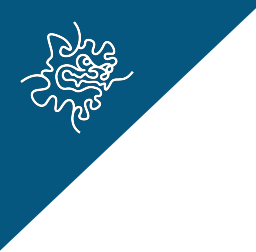




Simone Pigolotti (OIST)

Pattern formation and the physics of growing surfaces

YITP workshop “Frontiers in Nonequilibrium Physics”, 14/5/26



Physics of development

Egg



From larva to adult



Stade larvaire



De la métamorphose jusqu'au stade juvénile



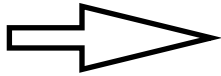
Stade adulte

© P. Salis, JE Randal



Physics of development

Egg



From larva to adult



Stade larvaire



De la métamorphose jusqu'au stade juvénile



Stade adulte

© P. Salis, JE Randal

- Can growth unlock new mechanisms for pattern formation?
- How are patterns maintained during growth?

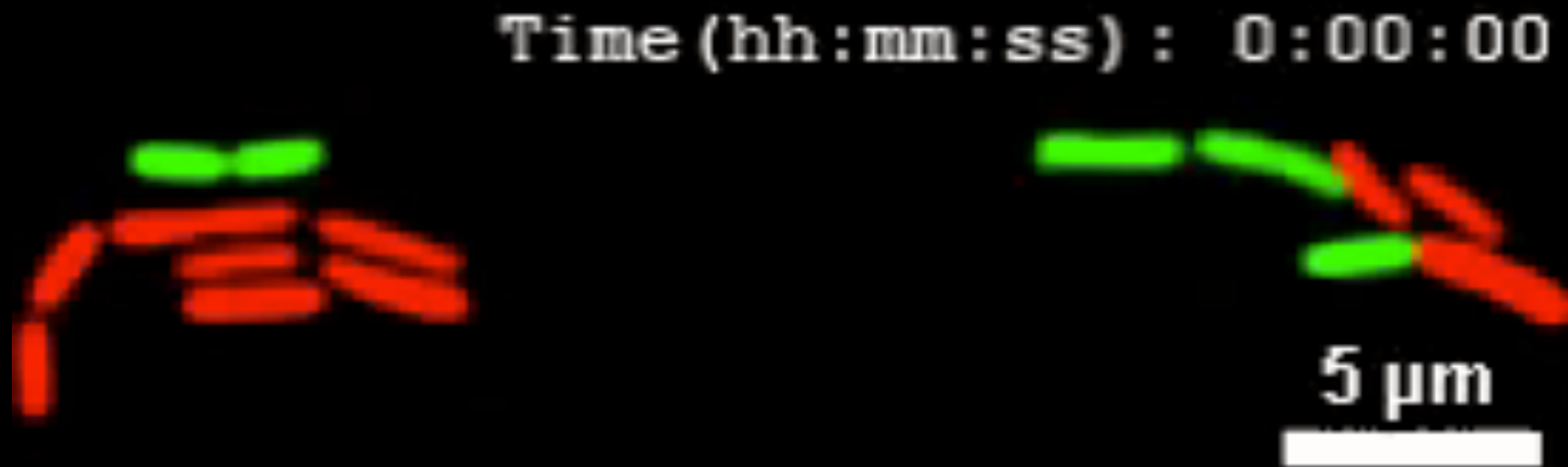


Outline

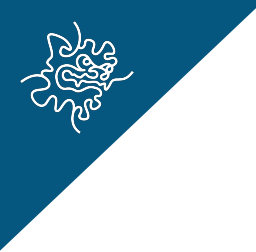
- Pattern formation in microchannels
- Disordered packing on growing surfaces
- Clownfish patterns as an interface problem



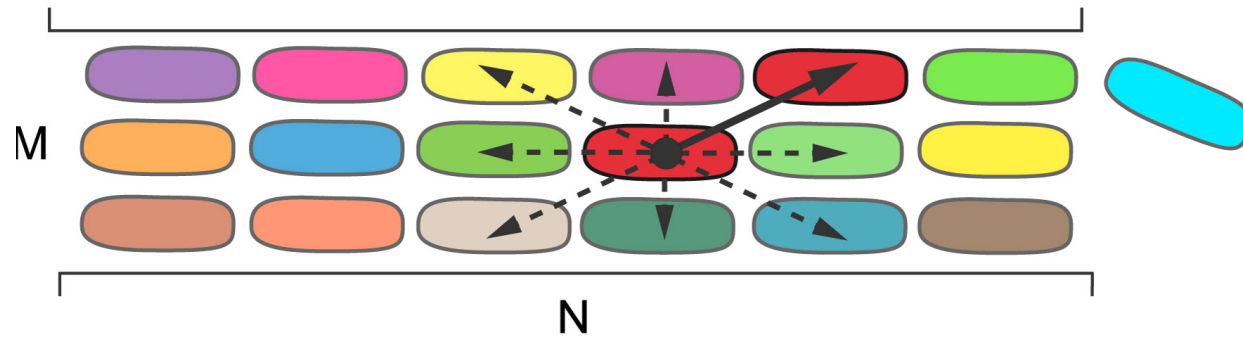
Competition in microchannels



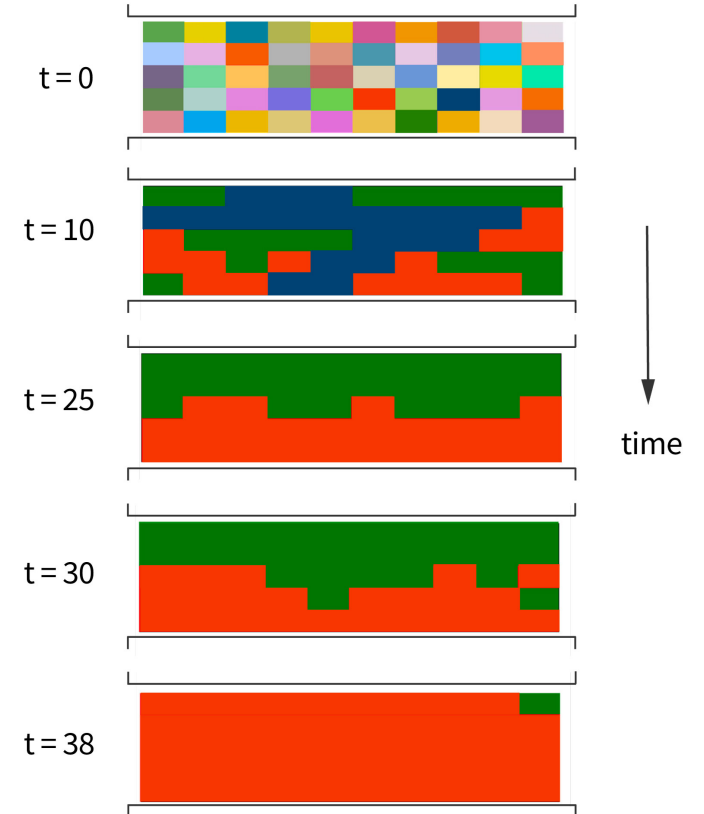
A. Koldaeva, P. Tsai, A. Shen, SP, Proc. Natl. Acad. Sci. (2022)



Model



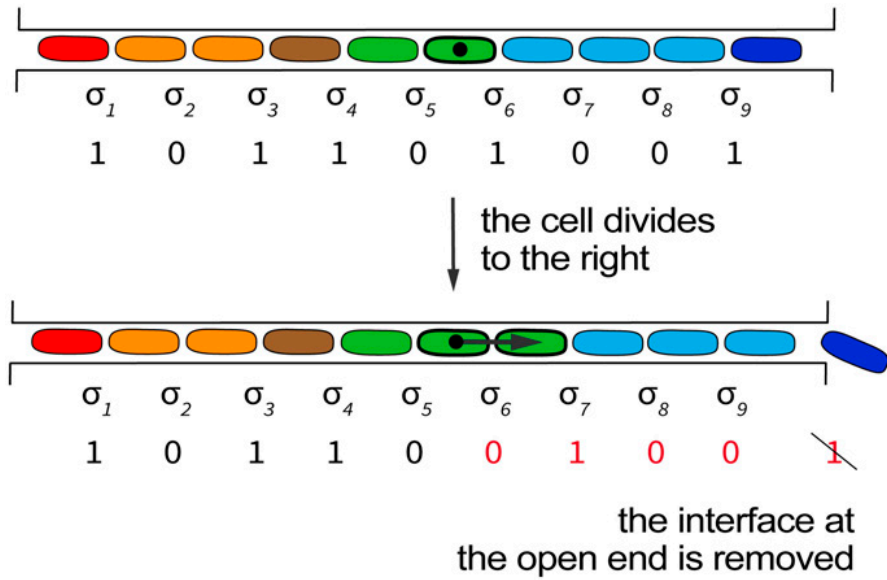
- Cells reproduce into neighboring sites with non-uniform rates.
- Each reproduction event pushes an entire lane of other cells, leading to an expulsion of another cell



A. Koldaeva, P. Tsai, A. Shen, SP, Proc. Natl. Acad. Sci. (2022)



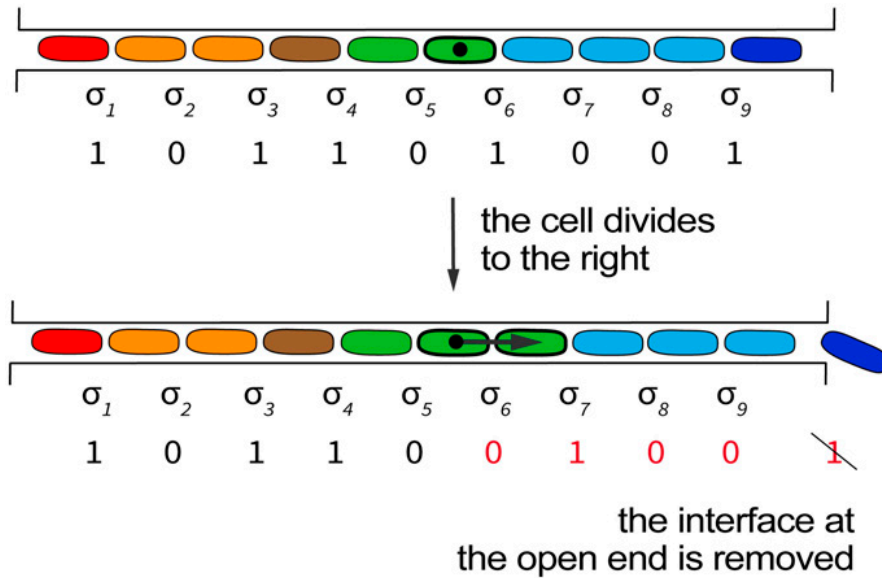
Interface density



A. Koldaeva, P. Tsai, A. Shen, SP, PNAS (2022)



Interface density



$$\hat{b}_r^i \vec{\sigma} = (\sigma_1, \sigma_2, \dots, \sigma_{i-1}, \sigma_{i+1}, \dots, \sigma_{N-1}, 0),$$

$$\hat{b}_l^i \vec{\sigma} = (0, \sigma_1, \sigma_2, \dots, \sigma_{i-1}, \sigma_{i+1}, \dots, \sigma_{N-1}),$$

$$\hat{c}_r^i \vec{\sigma} = (\sigma_1, \sigma_2, \dots, \sigma_{i-1}, \sigma_{i+1}, \dots, \sigma_{N-1}, 1),$$

