

Anomalous suppression of mean p_T fluctuations **potentially** signals QCD CEP

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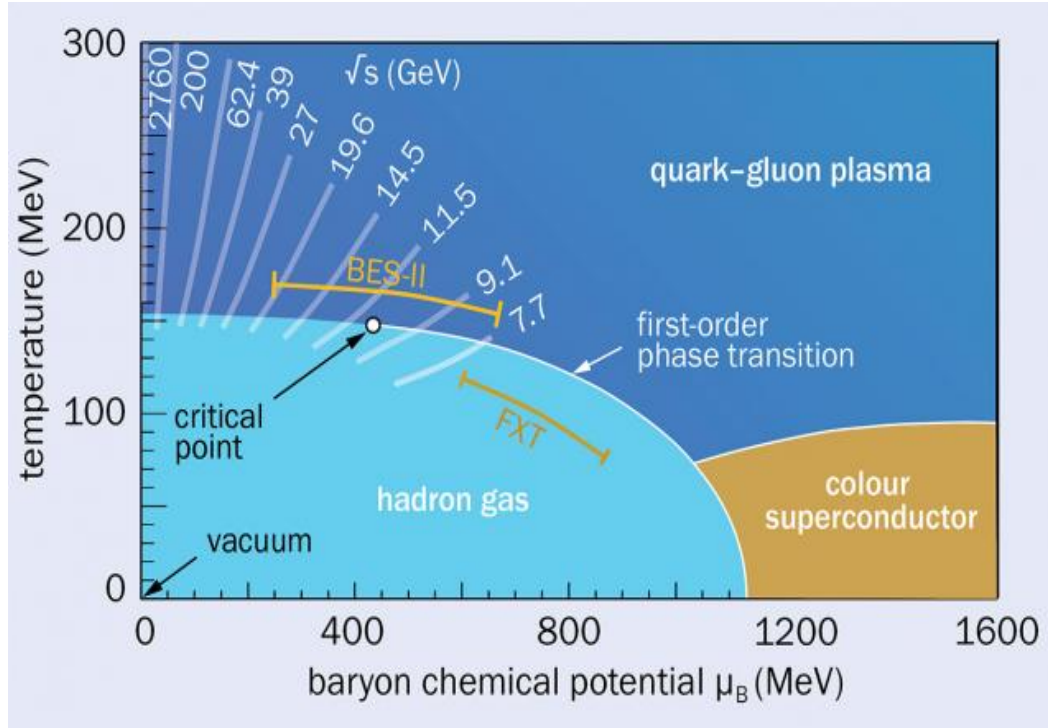
“QCD critical point and hydro evolution”

Yukawa Institute for Theoretical Physics, Kyoto, Jun. 2, 2026

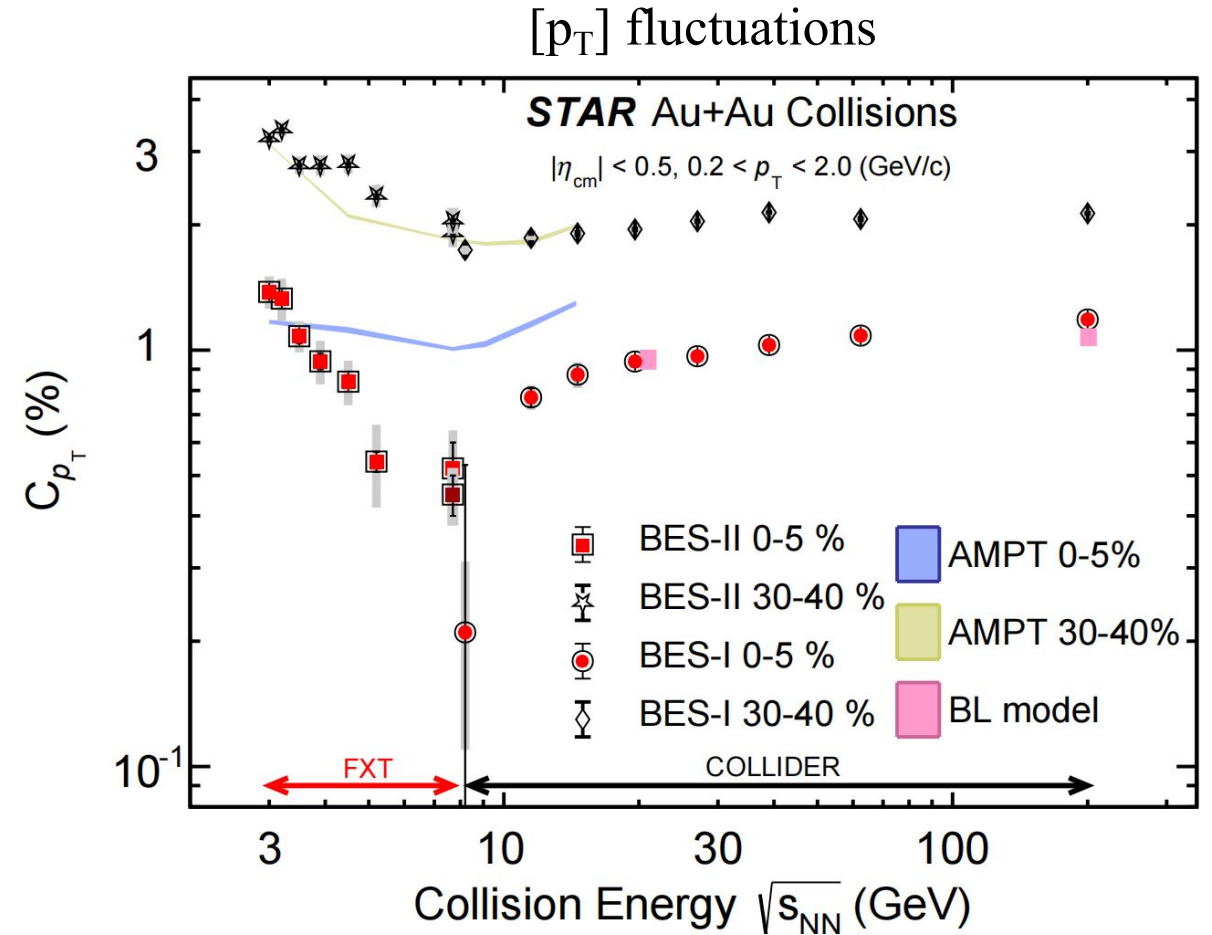


Introduction

- QCD phases and BES:



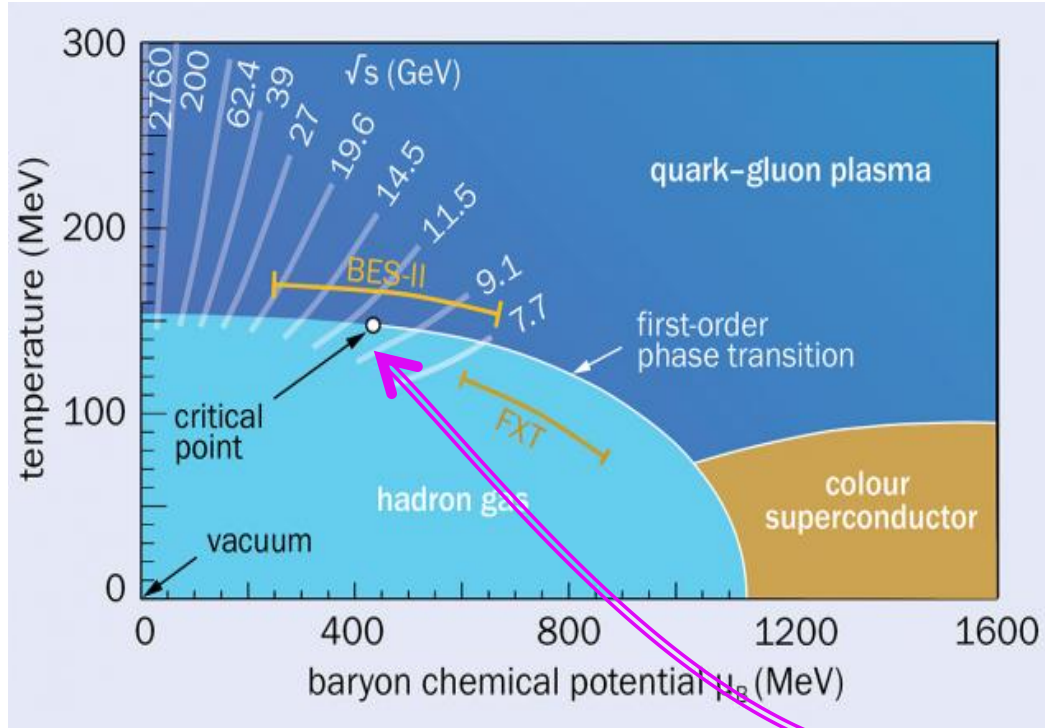
[STAR collaboration, 2604.06434]



