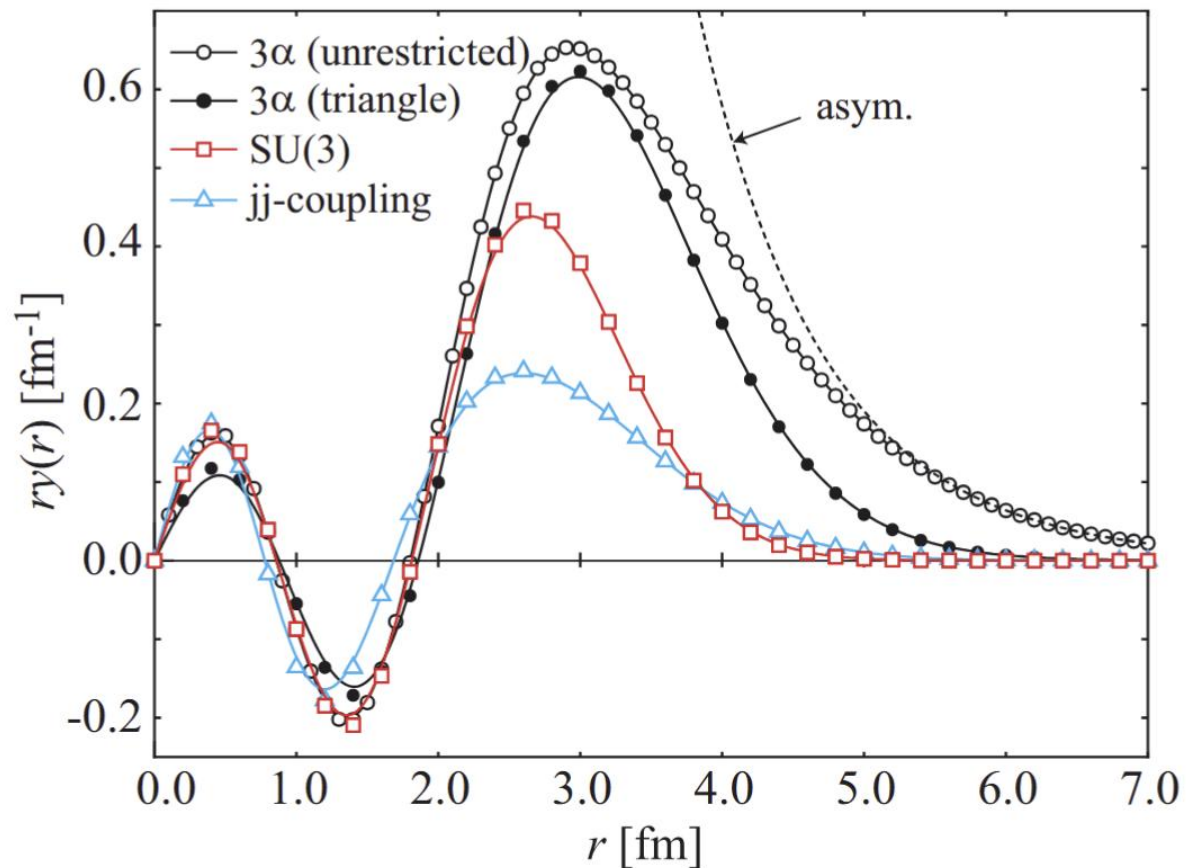
NOTE

All of them reproduce the observed radius of ^{12}C ,
 but yields different alpha particle preformation

^{12}C seems to be clustered!!

Numerical experiments.

The ordinary cluster model explains the data.
Fixed triangle somewhat undershoot.



Exp: J. Mabilia et al., PRC79, 054612 (2009)

Summary

(1) Shape of ^{12}C constructed from electron scattering data

Assumption: The 0^+ , 2^+ , 3^- and 4^+ states constitute “the ground band” sharing the same intrinsic state

⇒ An exotic triangular shape of ^{12}C emerges

(2) Alpha knockout from the ground state (ongoing)

Comparison of three different models

3α cluster model, Fixed 3α triangle, A simple shell model

⇒ Only 3α cluster model can explain the observation

⇒ More detailed analysis is ongoing for
 $^{12}\text{C}(p,p\alpha)^8\text{Be}(2^+)$, $^{16}\text{O}(p,p\alpha)^{12}\text{C}(0^+)$, etc...