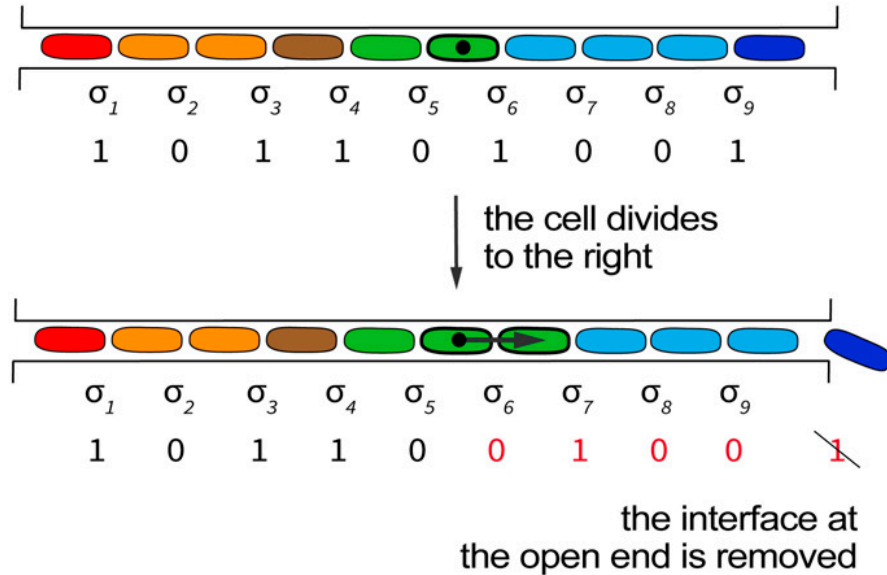
$$\hat{c}_l^i \vec{\sigma} = (1, \sigma_1, \sigma_2, \dots, \sigma_{i-1}, \sigma_{i+1}, \dots, \sigma_{N-1}).$$

Master equation for the interfaces:

$$\frac{dP_{\vec{\sigma}}}{dt} = -b \sum_{i=1}^{N-1} [q(i) + p(i+1)] P_{\vec{\sigma}}(t) + b \sum_{i=1}^{N-1} \delta_{\sigma_i, 0} [q(i) (P_{\hat{b}_r^i \vec{\sigma}} + P_{\hat{c}_r^i \vec{\sigma}}) + p(i+1) (P_{\hat{b}_l^i \vec{\sigma}} + P_{\hat{c}_l^i \vec{\sigma}})],$$



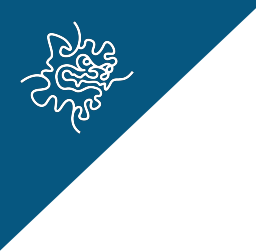
Interface density



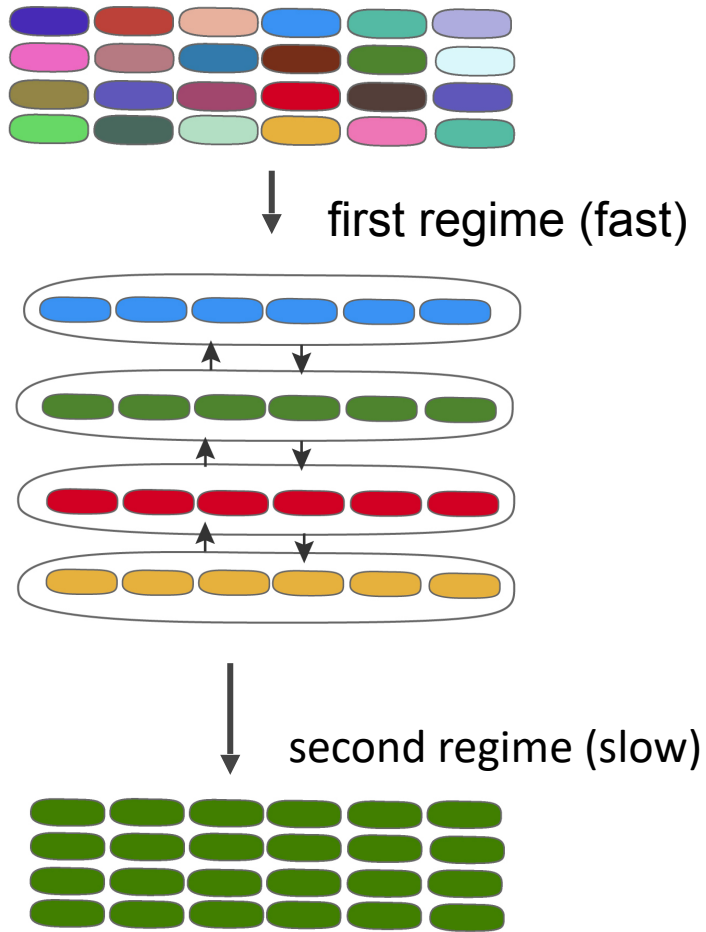
Results:

- $\{\sigma_i\}$ are i.i.d. with: $P(\sigma_i(t) = 1) = e^{-bt}$
- Diversity loss: $\langle A(t) \rangle = (N - 1)e^{-bt} + 1$
- Fixation time: $T_{N \rightarrow 1} = b^{-1} \sum_{A=2}^N \frac{1}{A-1} \approx b^{-1} [\log(N-1) + \gamma]$
- Fixation probabilities: approx. Gaussian with $\sigma \approx \sqrt{N}$

A. Koldaeva, P. Tsai, A. Shen, SP, PNAS (2022)



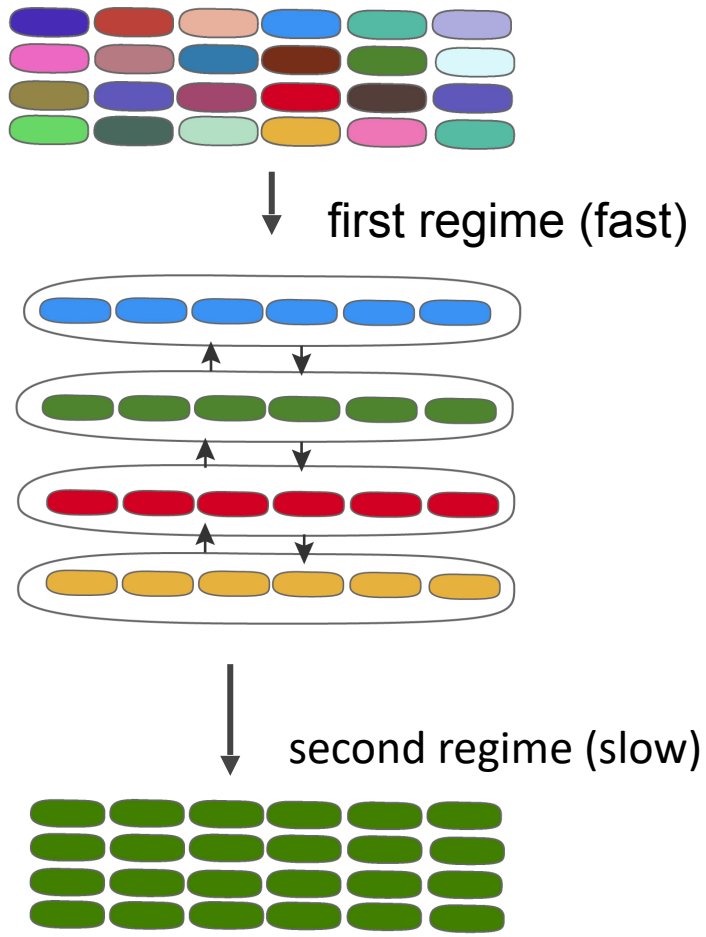
Fixation dynamics



A. Koldaeva, P. Tsai, A. Shen, SP, PNAS (2022)

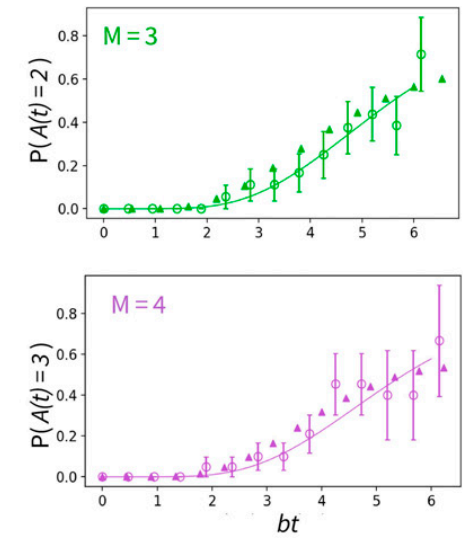
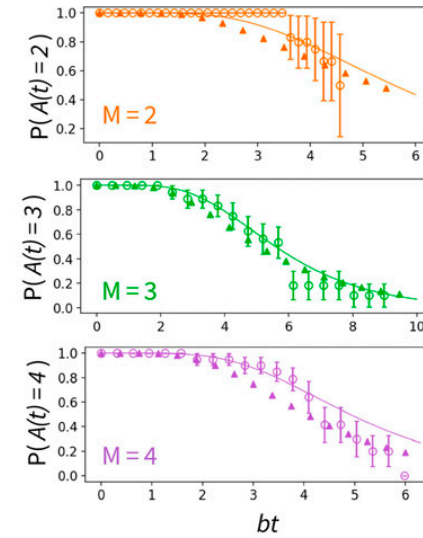
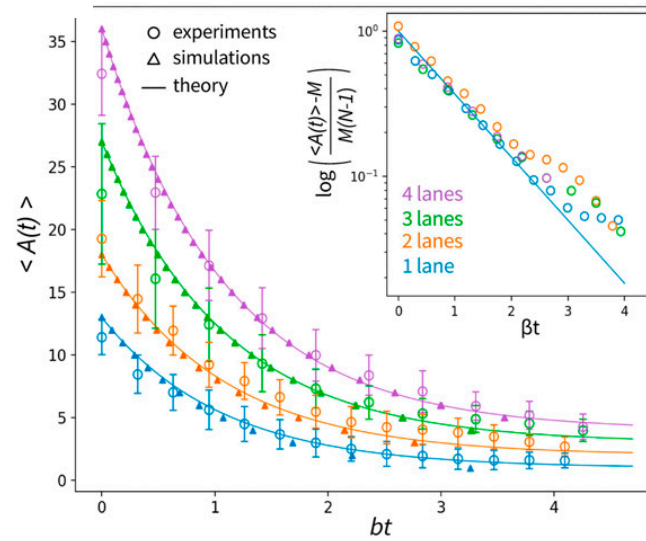