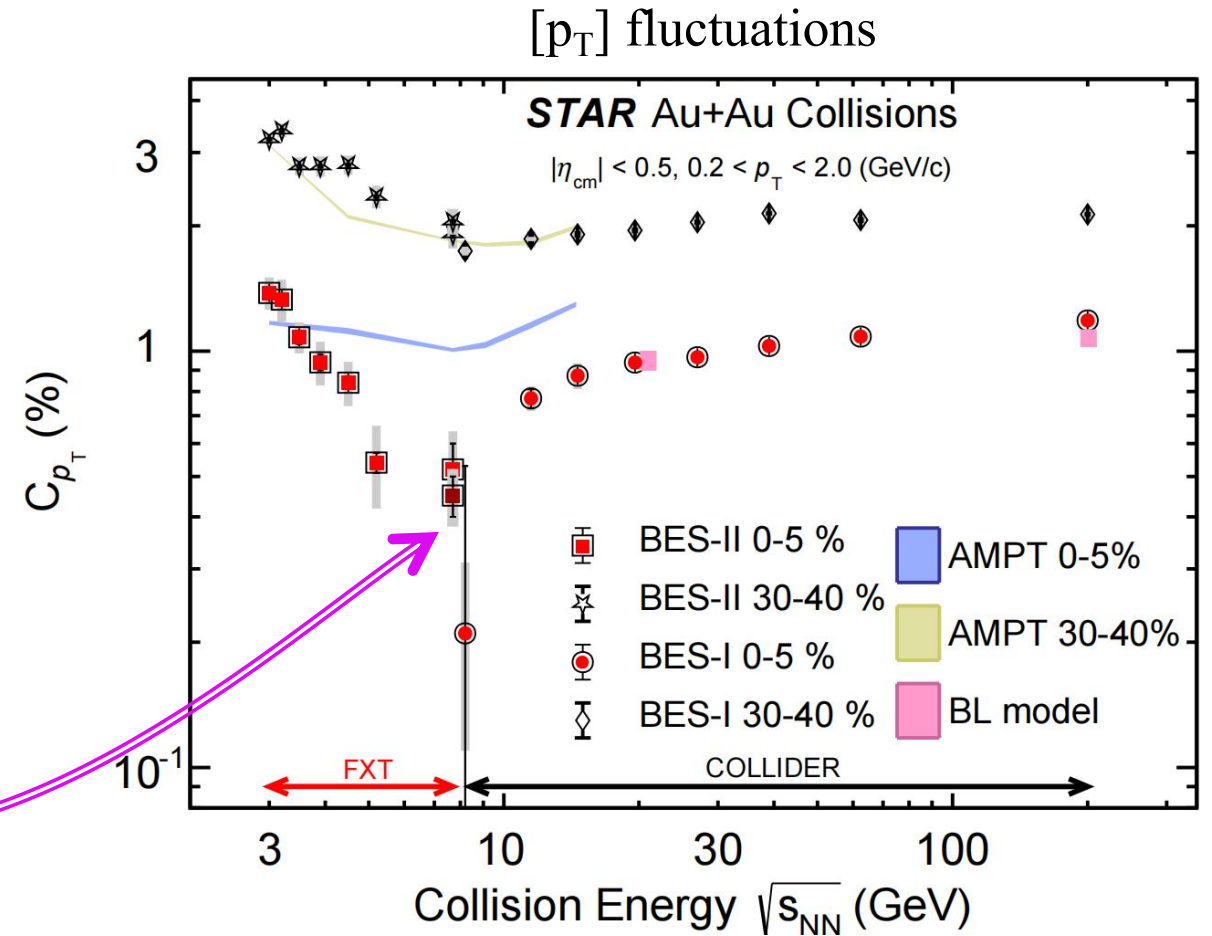
* Well identified non-monotonic signatures in BES are associated with QCD CEP.

Introduction

- QCD phases and BES:

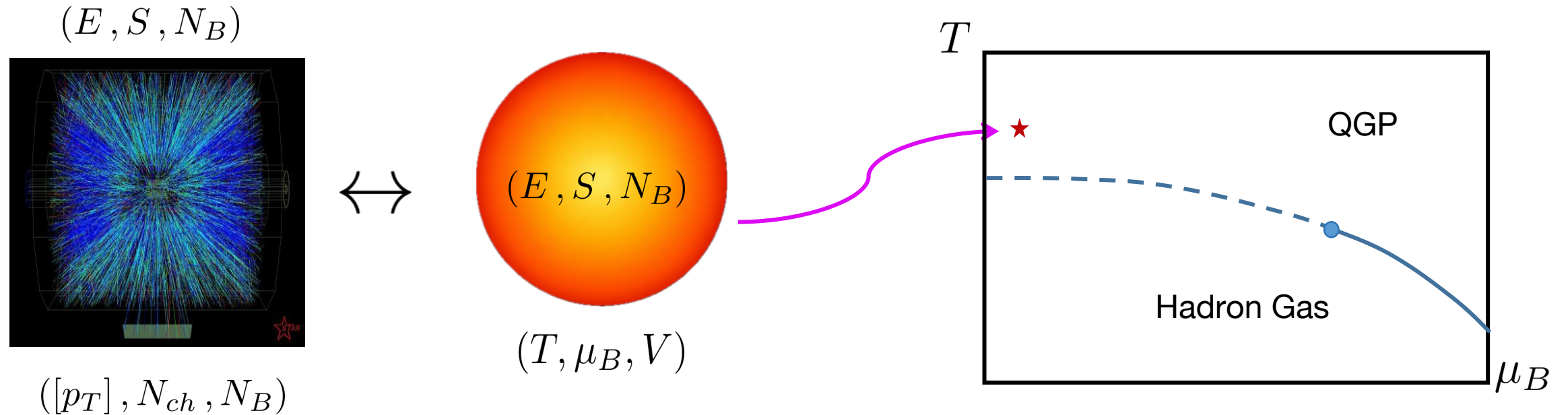


[STAR collaboration, 2604.06434]



* Well identified non-monotonic signatures in BES are associated with QCD CEP.

Particles from heavy-ion collisions = an effective QCD fireball



- In heavy-ion collisions, each collision produces a dynamical state: **Multi-Particle final State (MPS)**.
- Multi-Particle State is thermally equivalent to an effective smooth and uniform QCD fireball:

$$E^{\text{MPS}} = E^{\text{eff}} = \underbrace{e(T, \mu_B)}_{\text{QCD EoS}} V, \quad S^{\text{MPS}} = S^{\text{eff}} = \underbrace{s(T, \mu_B)}_{\text{QCD EoS}} V, \quad N_B^{\text{MPS}} = N_B^{\text{eff}} = \underbrace{n_B(T, \mu_B)}_{\text{QCD EoS}} V$$

MPS in ultra-central collisions (UCC) detect QCD thermodynamic response

E.g., speed of sound: ($\mu_B \approx 0$, LHC energies and top RHIC energies)

- Initial-state quantum fluct. introduce variations in energy from event to event: ΔE (or ΔS):

$$\text{const. } V \Rightarrow \underbrace{\Delta T}_{\text{resp.}} = \underbrace{C_V^{-1}}_{\text{thermal resp.}} \underbrace{(T\Delta S)}_{\text{func. source}} = \frac{c_s^2}{S} (T\Delta S), \quad C_V \sim \chi_{ee}$$

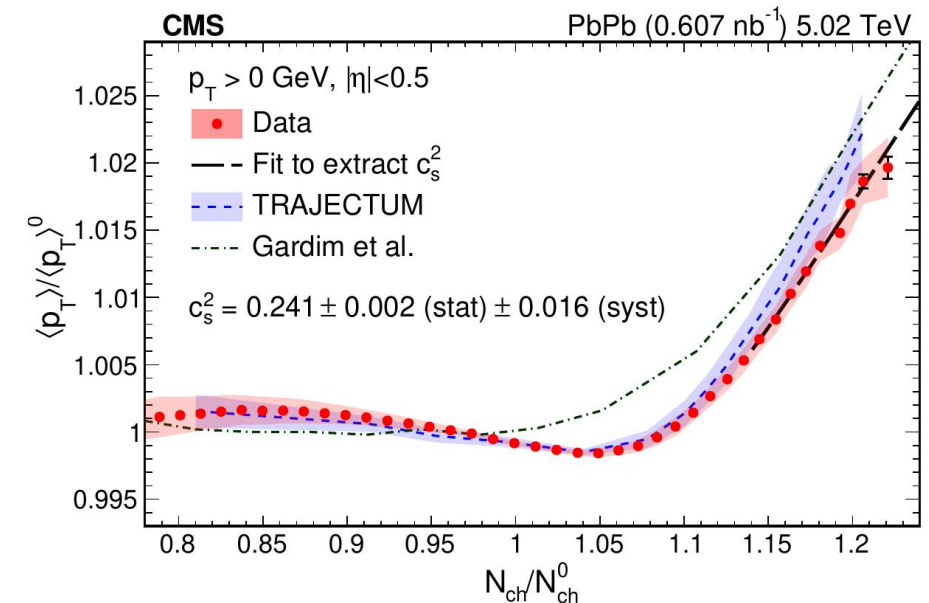
[CMS collaboration, 2401.06896]

