

# Fluctuation Signatures from the Equation of State to Proton Factorial Cumulants

*Jamie M. Karthein*

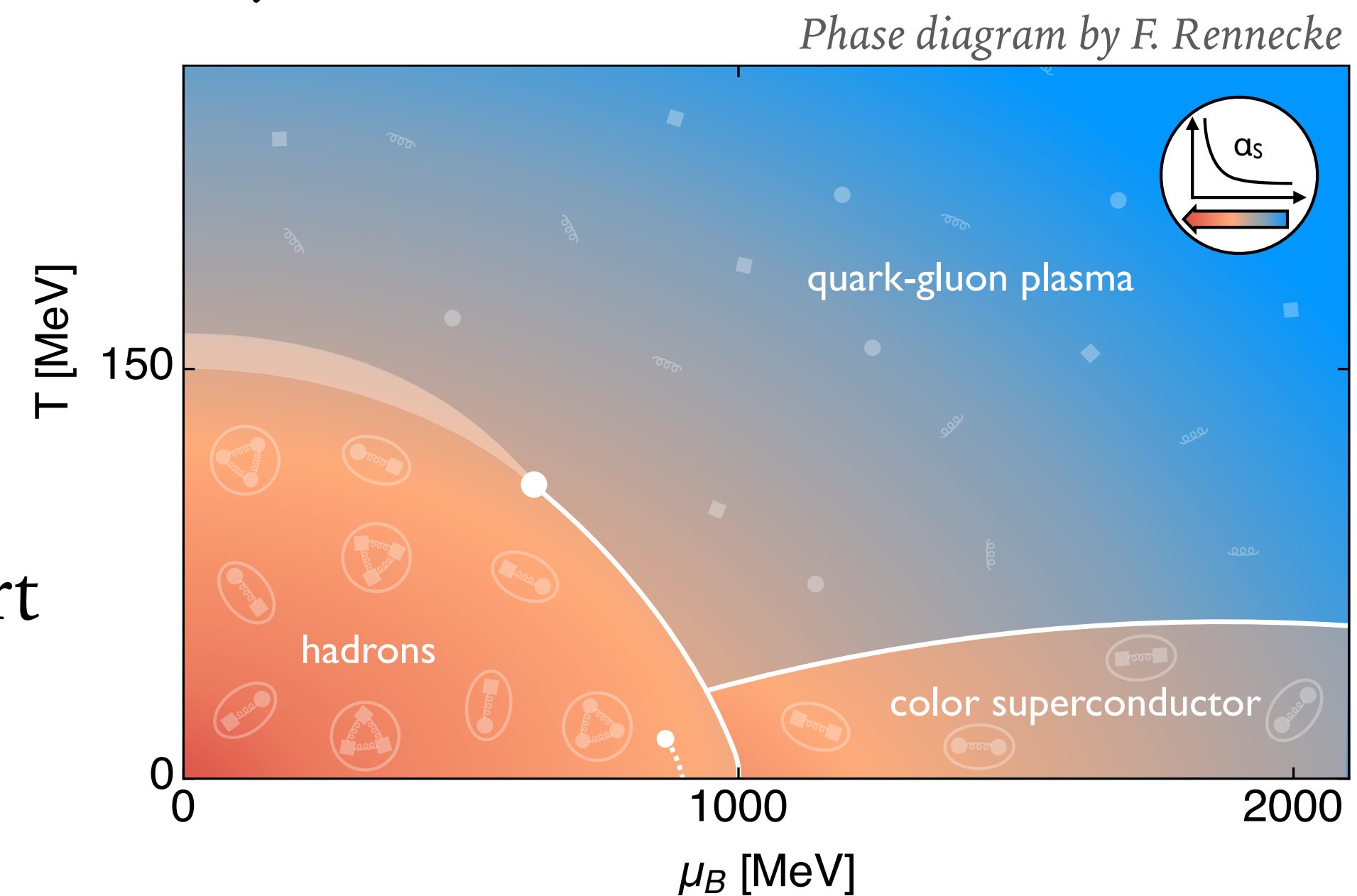


QCD Critical Point and Hydrodynamic Evolution at the  
Yukawa Institute for Theoretical Physics

# Phase Diagram (Sketch)

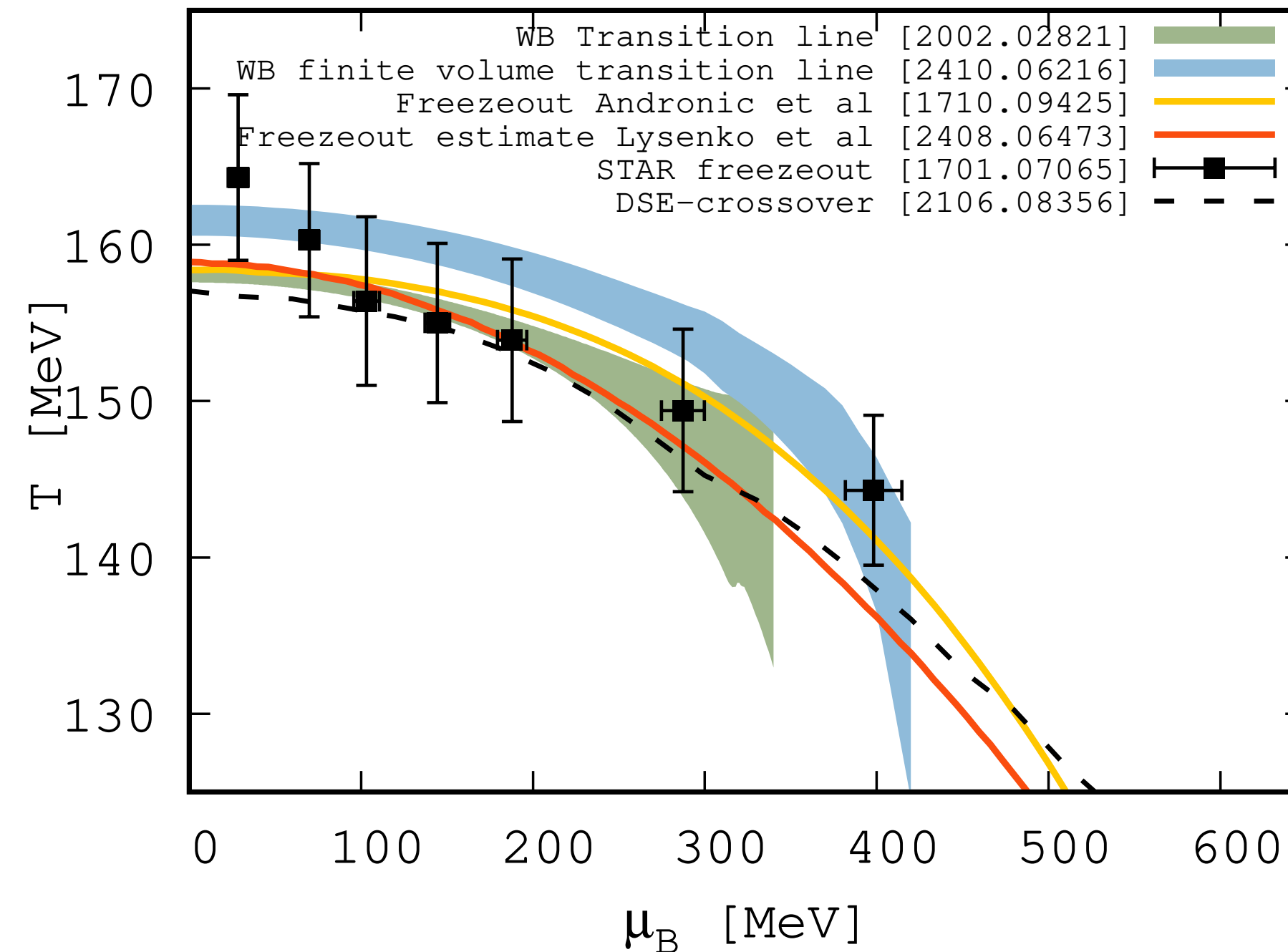


- Knowledge of the equation of state of strongly-interacting matter in equilibrium is crucial for:
  - Fluctuations, via derivatives of the pressure
  - The hadronic spectrum, i.e. the composition of the system in HICs, via thermal models
  - Hydrodynamic simulations
  - Hadronic transport simulations
  - Merger simulations
  - The behavior of the bulk viscosity & transport
  - The interior composition of neutron stars
  - ...



# (Less Sketchy) Phase Diagram

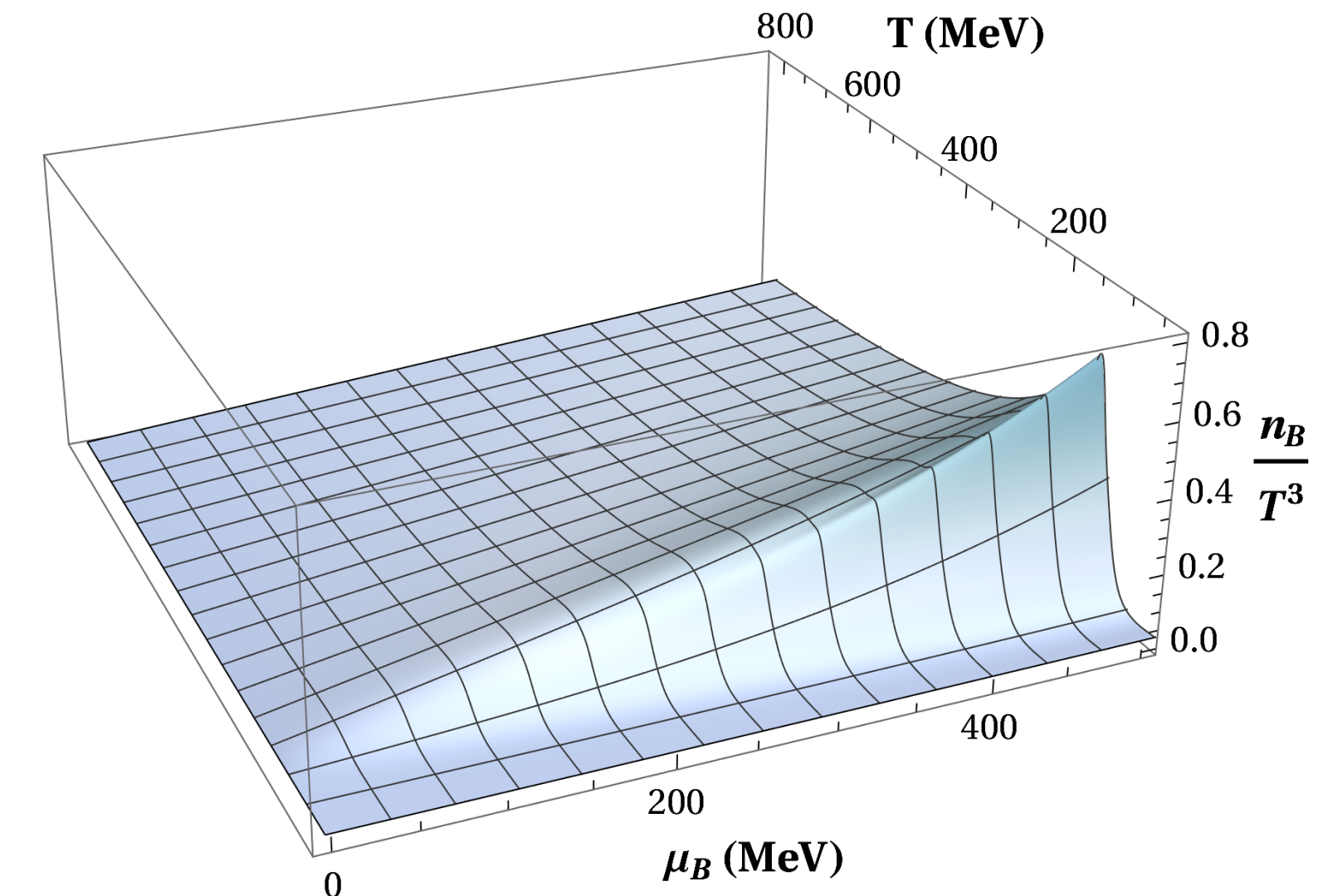
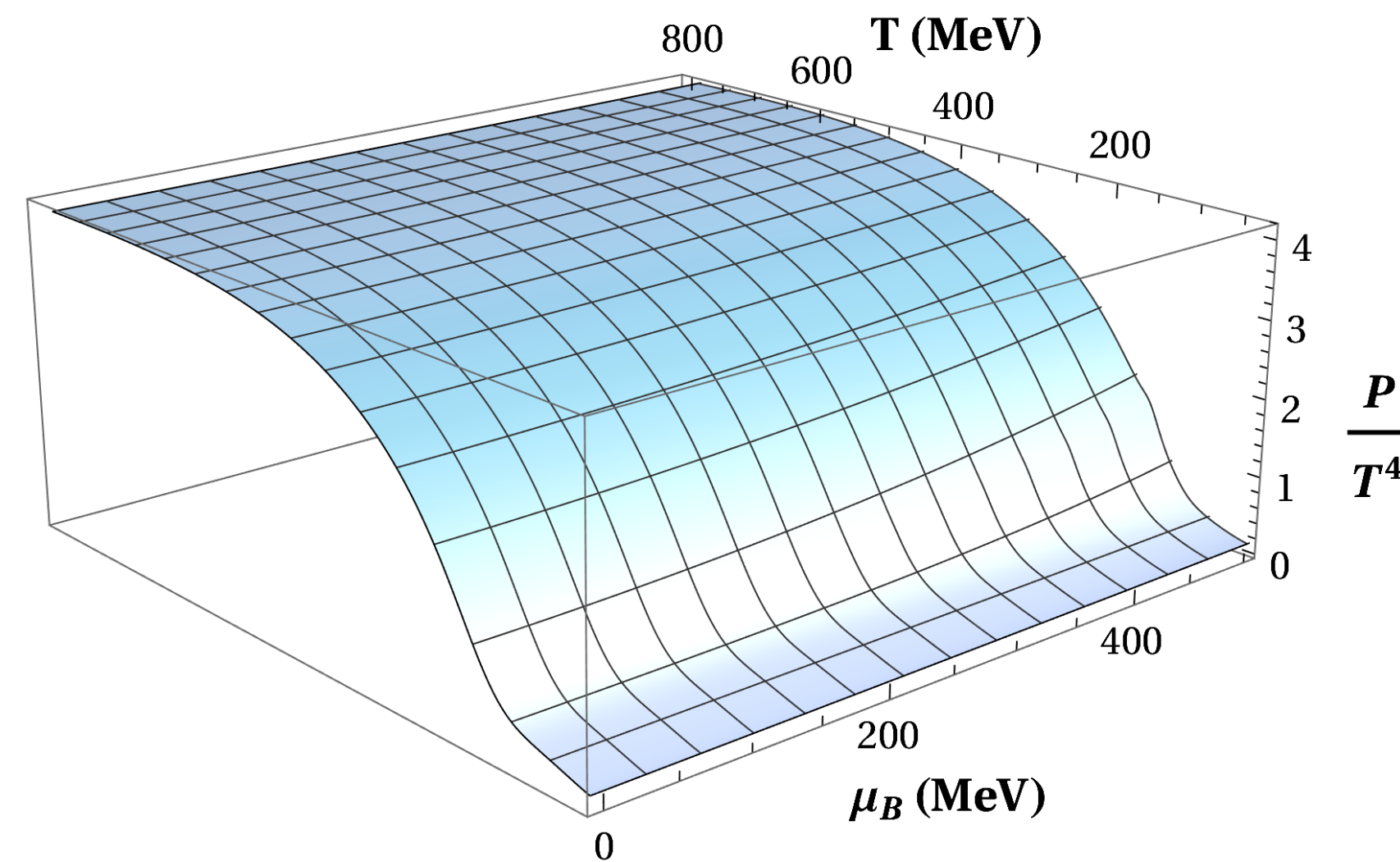
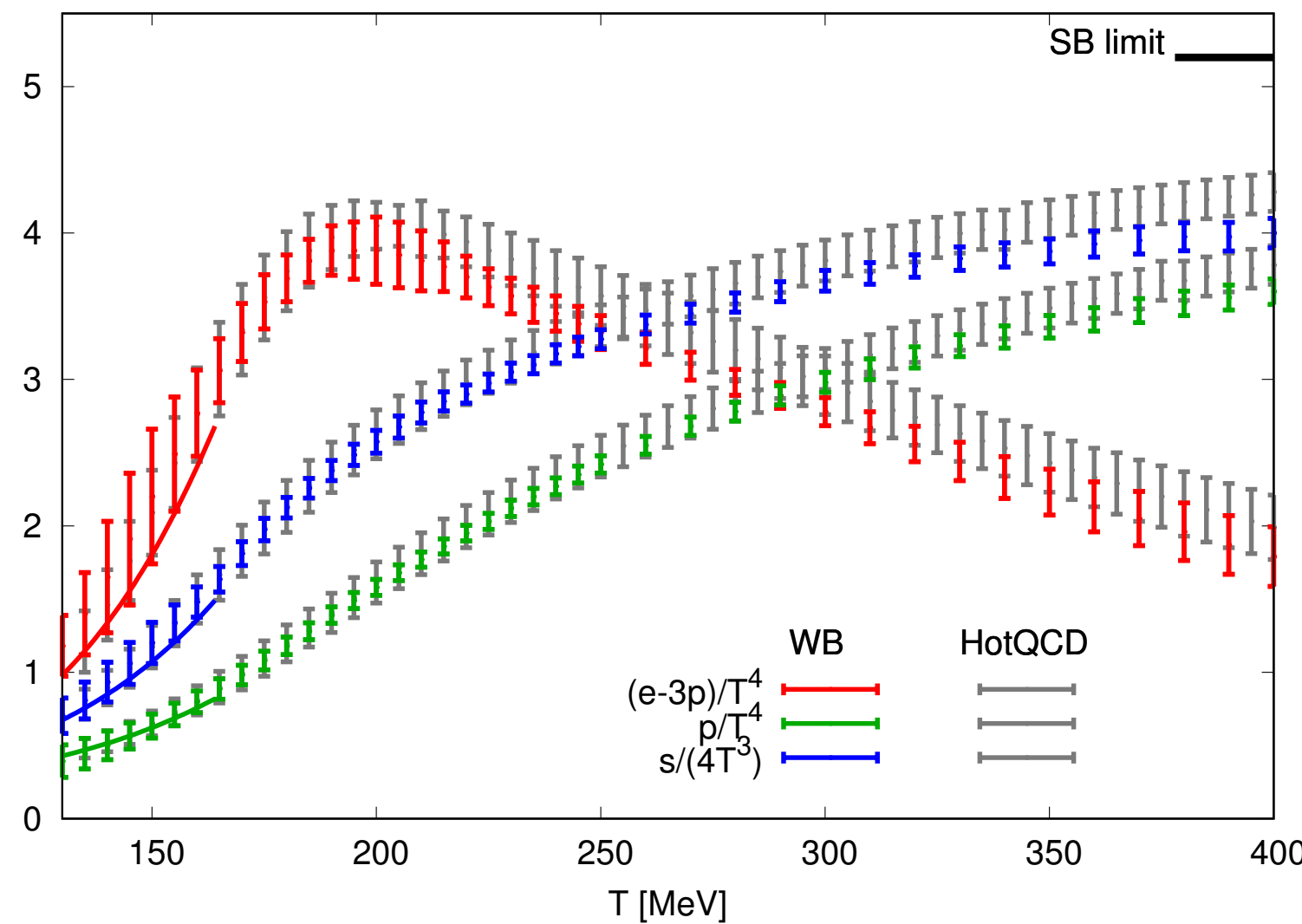
- Much of the current knowledge of the QCD phase diagram from *ab initio* theory and experiment from QCD crossover transition & freeze-out



# Lattice EoS at Finite $T$ & $\mu_B$

- Equilibrium thermodynamics calculated from first principles lattice QCD computations are well-established with good agreement amongst techniques

$$\frac{p(T, \mu_B)}{T^4} = \frac{p(T, 0)}{T^4} + \sum_{n=1}^{\infty} \frac{1}{(2n)!} \left. \frac{d^{2n}(p/T^4)}{d(\frac{\mu_B}{T})^{2n}} \right|_{\mu_B=0} \left(\frac{\mu_B}{T}\right)^{2n} = \sum_{n=1}^{\infty} \frac{1}{(2n)!} \chi_{2n}^B \left(\frac{\mu_B}{T}\right)^{2n}$$



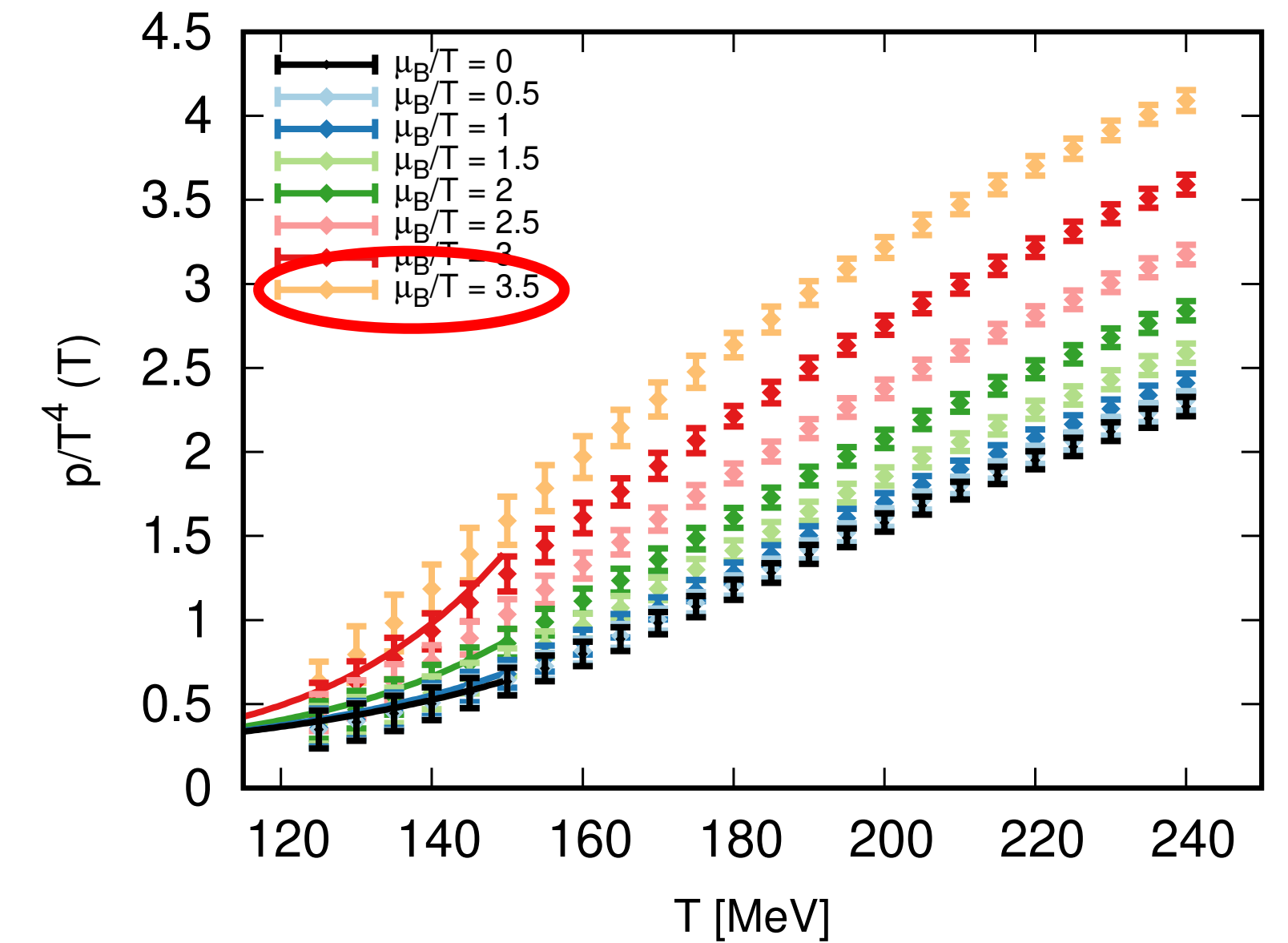
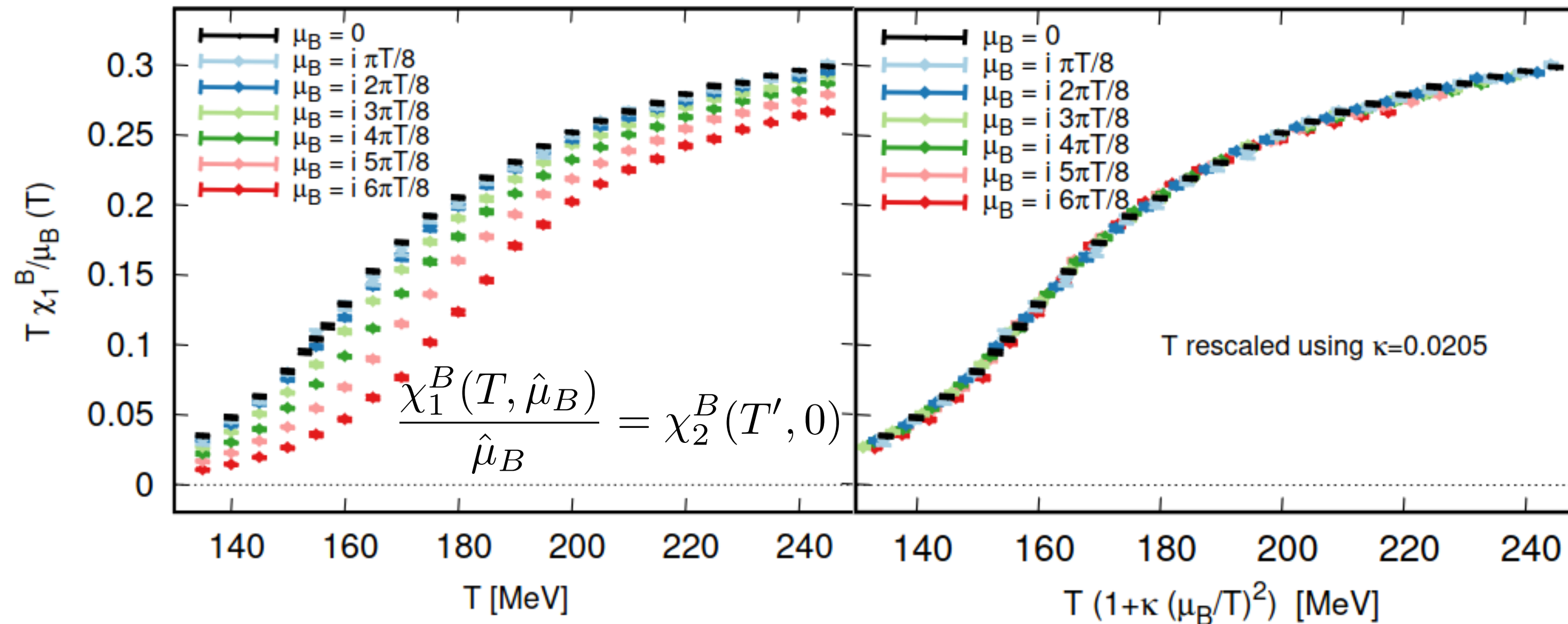
A. Bazavov PRD (2014), S. Borsanyi PLB (2014)

# Alternative Expansion Scheme: TExS

- Extend the coverage at finite density with the  $T'$ -Expansion Scheme (TExS) via a rescaled  $\mu_B$ -dependent temperature  $T'$

$$T'(T, \hat{\mu}_B) = T \left( 1 + \kappa_2^{BB}(T) \hat{\mu}_B^2 + \kappa_4^{BB}(T) \hat{\mu}_B^4 + \mathcal{O}(\hat{\mu}_B^6) \right)$$

$T \rightarrow T'$



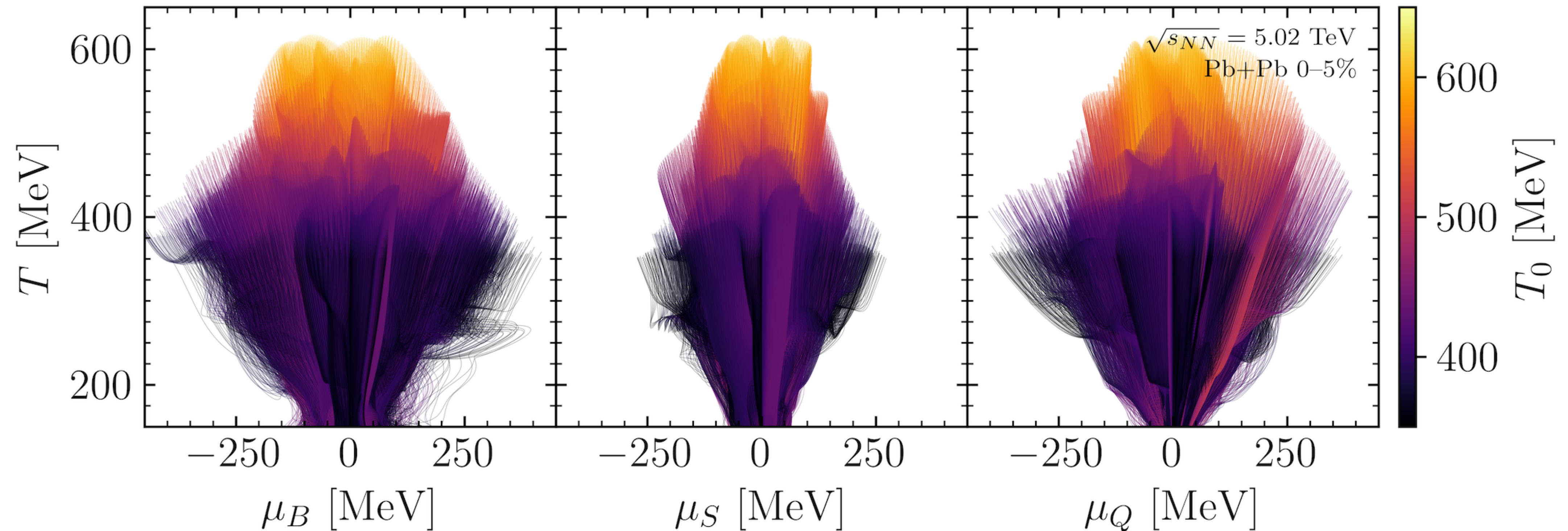
$T'$  expansion coefficients are reshuffling of Taylor terms:

$$\kappa_2^{BB}(T) = \frac{1}{6T} \frac{\chi_4^B(T)}{\chi_2^{B'}(T)} \quad \kappa_4^{BB}(T) = \frac{1}{360 \chi_2^{B'}(T)^3} \left( 3 \chi_2^{B'}(T)^2 \chi_6^B(T) - 5 \chi_2^{B''}(T) \chi_4^B(T)^2 \right)$$

# BQS Fluctuations at the LHC



- Hydrodynamic fluid cell trajectories show the range across which all three conserved charge chemical potentials can fluctuate for PbPb collisions at the LHC



*C. Plumberg, D. Almaalol, T. Dore et al, 2405.09648*

I. Lattice QCD-based equations of state

II. A first study of equilibrium proton factorial cumulants with MaxEnt

# I. Lattice QCD-based equations of state

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# Equation of State with Three Conserved Charges



- During HICs the system is not only confined to the  $T$ - $\mu_B$  plane: determine the equations that depend on  $\mu_B, \mu_Q, \mu_S$

$$\frac{P(T, \mu_B, \mu_Q, \mu_S)}{T^4} = \sum_{i,j,k} \frac{1}{i!j!k!} \chi_{ijk}^{BQS}(T) \left(\frac{\mu_B}{T}\right)^j \left(\frac{\mu_Q}{T}\right)^k \left(\frac{\mu_S}{T}\right)^i$$

where:

$$\chi_{ijk}^{BQS}(T) = \left. \frac{\partial^{i+j+k}(p/T^4)}{\partial(\frac{\mu_B}{T})^i \partial(\frac{\mu_Q}{T})^j \partial(\frac{\mu_S}{T})^k} \right|_{\mu_B, \mu_Q, \mu_S=0}$$