Fixation dynamics



first regime: $A(t) = MN \rightarrow M$

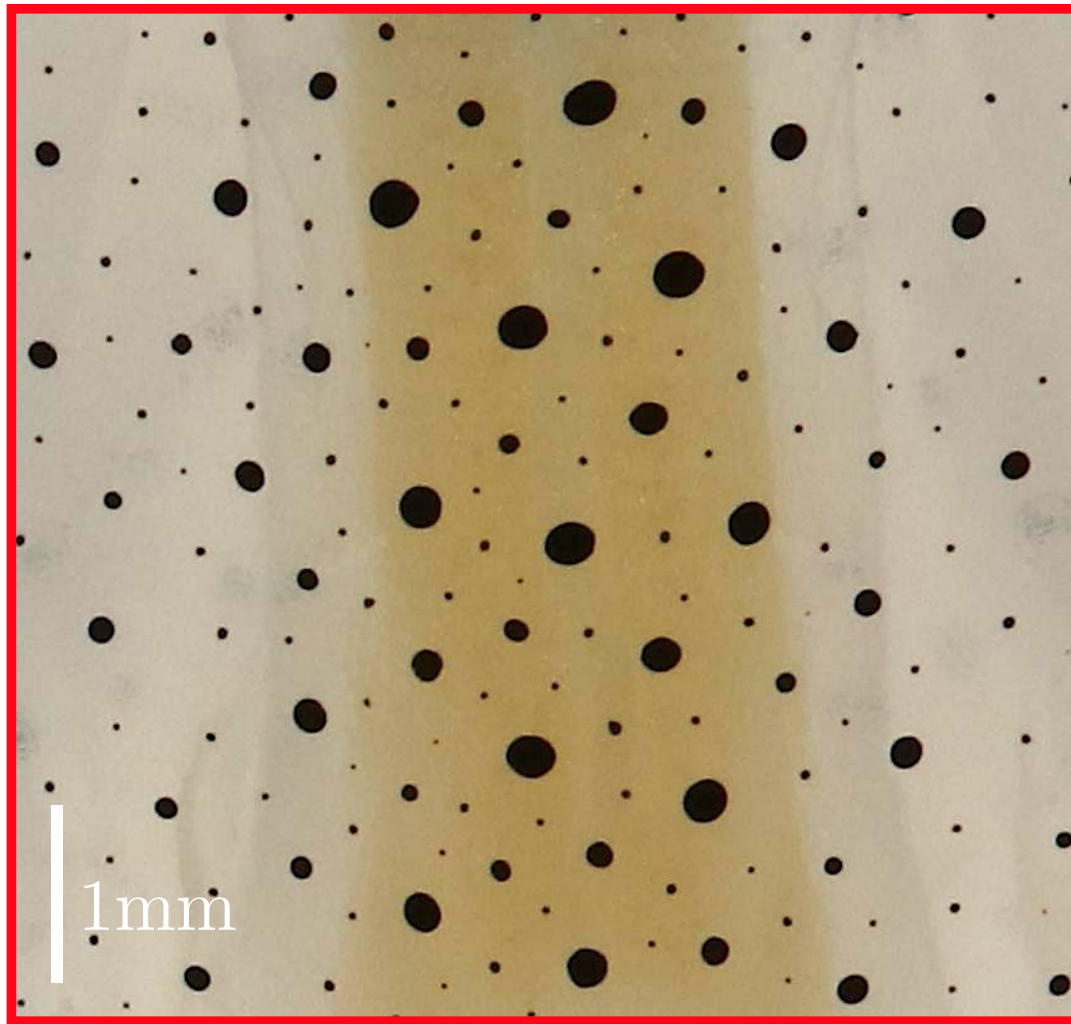
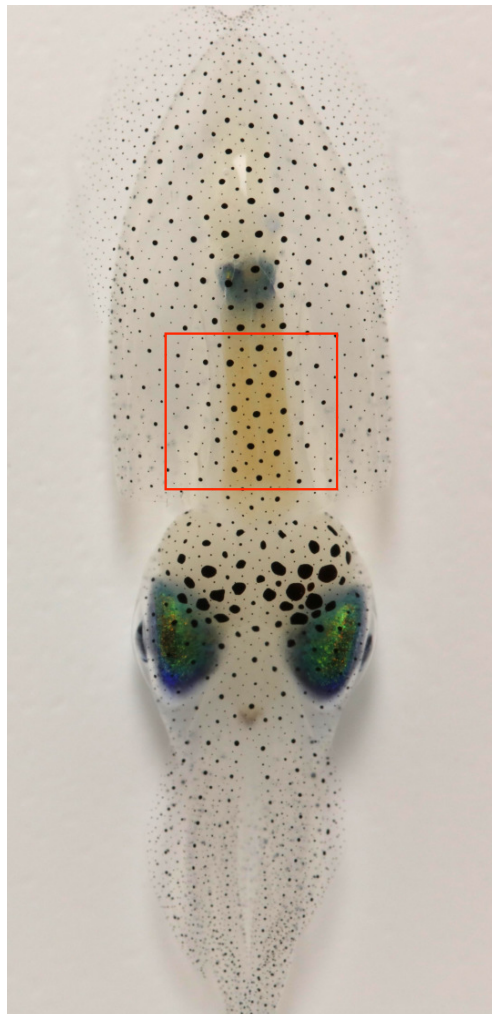
second regime: $A(t) = M \rightarrow 1$



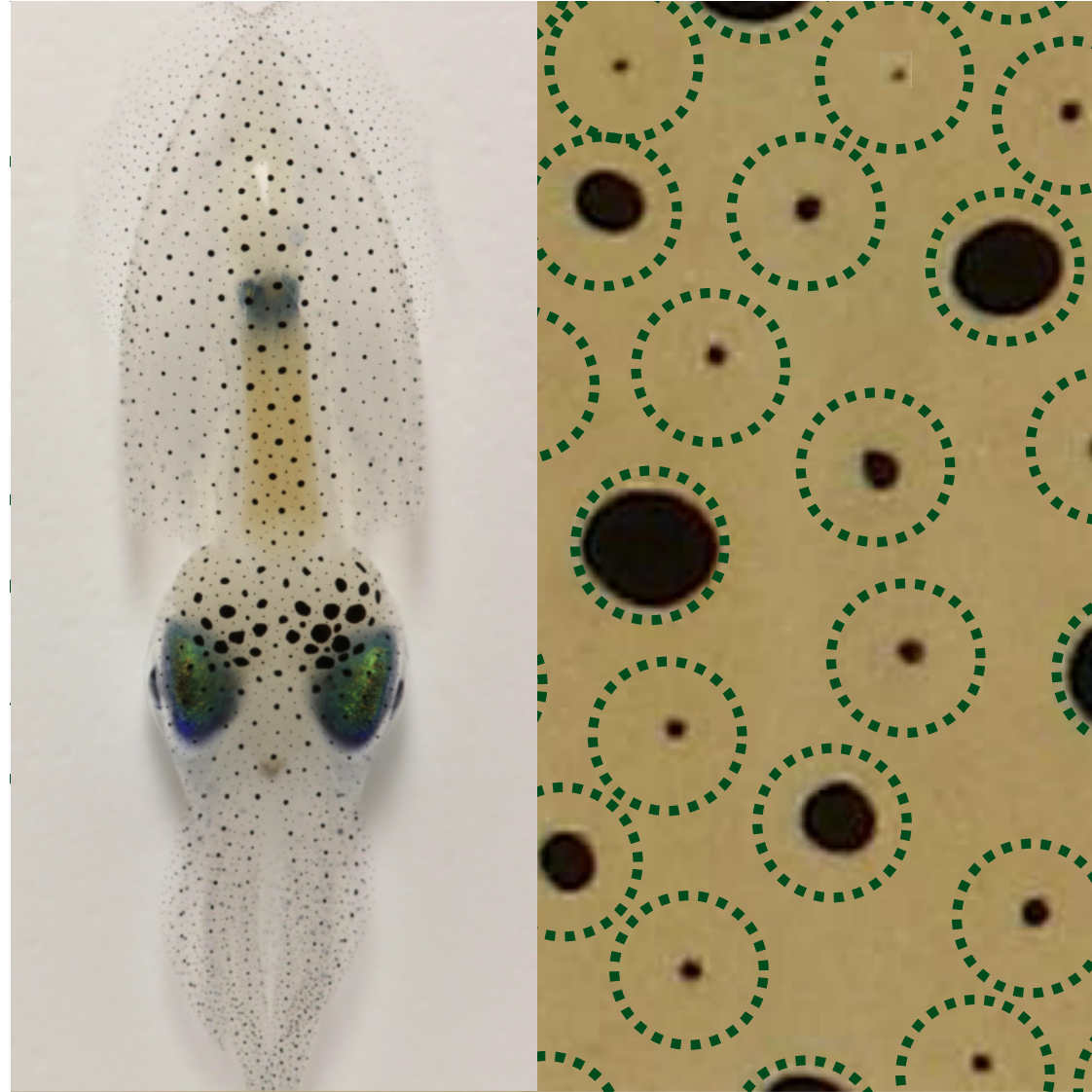
A. Koldaeva, P. Tsai, A. Shen, SP, PNAS (2022)



Ok, what about multicellular organisms?