- Further related to exp., $[p_T]$ vs N_{ch} in UCC measures c_s :

$$\Delta \log[p_T] = c_s^2 \Delta \log N_{ch}$$

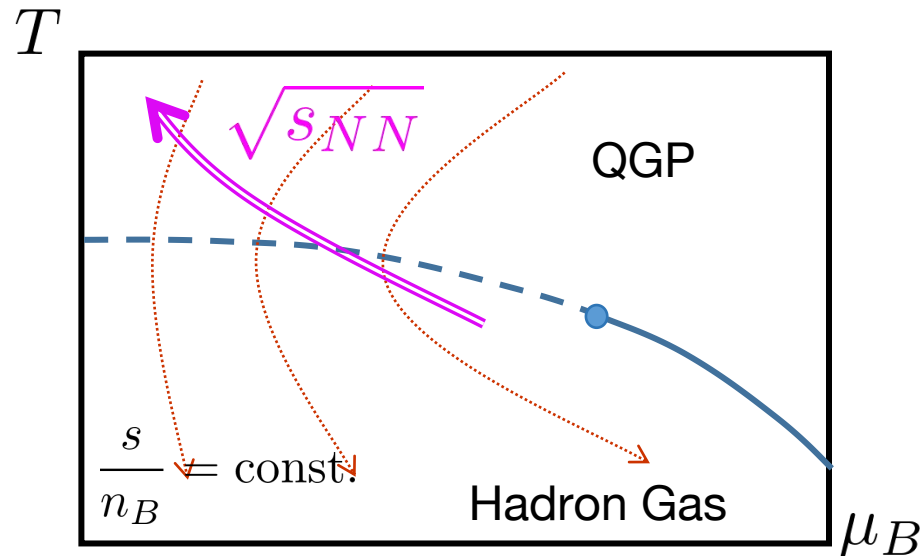
- Assumptions and conditions:

- Linear correlations: $[p_T] \sim T$, $N_{ch} \sim S$
- Saturated effective volume in UCC: $V = \text{const.}$



[F. Gardim et al, 19',]

Thermodynamical resp. with finite baryon chemical potential



- Isentropic condition: A QCD fireball close to local equilibrium (hydro) evolves approximately with a constant s/n_B

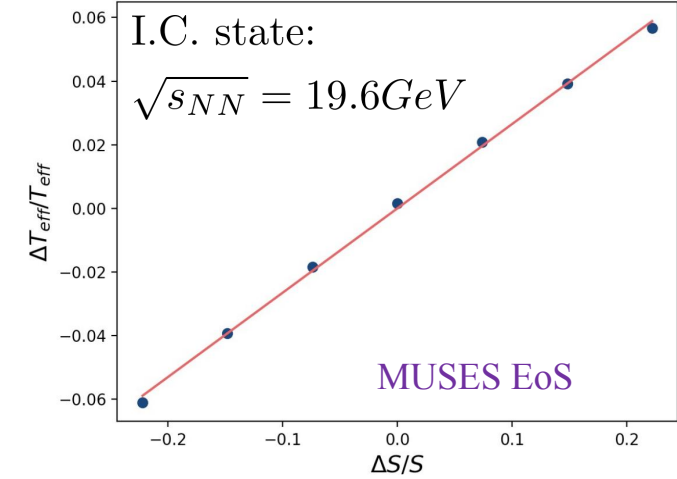
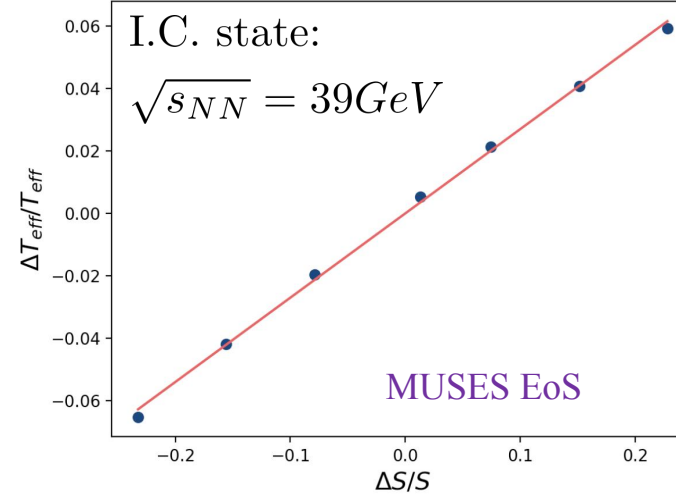
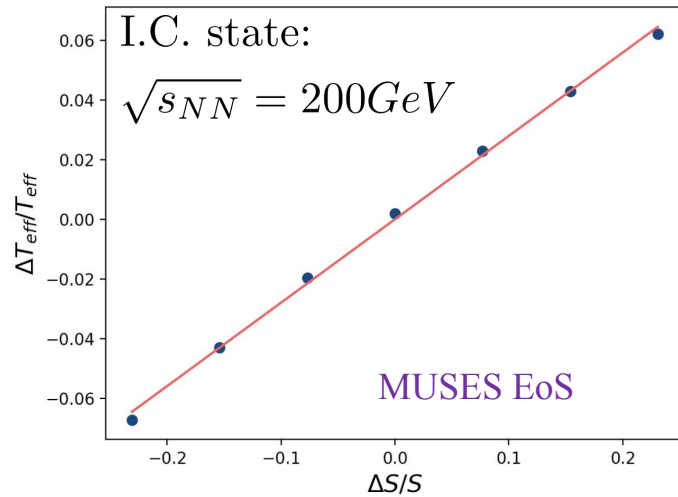
$$\sigma \equiv \frac{s}{n_B}$$

- Dissipations alter the condition, but not much.

- Fundamental thermodynamic relations and the constant s/n_B condition,

$$\begin{cases} de = Tds + \mu_B dn_B \\ dP = sdT + n_B d\mu_B \end{cases} \rightarrow Ts \left(1 + \underbrace{\left(\frac{\partial \mu_B}{\partial T} \right)_\sigma}_{\equiv \lambda} \sigma^{-1} \right) d \log T = c_s^2 Ts \left(1 + \frac{\mu_B}{T} \sigma^{-1} \right) d \log s$$

Thermodynamical resp. with finite chemical potential



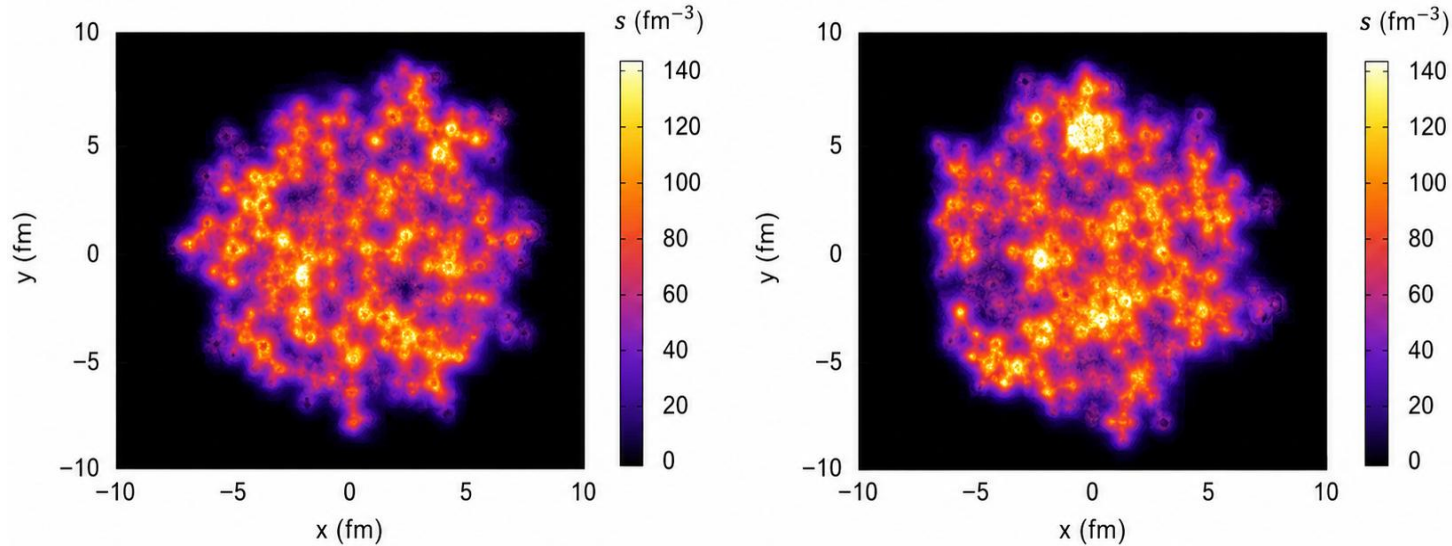
[MUSES EoS, 2402.0863]