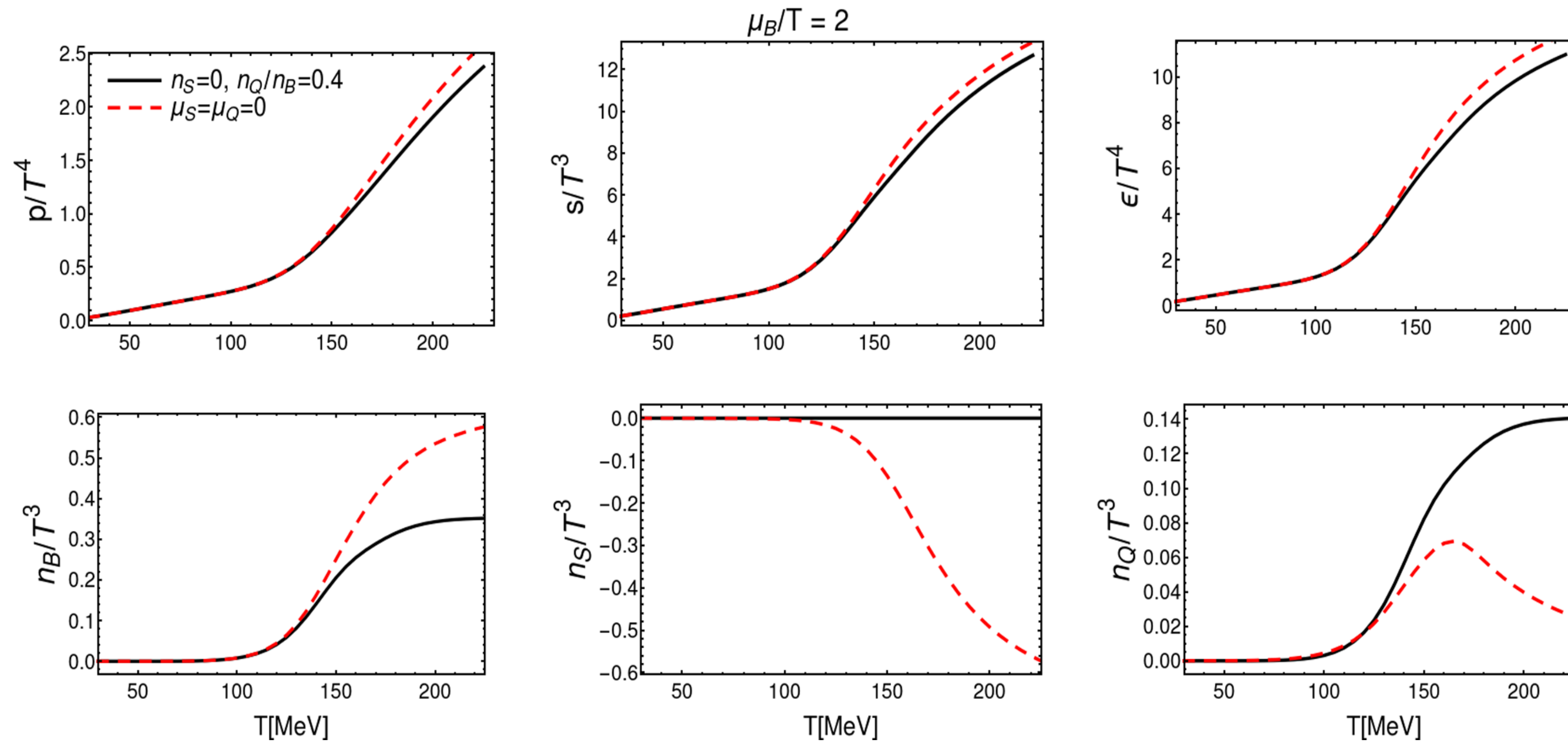
Lattice results only between  $T \sim 135 - 220$  MeV for all 22 coefficients

- Utilize HRG for low  $T$
- Impose Stefan-Boltzmann limit at high  $T$

# EoS with Conserved Charge Constraints



- Observe the effect of imposing strangeness neutrality and a fixed ratio of baryon number to electric charge as expected for the densities

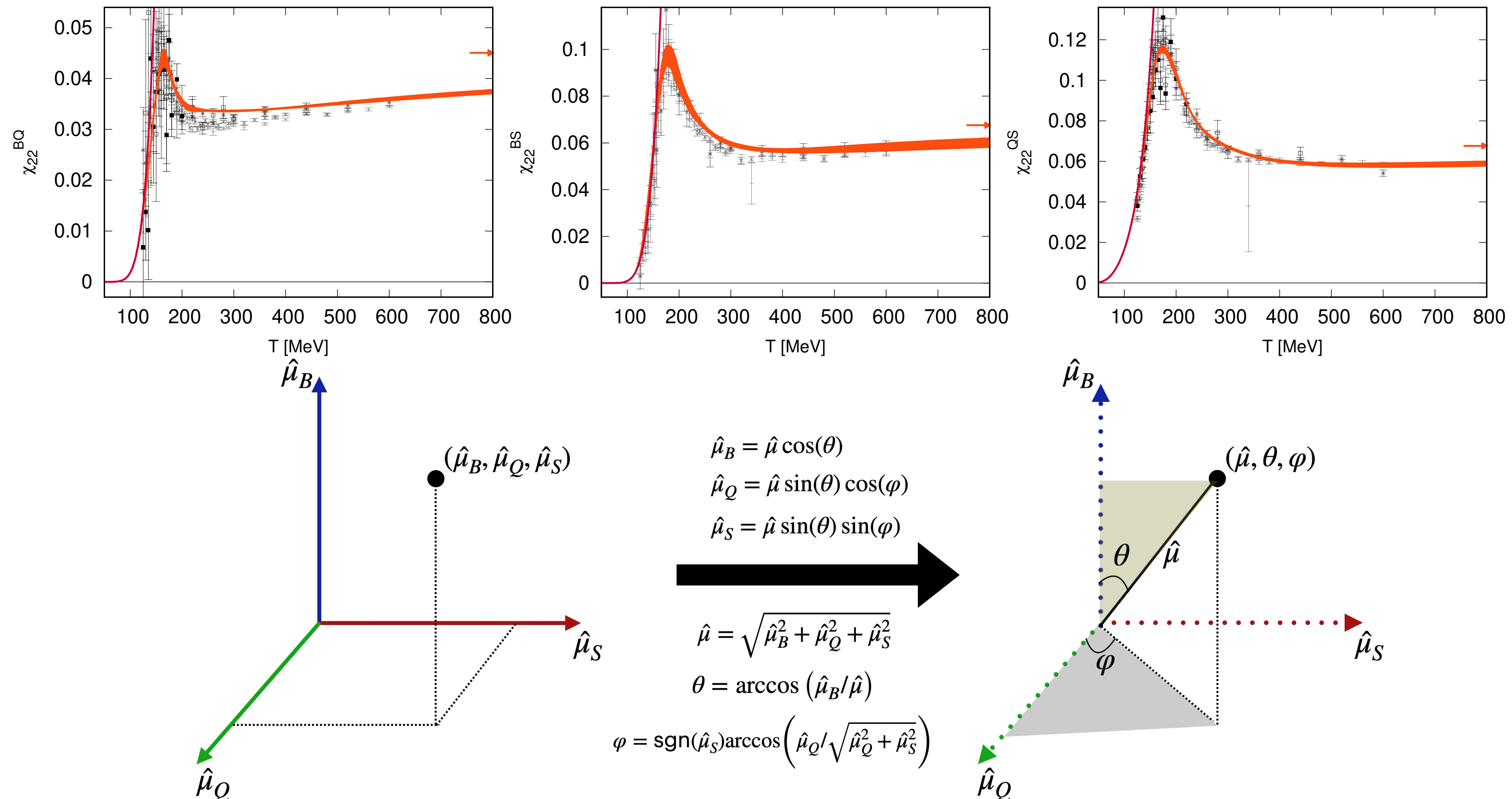


See also: A. Monnai et al, PRC (2019)

J. Noronha-Hostler, JS=JMK et al, PRC (2019)

# New Lattice QCD EoS in $T-\mu_B-\mu_Q-\mu_S$ : 4D TExS

- New continuum estimated lattice results in full 4D phase diagram of temperature and conserved charges by transforming to spherical coordinates



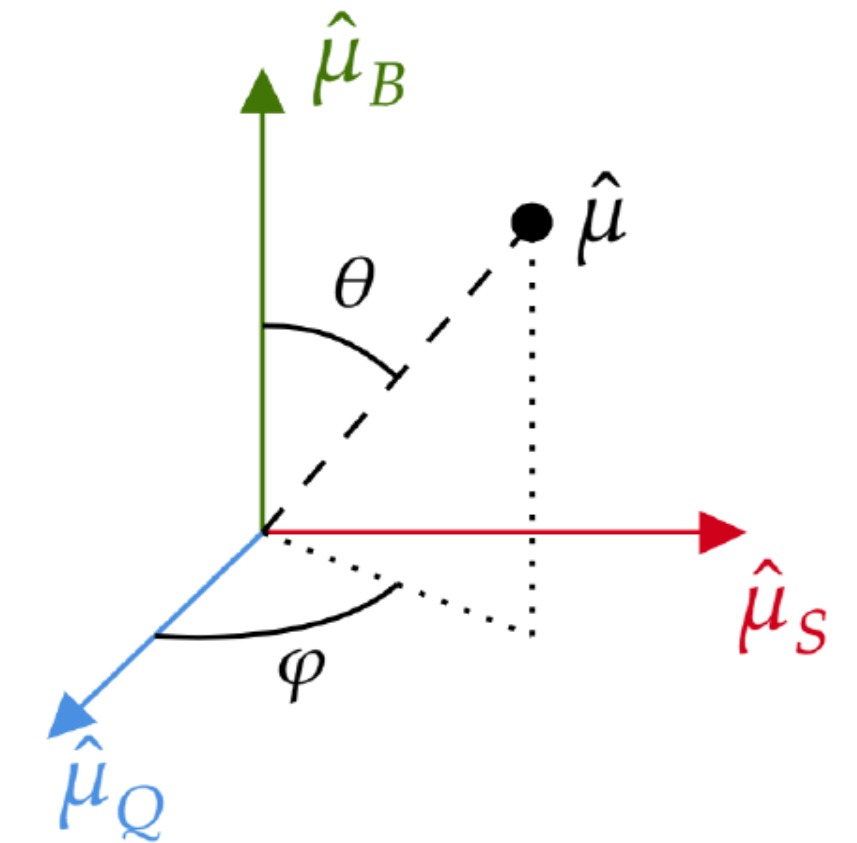
A. Abuali et al, PRD (2025)

# New Lattice QCD EoS in $T-\mu_B-\mu_Q-\mu_S$ : 4D TExS

- Extended coverage via TExS along lines of constant  $\mu/T$  at fixed  $(\theta, \phi)$  and define the generalized susceptibilities

$$X_2^{\theta,\phi}(T) = c_\theta^2 \cdot \chi_2^B(T) + s_\theta^2 c_\phi^2 \cdot \chi_2^Q(T) + s_\theta^2 s_\phi^2 \cdot \chi_2^S(T) + \dots$$

$$X_4^{\theta,\phi}(T) = c_\theta^4 \cdot \chi_4^B(T) + s_\theta^4 c_\phi^4 \cdot \chi_4^Q(T) + s_\theta^4 s_\phi^4 \cdot \chi_4^S(T) + \dots$$



to determine the expansion with constrained high  $T$  by approach to SB limit with coefficient  $\lambda$ :

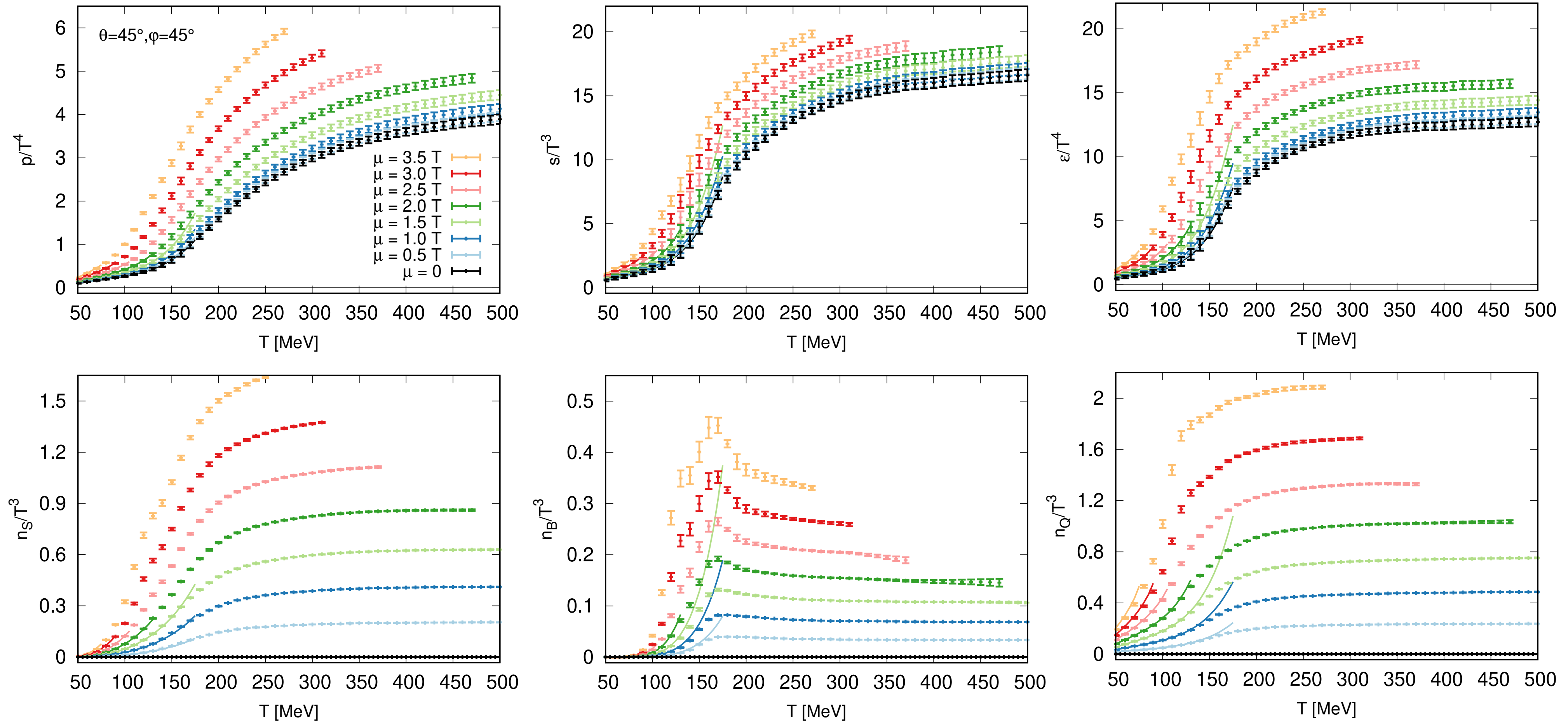
$$T'_F(T, \hat{\mu}_B) = T (1 + \lambda_{2,F}(T) \hat{\mu}_B^2 + \dots)$$

$$\lambda_2^{\theta,\phi}(T) = \frac{1}{6T} \frac{1}{X_2^{\theta,\phi}(T)} \times \left( X_4^{\theta,\phi}(T) - \frac{\bar{X}_4^{\theta,\phi}(0)}{\bar{X}_2^{\theta,\phi}(0)} X_2^{\theta,\phi}(T) \right)$$

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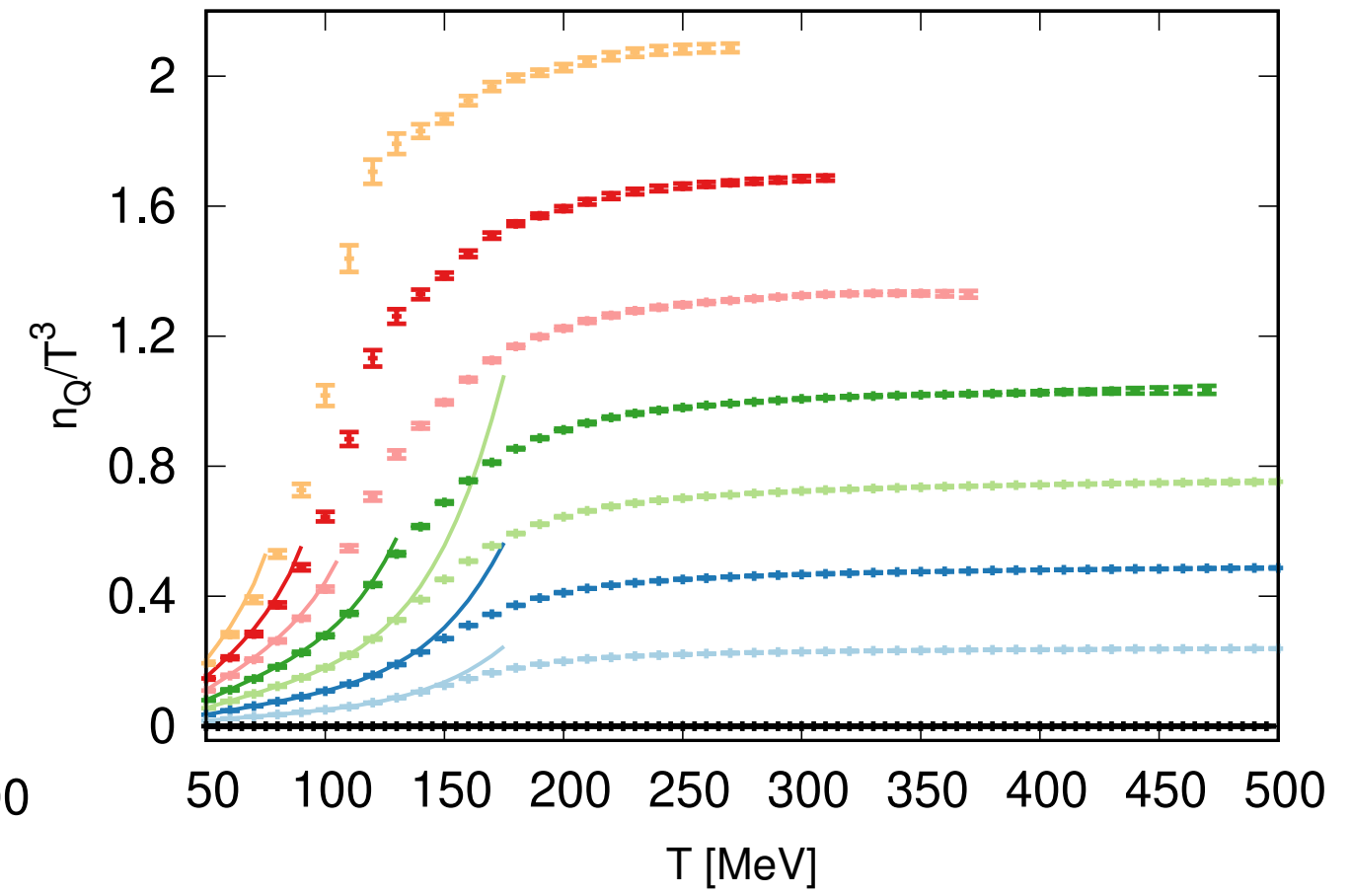
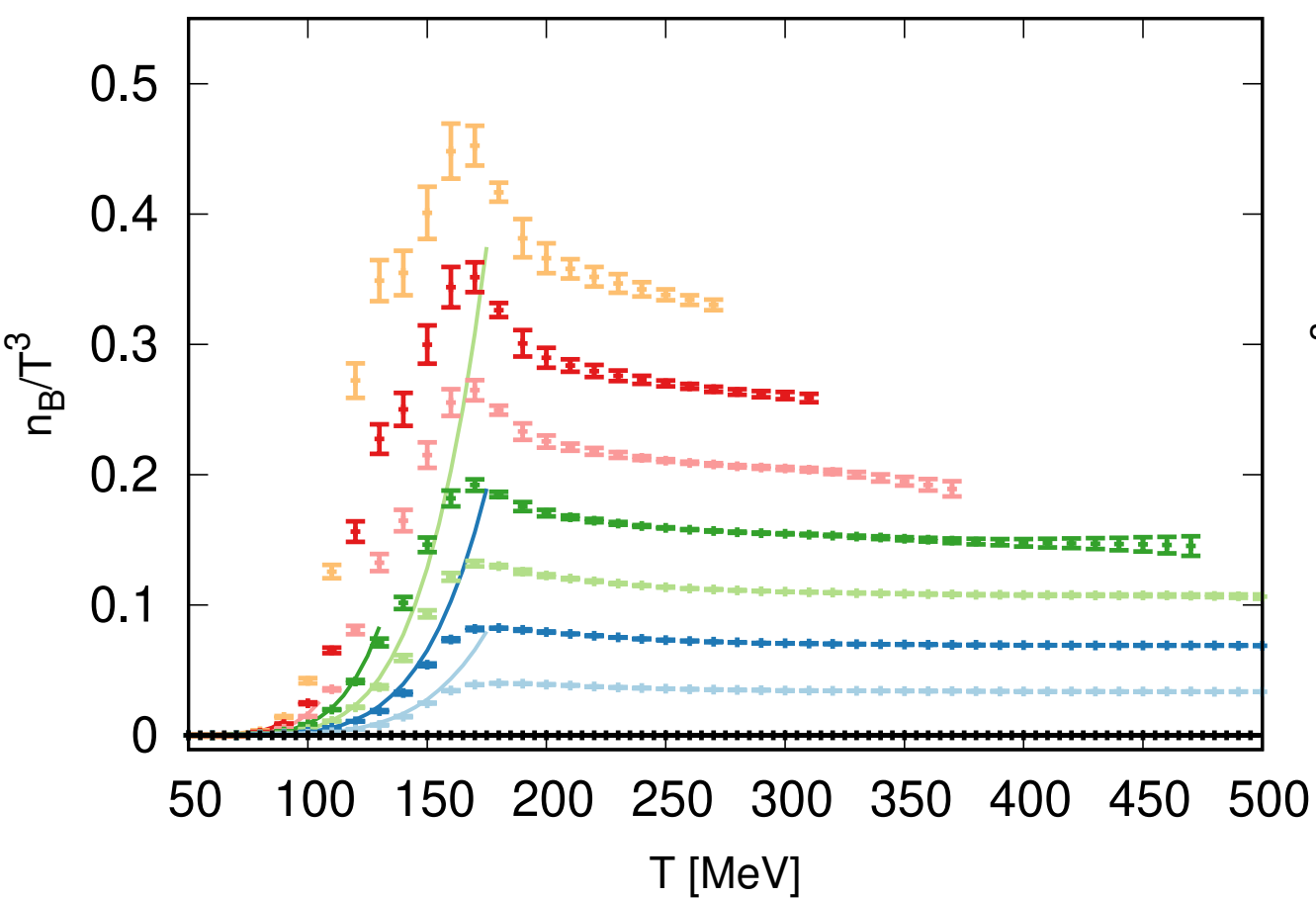
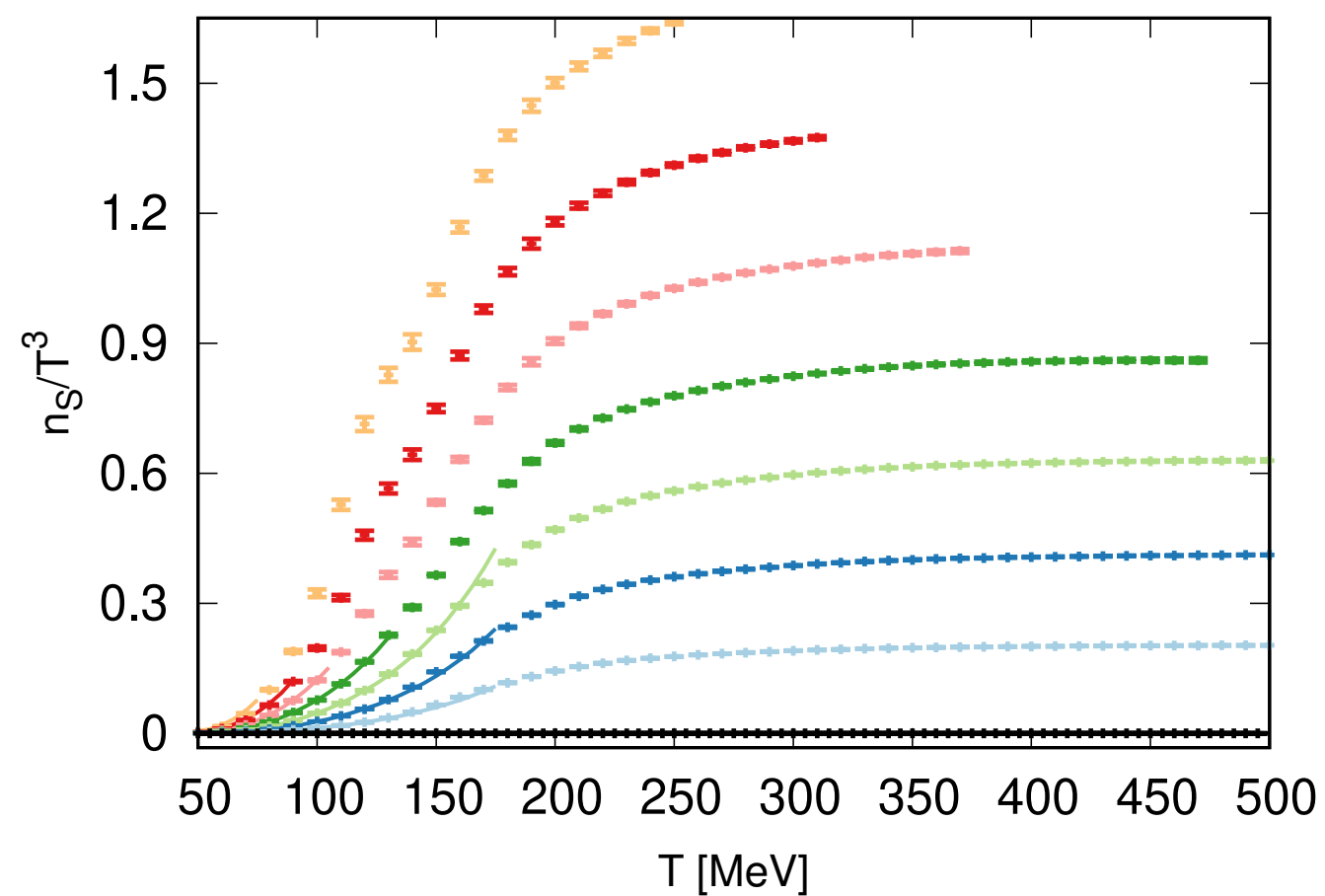
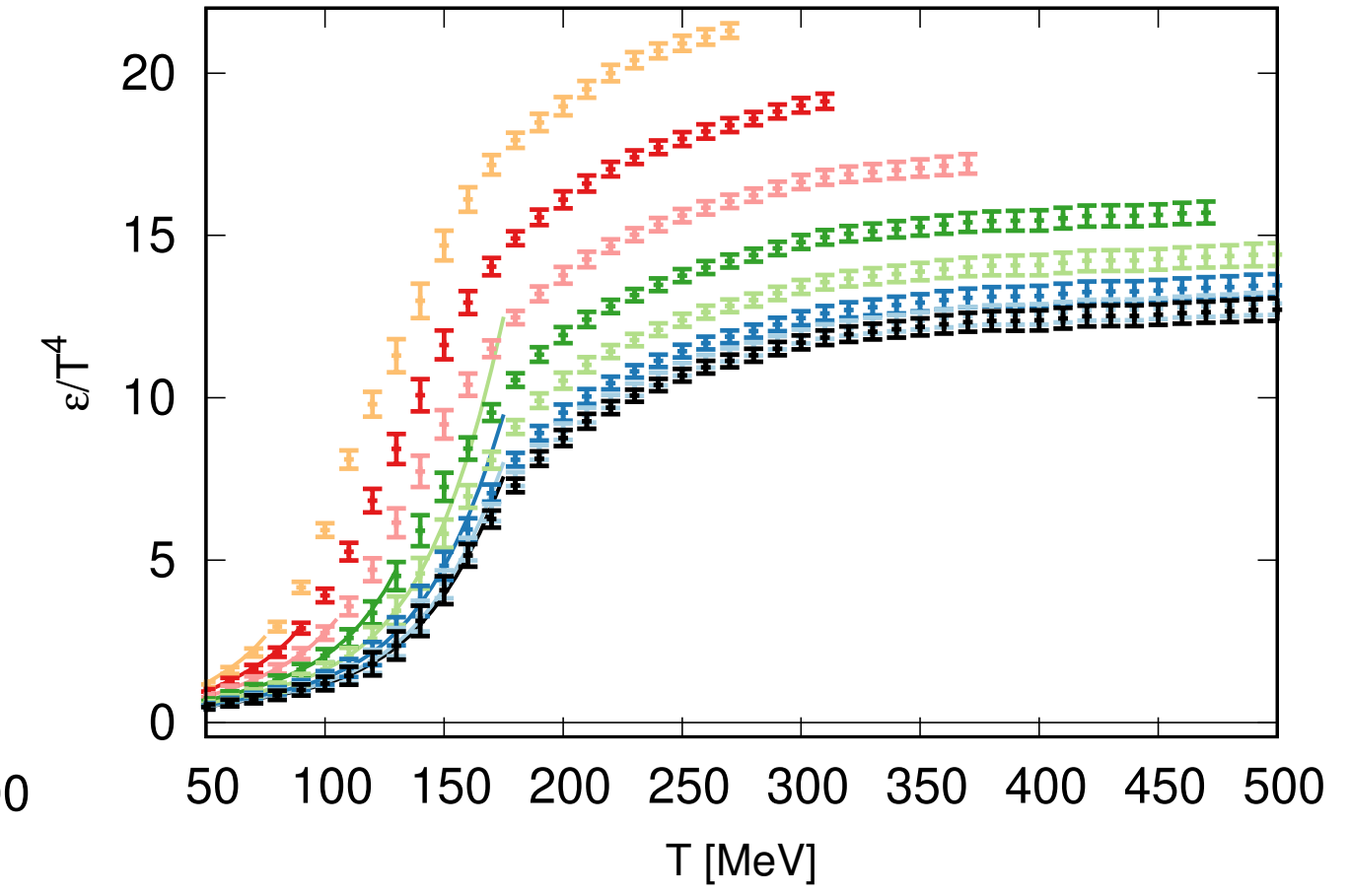
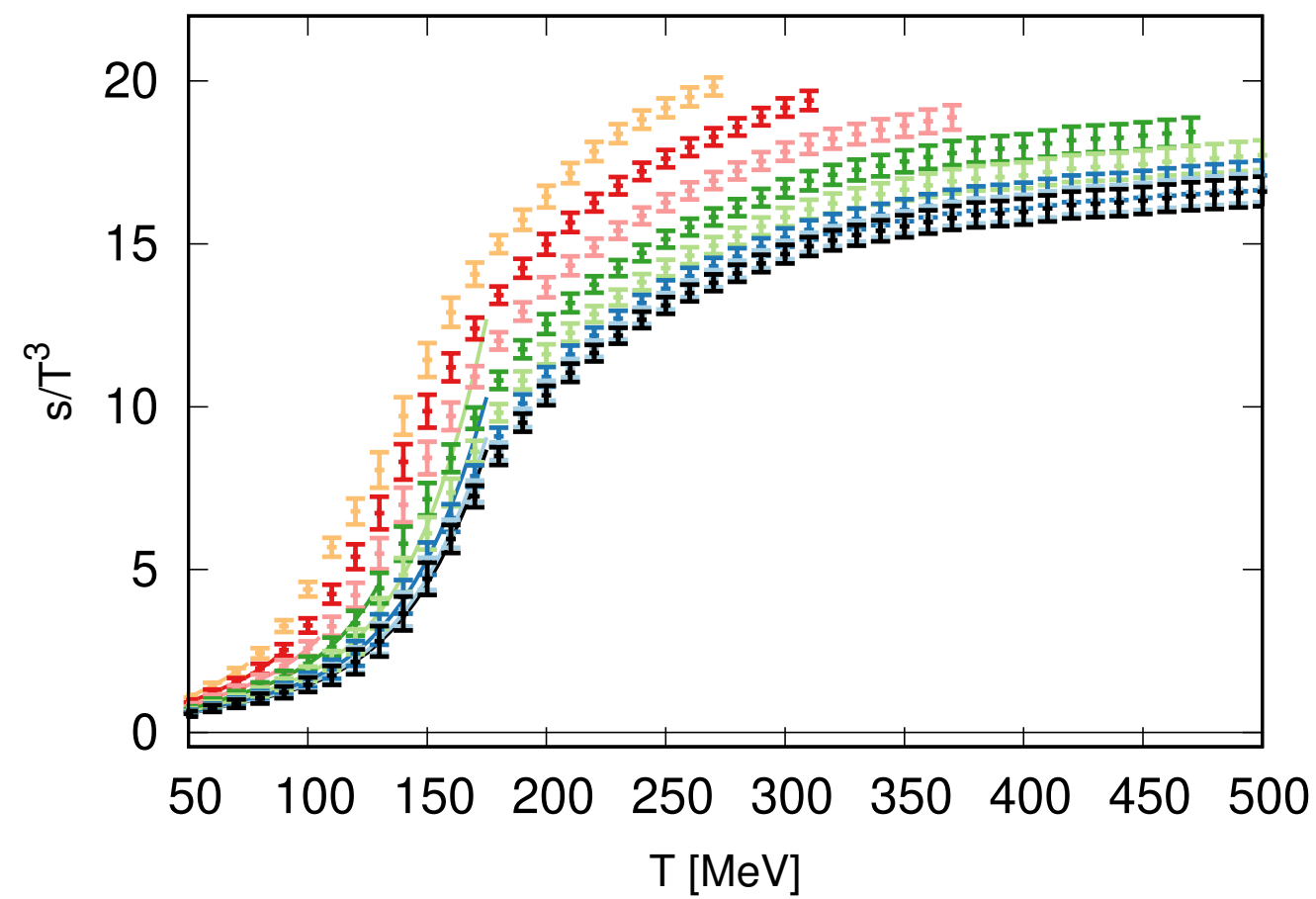
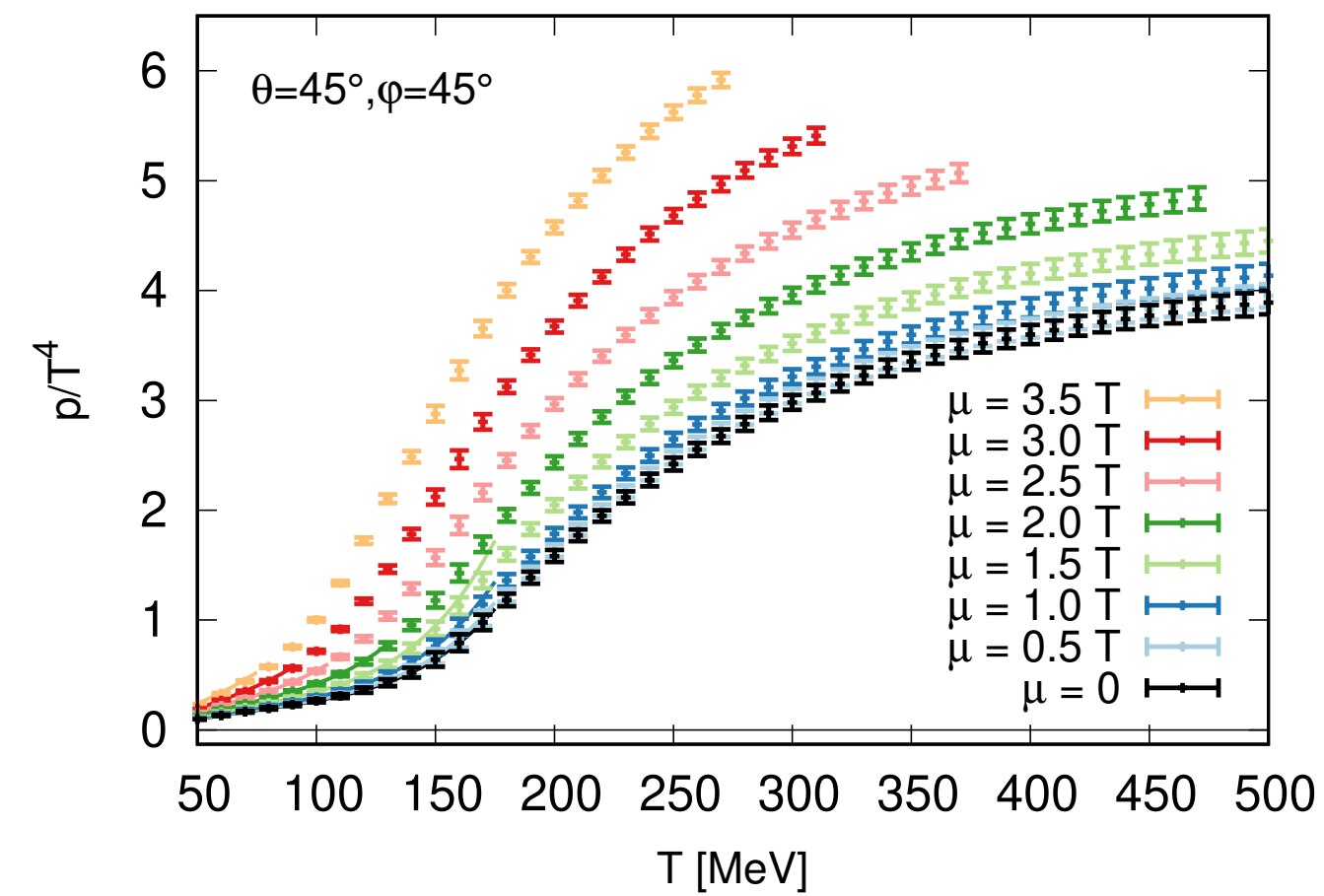


*A. Abuali et al, PRD (2025)*

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What about the critical point?