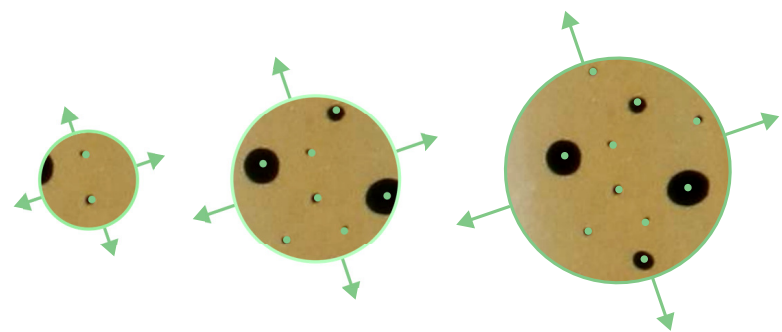


R. Ross SP, Phys. Rev. X (2025)

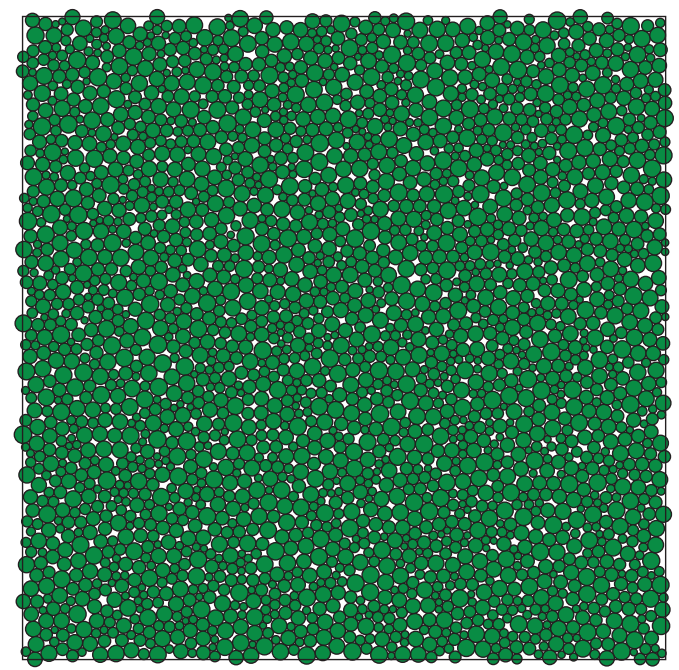
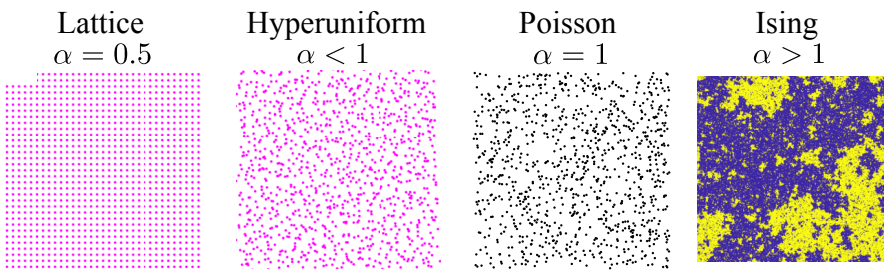




Hyperuniformity?



$$\langle N^2 \rangle - \langle N \rangle^2 \sim \langle N \rangle^\alpha$$



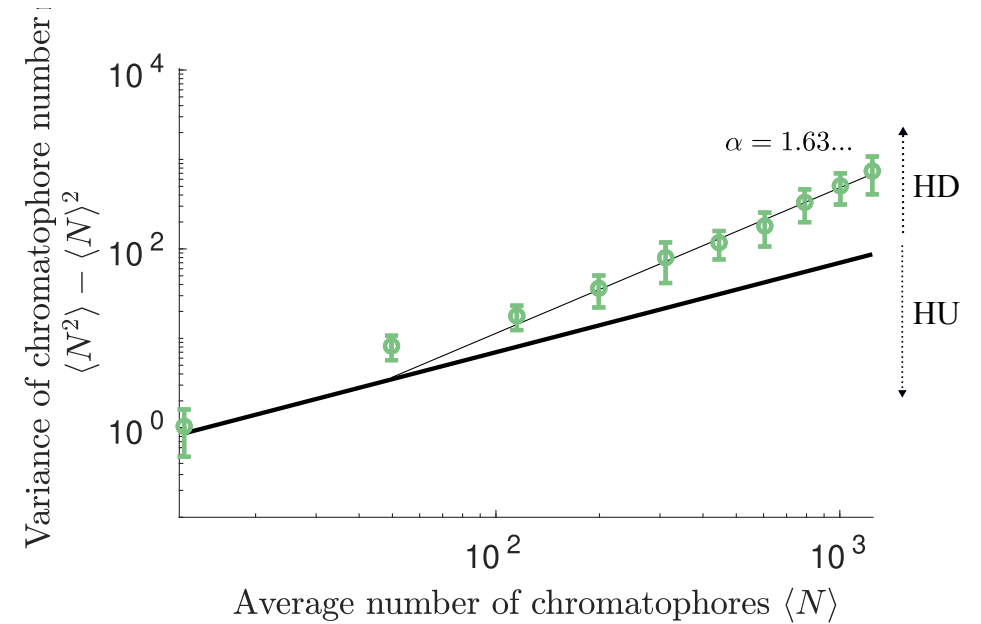
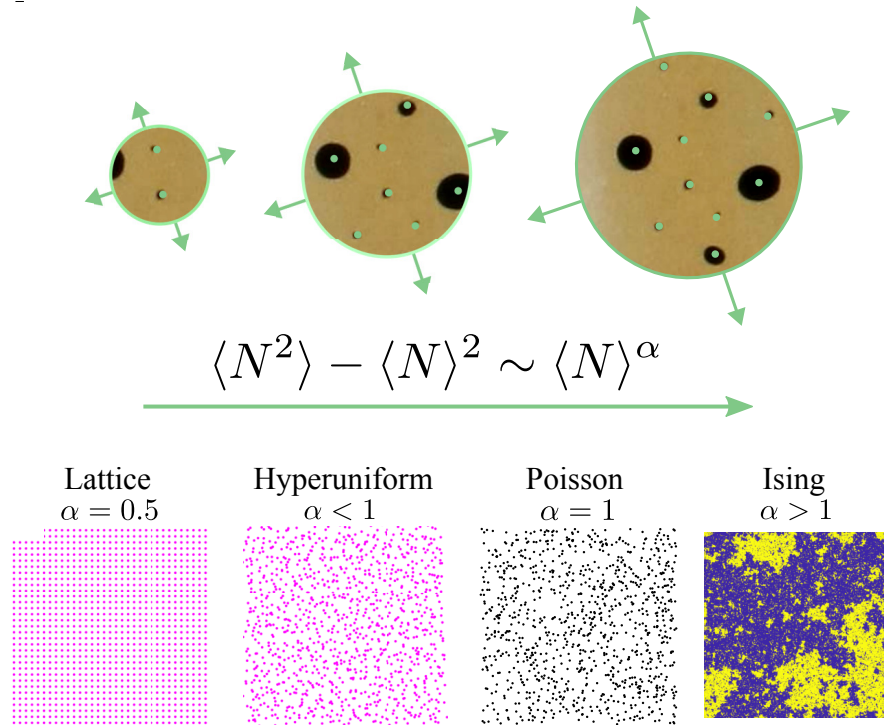
Zachary, Jiao, Torquato, PRE (2011)

Random jammed systems are hyperuniform ($\alpha < 1$)