- Accounting for the condition of constant volume: $s \rightarrow S/V$, $\sigma \rightarrow S/N_B$

$$d \log T = \frac{c_s^2 \left(1 + \frac{\mu_B}{T} \sigma^{-1}\right)}{1 + \lambda \sigma^{-1}} d \log S \quad \Leftrightarrow \quad dT = \underbrace{\frac{c_s^2}{S} \frac{1 + \frac{\mu_B}{T} \sigma^{-1}}{1 + \lambda \sigma^{-1}}}_{C_\sigma^{-1}} T dS$$

where C_σ is the generalized heat capacity with finite chemical potential wrt the isentropic condition

The effect of initial-state fluctuations



Au-Au: $\sqrt{s_{NN}} = 200 \text{ GeV}$

$b = 0 \text{ fm}$

$\Delta S \approx 0$

- Initial-state fluctuations are of quantum origin, and provide the **dominant** fluct. in the fireball.
- Initial-state fluctuations are associated with S , they appear in the resp. relation accordingly,

$$d \log T = \frac{c_s^2 \left(1 + \frac{\mu}{T} \sigma^{-1}\right)}{1 + \lambda \sigma^{-1}} d \log S + \frac{\delta_0}{1 + \lambda \sigma^{-1}}$$

* The effect of thermal fluct. are expected sub-leading, and are not considered in this work.

The thermodynamic resp. with finite μ_B in $[p_T]$ vs N_{ch}

- Accounting for the linear correlations ($[p_T] \sim T$ and $N_{ch} \sim S$), we find in UCC,

$$\Delta \log[p_T] = \frac{c_s^2 \left(1 + \frac{\mu}{T} \sigma^{-1}\right)}{1 + \lambda \sigma^{-1}} \Delta \log N_{ch} + \frac{\delta}{1 + \lambda \sigma^{-1}} + \underbrace{\delta'}_{\text{All other fluct.}}$$

- Observations:**

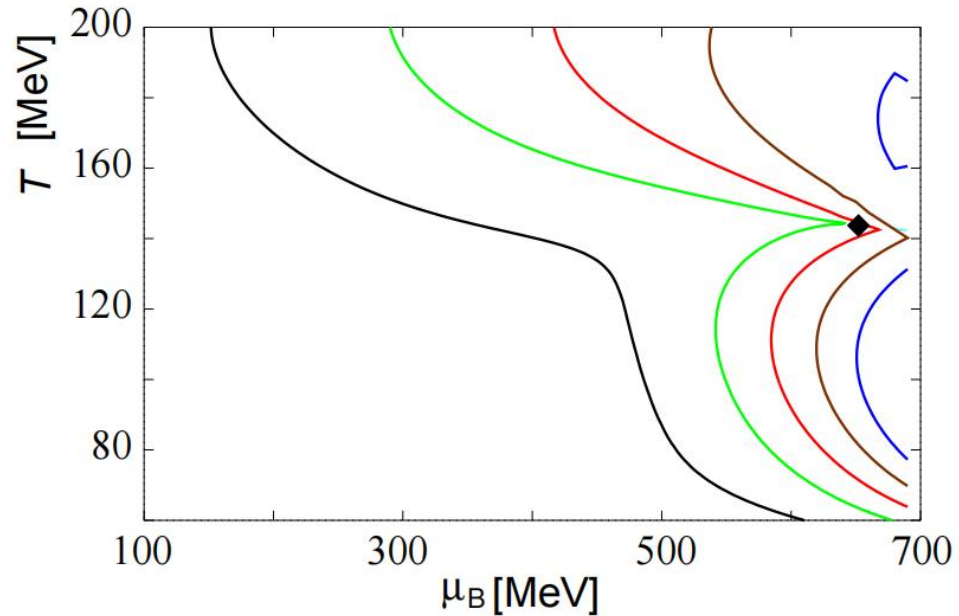
- Linear correlation is determined by $1/C_\sigma$. (NB: approaching QCD CEP, $C_\sigma \propto C_V \sim \xi^{0.17}$)
- Fluctuations (from both ΔN_{ch} and δ) are suppressed by $1/(1 + \lambda/\sigma)$. (NB: λ is singular around QCD CEP.)

- Predictions for UCC:**

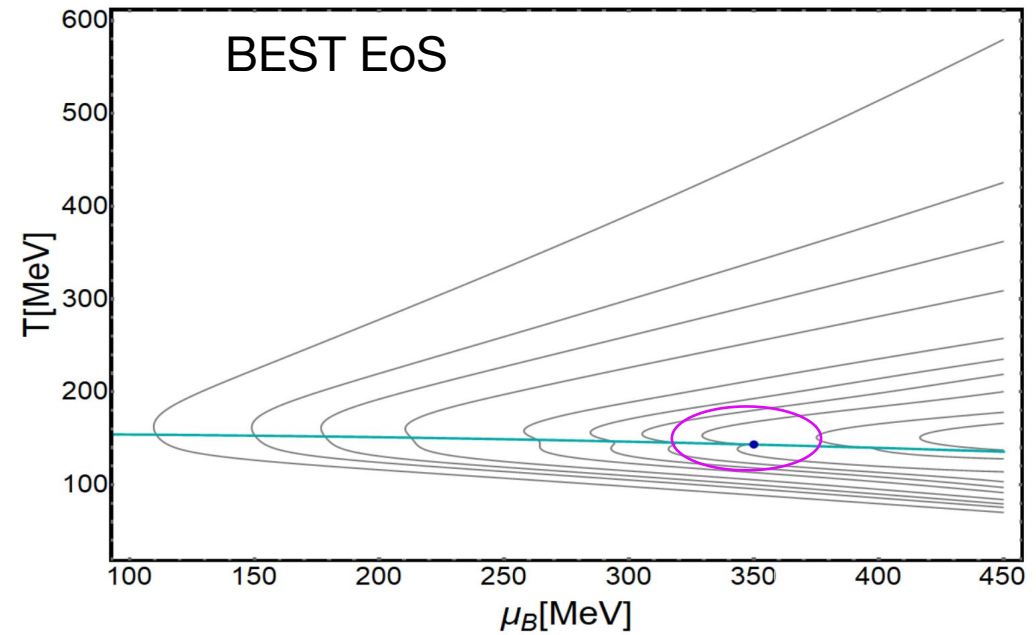
- Linear correlation between $[p_T]$ and N_{ch} is reduced around QCD CEP.
- Fluctuations of $[p_T]$ are suppressed around QCD CEP due to $1/(1 + \lambda/\sigma)$.

Critical properties of λ

[C. Nonaka and M. Asakawa, nuch-th/0410078]



[BEST collaboration, 1805.05259]