*A. Abuali et al, PRD (2025)*

- Original work from the BEST collaboration to implement critical features into EoS for heavy-ion collisions:

$$P(\mu, T) = P^{\text{reg}}(\mu, T) + P^{\text{sing}}(\mu, T)$$

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Matched to  
lattice QCD  
Taylor coefficients

3D Ising model

$$T^4 c_n^{\text{LAT}}(T) = T^4 c_n^{\text{Non-Ising}}(T) + c_n^{\text{Ising}}(T)$$

# BEST Critical EoS



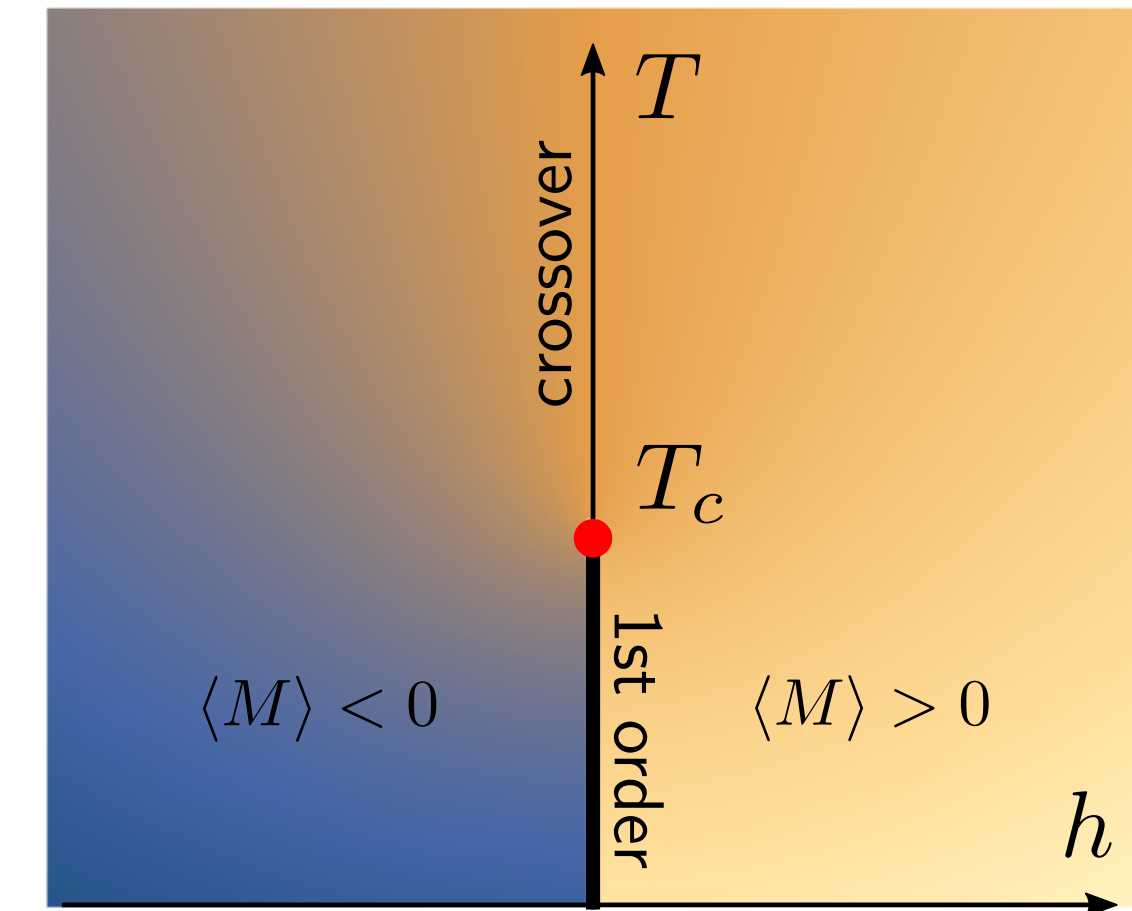
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$$T^4 c_n^{\text{LAT}}(T) = T^4 c_n^{\text{Non-Ising}}(T) + c_n^{\text{Ising}}(T)$$



$$h = h_0 R^{\beta\delta} H(\theta)$$

$$r = \frac{T - T_c}{T_c} = R(1 - \theta^2)$$

$$G = h_0 M_0 R^{2-\alpha} [g(\theta) - \theta H(\theta)] = -P$$

*P. Parotto et al, PRC (2020)*

*Fig: A. Bzdak et al, Phys.Rept. (2020)*

# BEST Critical EoS

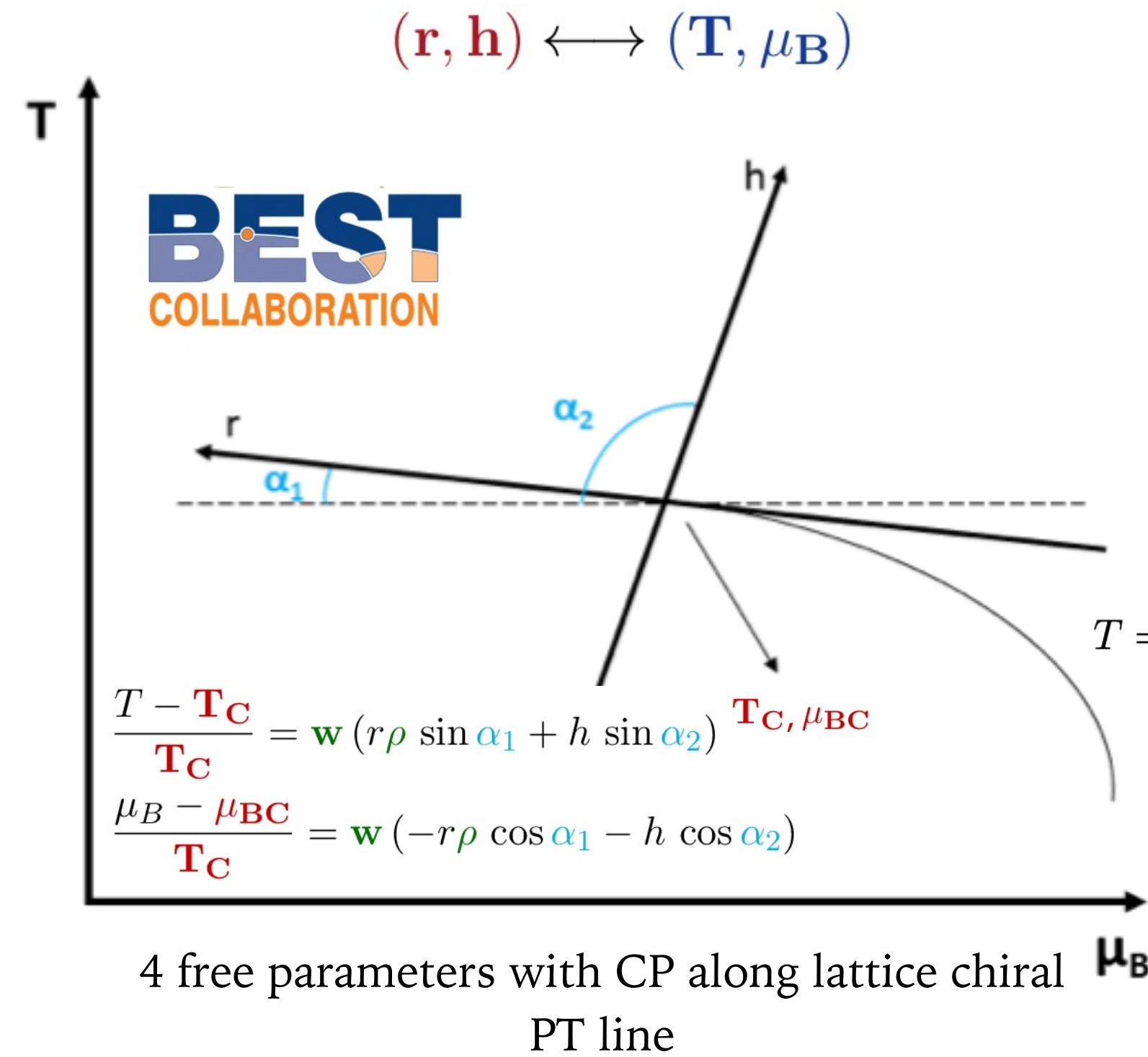


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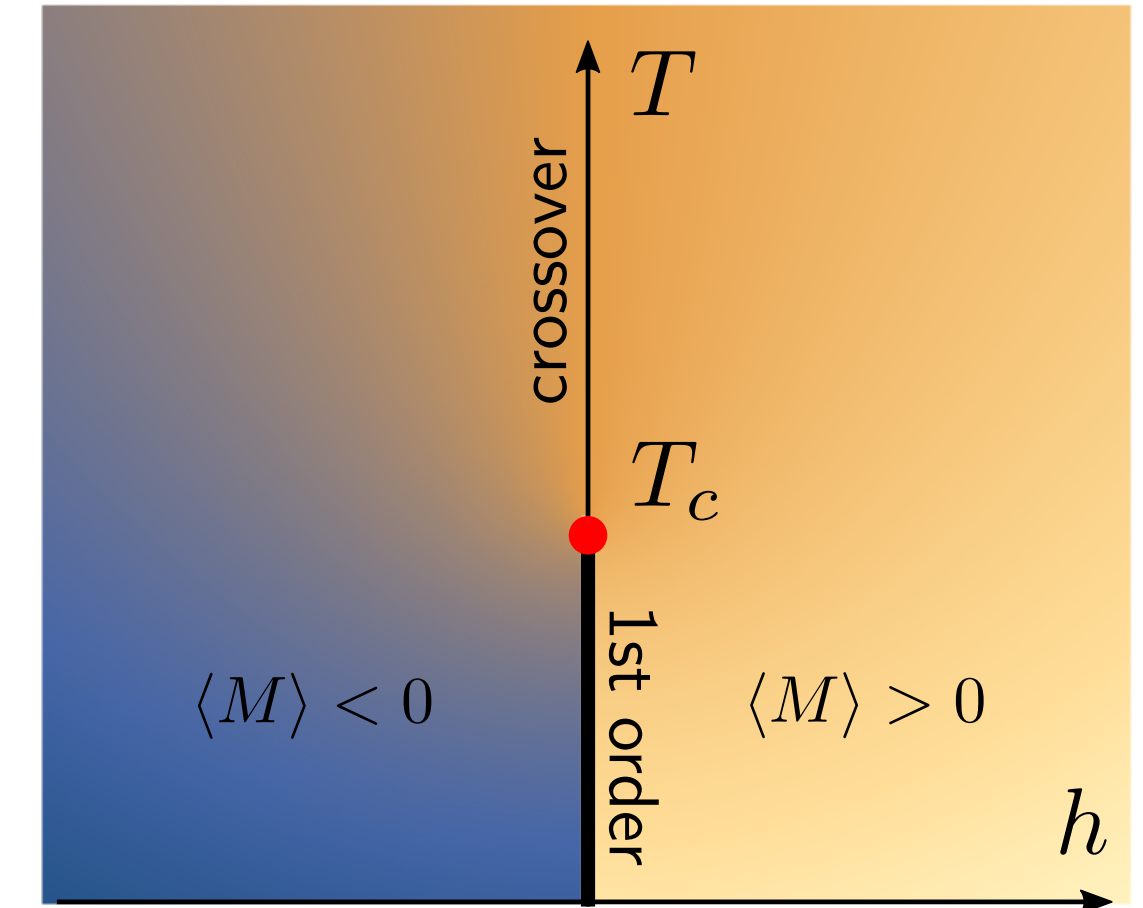
3D Ising model



$$T^4 c_n^{\text{LAT}}(T) = T^4 c_n^{\text{Non-Ising}}(T) + c_n^{\text{Ising}}(T)$$

$$T = T_0 + \kappa T_0 \left( \frac{\mu_B}{T_0} \right)^2 + O(\mu_B^4)$$

$$\alpha_1 = \tan^{-1} \left( 2 \frac{\kappa}{T_0} \mu_{BC} \right)$$



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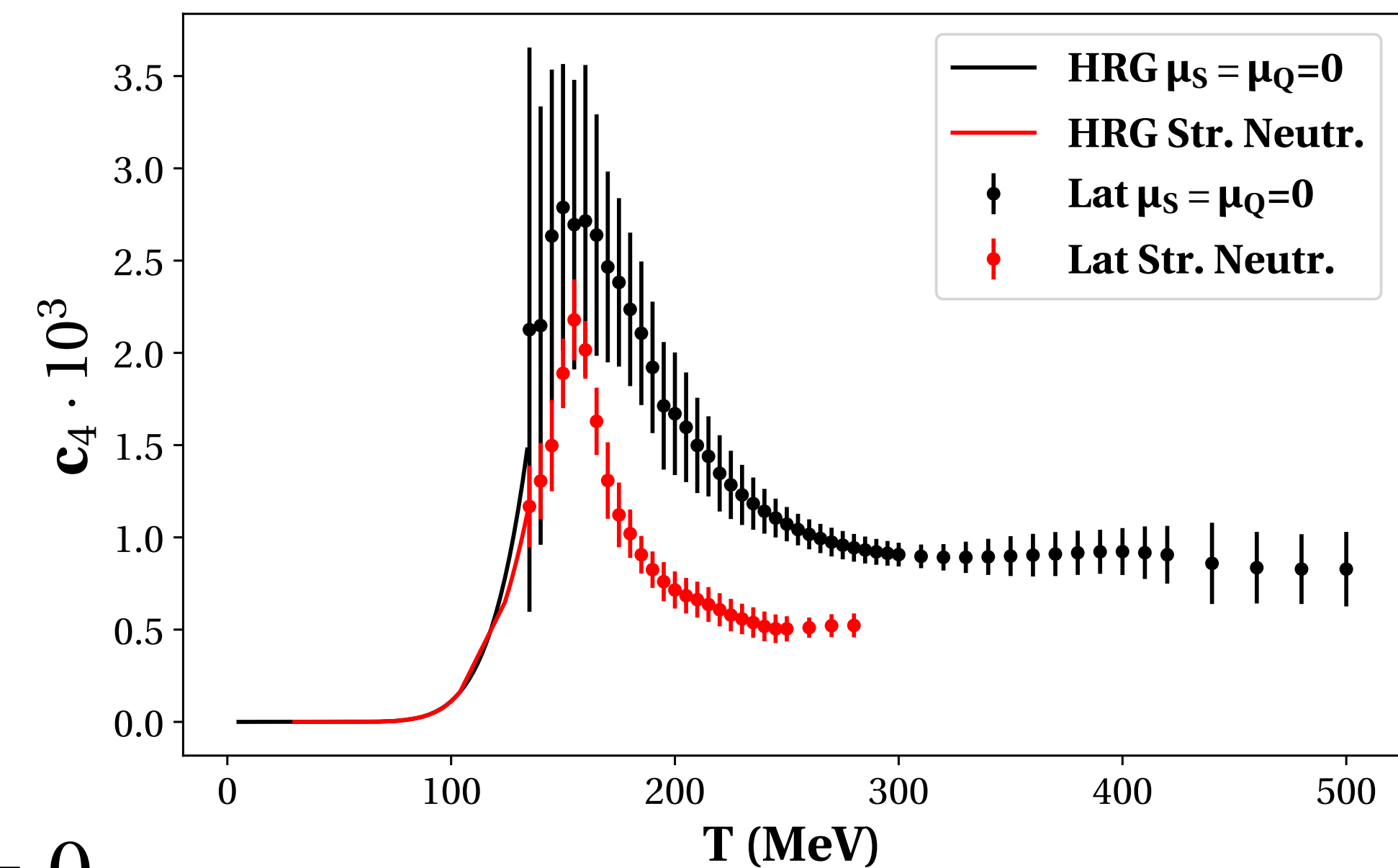
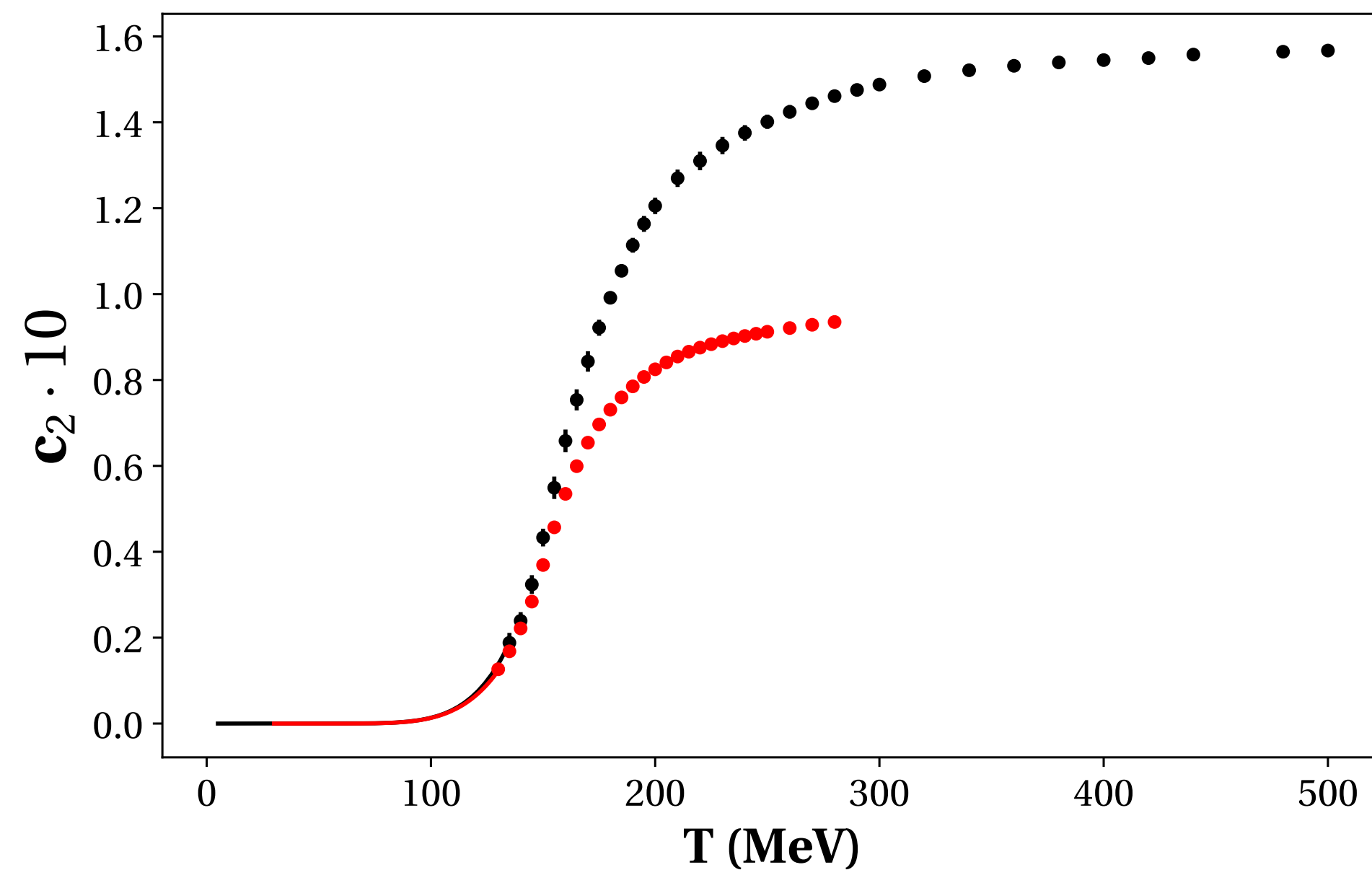
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*P. Parotto et al, PRC (2020)*  
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# Taylor Coefficients from LQCD

- Lattice results for Taylor expansion of pressure around  $\mu_B = 0$  up to  $\mathcal{O}(\mu_B^4)$  are the backbone of the procedure for creating this equation of state

$$\frac{P(T, \mu_B)}{T^4} = \sum_n c_{2n}(T) \left(\frac{\mu_B}{T}\right)^{2n}$$



- $\mu_Q = \mu_S = 0$
- strangeness-neutral

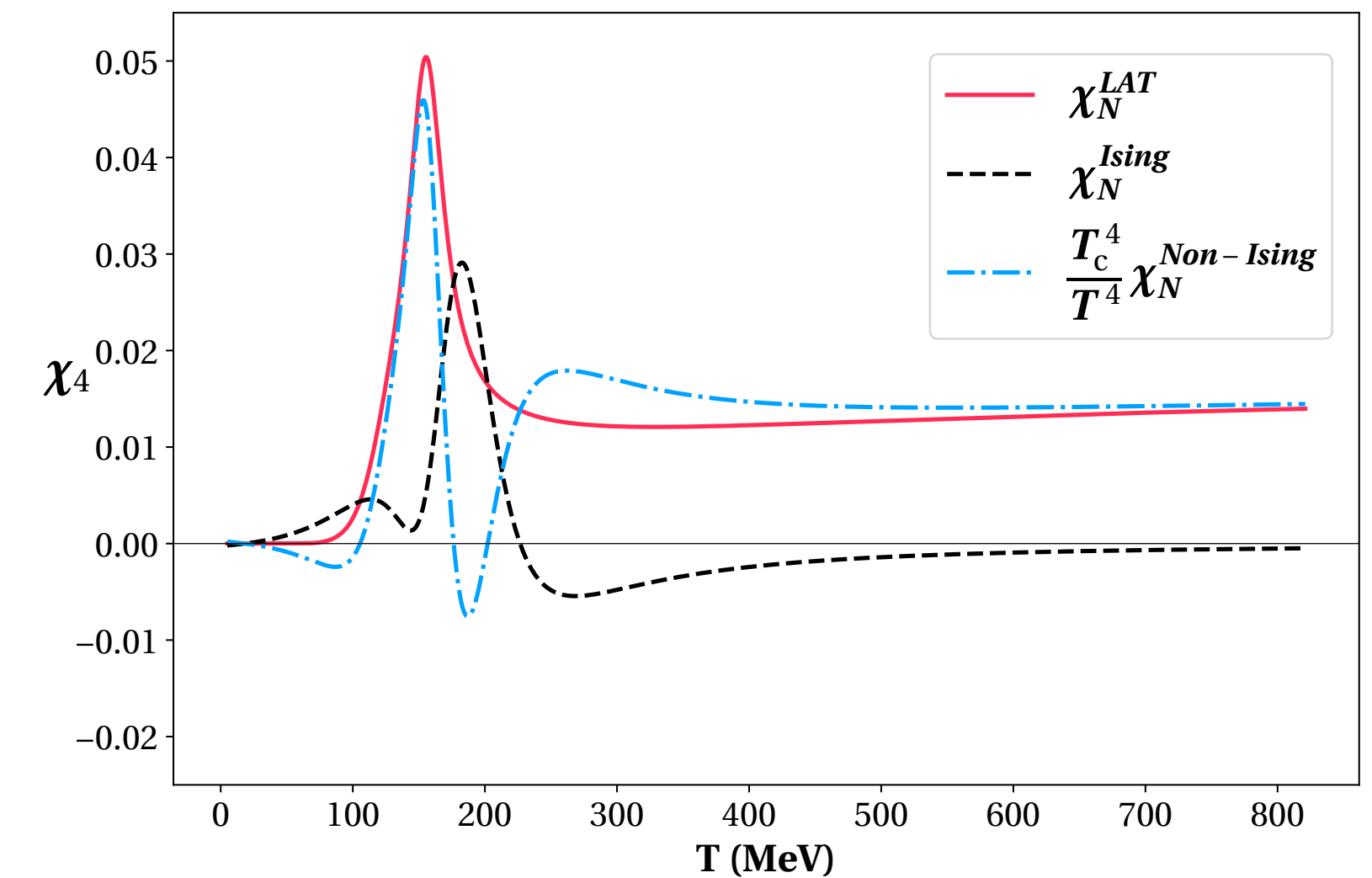
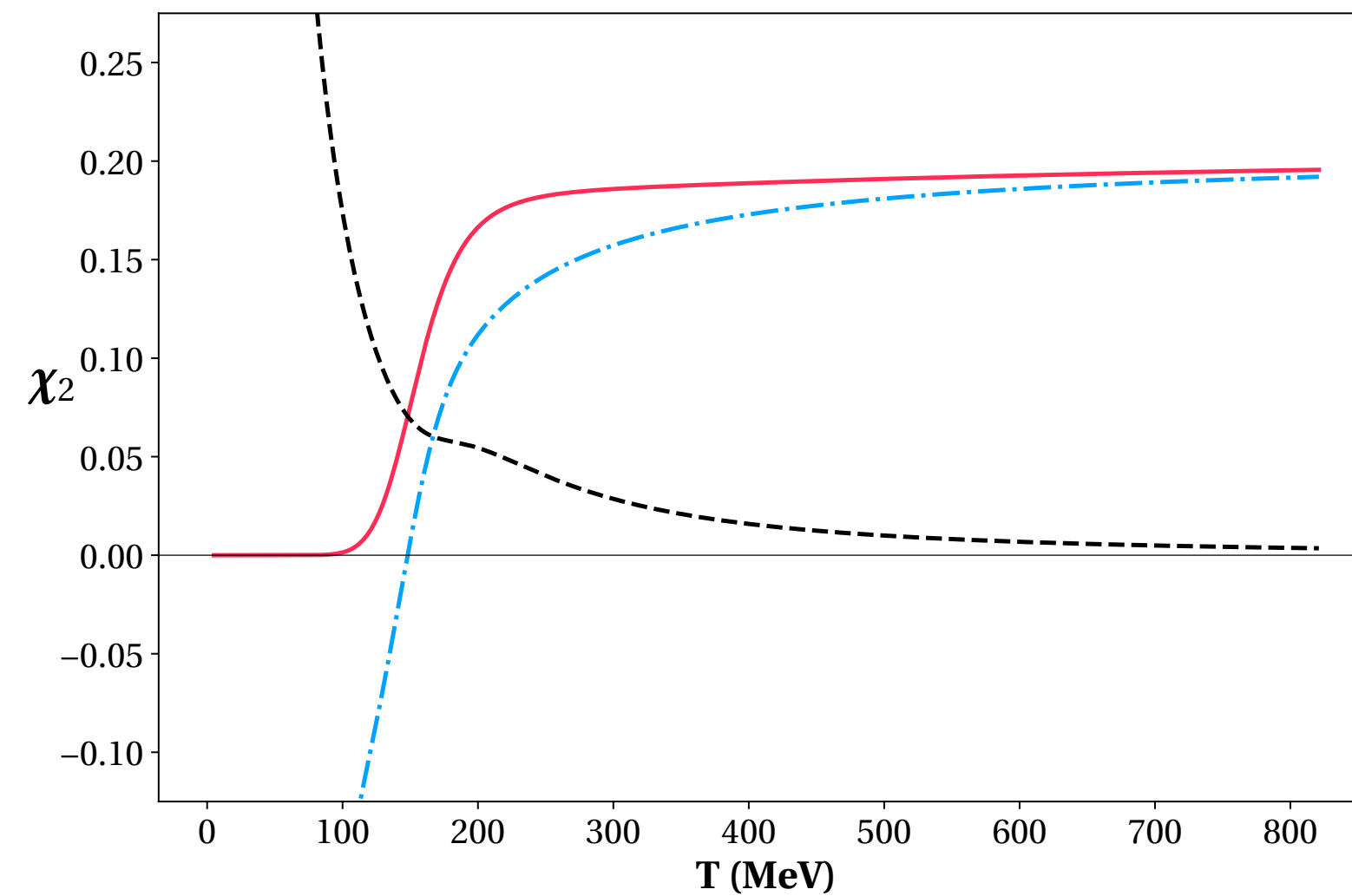
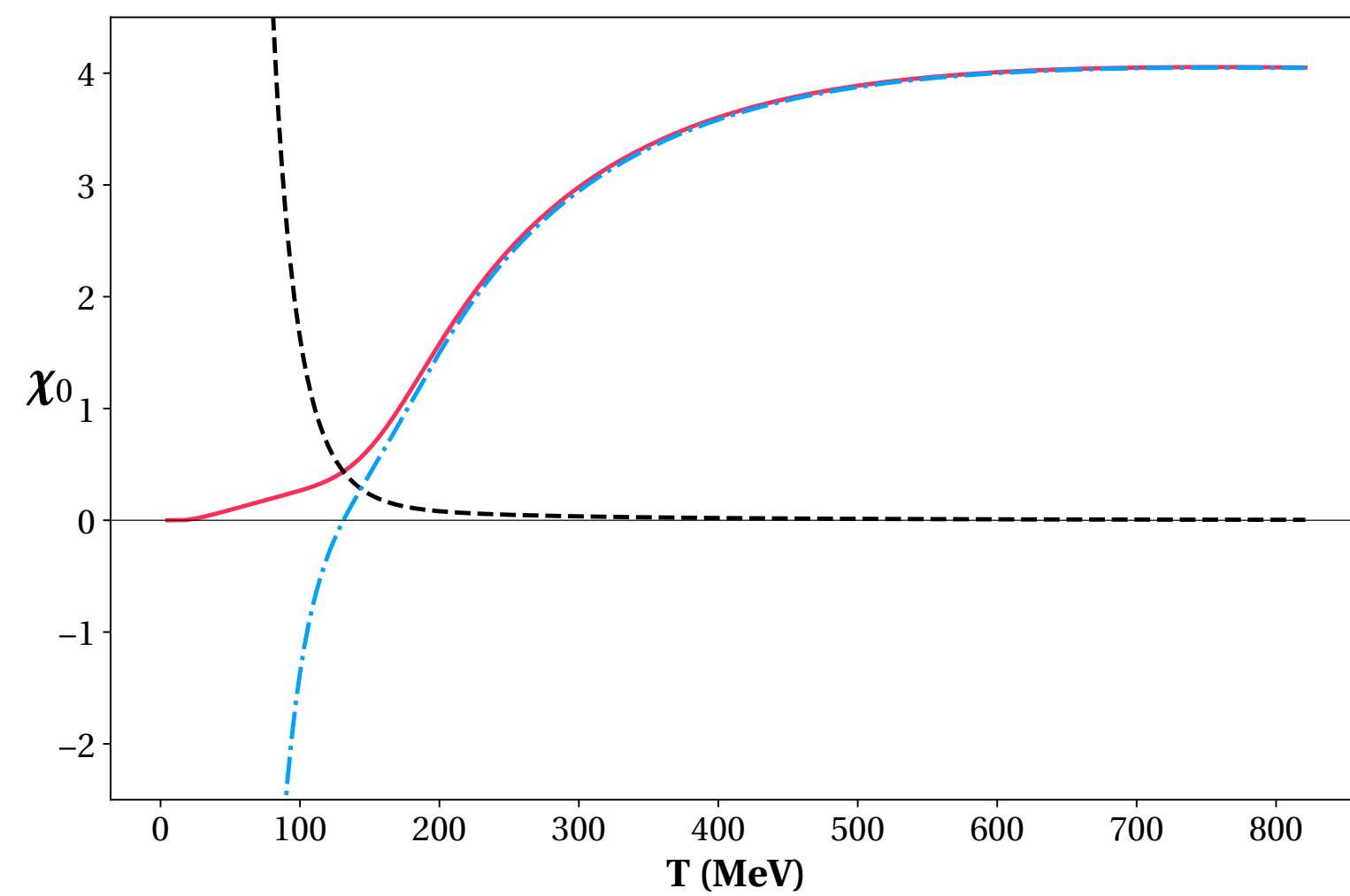
*JMK et al, EPJ+ (2021)*  
*J. Guenther et al, NPA (2017)*  
*R. Bellweid et al, PRD (2015)*  
 See also: *A. Bazavov et al, PRD (2017)*