R. Ross SP, Phys. Rev. X (2025)
S. Torquato, Phys. Rep. (2018)



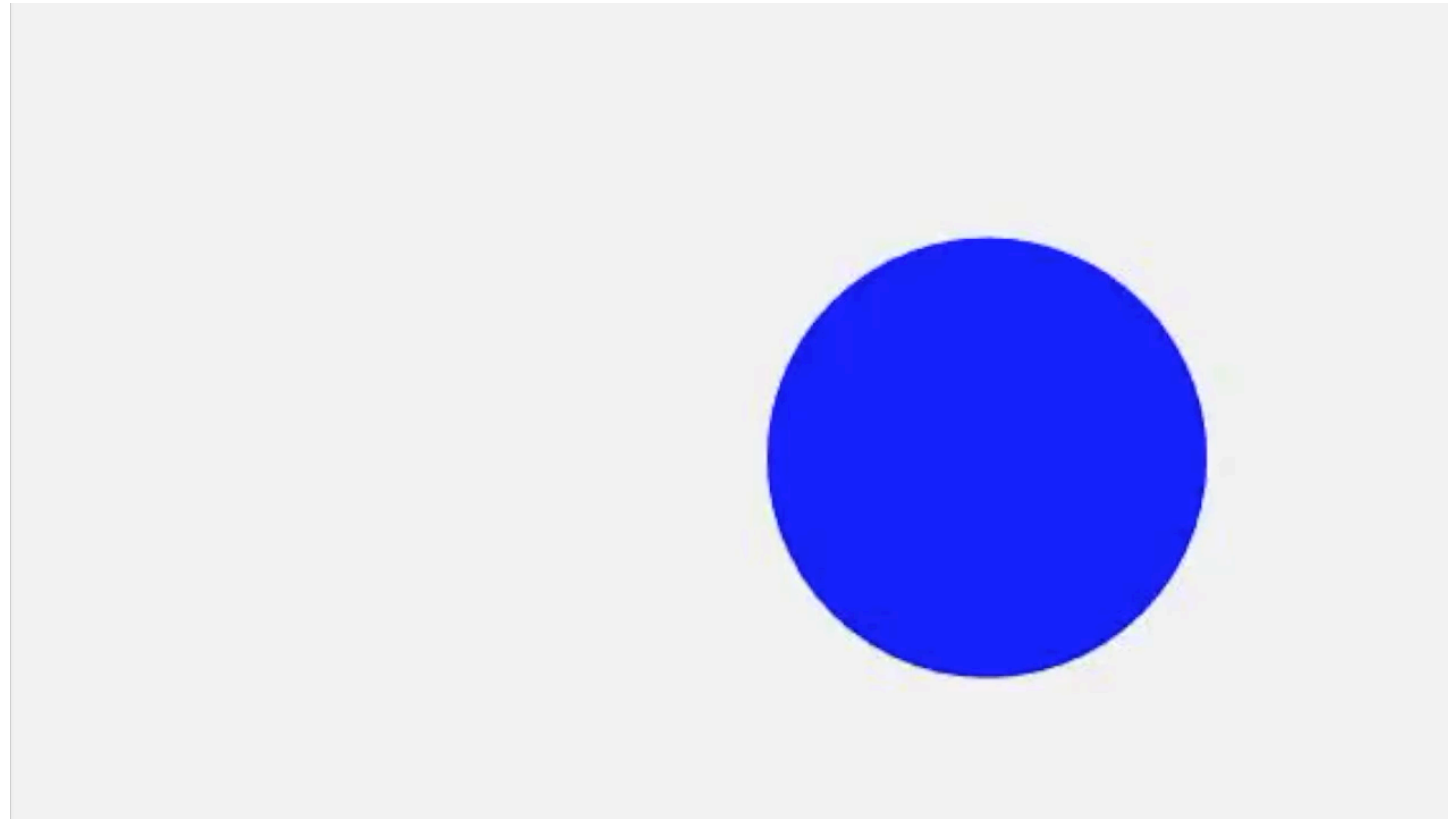
Hyperdisordered behavior



R. Ross SP, Phys. Rev. X (2025)



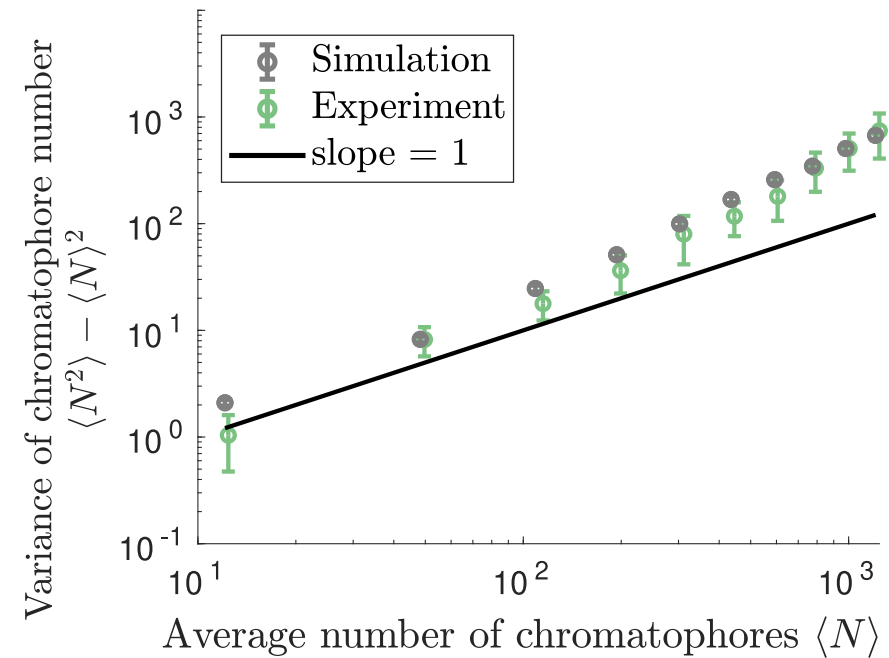
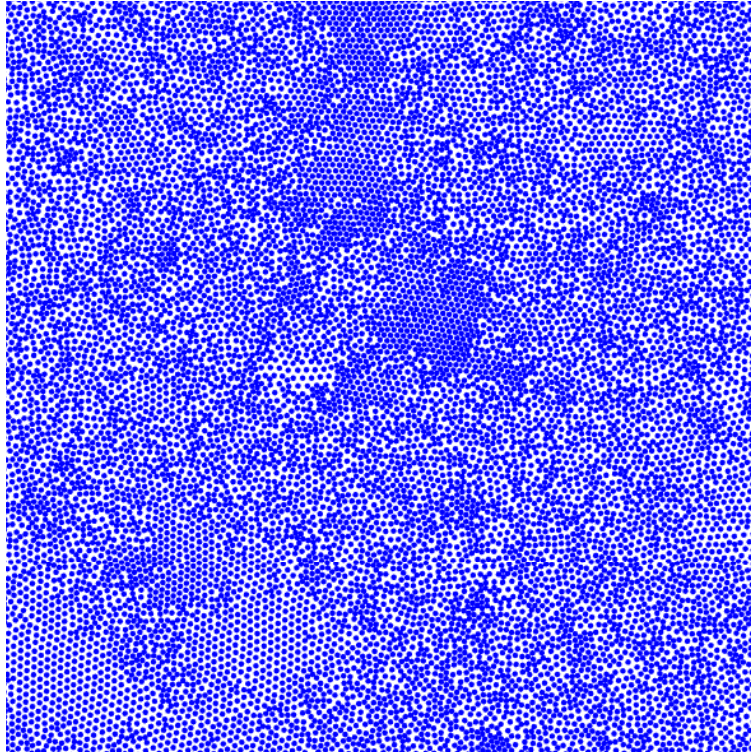
Packed disks on growing domains



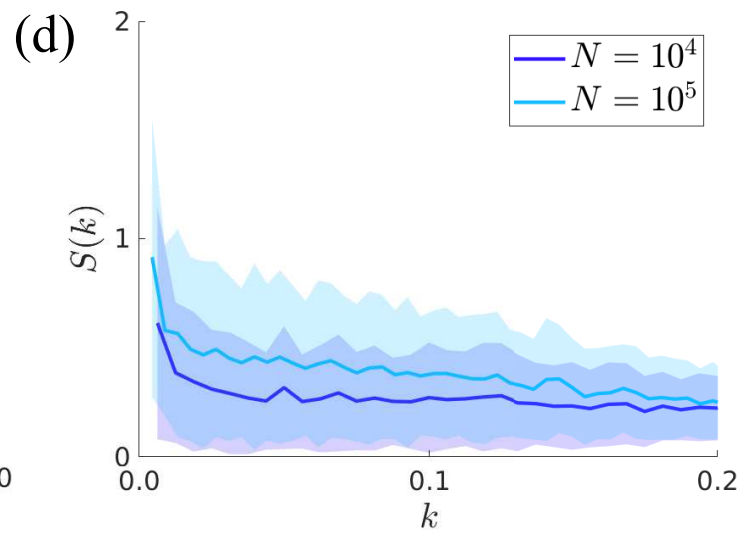
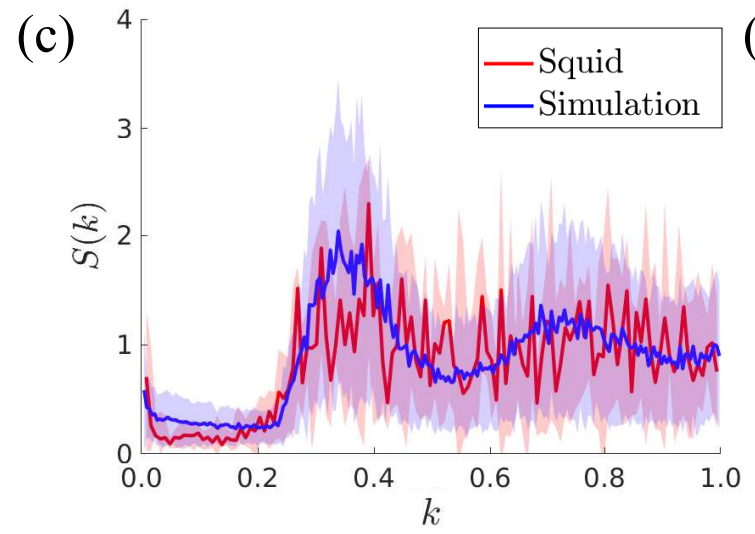
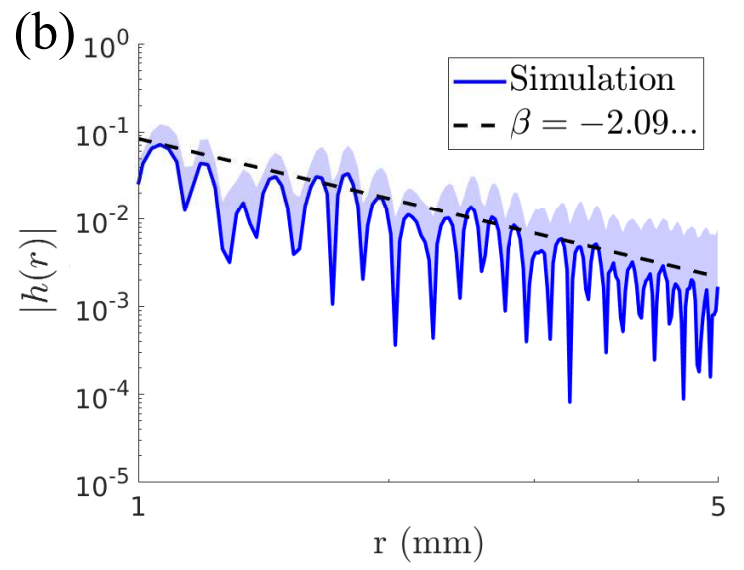
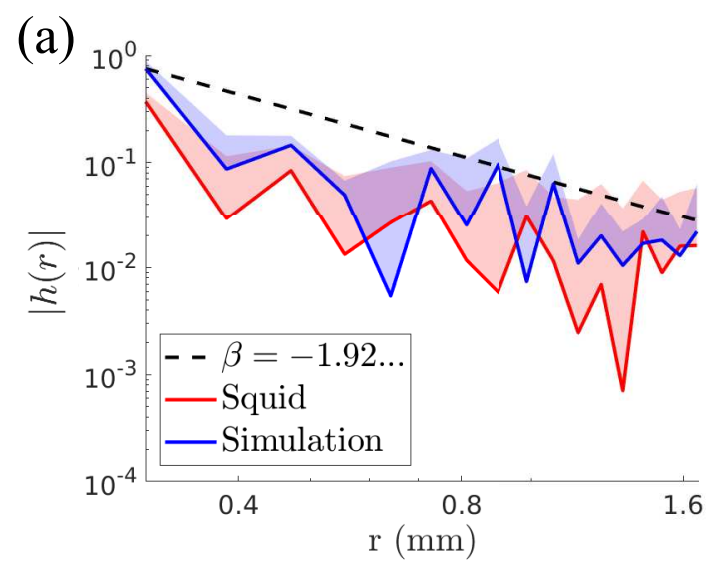
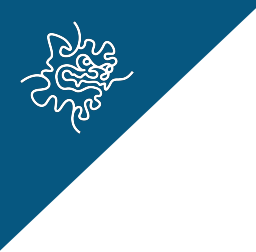
R. Ross SP, Phys. Rev. X (2025)