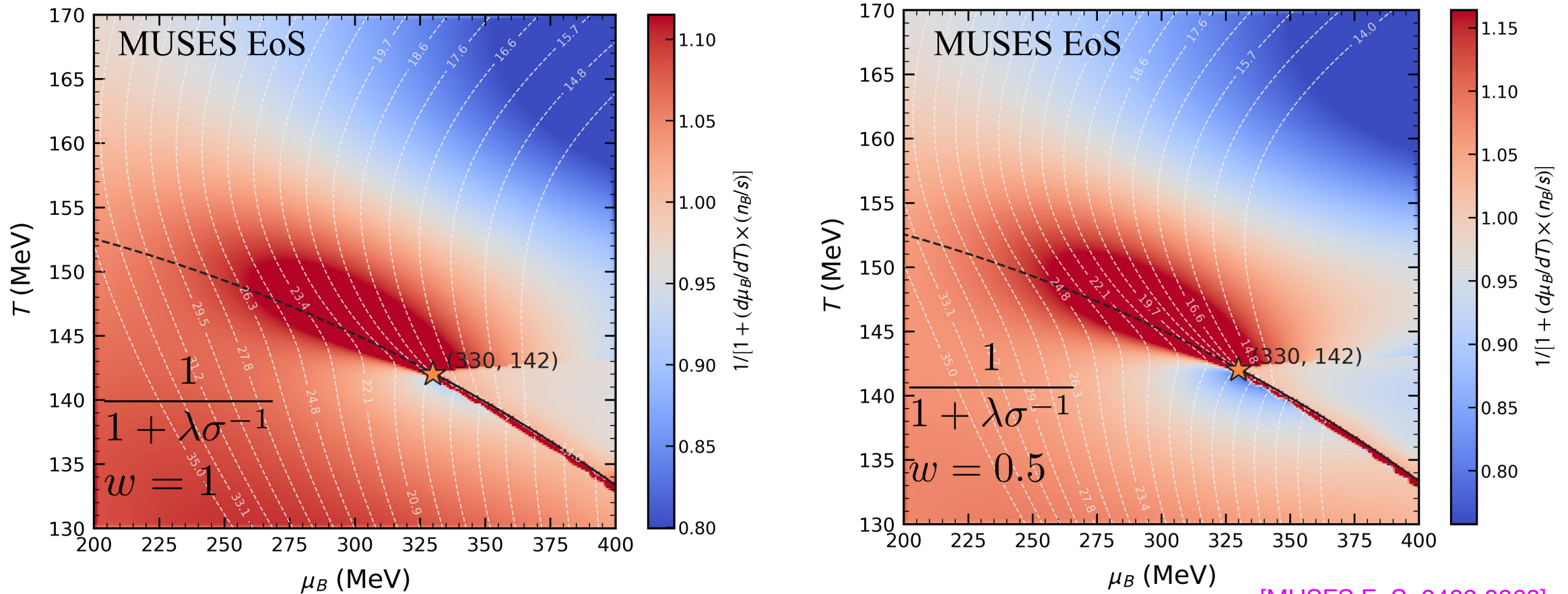


- Naive understanding: $|\lambda|$ maximizes due to the “focusing effect” of s/n_B lines around CEP. (Qualitative)

$$\lambda = \left(\frac{\partial \mu_B}{\partial T} \right)_\sigma \Leftrightarrow \text{slope along the const. } \frac{s}{n_B} \text{ line}$$

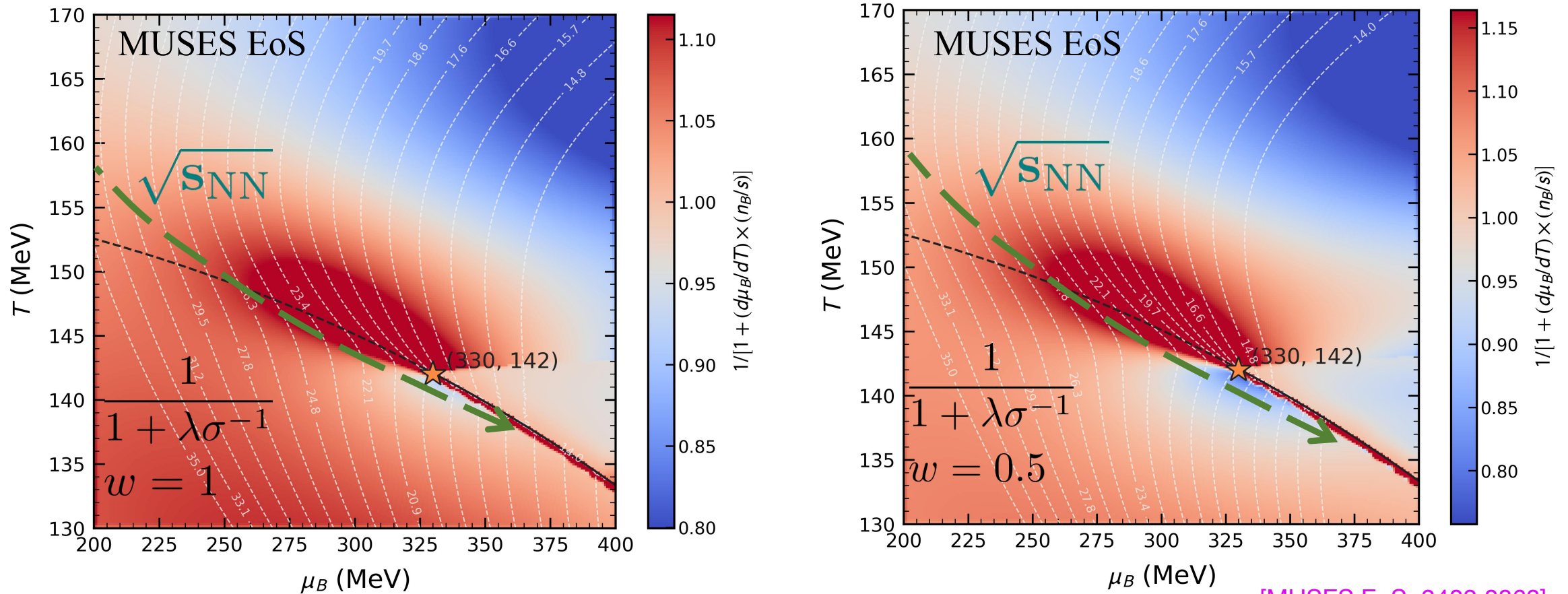
- Precise value of λ relies on non-universal behavior near CEP. (Quantitative)

The suppression factor $1/(1 + \lambda/\sigma)$ in QCD phases: plateau and basin



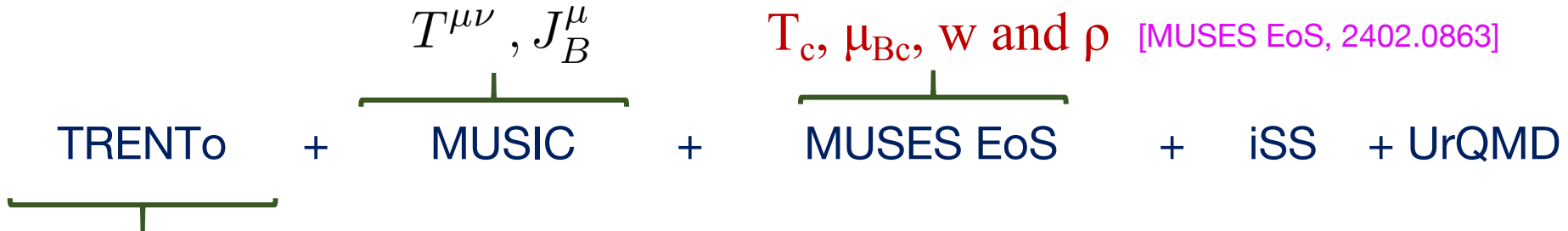
- Location of CEP is chosen regarding S/N_B at 7.7 GeV.
- Criticality of the CEP relies on parameter w in MUSES EoS.

The suppression factor $1/(1 + \lambda/\sigma)$ in QCD phases: plateau and basin



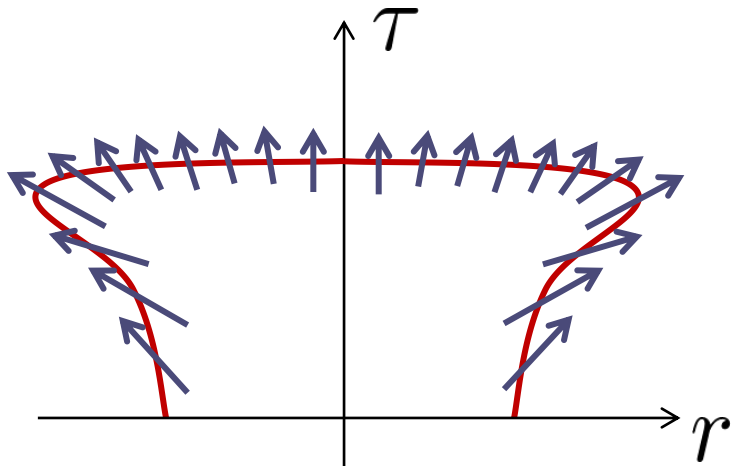
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Validation using Hybrid Hydro modeling (2+1D)



sample s and n_B independently, $\tau_0 \sim$ nucleus crossing time, k increases monotonically

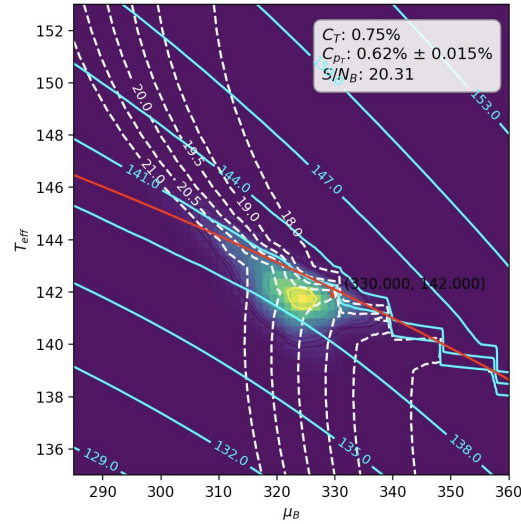
- Extract T and μ_B from the hydro freeze-out surface