# Singular and Non-singular Contributions



- Require that the total free energy (pressure) is the one from the lattice, so order-by-order we have:

$$\chi_N^{Lat}(T) = \chi_N^{Ising}(T) + \chi_N^{Non-Ising}(T)$$

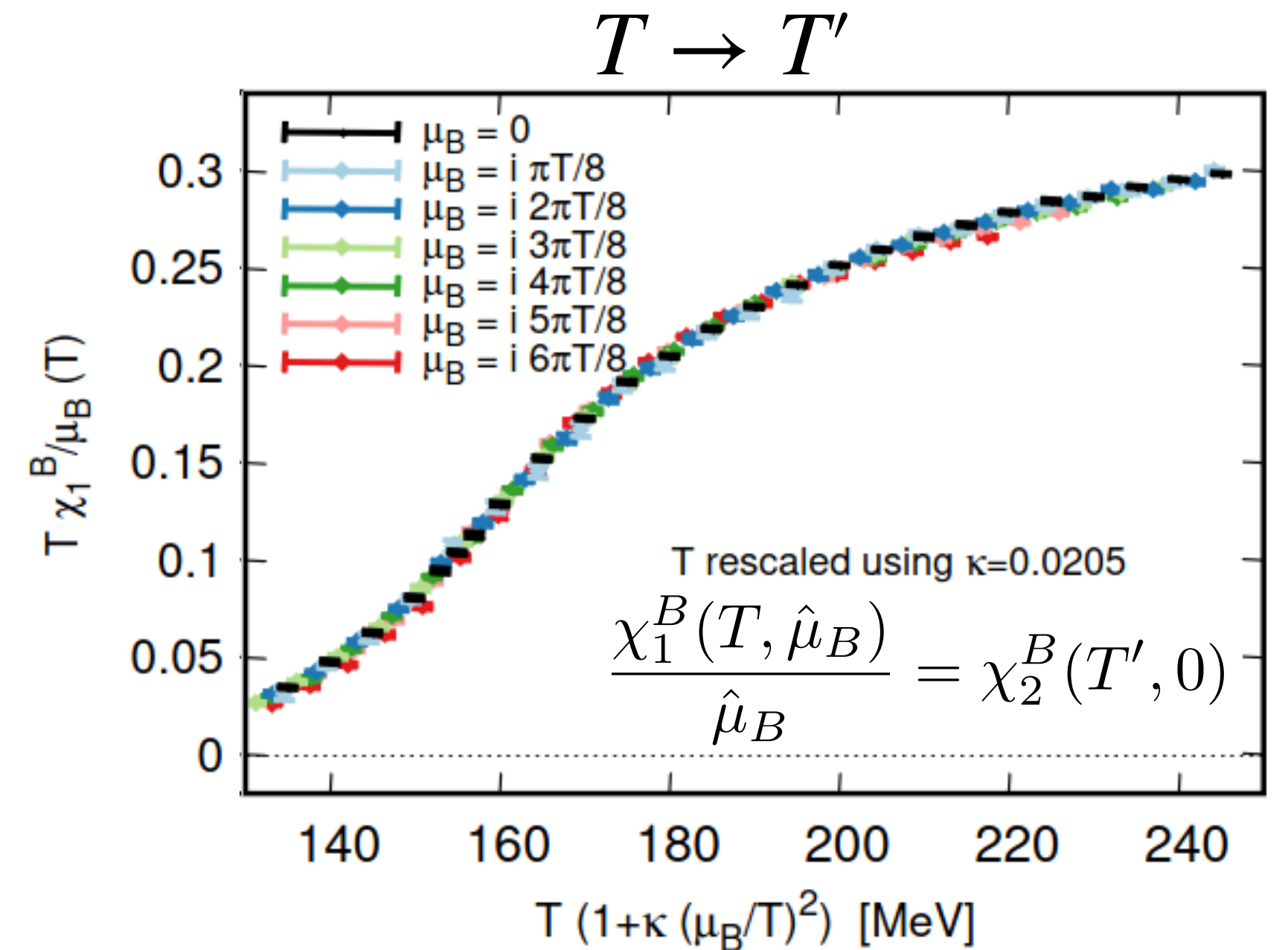


$$P(T, \mu_B) = T^4 \sum_{n=0}^2 c_{2n}^{\text{non-Ising}}(T) \left(\frac{\mu_B}{T}\right)^{2n} + T_C^4 P_{\text{symm}}^{\text{Ising}}(T, \mu_B)$$

- Recent development of critical EoS to extend coverage to  $\mu_B \simeq 700$  MeV in BES-II range: Ising TExS EoS is successor to BEST EoS
- Utilize lattice results from  $T'$ -Expansion Scheme (TExS) and incorporate singularity in  $n_B$  via a singularity in  $T'$

$$T'(T, \mu_B) = \underbrace{T'_{\text{lat}}(T, \mu_B)}_{\text{lowest orders in } (\mu_B/T)} + \underbrace{T'_{\text{crit}}(T, \mu_B) - \text{Taylor}_{n \leq 2}[T'_{\text{crit}}(T, \mu_B)]}_{\text{higher orders in } (\mu_B/T)}$$

which has the same singularity as  $T'_{\text{crit}}$  and the same truncated Taylor expansion as  $T'_{\text{lat}}$



# Ising TExS EoS

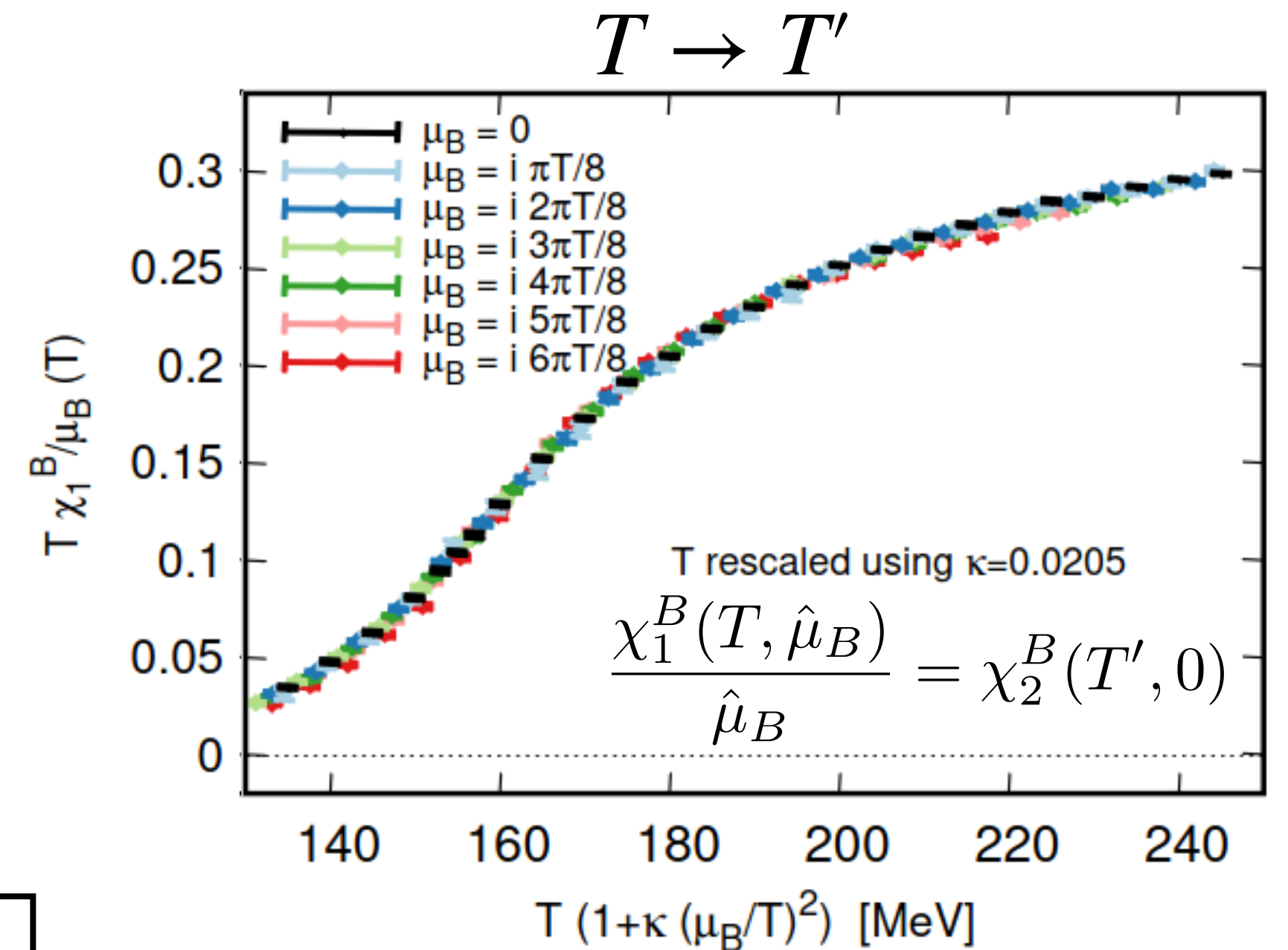


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Incorporate critical physics into higher orders in  $\mu_B$  by subtracting the 2nd order Taylor term



# Ising TExS EoS

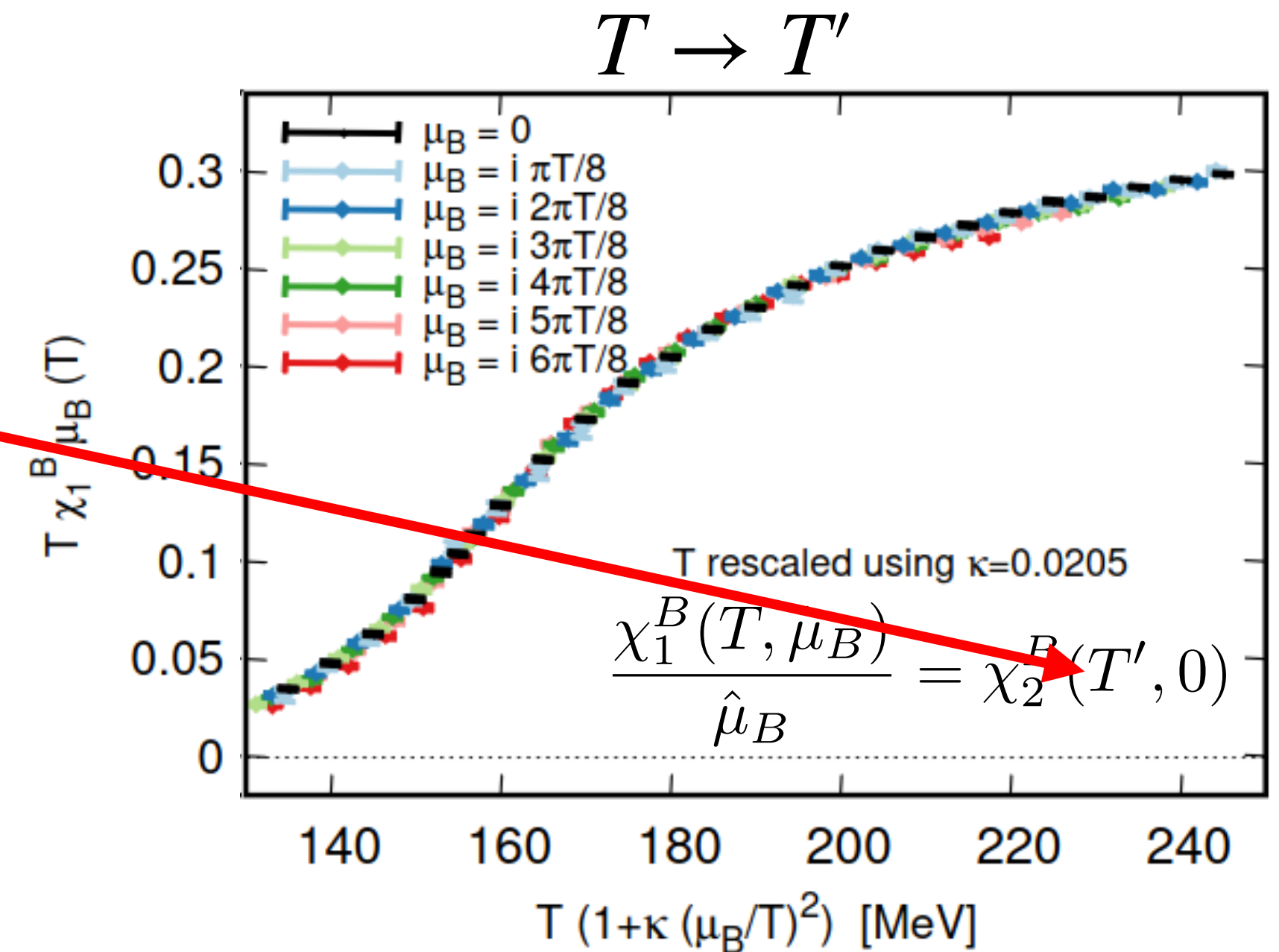


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$T'$  completely defines the baryon density with a critical point  $\rightarrow$  converges faster



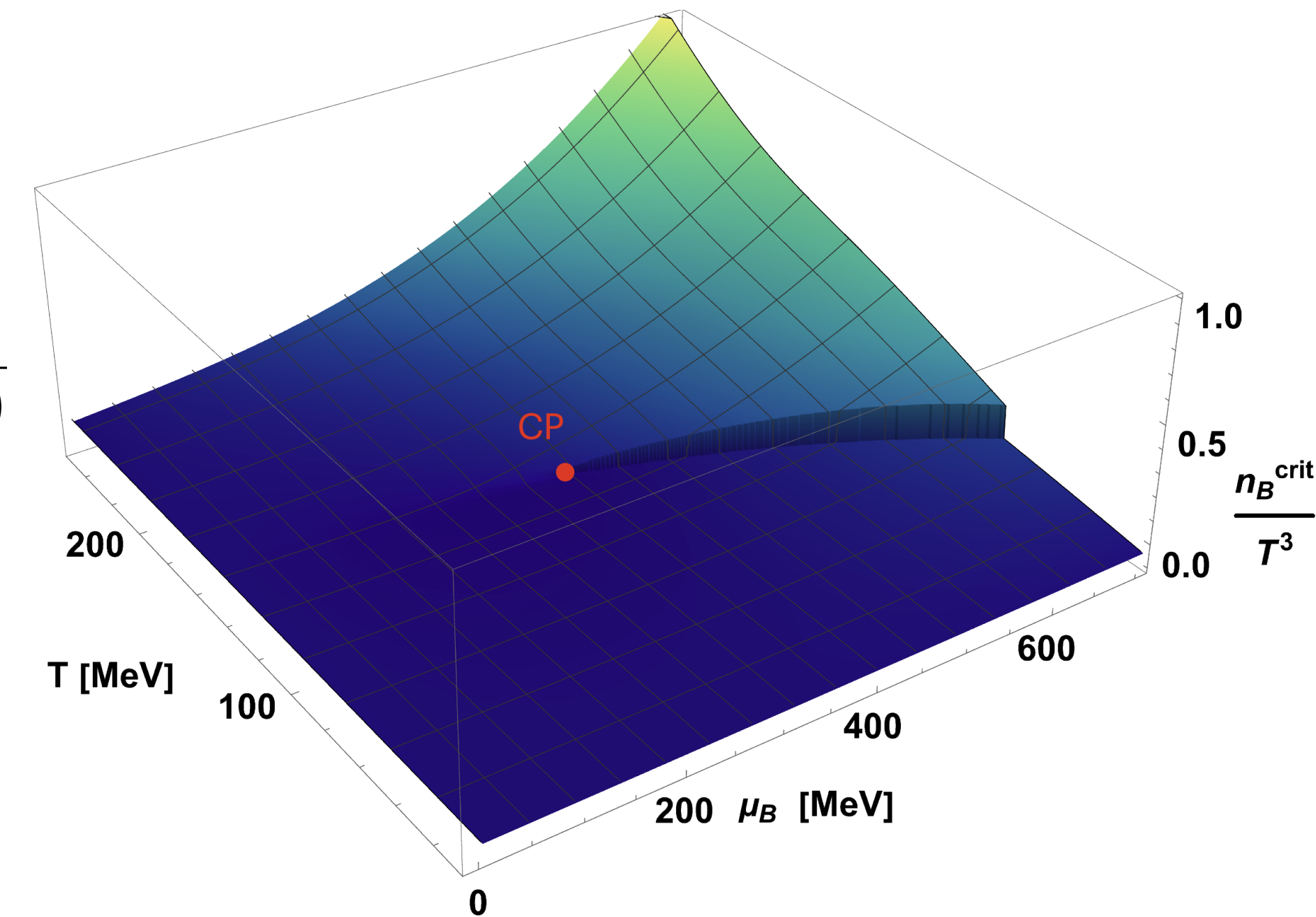
M. Kahangirwe et al, PRD (2024)

- Recent development of critical EoS to extend coverage to  $\mu_B \simeq 700$  MeV in BES-II range: Ising TExS EoS is successor to BEST EoS
- Map Ising transition line quadratically along chiral phase transition rather than tangent to the critical point with a linear mapping

$$(\mathbf{r}, \mathbf{h}) \longleftrightarrow (\mathbf{T}, \mu_B) : \begin{aligned} h(\mu, T) &= -\frac{\Delta T' \cos \alpha_1}{T_c w \sin(\alpha_1 - \alpha_2)} \\ r(\mu, T) &= -\frac{\mu^2 - \mu_c^2}{2\mu_c T_c \rho w \cos \alpha_1} + \frac{\Delta T' \cos \alpha_2}{T_c \rho w \sin(\alpha_1 - \alpha_2)} \end{aligned}$$

$$\Delta T'(\mu, T) \equiv \frac{T'(\mu, T) - T_0}{(\partial T' / \partial T)_{\text{CP}}}$$

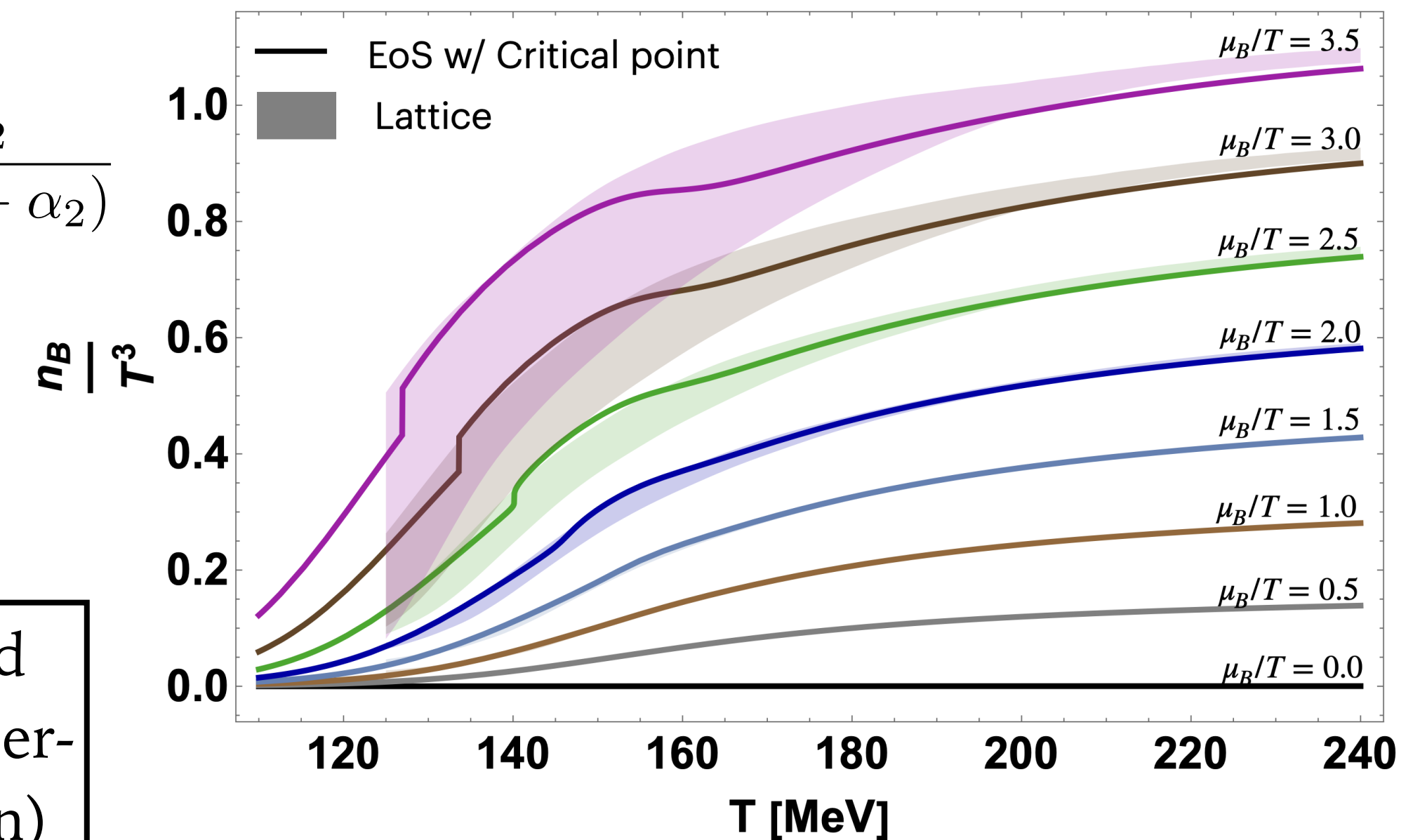
New quadratic mapping puts Ising phase transition line directly on chiral crossover from lattice



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Full EoS agrees with lattice and includes critical features with user-chosen parameters (e.g. location)



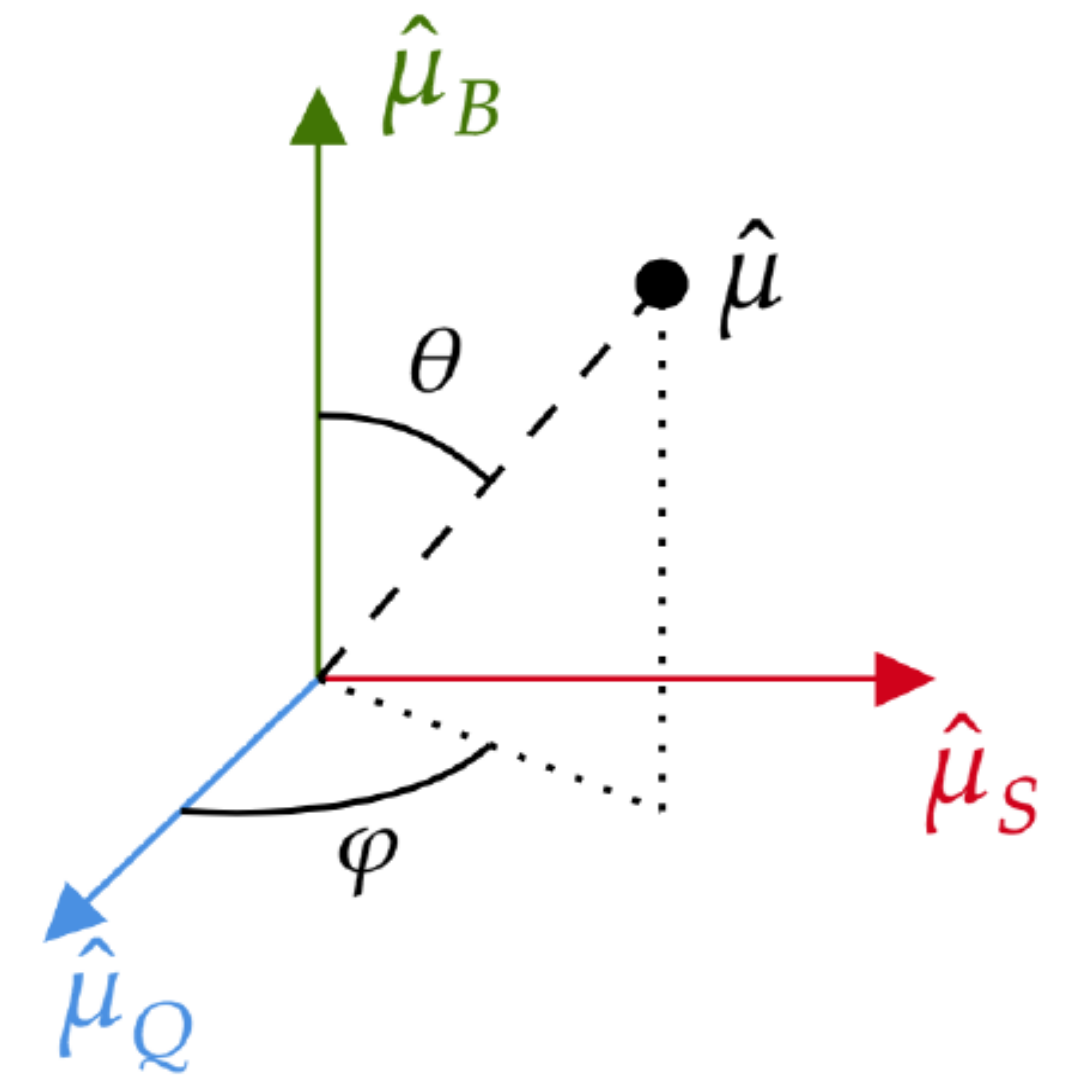
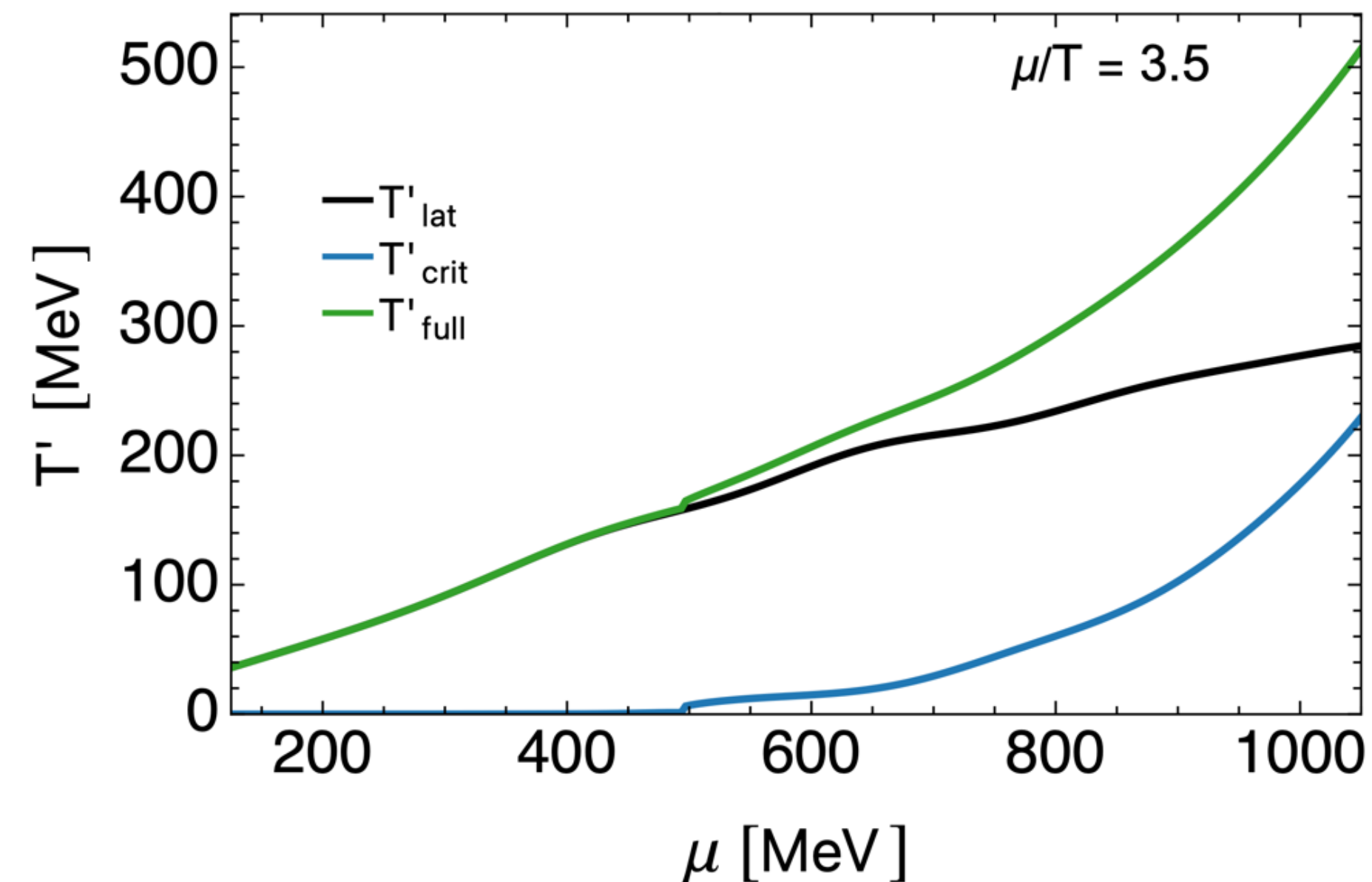
*M. Kahangirwe et al, PRD (2024)*