Packed disks on growing domains



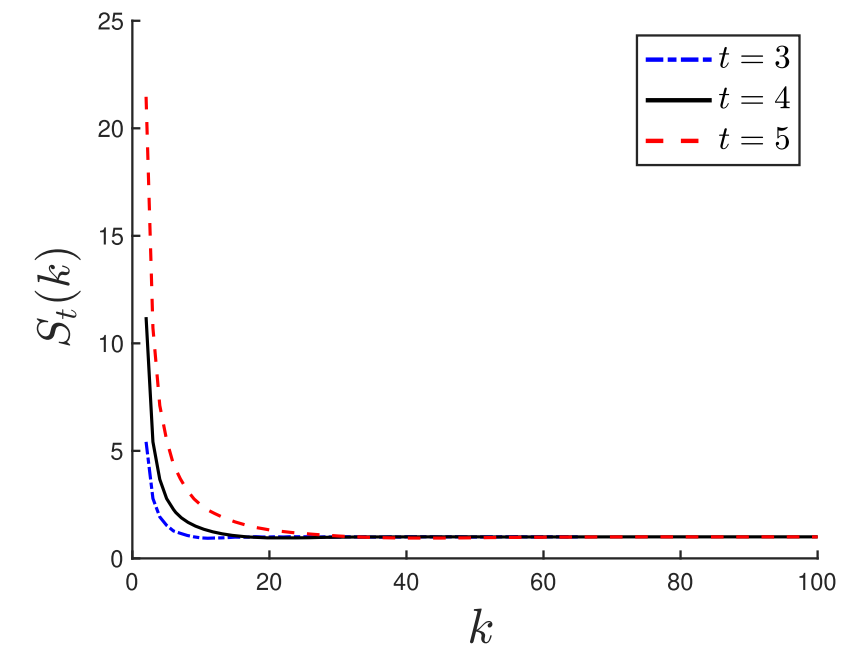
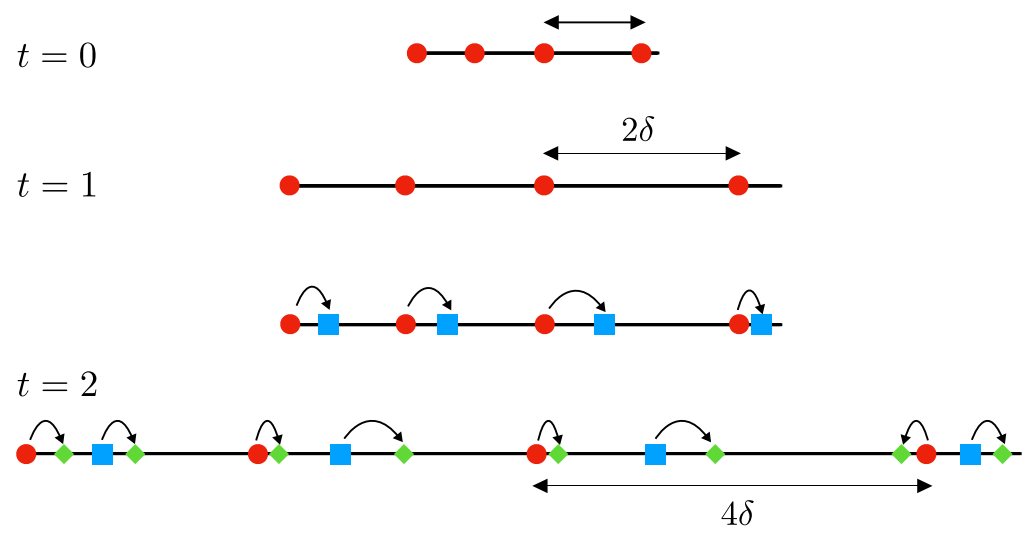
R. Ross SP, Phys. Rev. X (2025)



R. Ross SP, Phys. Rev. X (2025)



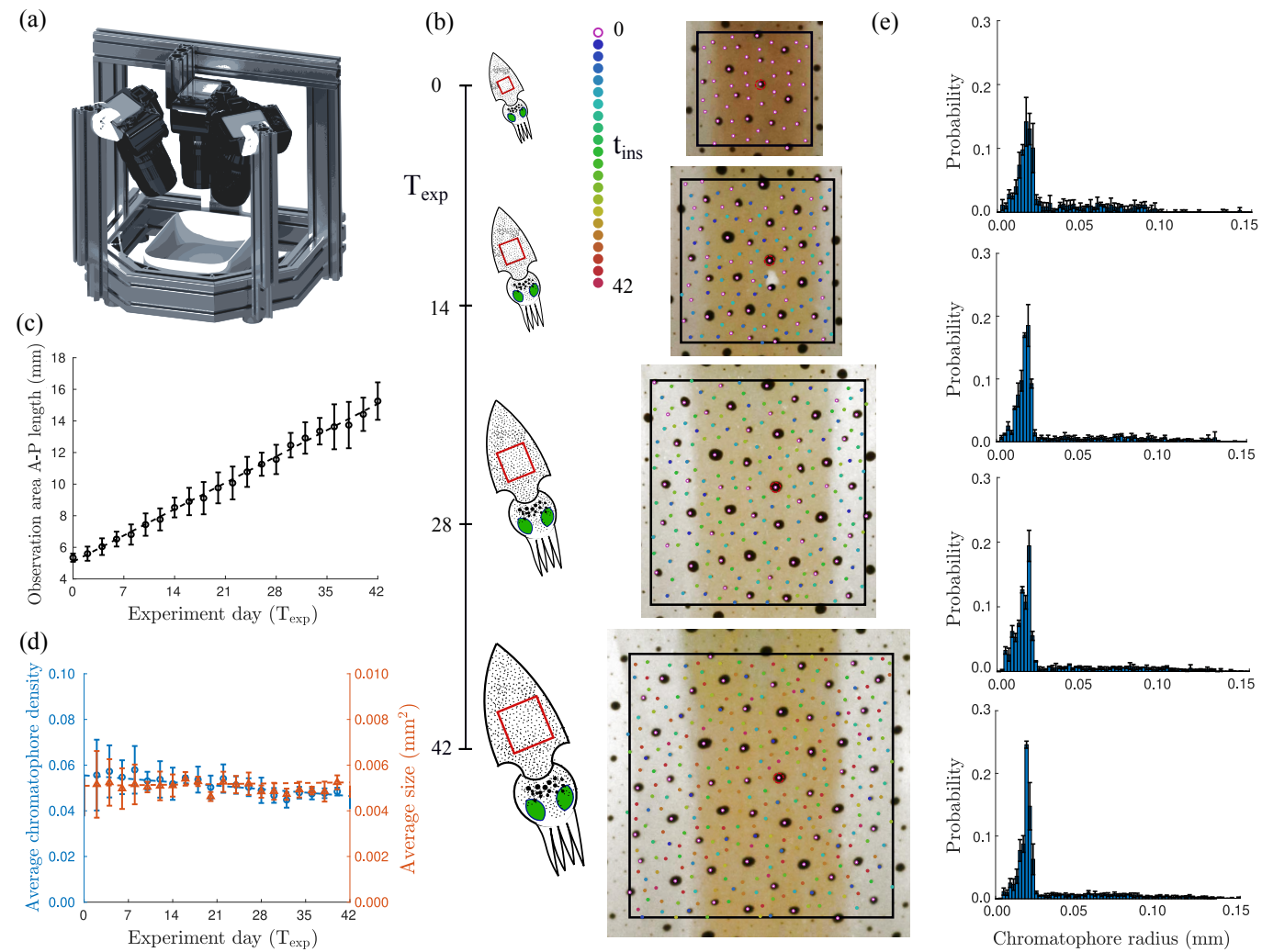
Hierarchical toy model (1D)



Density fluctuations are exported to larger and larger scales



Stationary distribution of chromatophore sizes



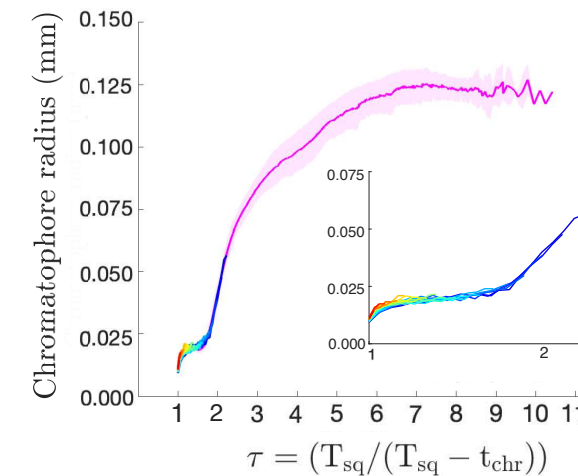
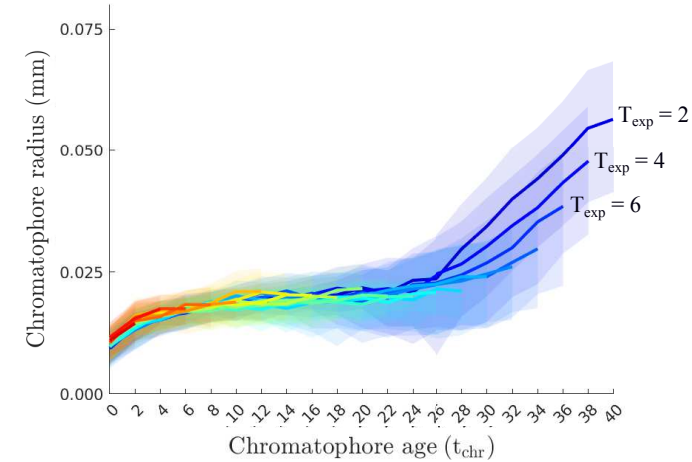
R. Ross SP, Phys. Rev. X (2025)



Stationary distribution of chromatophore sizes

$$p(R) = \int_0^\infty dt_{\text{chr}} p(R|t_{\text{chr}}, T_{\text{squid}}) \frac{2(T_{\text{squid}} - t_{\text{chr}})}{T_{\text{squid}}^2}$$