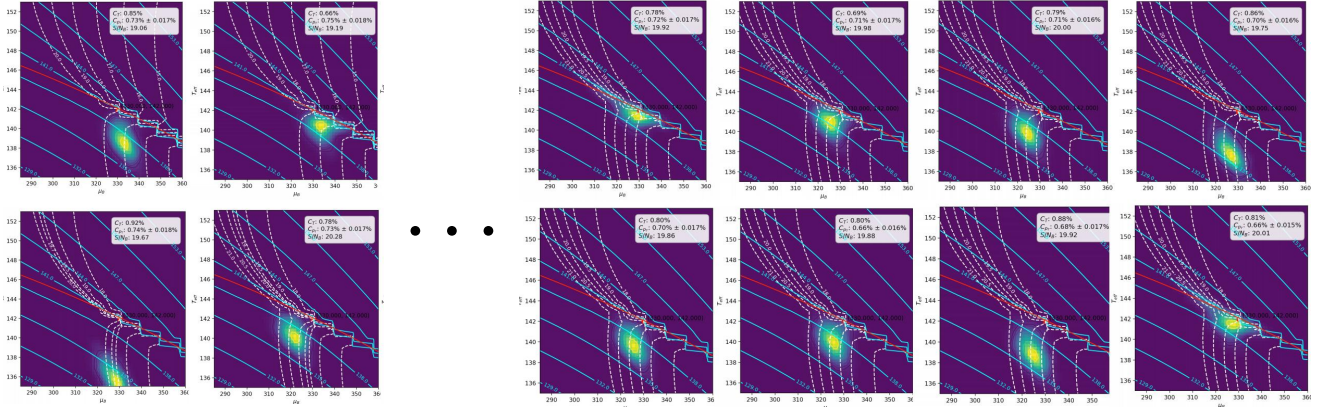
$$\begin{cases} E = \int_{\Sigma} d\sigma_{\mu} T^{0\mu} = e^{EOS}(T, \mu_B)V \\ S = \int_{\Sigma} d\sigma_{\mu} s^{\mu} = s^{EOS}(T, \mu_B)V \\ N_B = \int_{\Sigma} d\sigma_{\mu} n_B^{\mu} = n_B^{EOS}(T, \mu_B)V \end{cases}$$

* We consider (at least transient) thermalization of QCD system above 3.5 GeV in BES.

Assuming 7.7 GeV is special ...



$\times N \rightarrow$

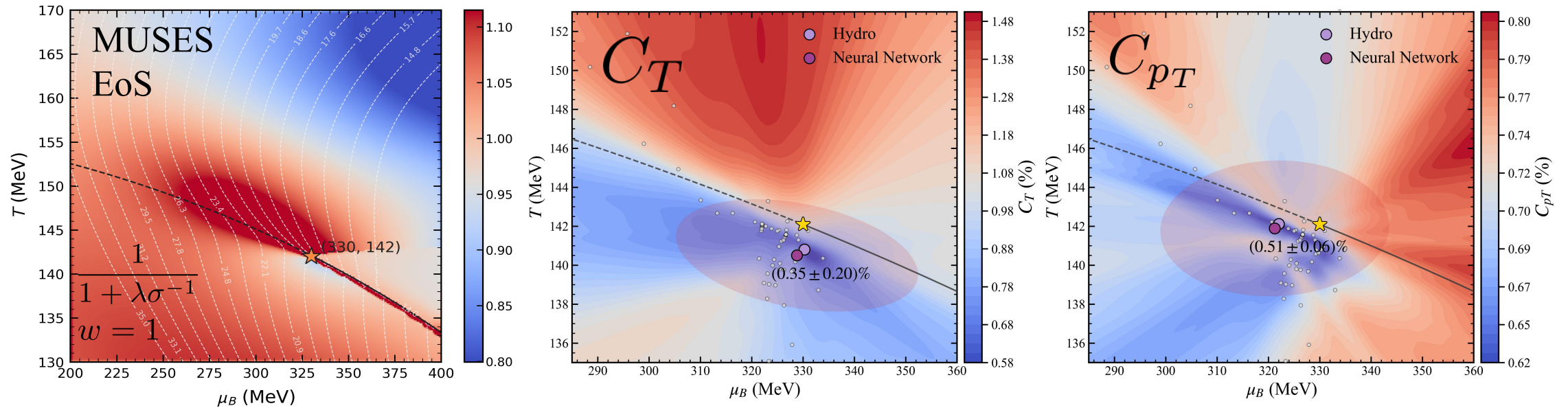


- Scan QCD phases around 7.7 GeV: with $\langle N_{ch} \rangle$, $\langle N_B \rangle$ and $\langle [p_T] \rangle$ tuned according to STAR exp.
- Accounting for errors, we roughly have for 7.7 GeV 0-5% events, $S/N_B \approx 20 \pm 1$
 for $w = 1.0$ in MUSES $\rightarrow T_c \sim 142\text{MeV}$, $\mu_{Bc} \sim 330\text{ MeV}$
- EbE simulations of 0-5% central Au-Au collision EbE (1k events, $N \sim 70$ sets) and def.

$$C_T \equiv \sqrt{\langle (\delta T)^2 \rangle} / \langle T \rangle$$

$$C_{p_T} \equiv \sqrt{\langle (\delta [p_T])^2 \rangle} / \langle [p_T] \rangle$$

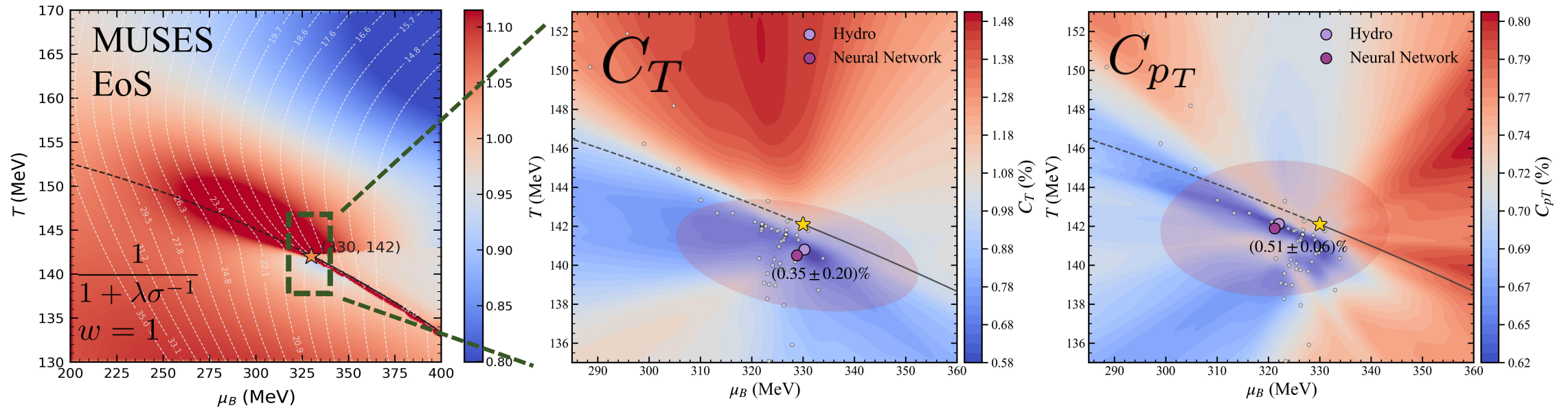
Reconstruct the “plateau and basin” from hydro + EoS



- White symbols correspond to raw hydro results.
- Distribution reconstructed using NN to interpolate hydro results and estimate the minimum.
- **Distribution of $C_T \neq$ Distribution of C_p ,**