# Critical EoS in $T$ - $\mu_B$ - $\mu_Q$ - $\mu_S$ : 4D Ising TExS

- Extend to all relevant conserved charges based on the Ising TExS critical EoS approach and the 4D TExS EoS in spherical coordinates:

$$T'_{\text{full}}(T, \hat{\mu}, \theta, \varphi) = T'_{\text{lat}}(T, \hat{\mu}, \theta, \varphi) + T'_{\text{crit}}(T, \hat{\mu}, \theta, \varphi)$$

$$T'_{\text{crit}}(T, \hat{\mu}, \theta, \varphi) = \left( \frac{\partial X_2}{\partial T} \Big|_{T_0} \right)^{-1} \cdot \frac{n_{\text{Ising}}(T, \hat{\mu}, \theta, \varphi)}{\hat{\mu}} - T'_{\text{Taylor}}^{\mathcal{O} \leq 2}$$



# Critical EoS in $T-\mu_B-\mu_Q-\mu_S$ : 4D Ising TExS

- Adding additional dimensions in conserved charges promotes the critical point to a critical surface: first consider two chemical potentials - a critical ellipse

Parametrize the location of the critical ellipse by:

- $\{\tau, \mu_{B,c}, \mu_{S,c}\}$
- $\{\mu_c, d\mu/d\theta, d^2\mu/d\theta^2\}$  along strangeness neutrality

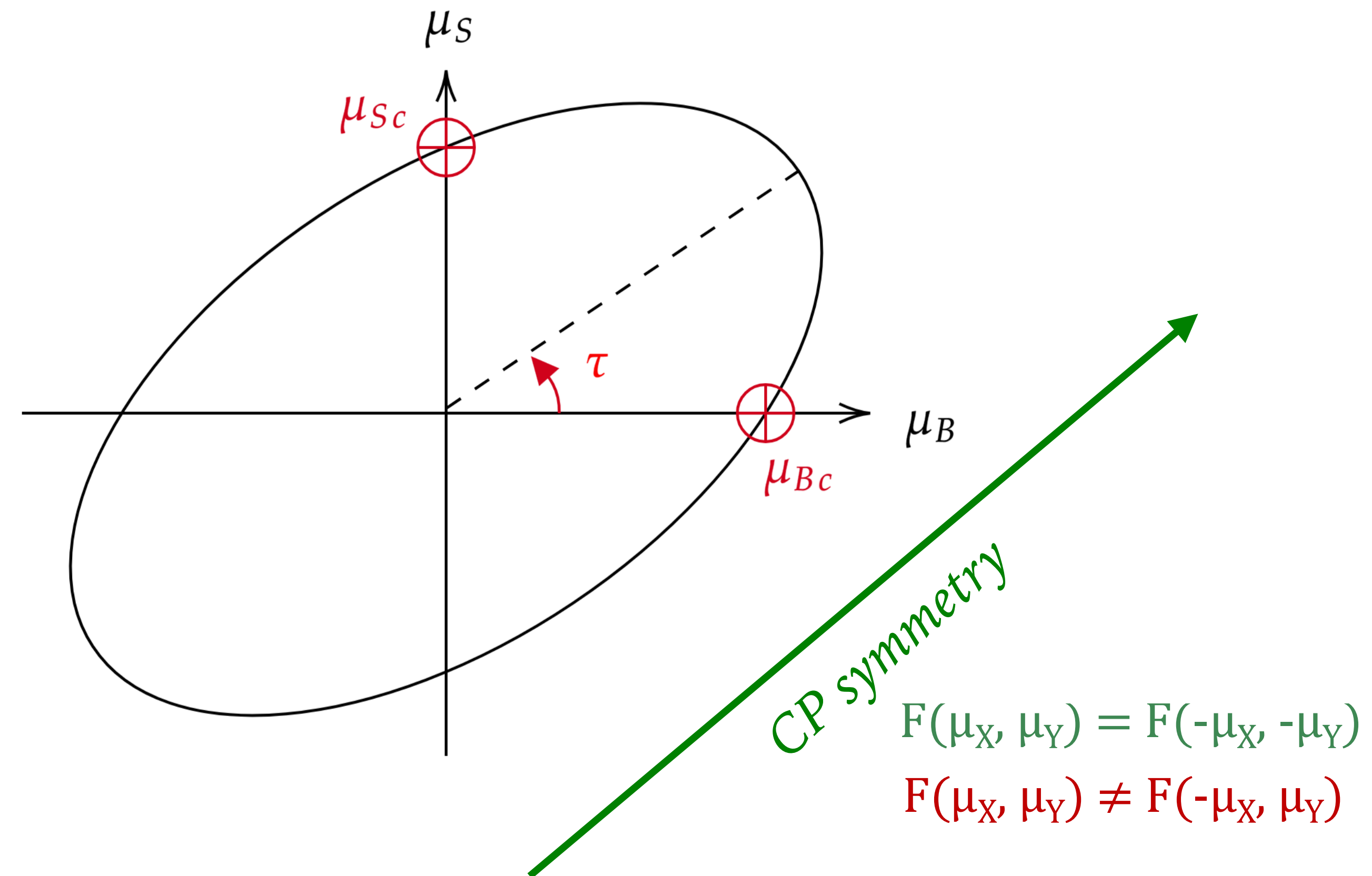


Figure adapted from F. DiClemente CPOD 2026

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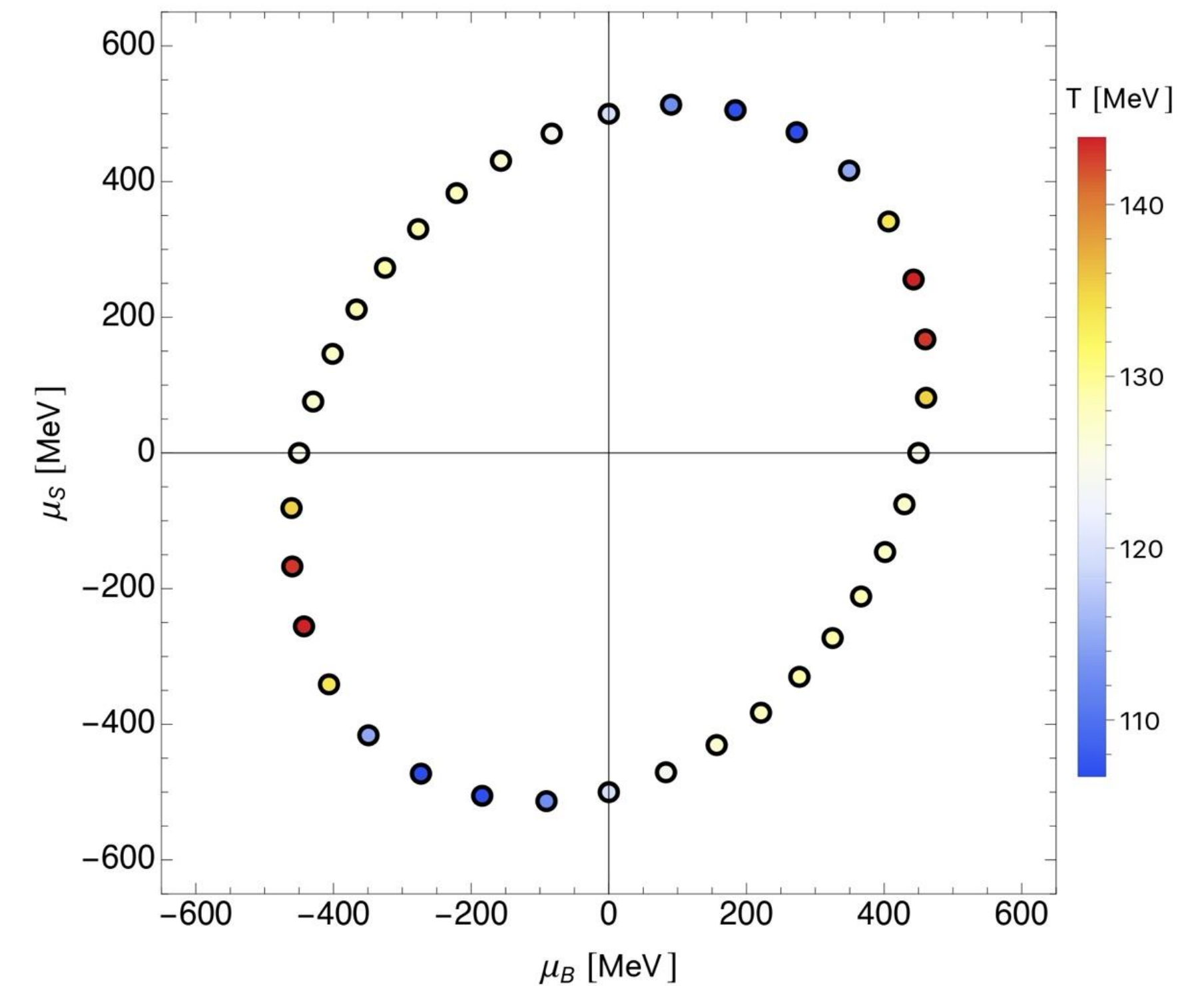
An example critical ellipse:

$$\tau = 57^\circ, \mu_{B,c} = 450 \text{ MeV}, \mu_{S,c} = 500 \text{ MeV}$$

Temperature constrained by  $\Delta T' = 0$

$$\Delta T'(\mu, T) \equiv \frac{T'(\mu, T) - T_0}{(\partial T' / \partial T)_{cp}}$$

$$T_0 = T_c \cdot [1 + \lambda_2(T_c) \cdot \hat{\mu}_c^2] \quad \longrightarrow \quad \hat{\mu}_c = \sqrt{\frac{T_0/T_c - 1}{\lambda_2(T_c)}}$$



# Critical EoS in $T-\mu_B-\mu_Q-\mu_S$ : 4D Ising TExS

- Preliminary results given an example choice of parameters and trajectory in hyperplane

$$\mu_{B,c} = 450 \text{ MeV}$$

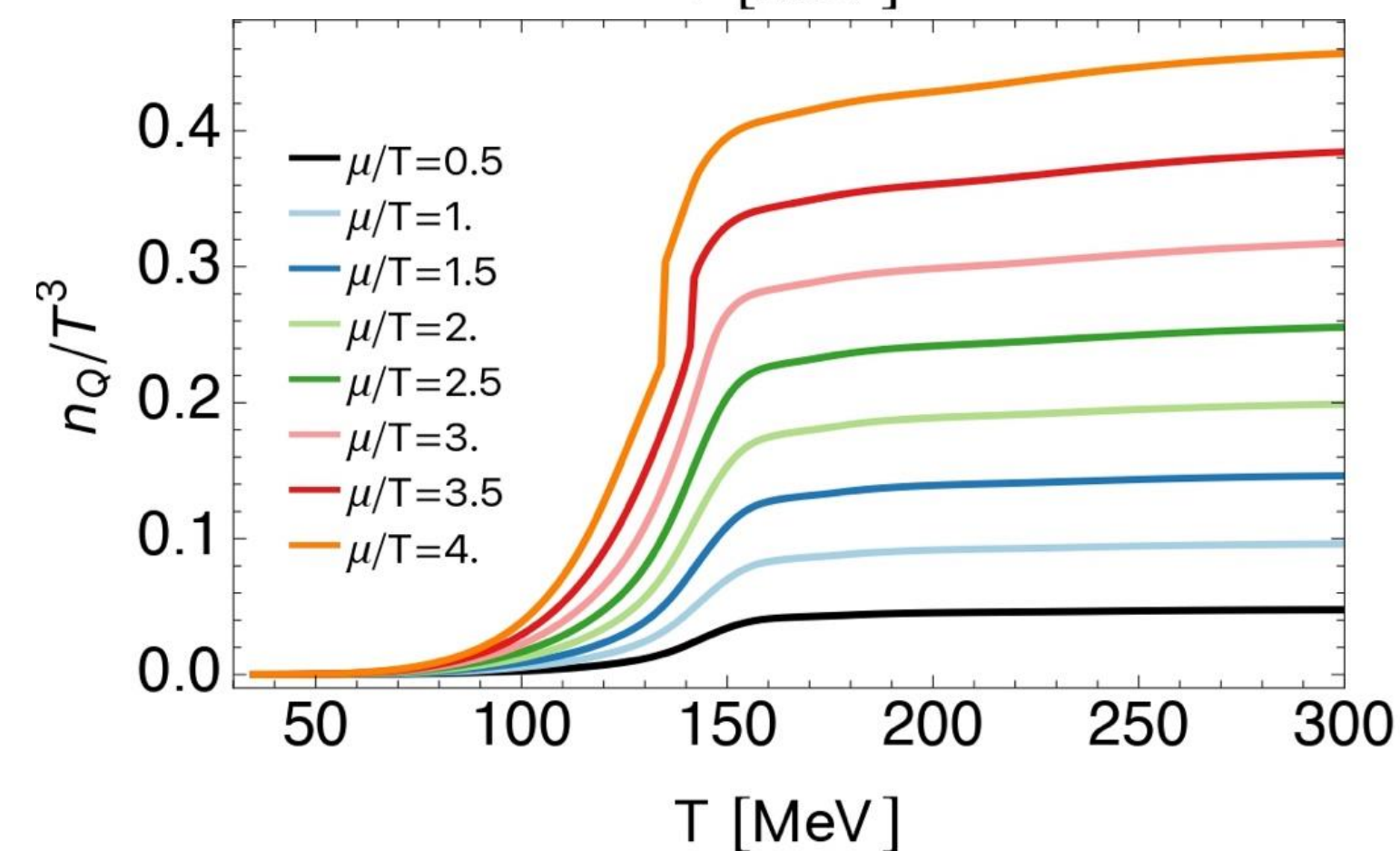
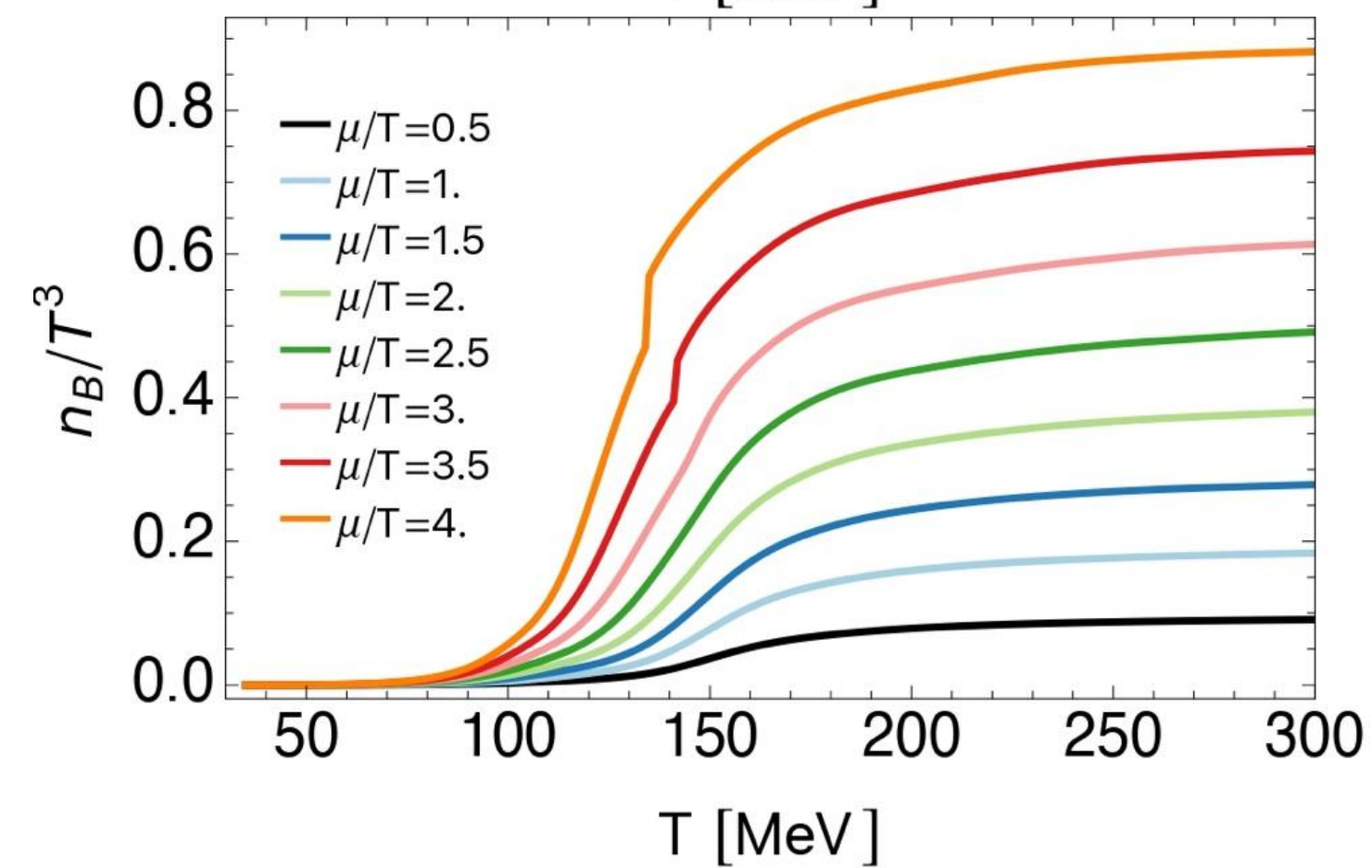
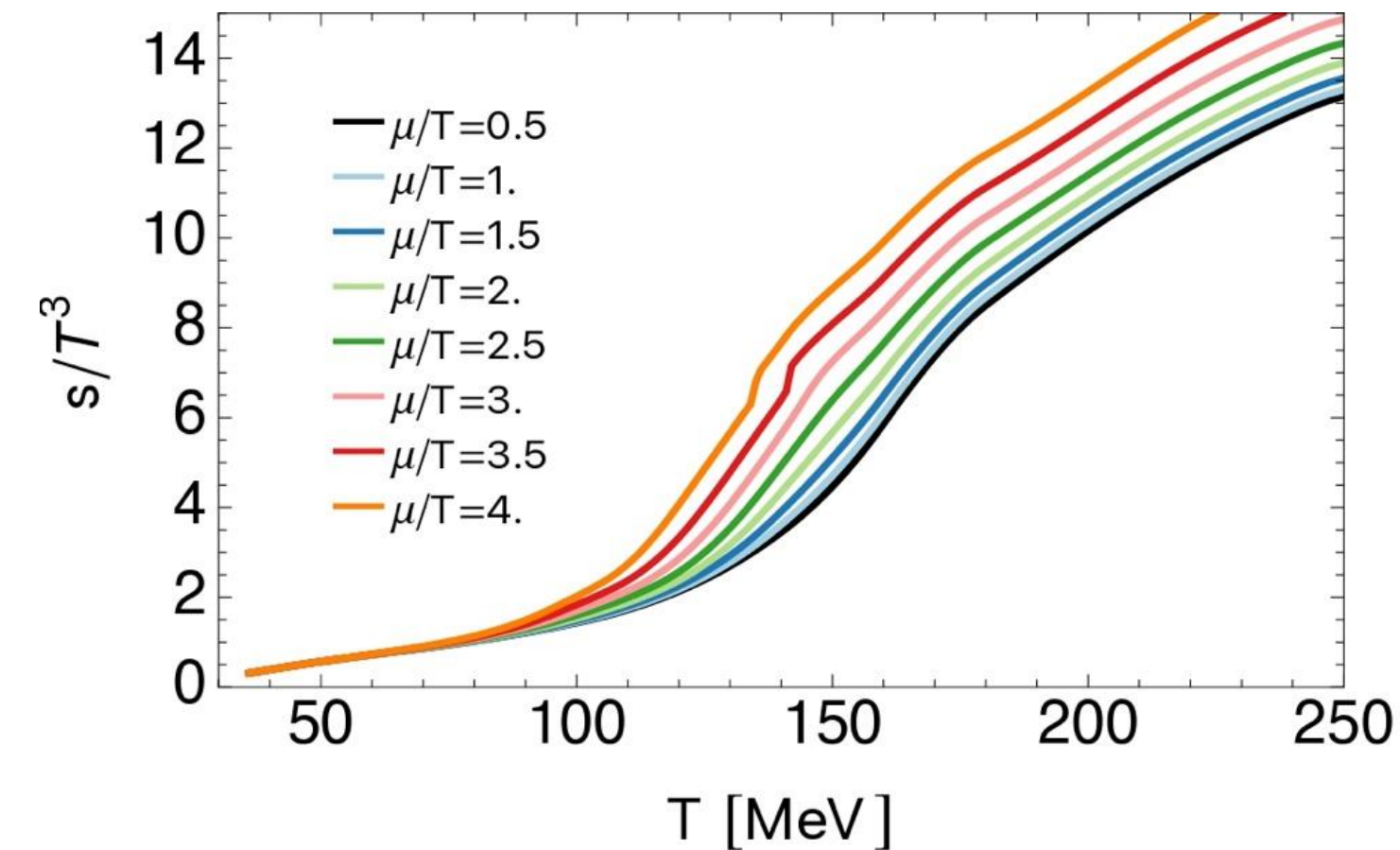
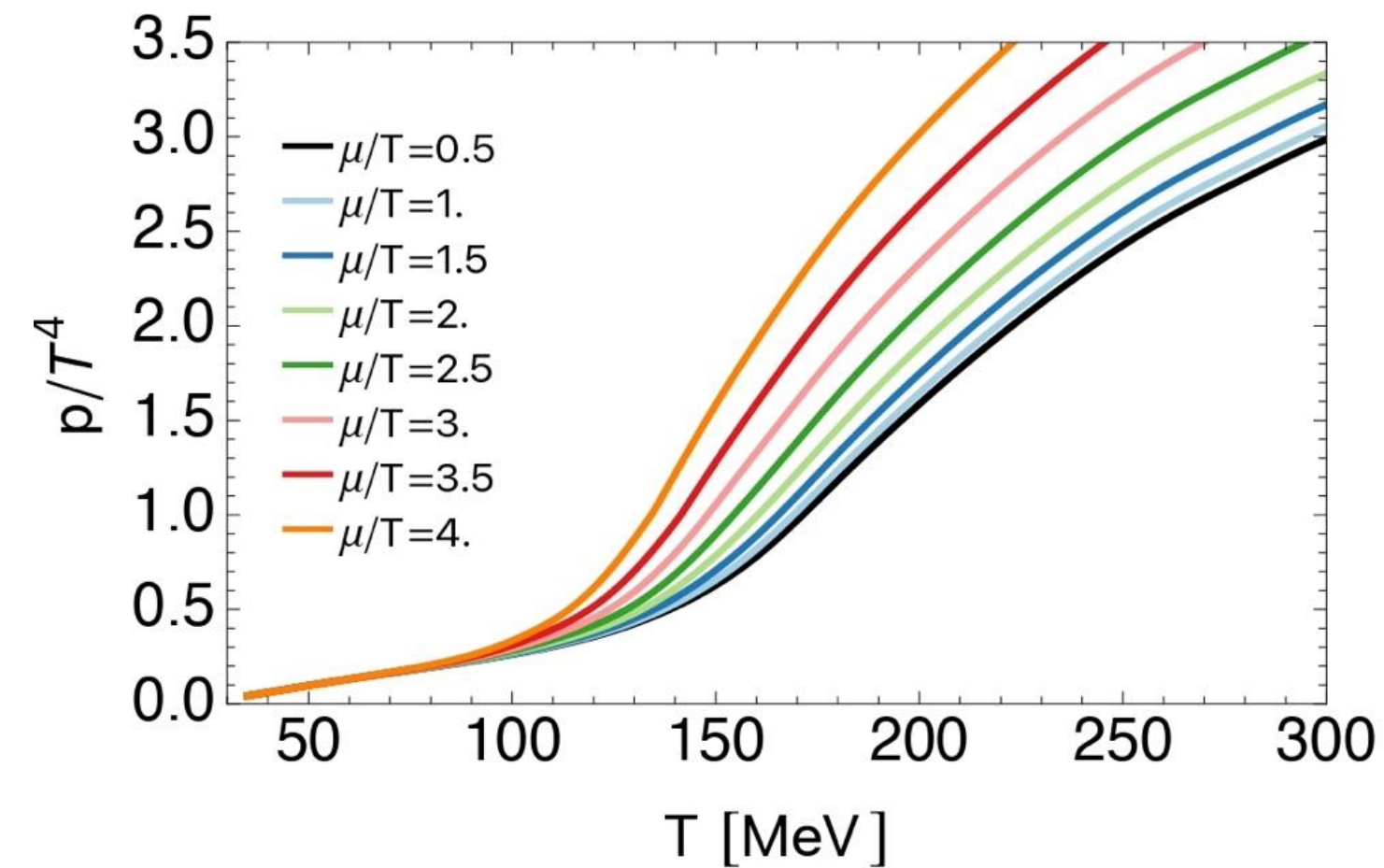
$$\mu_{S,c} = 500 \text{ MeV}$$

$$\alpha_{12} = \alpha_1 - \alpha_2 = 14.6^\circ$$

$$w = 5$$

$$\rho = 0.3$$

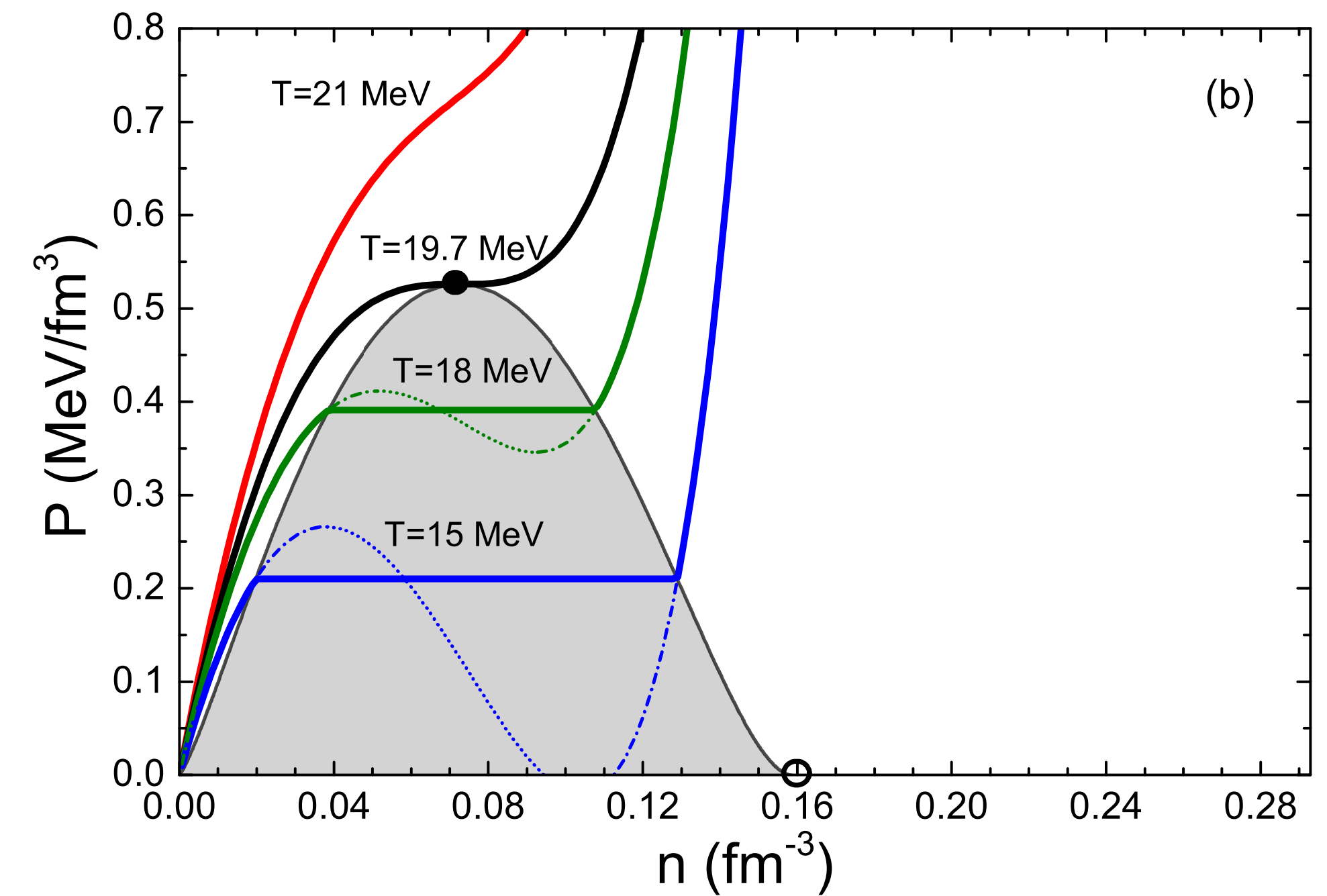
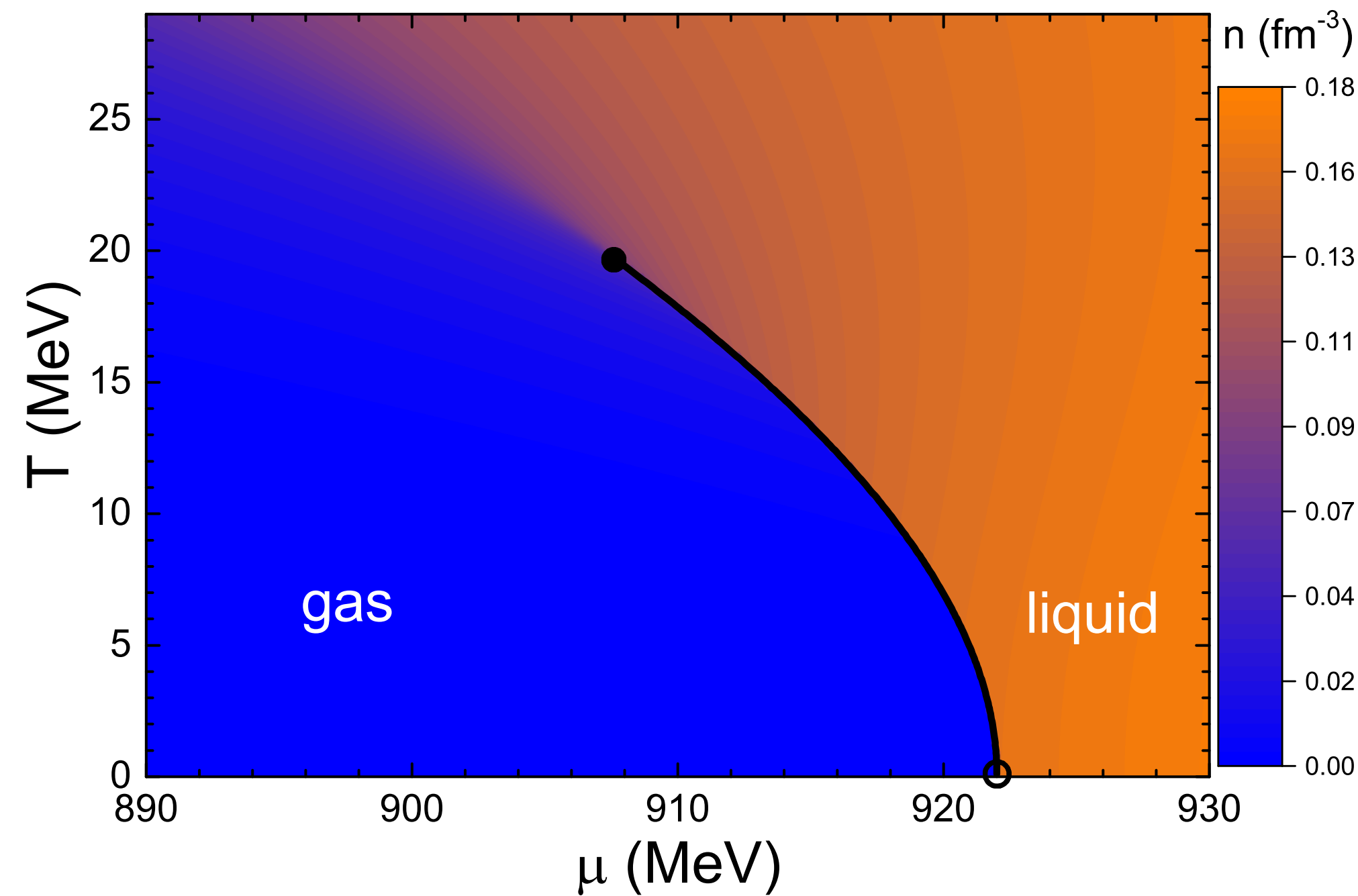
$$\mu_S = \mu_B/3$$



# First-order Phase Transition Features



- Besides the critical point itself, the features of the first order phase transition can tell us about the existence of a critical point
  - Coexistence region & spinodals