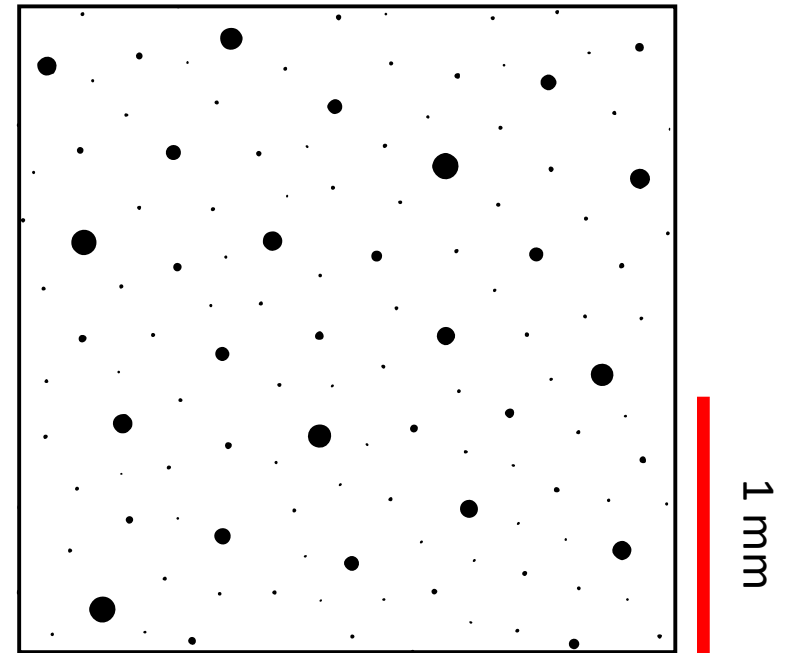
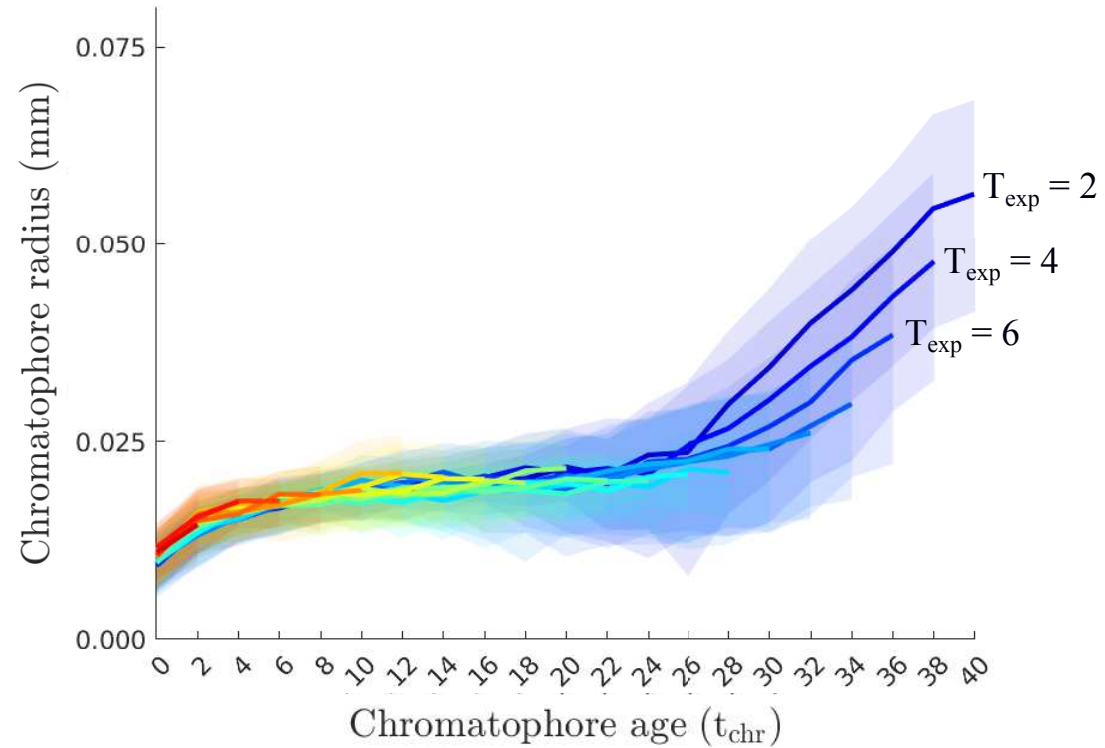
Chromatophore dynamics must depend on the “absolute” time



R. Ross SP, Phys. Rev. X (2025)



Chromatophore sizes

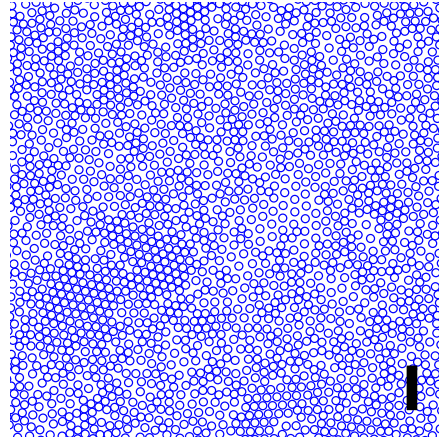


R. Ross SP, Phys. Rev. X (2025)

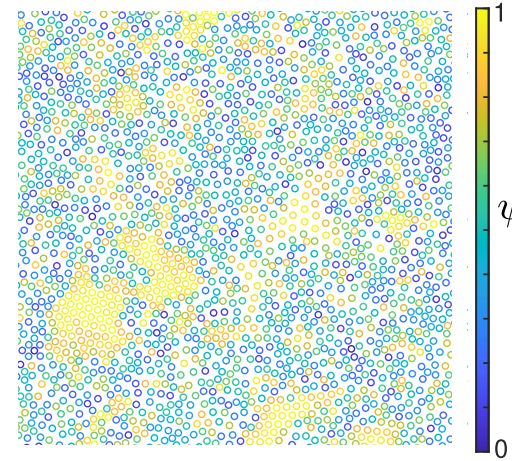


A critical growing state?

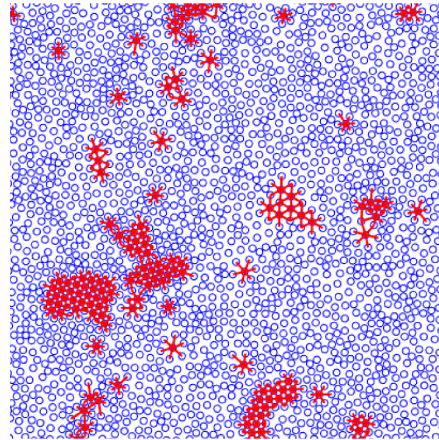
(c) Simulation output



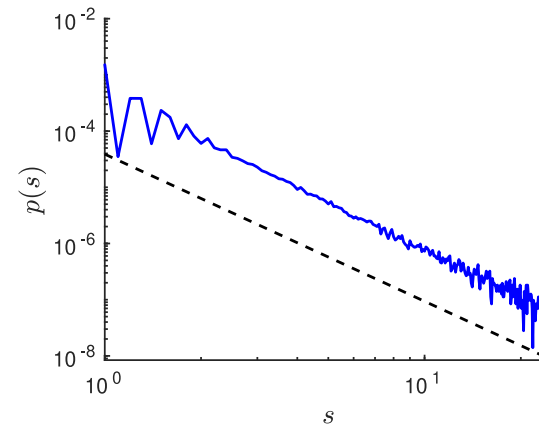
(d) Hexatic order parameter



(e) Connected components

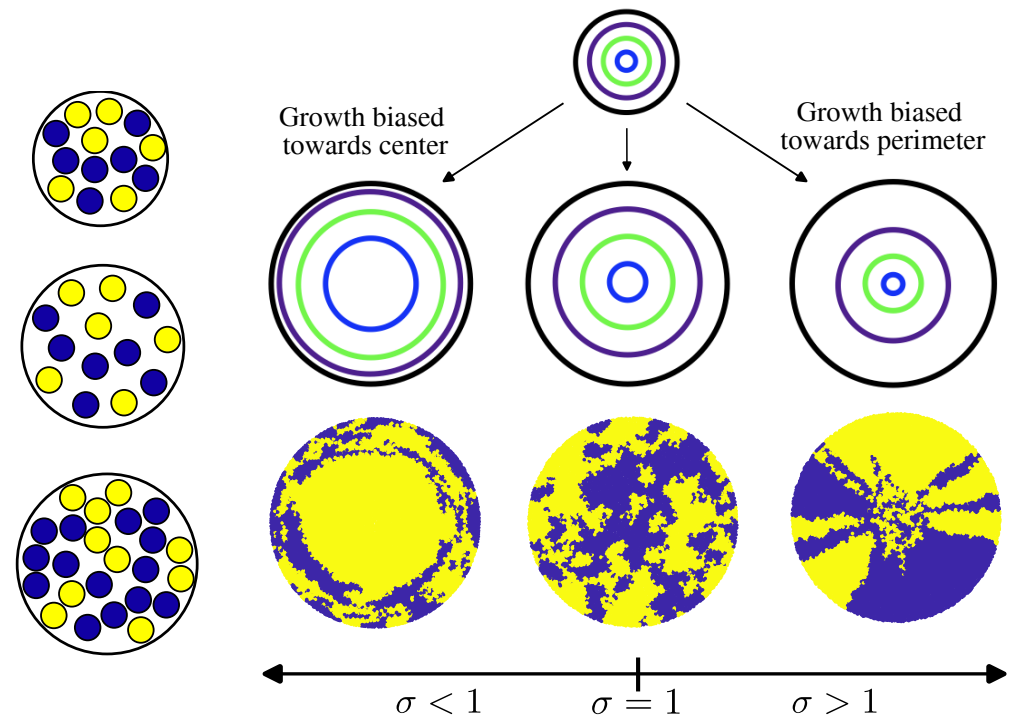


(f)