$$[p_T] = aT \left(1 + b \frac{\mu_B}{T} + \dots \right)$$

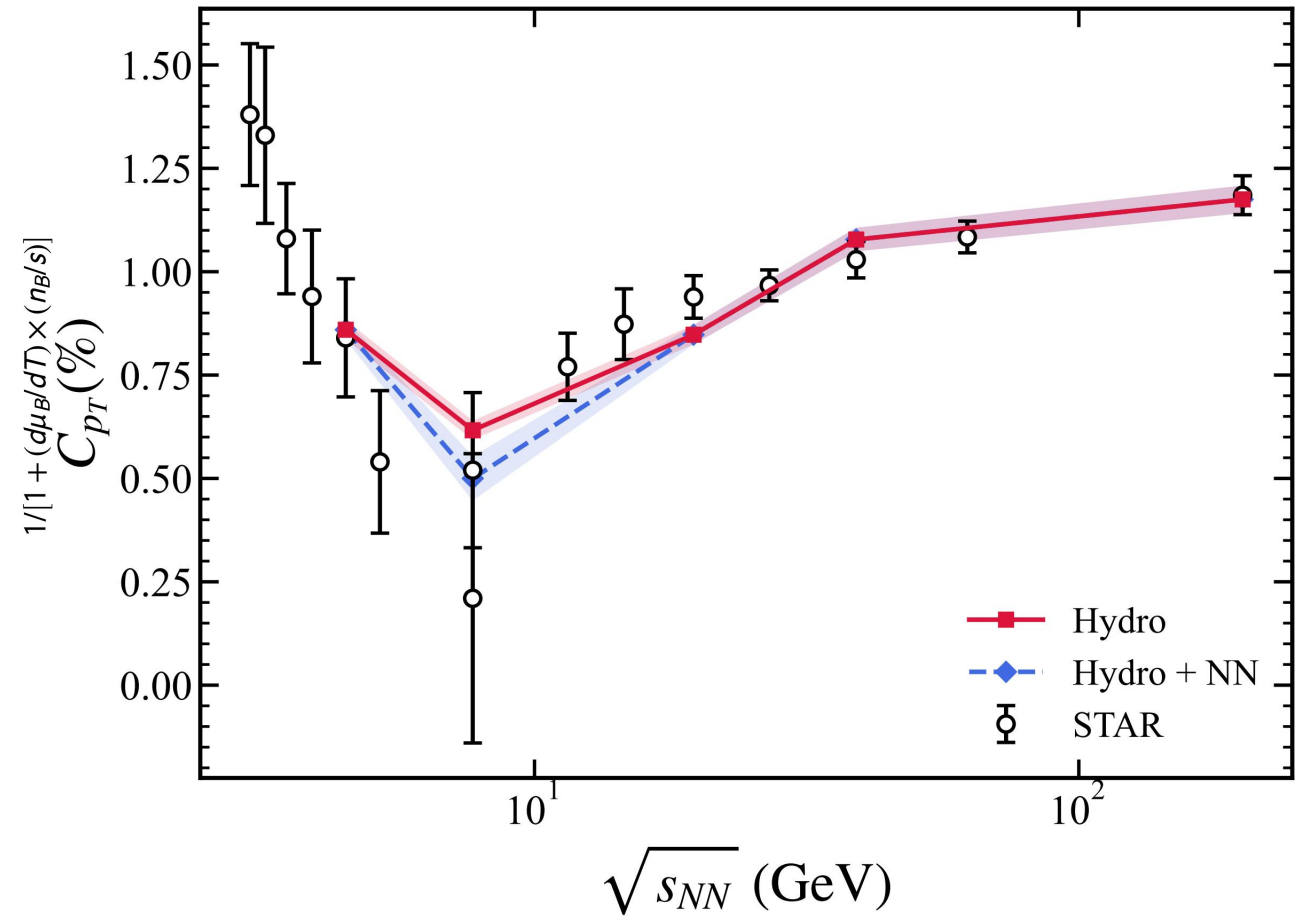
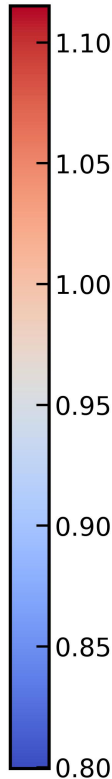
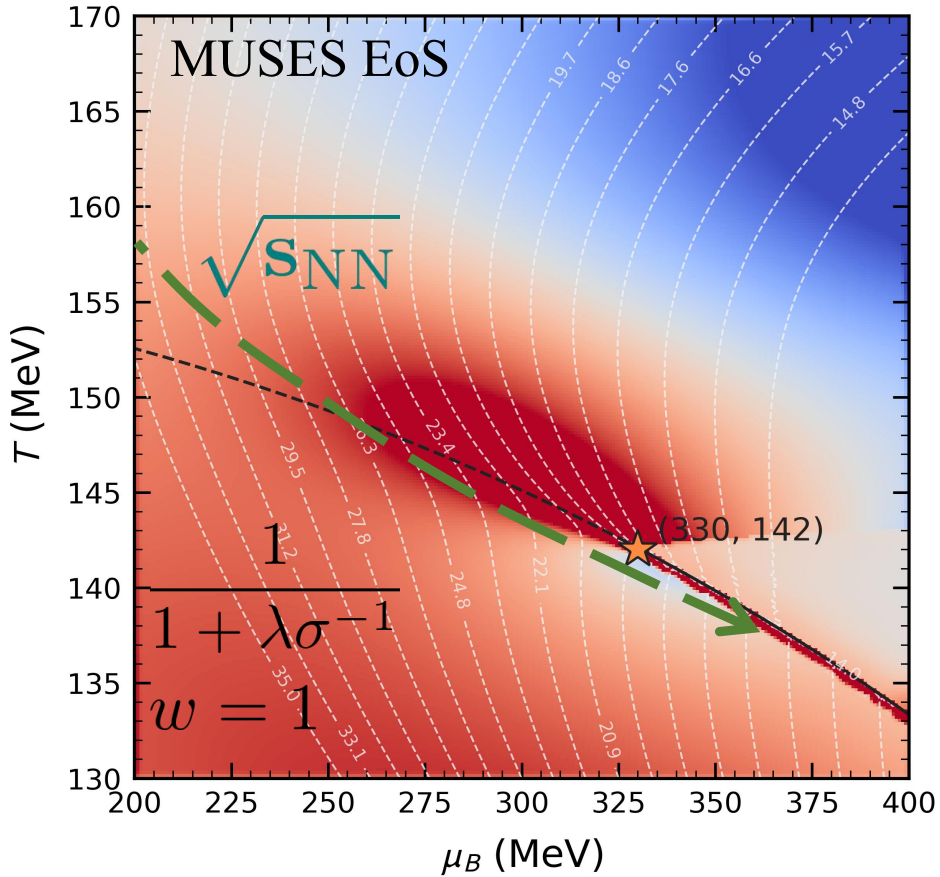
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Extend to other energies

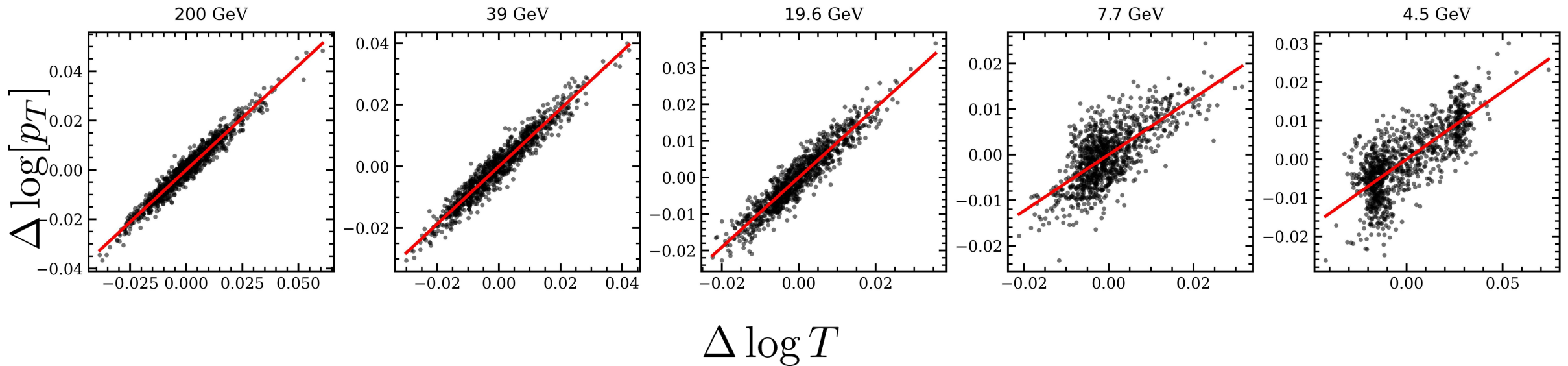


Hydro parameters are chosen such that they vary monotonically with respect to $\sqrt{S_{NN}}$.

Summary

- Anomalous suppression of mean p_T fluctuations **potentially** signals QCD CEP ==> Qualitative True
- Quantitative analyses, such the location of CEP, requires more rigorous treatment of hydro modeling and model parameters (Bayesian analyses), effect of thermal fluct., etc.