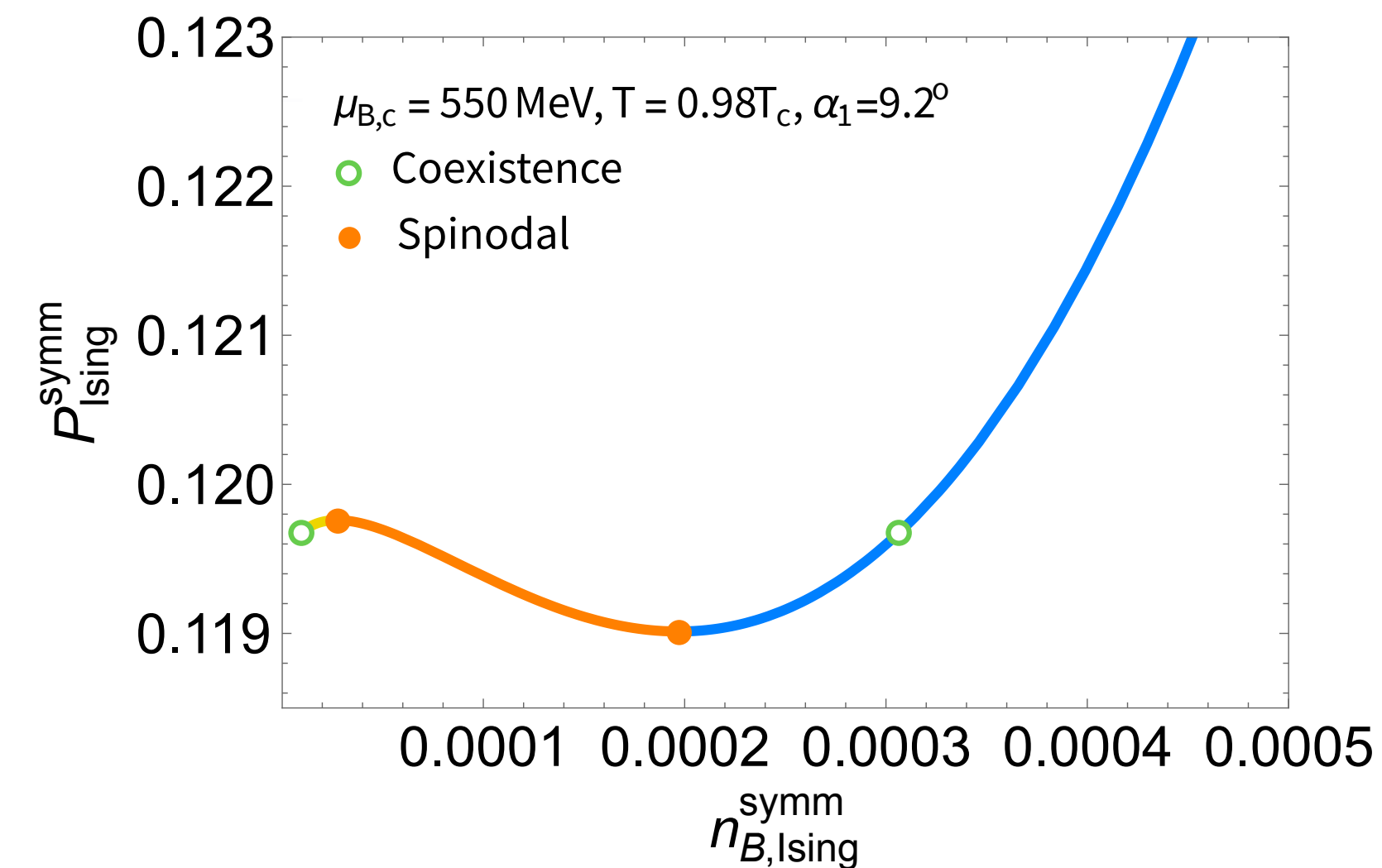
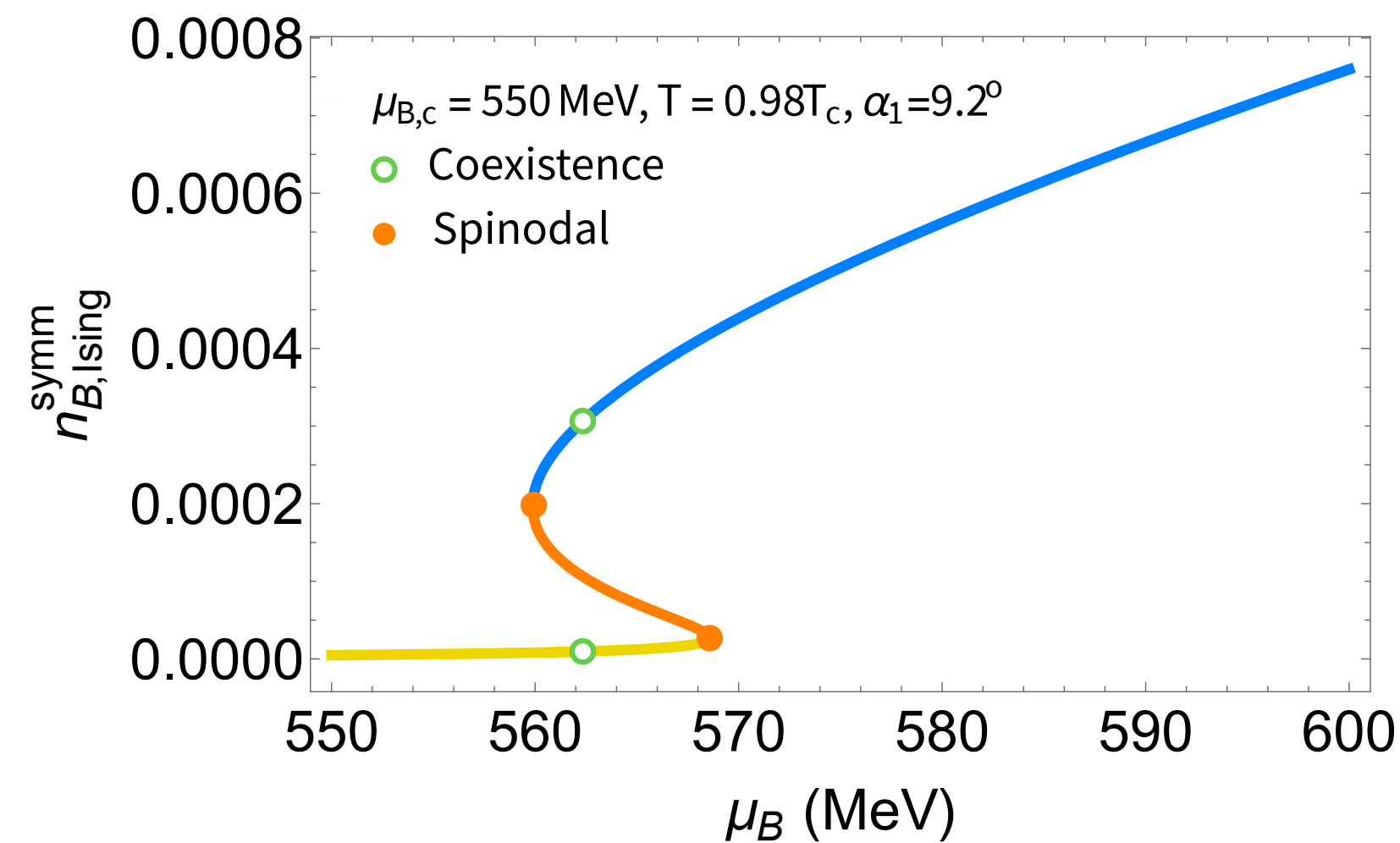


*V. Vovchenko, D.V. Anchishkin, M.I. Gorenstein, PRC (2015)*

# EoS for First Order Regime

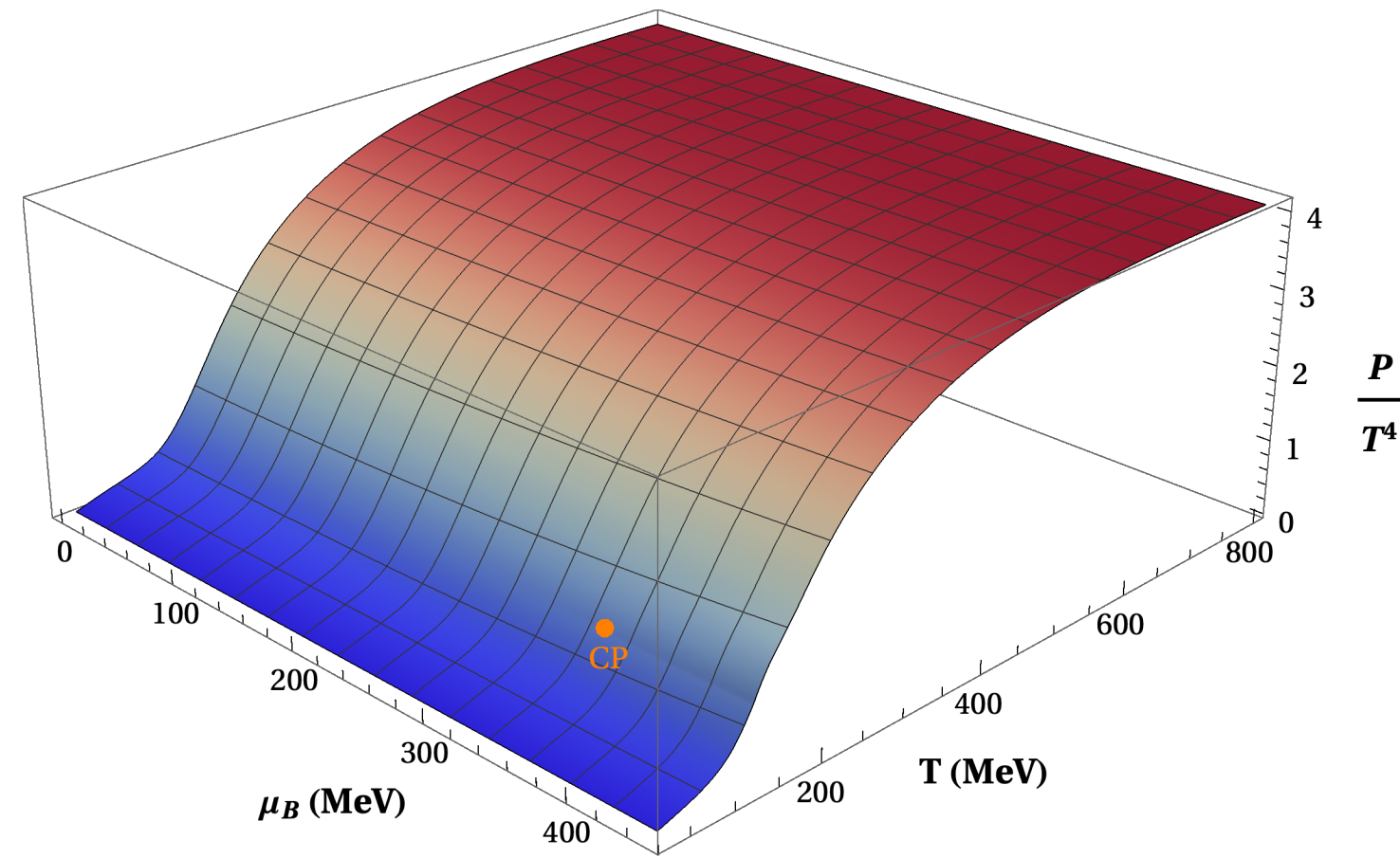
- Utilize mean field Ising model to incorporate first order features, e.g. spinodals, not present in 3D Ising due to shift into complex plane:

$$\Delta\phi = \pi \left( \beta\delta - \frac{3}{2} \right) \quad \text{X. An, D. Mesterhazy, M. Stephanov, J. Stat. Mech (2018)}$$

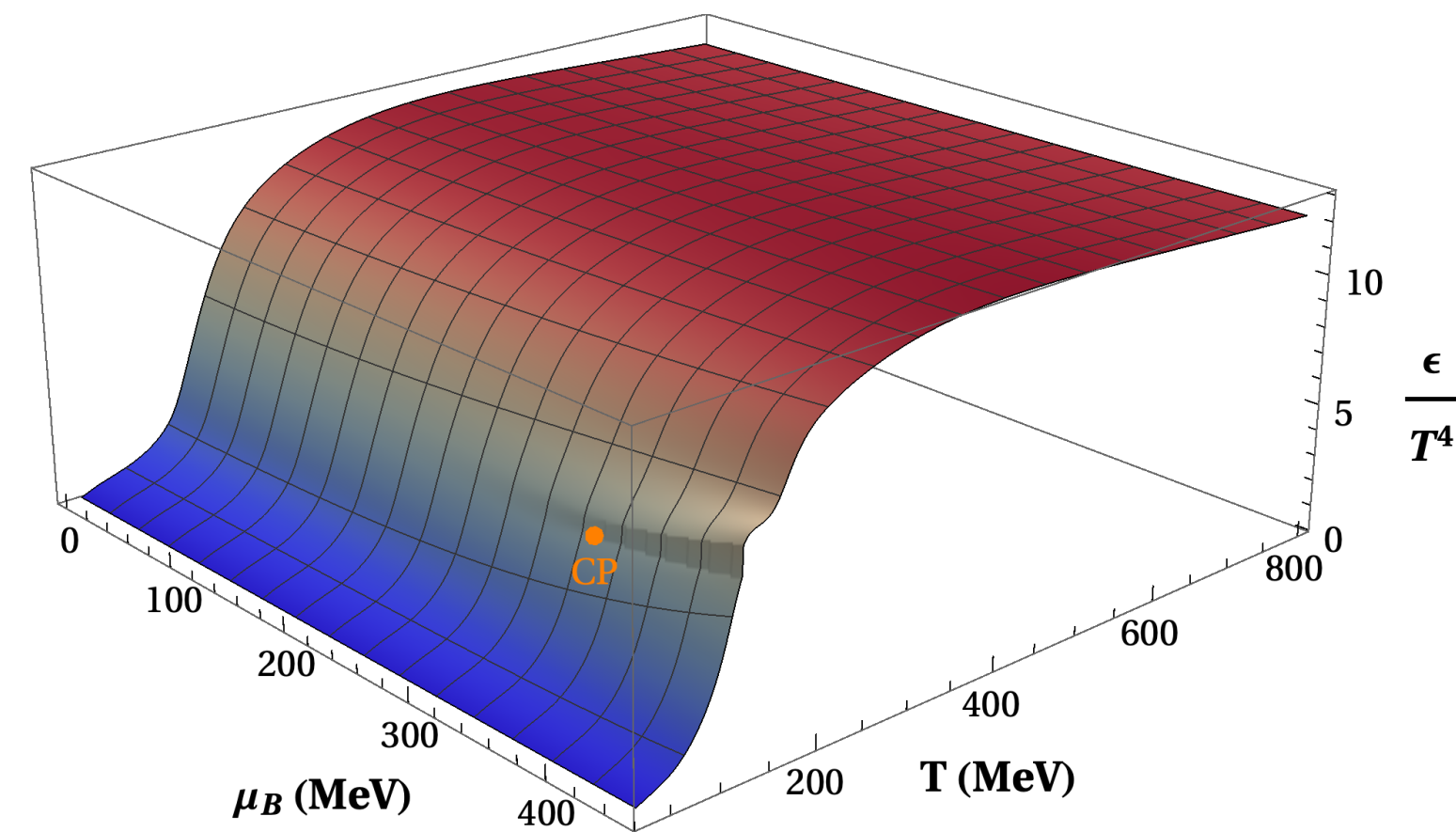


Strong dependence of the features on the angle  $\alpha_1$   
 → future work will use new quadratic mapping

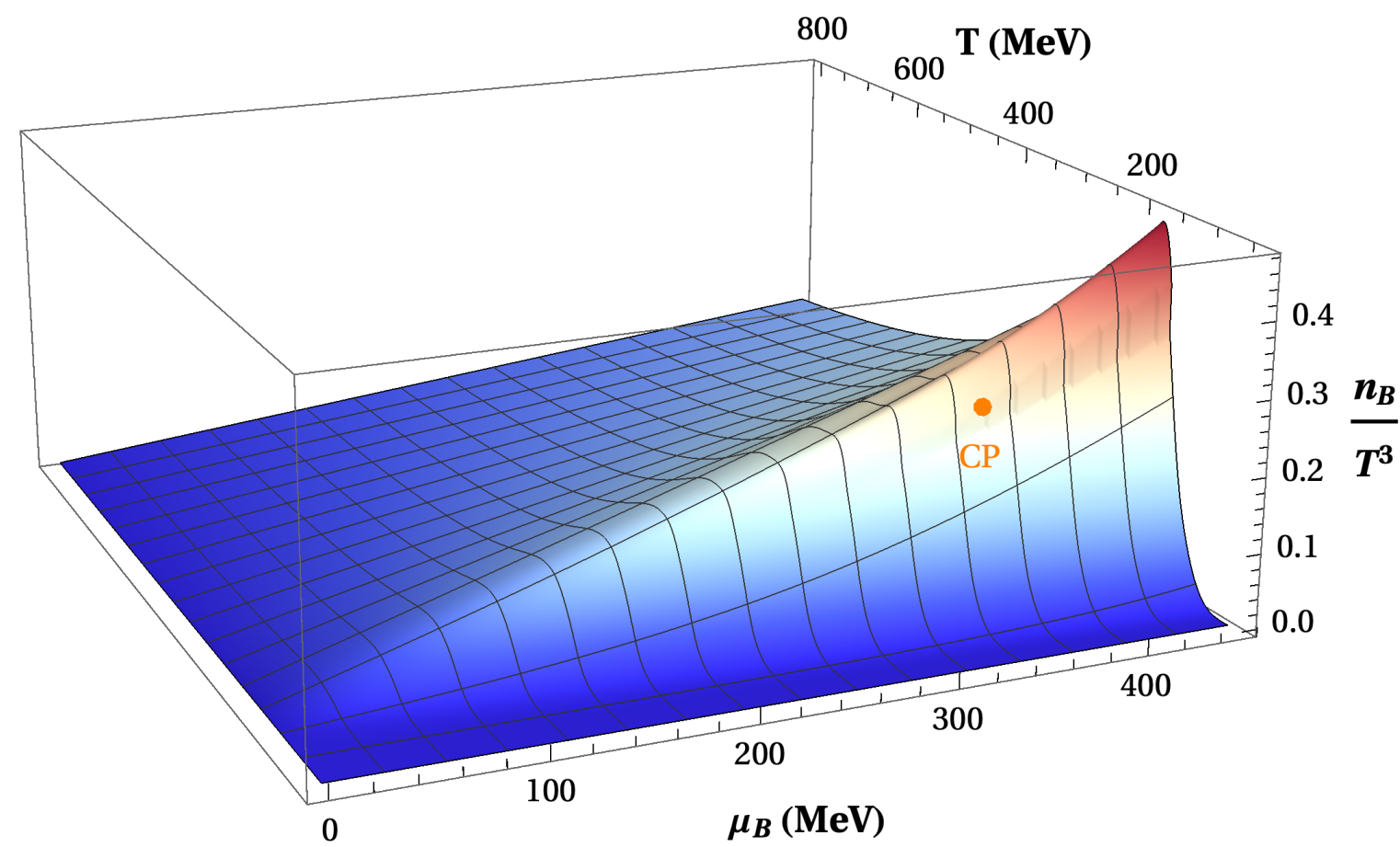
# How to use EoS to connect to experiment?



Pressure



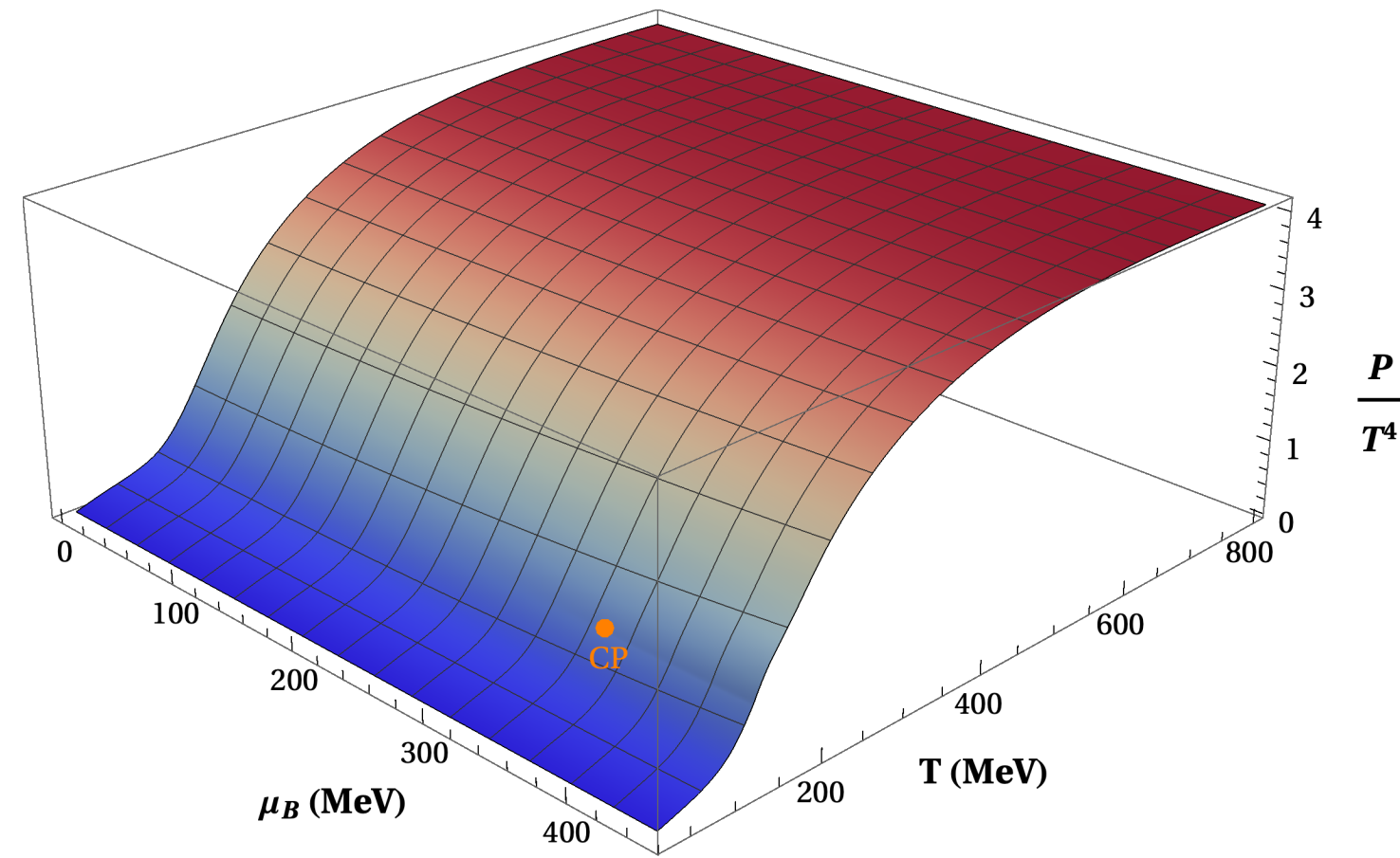
Energy density



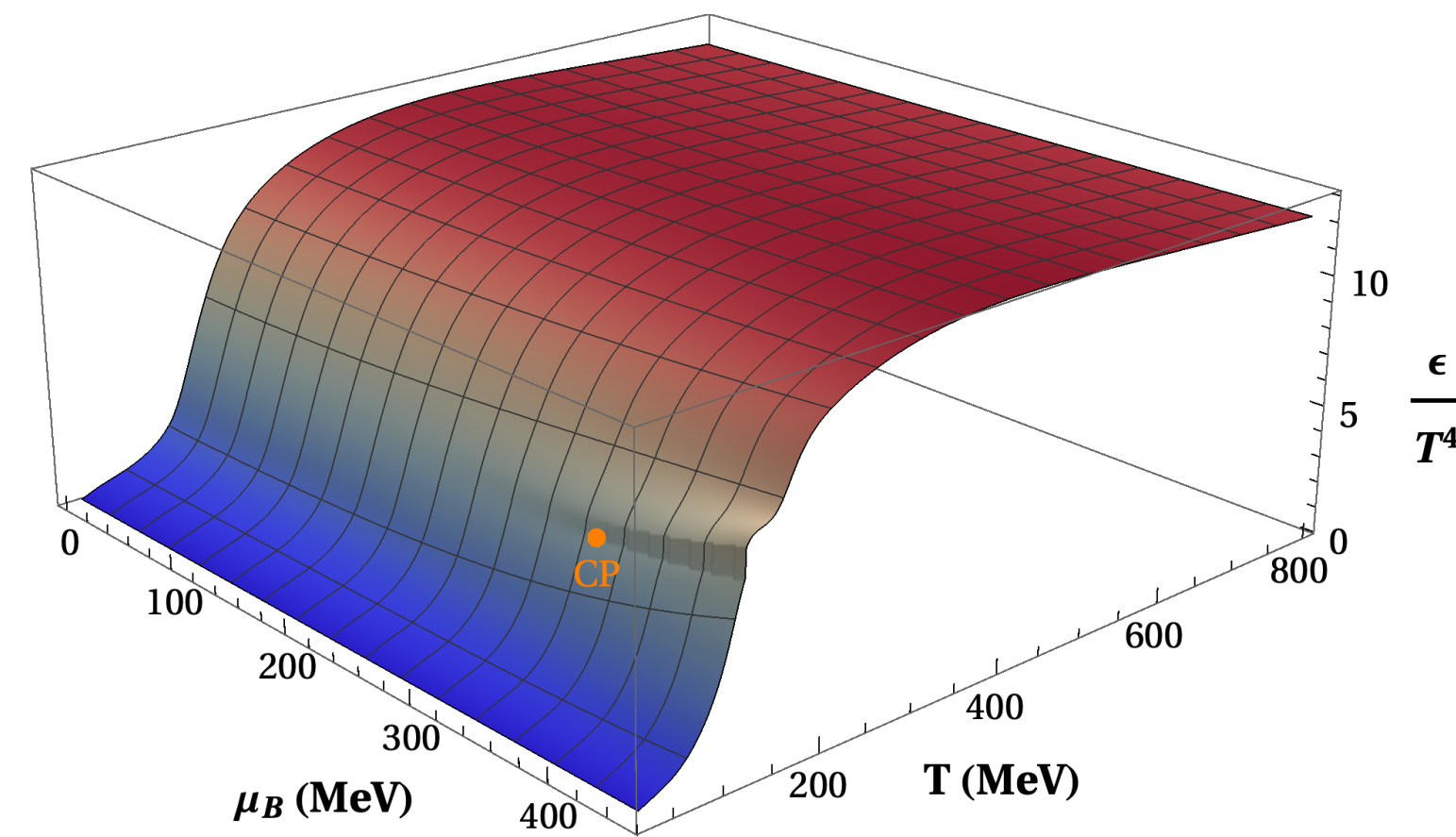
Baryon density  *EoS with parametrized CP*

JMK et al, EPJ+ (2021)

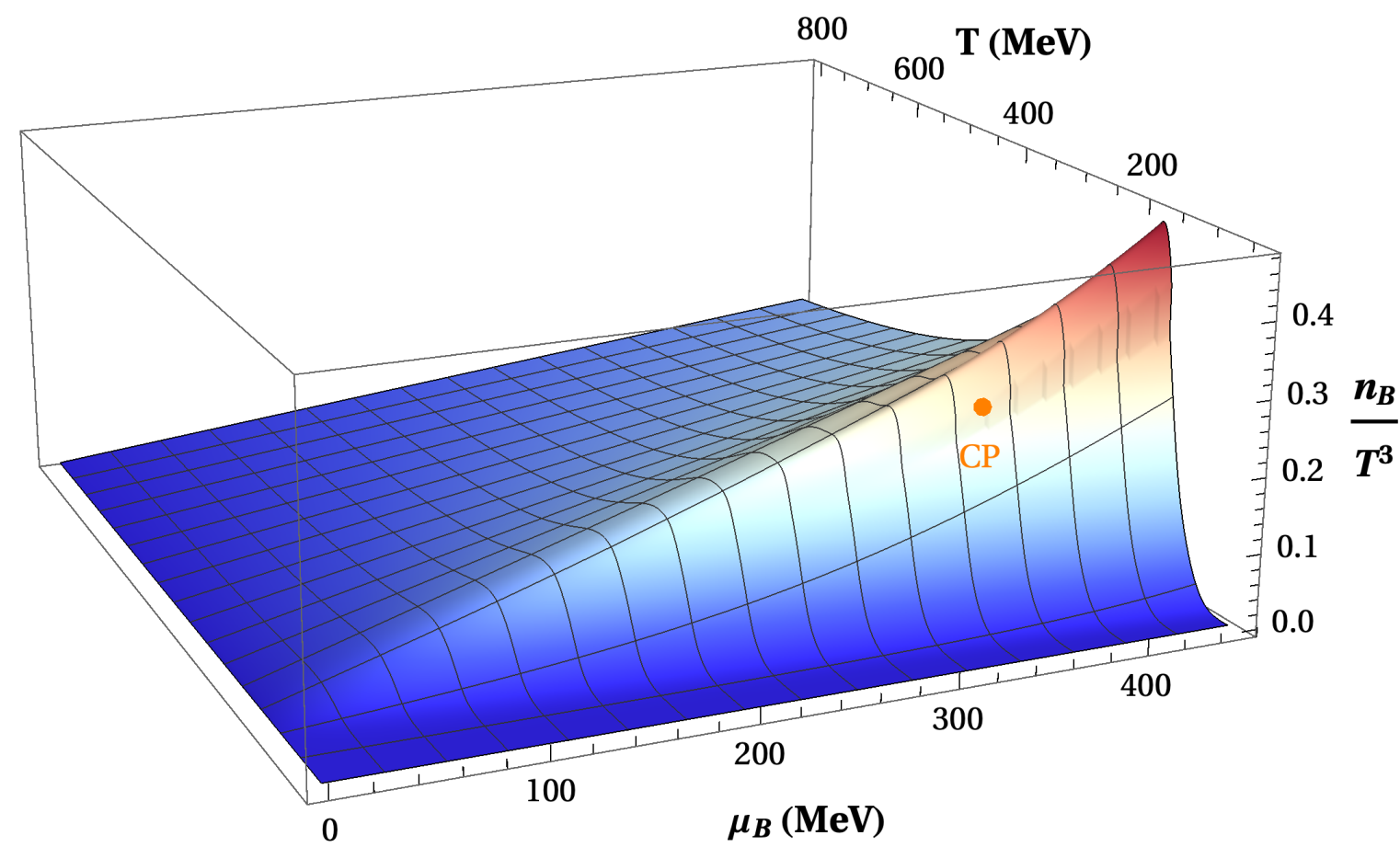
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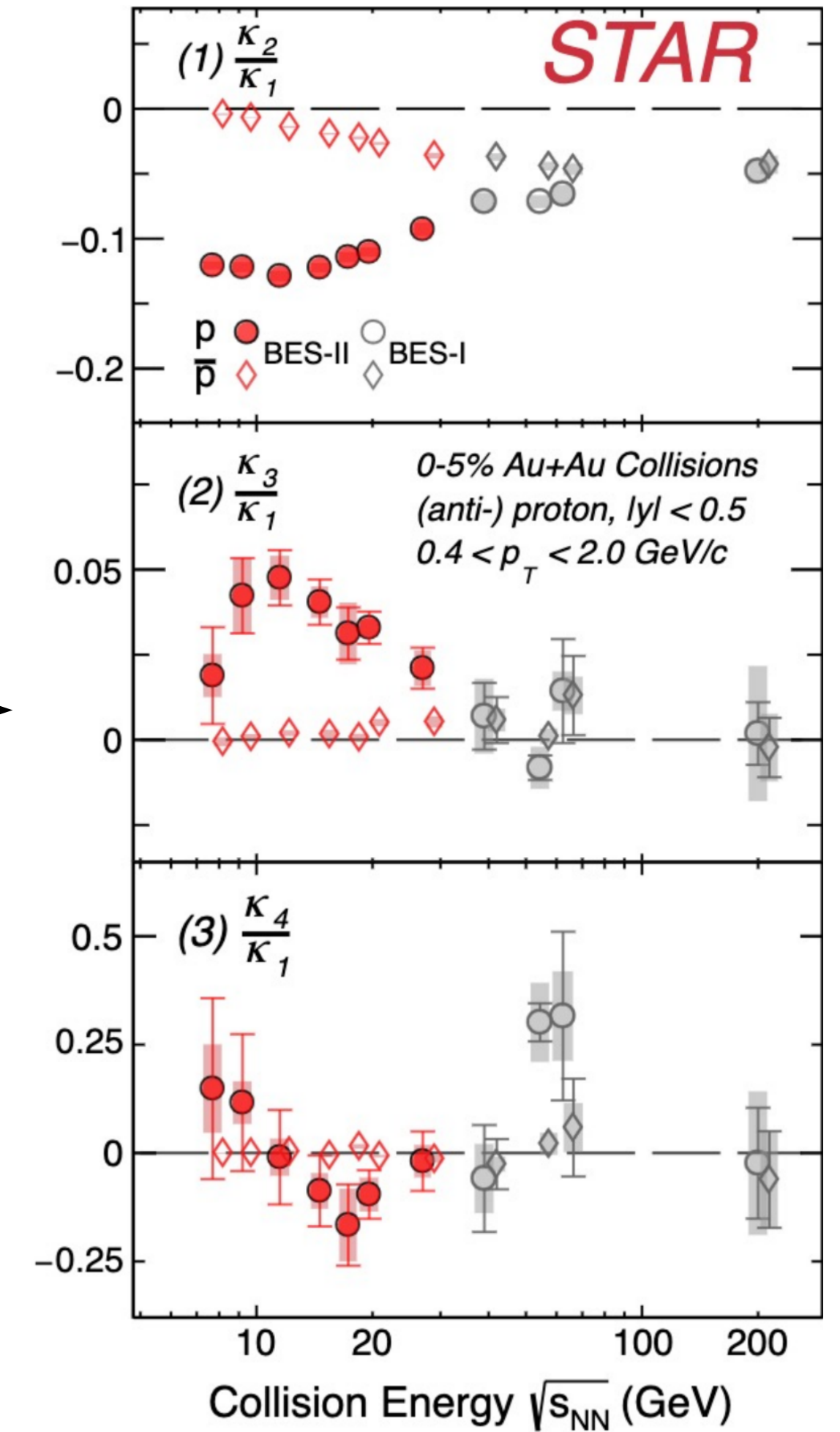
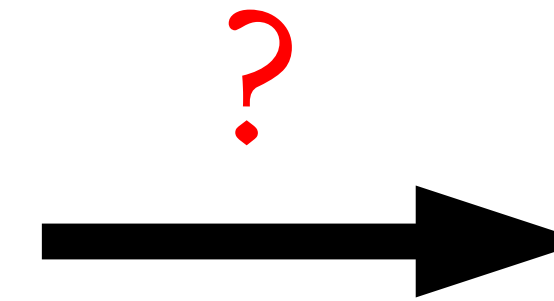
Pressure