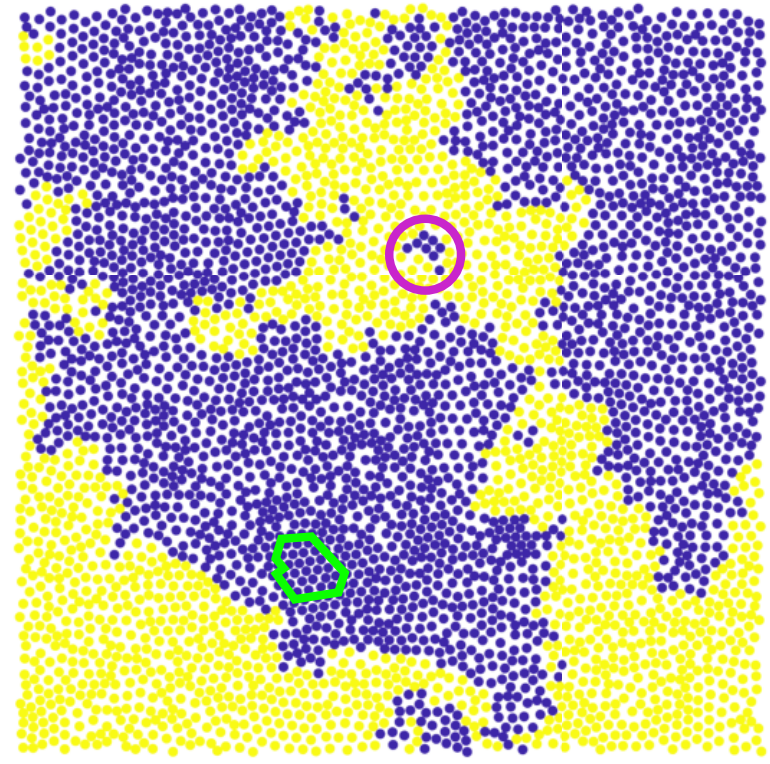




The growing voter model



Displacement field $\lambda = [r/R(t)]^\sigma$

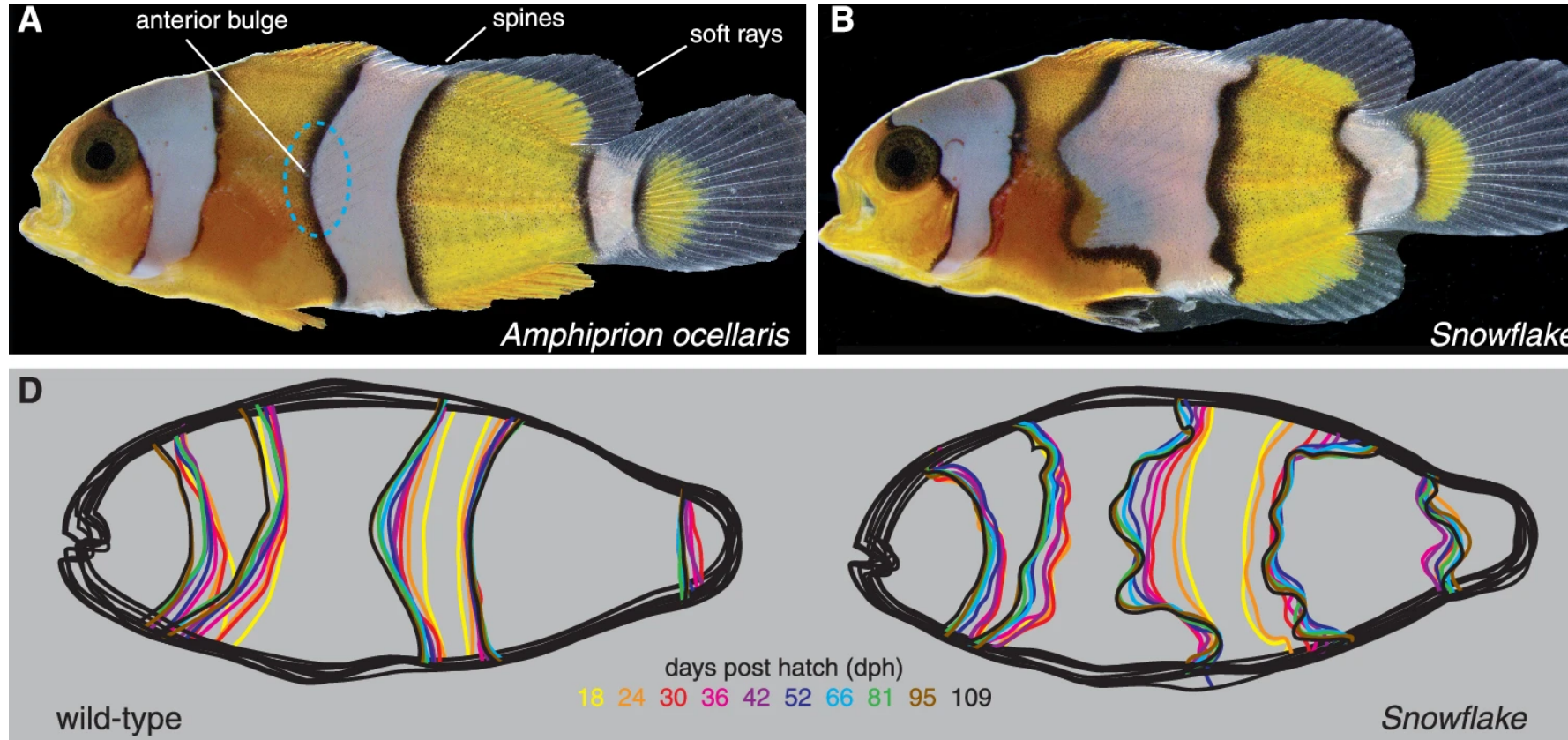


Phenomenological theory for the critical exponents

R. Ross, SP, arXiv:2411.09172



The snowflake mutation of clownfish

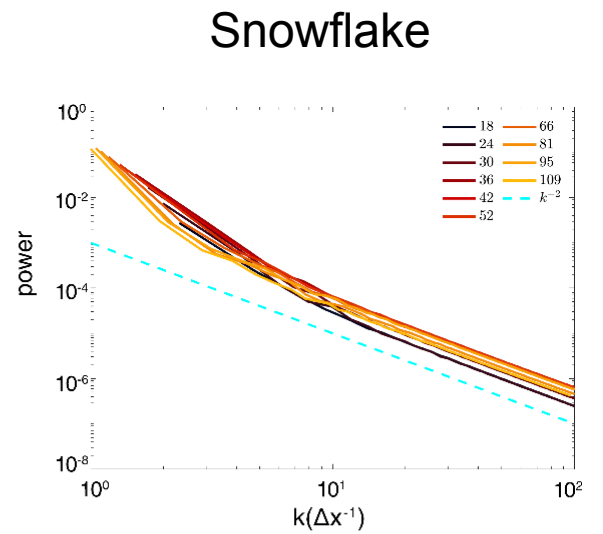
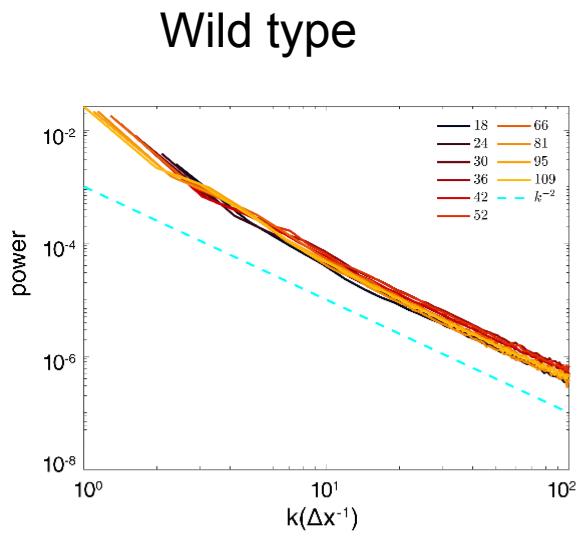
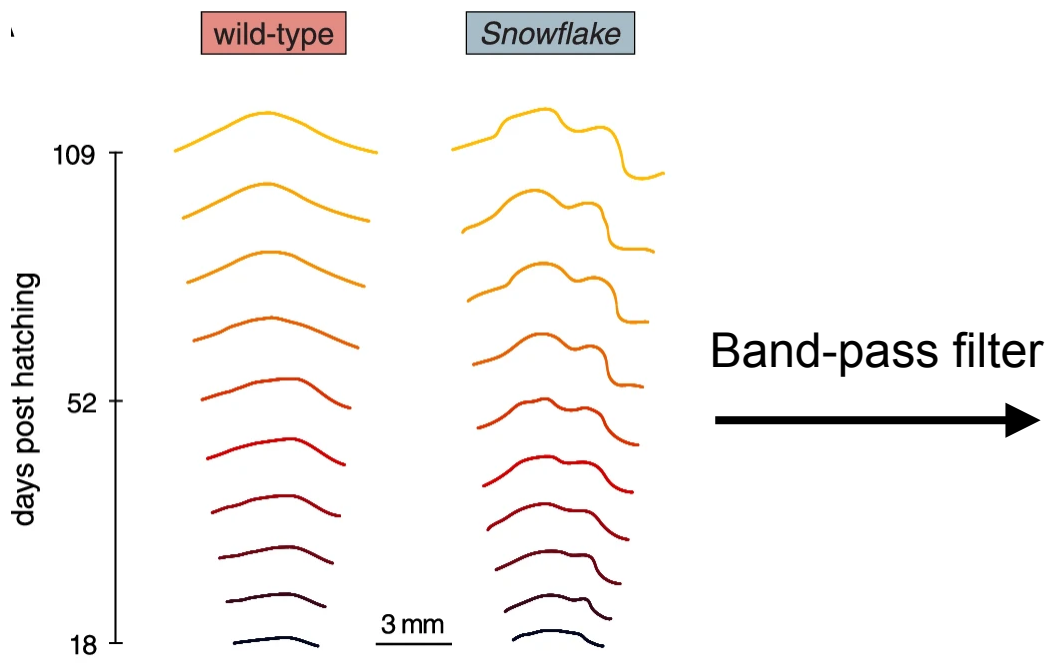


M. Klann ... SP ... V Laudet, Nature Communications (2026)



Edwards-Wilkinson model

$$\partial_t h(x, t) = D \nabla^2 h + \sigma \xi(x, t)$$



- Spectrum compatible with EW model for both wild type and Snowflake
- Either larger noise amplitude or smaller surface tension in Snowflake

M. Klann ... SP ... V Laudet, Nature Communications (2026)



Conclusions

- Growth is ubiquitous in biological development
- It gives rise to new physics
- New universality classes?



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Laudet Unit (Clownfish)

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