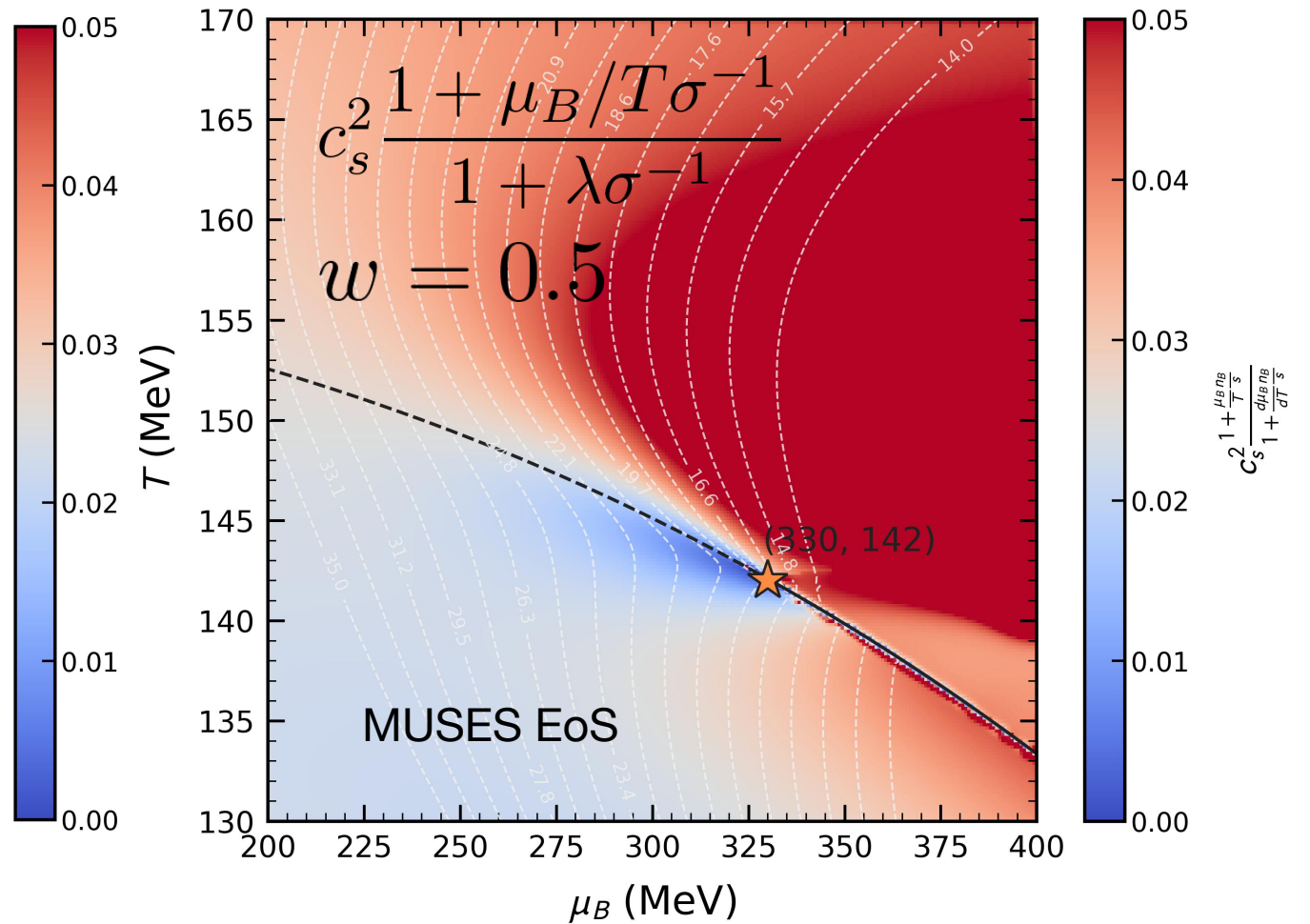
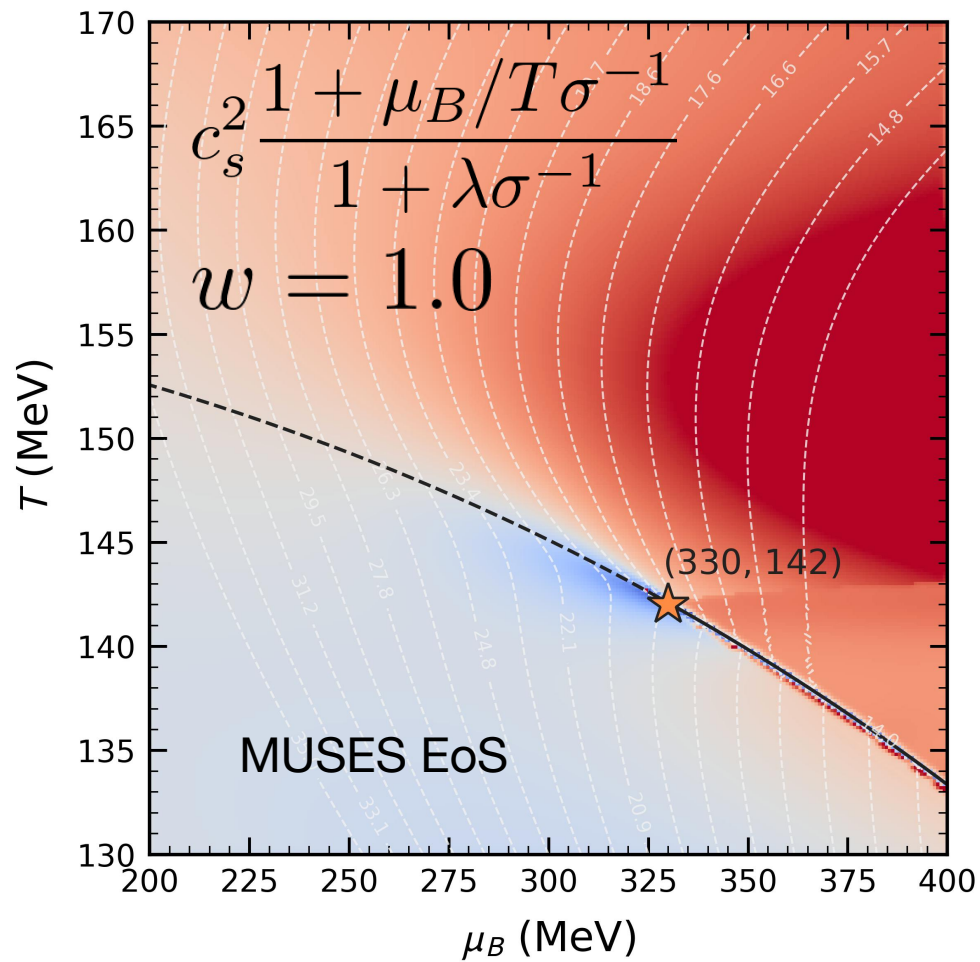
Linear correlations between $[p_T]$ and T



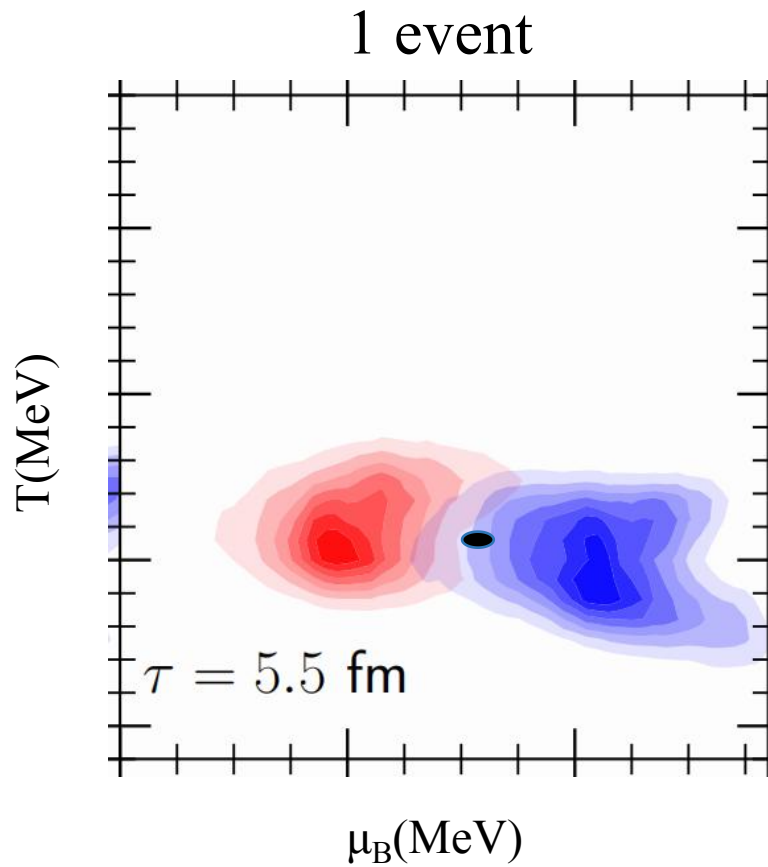
$$[p_T] = aT + b \quad \rightarrow \quad \Delta \log T = \Delta \log [p_T]$$

$$[p_T] = aT \left(1 + b \frac{\mu_B}{T} + \dots \right) \quad \nrightarrow \quad \Delta \log T = \Delta \log [p_T]$$

Linear coefficient $\sim 1/C_\sigma$



Why thermal fluctuations are sub-leading



VS