Energy density



Baryon density



JMK et al, EPJ+ (2021)

A. Pandav (STAR collaboration), CPOD 2024

# Effective Model for Critical Fluctuations

- ▶ Linear sigma model predicts a critical point at which the sigma mass vanishes & correlation length diverges

$$\Omega = \int d^3 \mathbf{x} \left[ \frac{(\nabla \sigma)^2}{2} + \frac{m_\sigma^2}{2} \sigma^2 + \frac{\lambda_3}{3} \sigma^3 + \frac{\lambda_4}{4} \sigma^4 + \dots \right]$$

- ▶ Historically used to calculate particle multiplicity fluctuations

$$V \langle \delta n_{\mathbf{p}} \delta n_{\mathbf{k}} \rangle = \frac{g^2}{m_\sigma^2 T} \frac{4m^2}{E_{\mathbf{p}} E_{\mathbf{k}}} \left[ n_{\mathbf{p}}^+ (1 - n_{\mathbf{p}}^+) - n_{\mathbf{p}}^- (1 - n_{\mathbf{p}}^-) \right] \\ \times \left[ n_{\mathbf{k}}^+ (1 - n_{\mathbf{k}}^+) - n_{\mathbf{k}}^- (1 - n_{\mathbf{k}}^-) \right]$$

- ▶ Unknown couplings (particle,  $g$ , and higher point,  $\lambda$ ), correlation length ( $\xi = 1/m_\sigma$ )

*M. Gell-Mann and M. Levy, Nuovo Cim. (1960)*

*M. Stephanov, K. Rajagopal, E. Shuryak, PRL (1998)*

*O. Scavenius et al, PRC (2001)*

*Y. Hatta and M.A. Stephanov PRL (2003)*

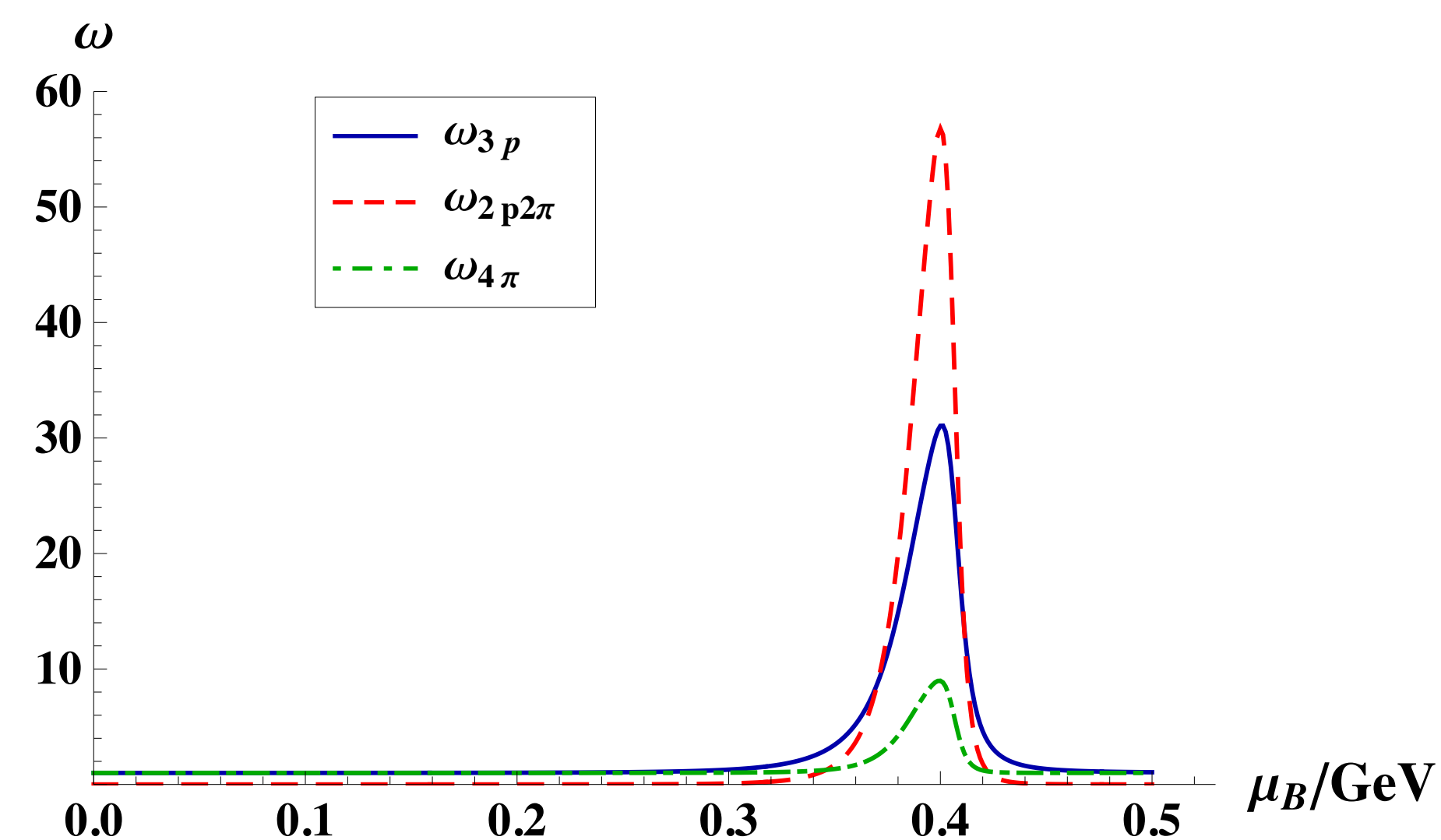
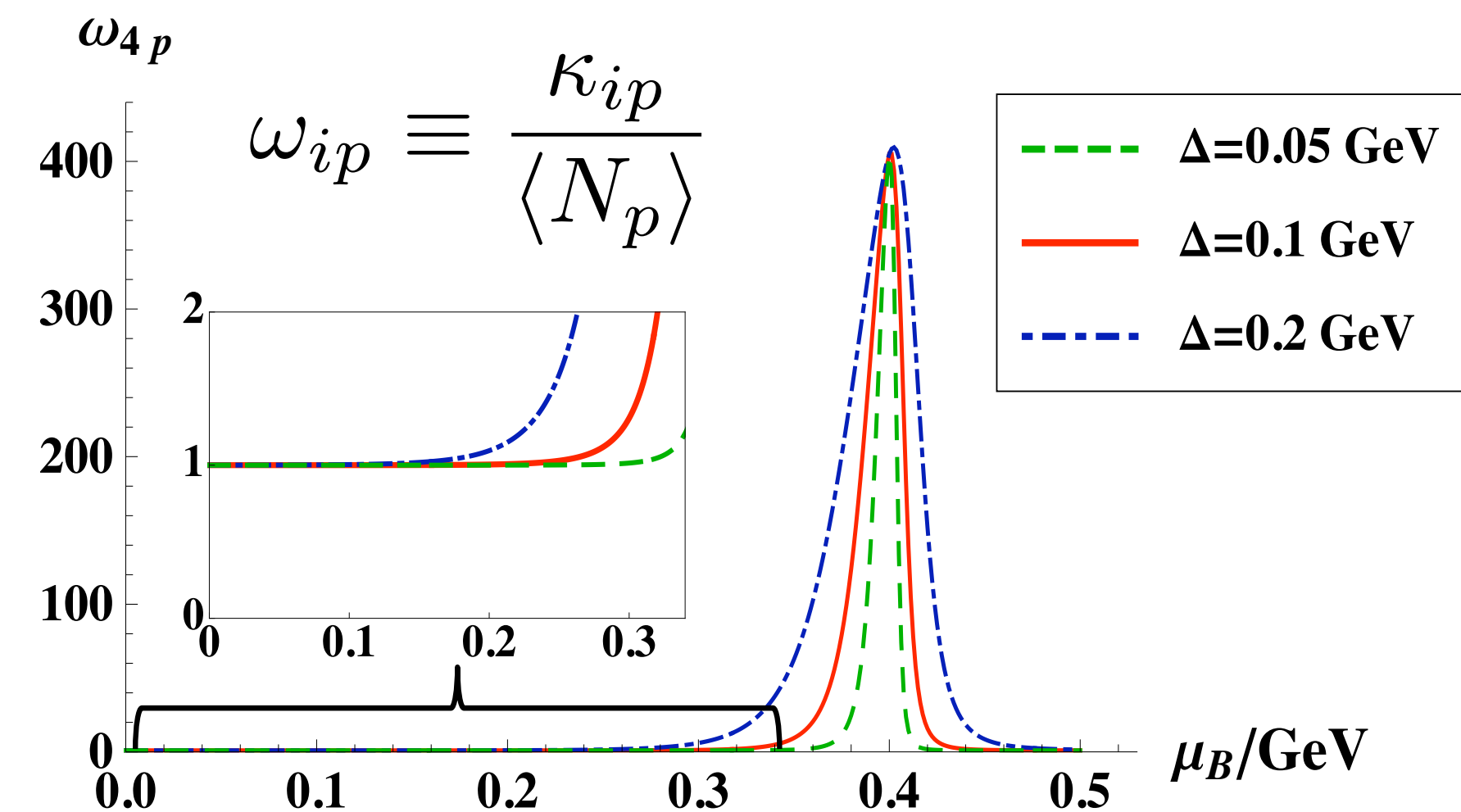
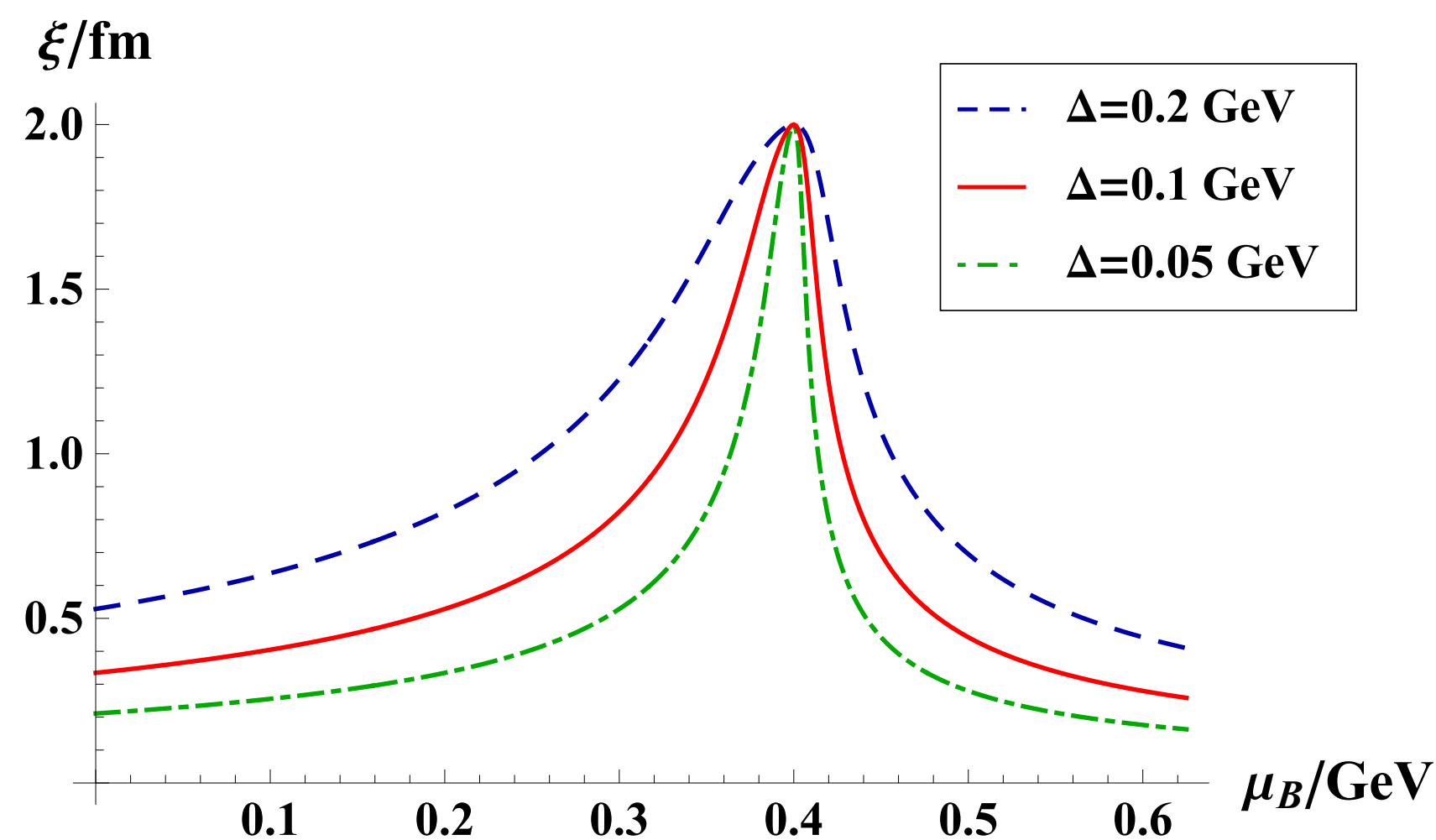
*... + many more*

# Early Estimates of Equilibrium Fluctuations



- Order-of-magnitude predictions of volume-independent normalized cumulants from 2010 relied on these unknown parameters
- Original estimates used parametrized correlation length with width  $\Delta$

$$\xi(\mu_B) = \frac{\xi_{\max}}{\left[1 + \frac{(\mu_B - \mu_B^c)^2}{W(\mu_B)^2}\right]^{1/3}}$$

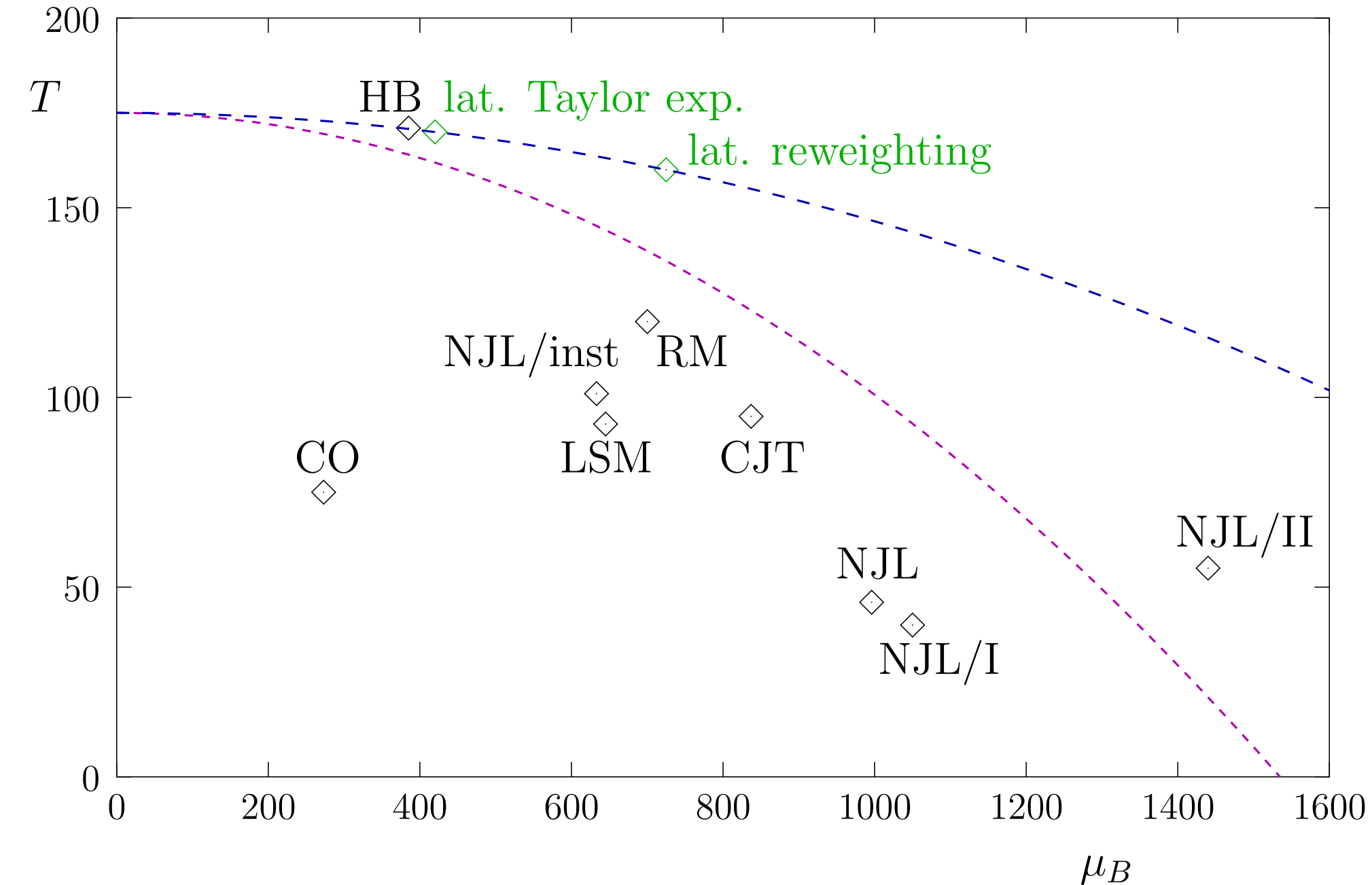


*C. Athanasiou, K. Rajagopal, M. Stephanov, PRD (2010)*

# Brief Aside: Critical Point from LSM



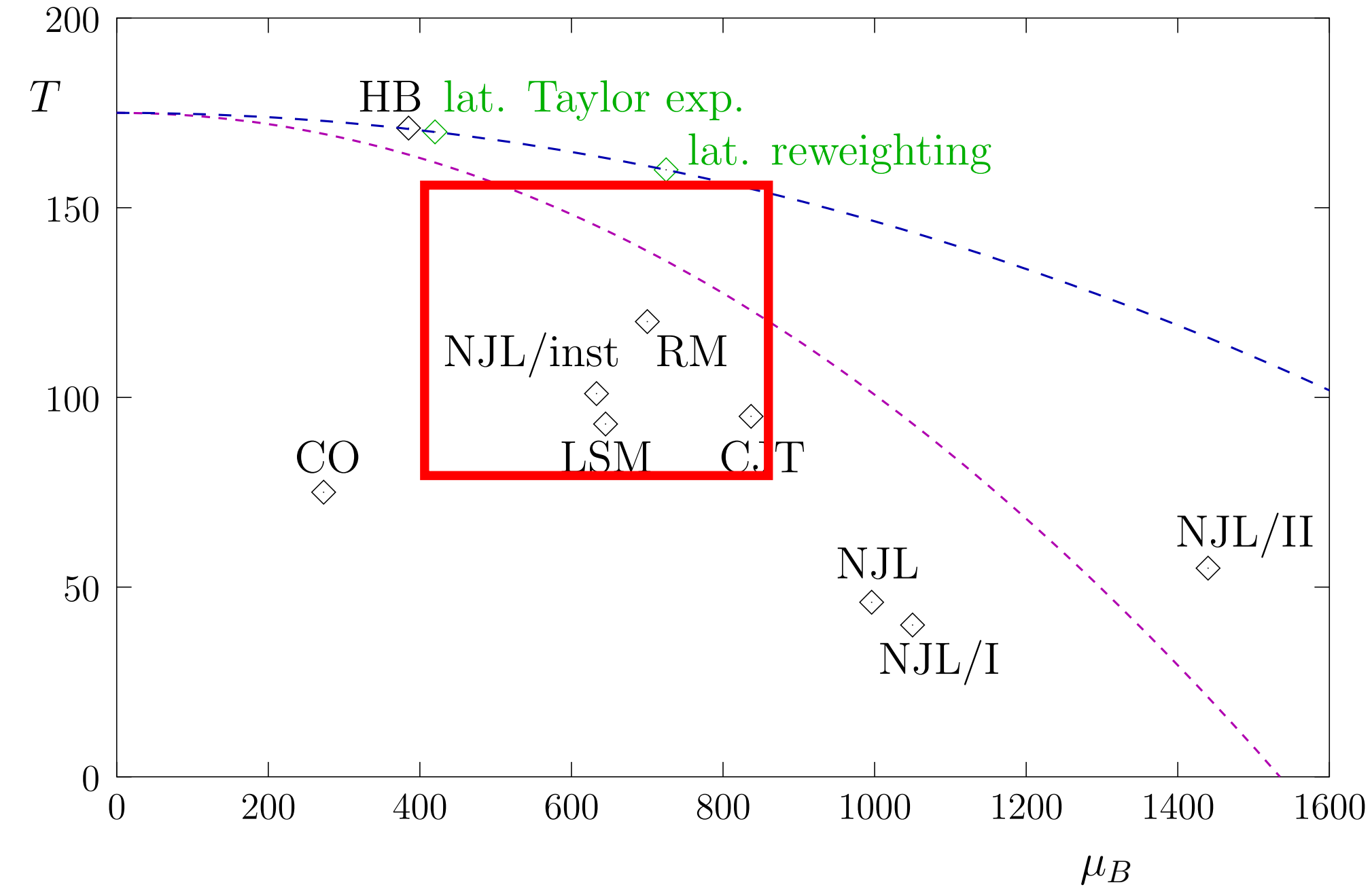
- Early calculations with two light flavored quarks obtained critical point location within linear sigma model



# Brief Aside: Critical Point from LSM



- Early calculations with two light flavored quarks obtained critical point location within linear sigma model

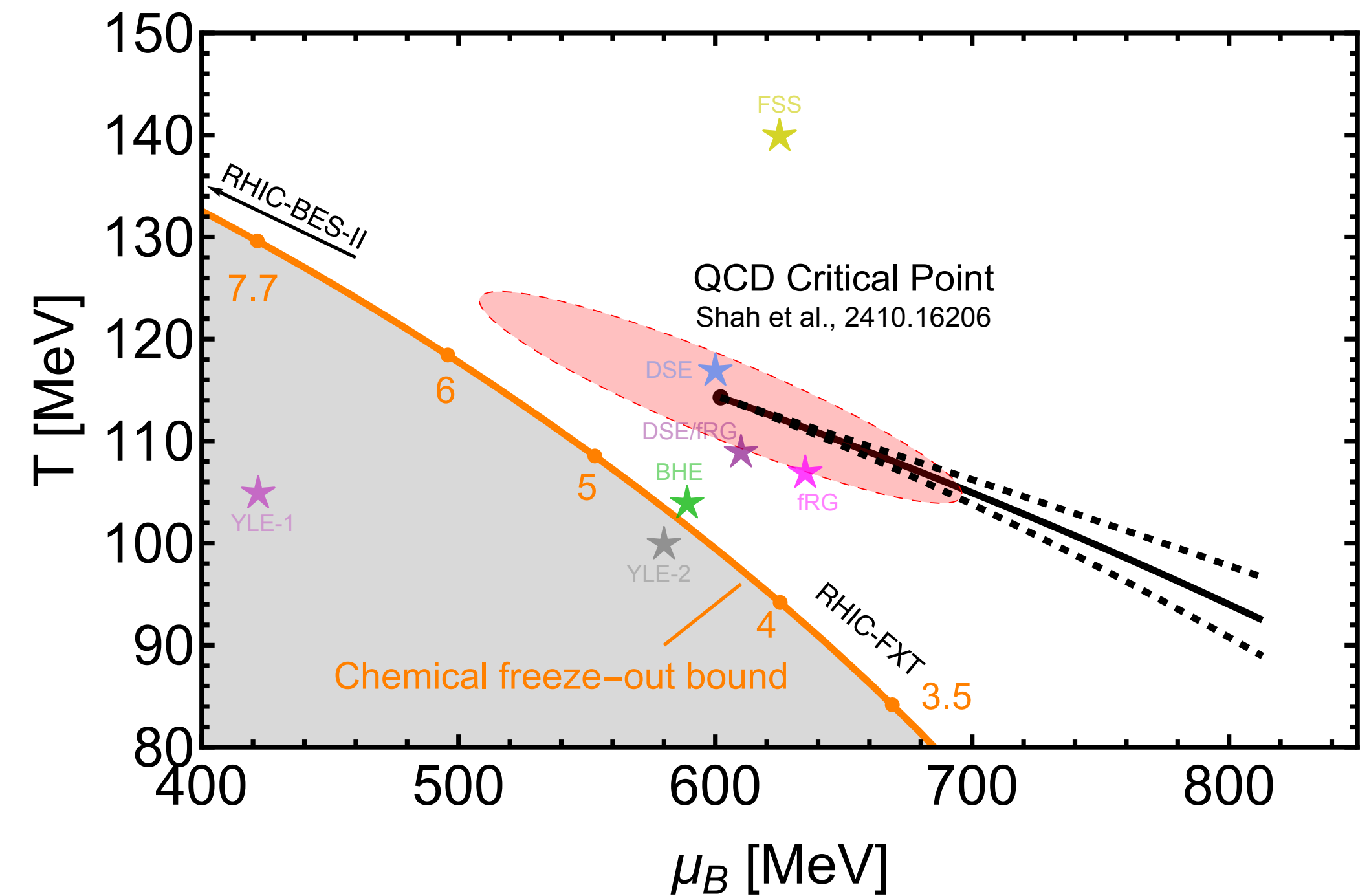
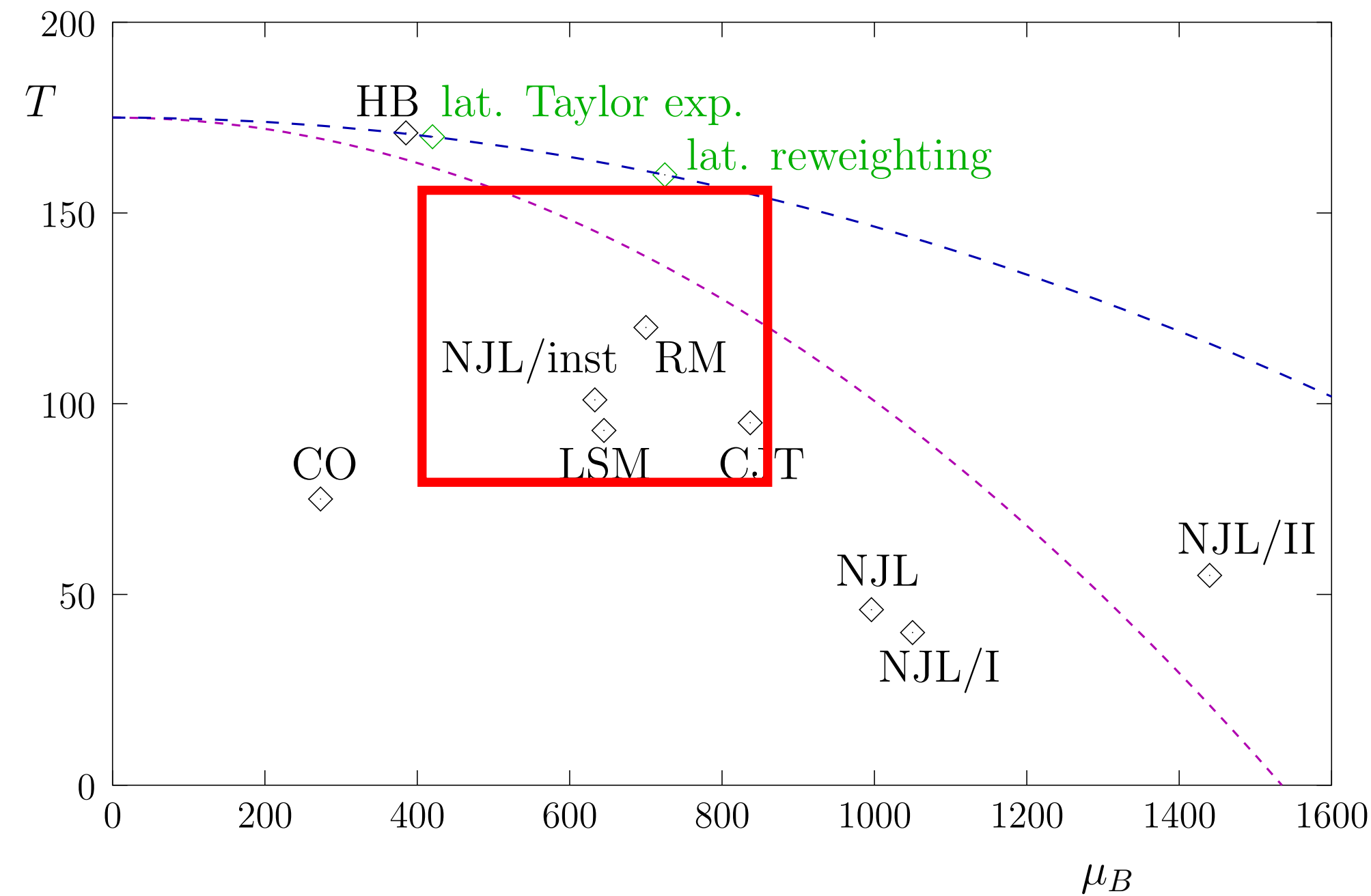


*M. Stephanov, Prog.Theor.Phys.Suppl. (2004)*  
*O. Scavenius et al, PRC (2001)*

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*M. Stephanov, Prog.Theor.Phys.Suppl. (2004)*

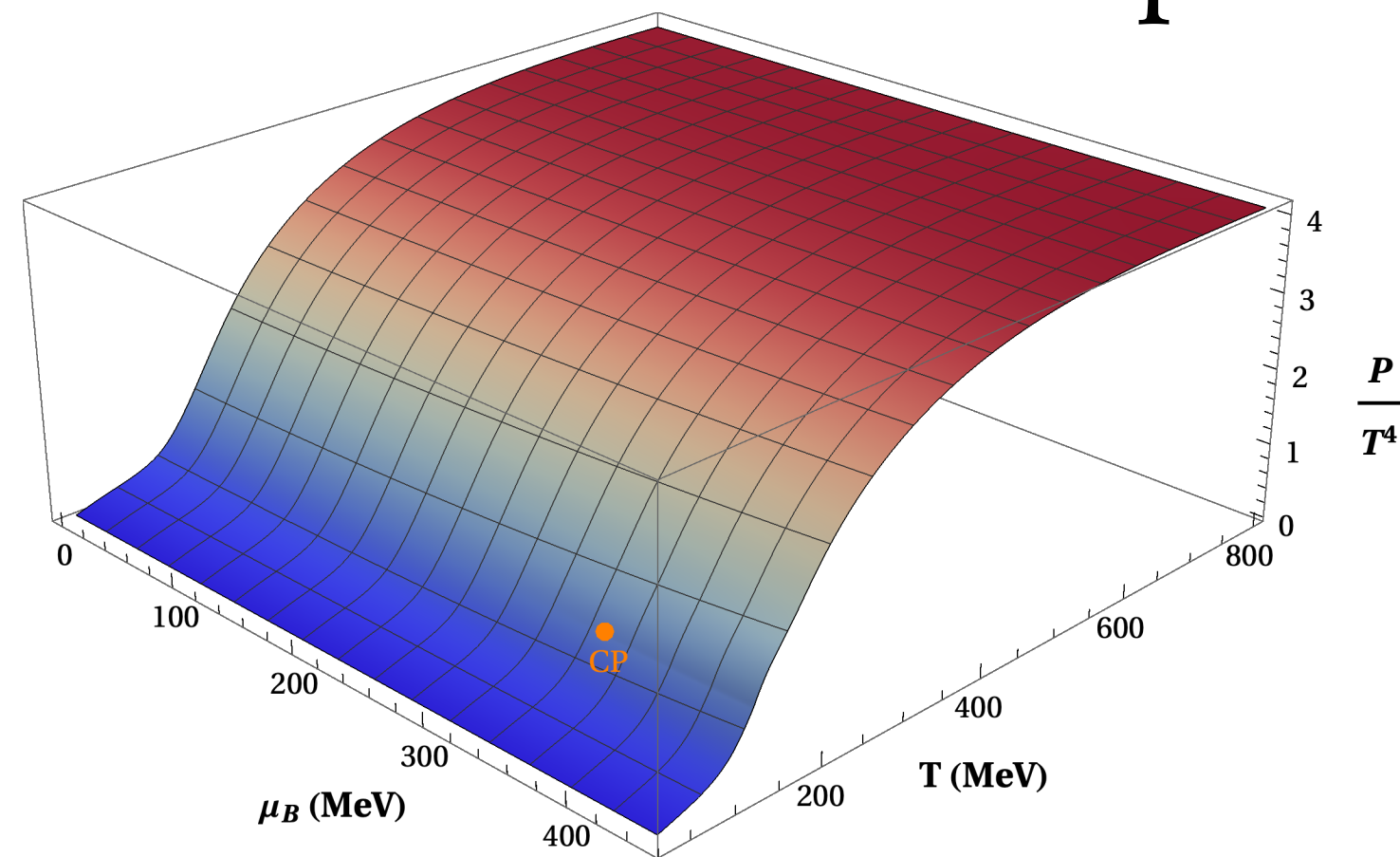
*O. Scavenius et al, PRC (2001)*

*V. Vovchenko and V. Koch, J.Subatomic Part.Cosmol. (2025)*

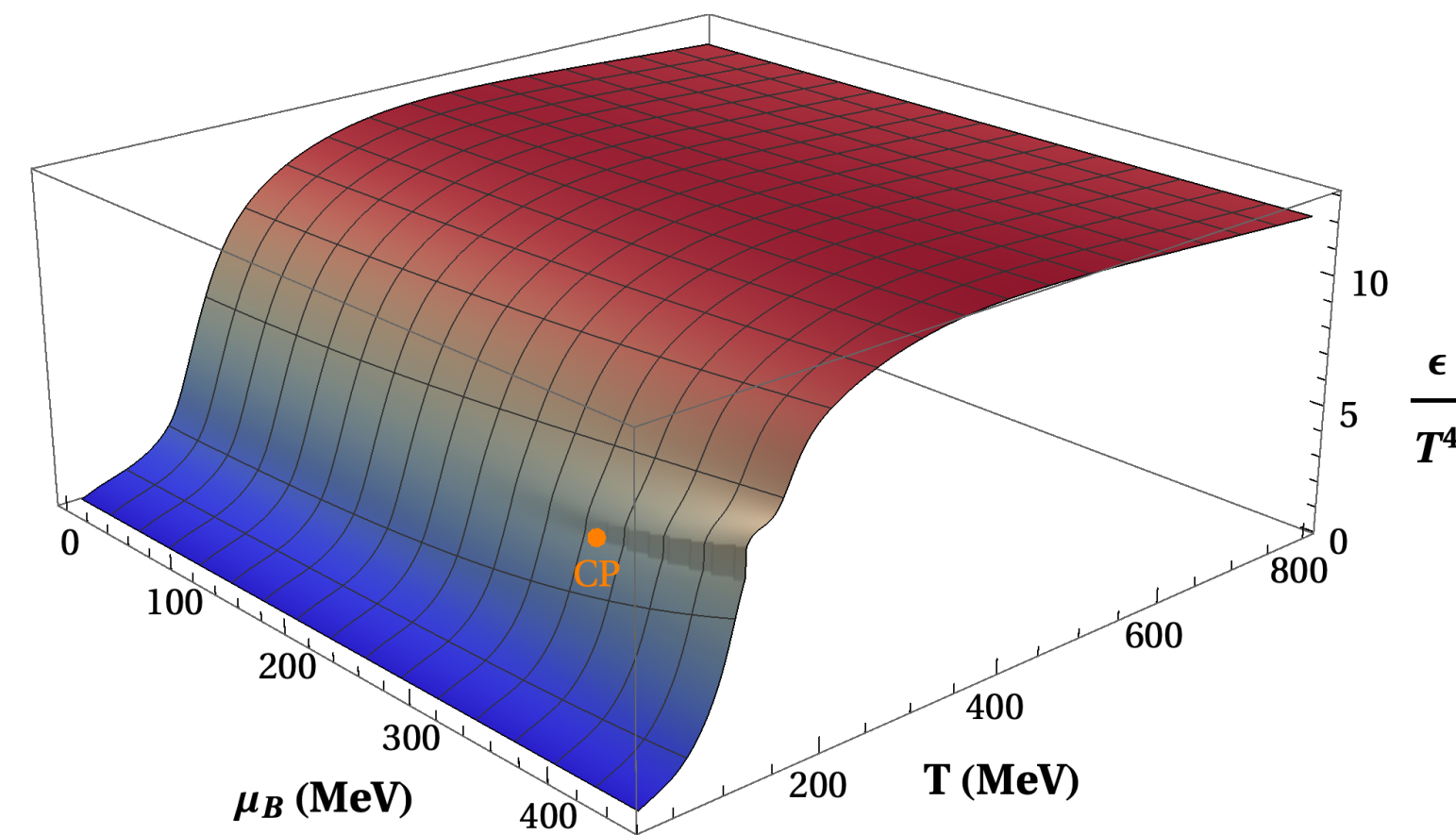
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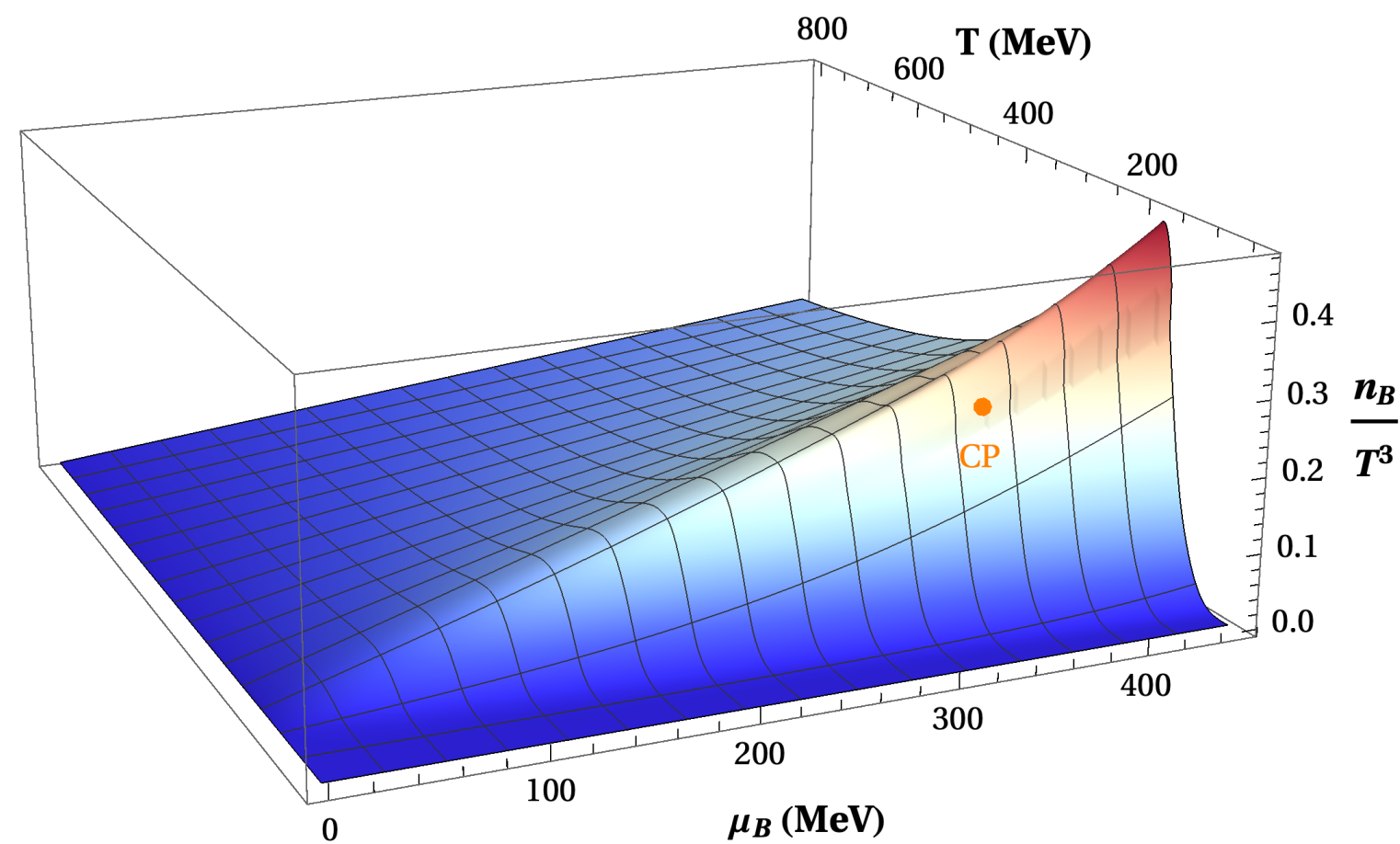
... in a model independent way



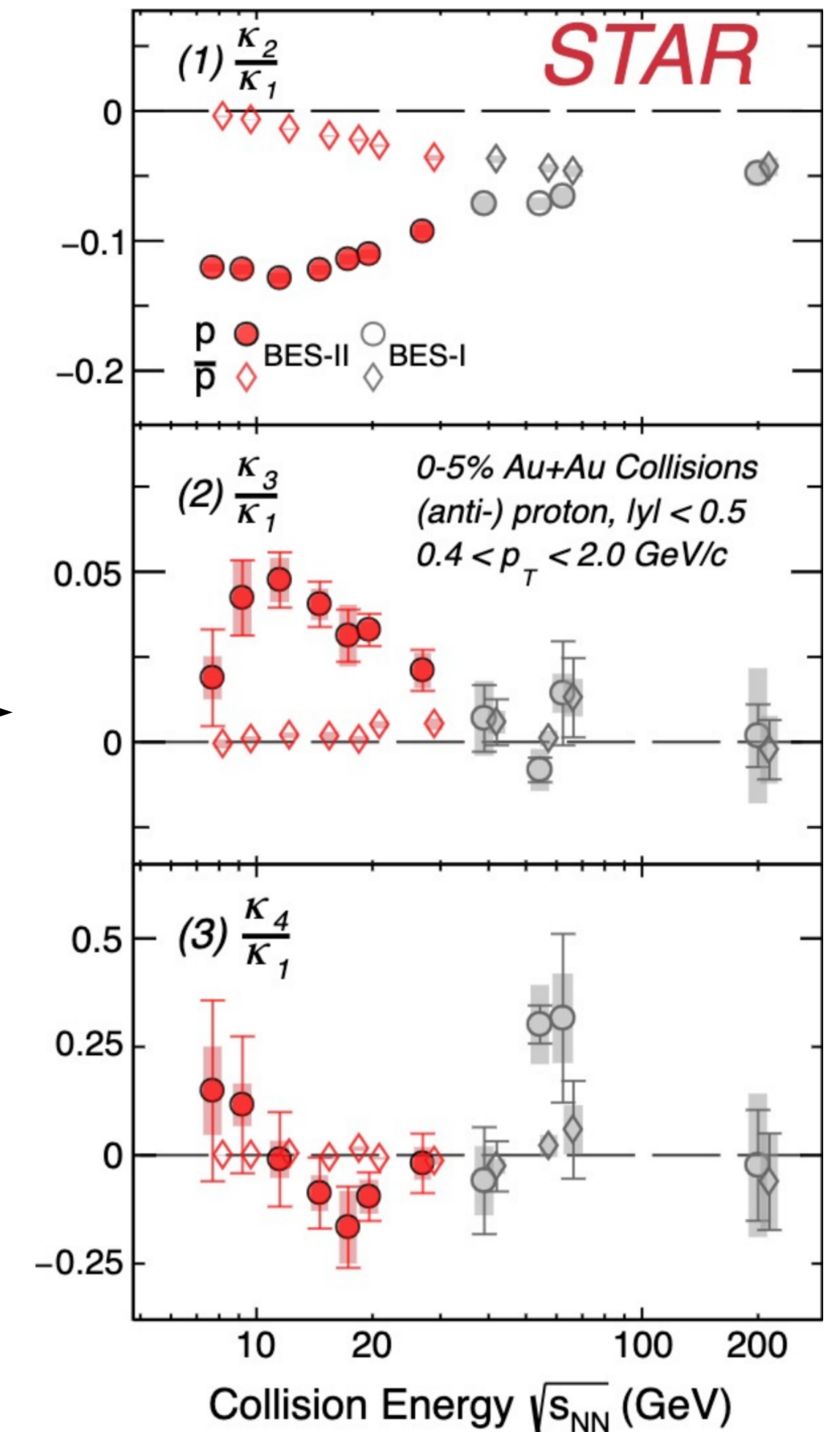
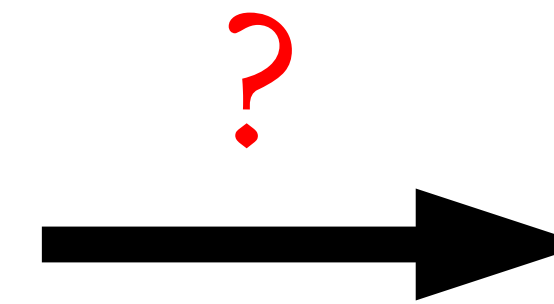
Pressure



Energy density



Baryon density



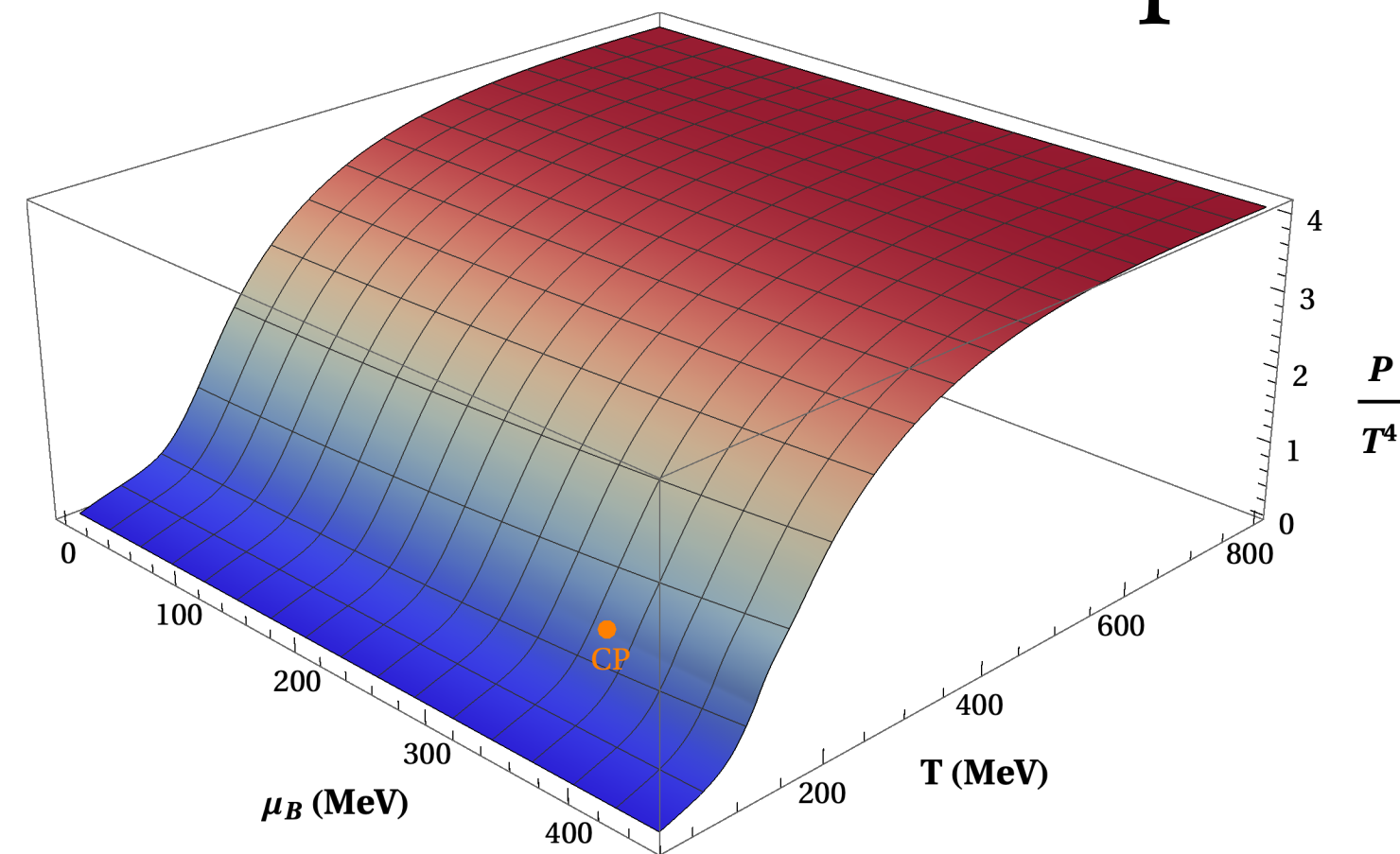
JMK et al, EPJ+ (2021)

A. Pandav (STAR collaboration), CPOD 2024

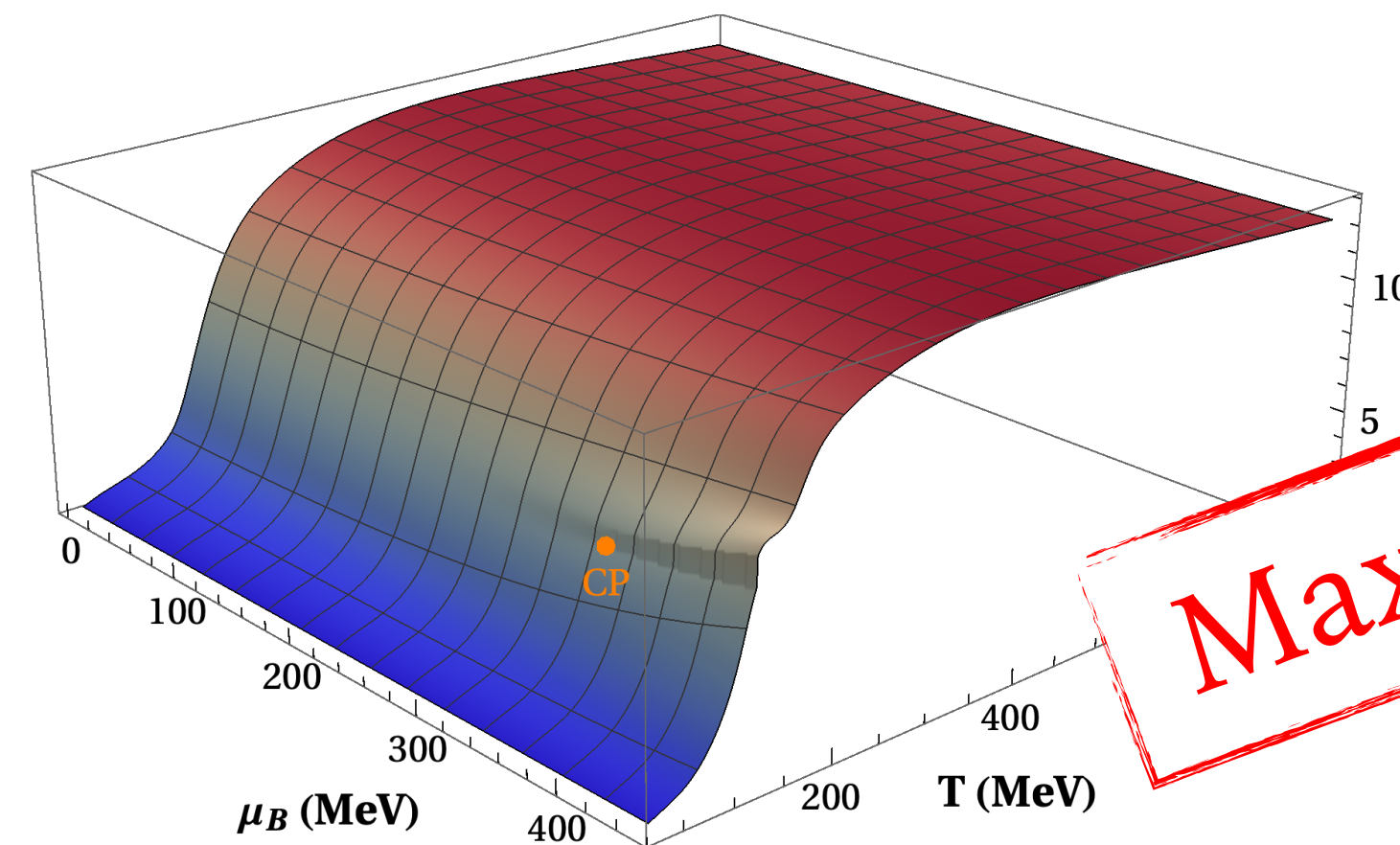
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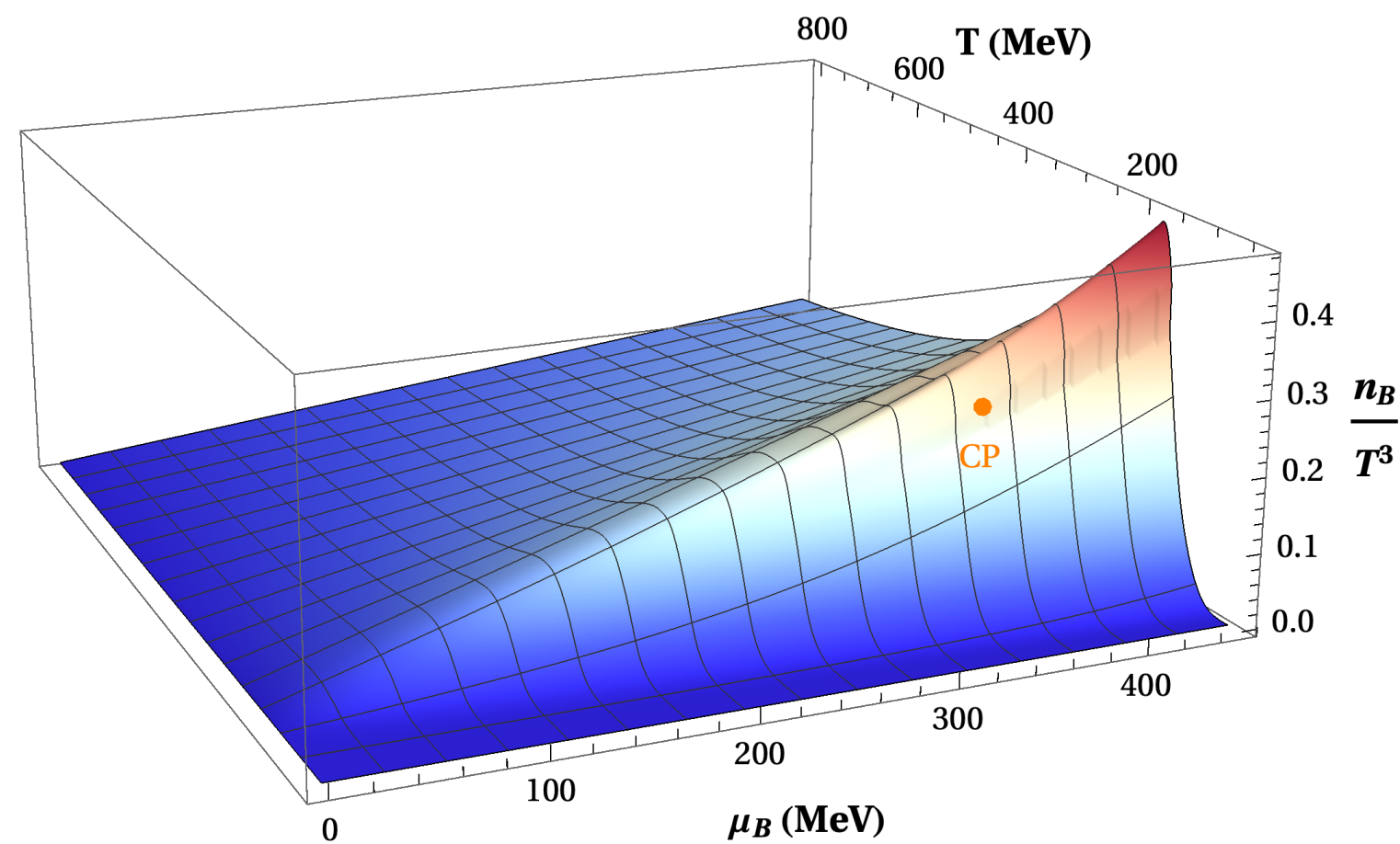
... in a model independent way



Pressure

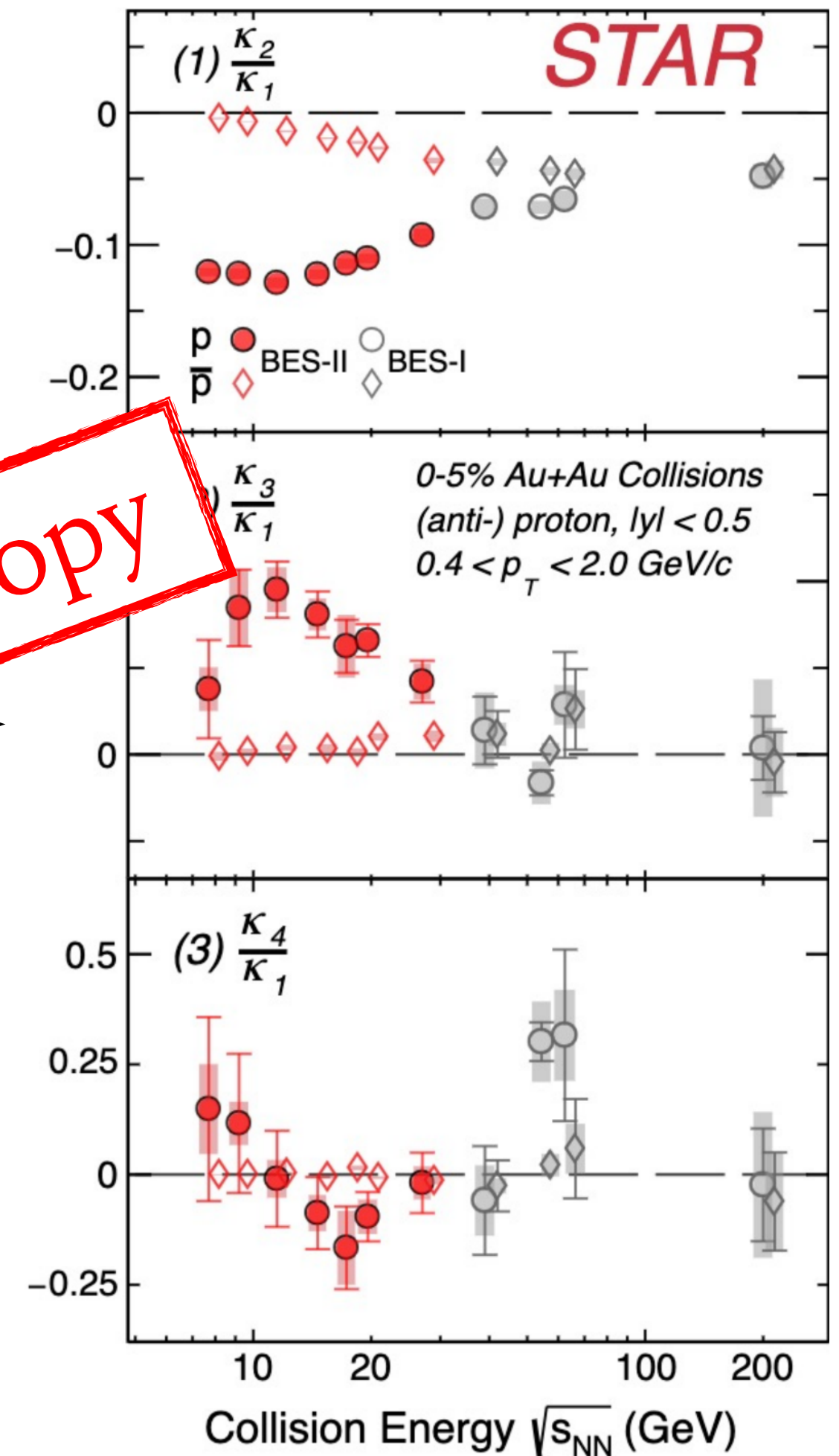


Energy density



Baryon density

**Maximum Entropy**



JMK et al, EPJ+ (2021)  
A. Pandav (STAR collaboration), CPOD 2024

## **II. A first study of equilibrium proton factorial cumulants with MaxEnt**

---

# Connecting Hydro & Particle Fluctuations



## Maximum entropy freezeout of hydrodynamic fluctuations

Maneesha Sushama Pradeep<sup>1</sup> and Mikhail Stephanov<sup>1</sup>

<sup>1</sup>*Department of Physics, University of Illinois, Chicago, IL 60607, USA*

(Dated: May 17, 2023)

We propose a general approach to freezing out *fluctuations* in heavy-ion collisions using the principle of maximum entropy. We find the results naturally expressed as a direct relationship between the *irreducible relative correlators* quantifying the deviations of hydrodynamic as well as hadron gas fluctuations from the ideal hadron gas baseline. The method also allows us to determine heretofore unknown parameters crucial for the freezeout of fluctuations near the QCD critical point in terms of the QCD equation of state.

$$\hat{\Delta}G_{A_1 \dots A_k} = \hat{\Delta}\mathcal{H}_{a_1 \dots a_k} \prod_{i=1}^k P_{A_i}^{a_i}$$

↑ Irreducible relative cumulant (IRC)
 ↑ (Particle) Gas
 ↑ Hydro (fields)
 ↑ Matching conditions projectors

Least-biased determination of particle multiplicity fluctuations consistent with hydrodynamic correlations

*M. Pradeep and M. Stephanov, PRL (2023)*

# Connecting Hydro & Particle Fluctuations



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*M. Pradeep and M. Stephanov, PRL (2023)*

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$$\times \left[ n_{\mathbf{k}}^+ (1 - n_{\mathbf{k}}^+) - n_{\mathbf{k}}^- (1 - n_{\mathbf{k}}^-) \right]$$

Extract unknown coupling in terms of EoS properties:

$$g_A = \sqrt{Z} \frac{E_A}{m_A} (\bar{H}^{-1})_{mm} \frac{w_c}{n_c} \left( \frac{E_A}{w_c} - \frac{q_A}{n_c} \right), \quad \sqrt{Z} \propto \frac{\sin(\alpha_1)}{w \sin(\alpha_1 - \alpha_2)}$$

*M. Pradeep and M. Stephanov, PRL (2023)*

# Connecting Hydro & Particle Fluctuations



## Maximum entropy freezeout of hydrodynamic fluctuations

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what about correlation length? ( $\xi = 1/m_\sigma$ )

Extract unknown coupling in terms of EoS properties:

$$g_A = \sqrt{Z} \frac{E_A}{m_A} (\bar{H}^{-1})_{mm} \frac{w_c}{n_c} \left( \frac{E_A}{w_c} - \frac{q_A}{n_c} \right), \quad \sqrt{Z} \propto \frac{\sin(\alpha_1)}{w \sin(\alpha_1 - \alpha_2)}$$

M. Pradeep and M. Stephanov, PRL (2023)

# Correlation Length



- Derive correlation length dependence on mapping parameters using universal scaling form with approximate rational critical exponents

plot  $\xi^2/w^2$  to relate to  $p$  fluctuations:

$$\langle \delta n_p \delta n_p \rangle =$$

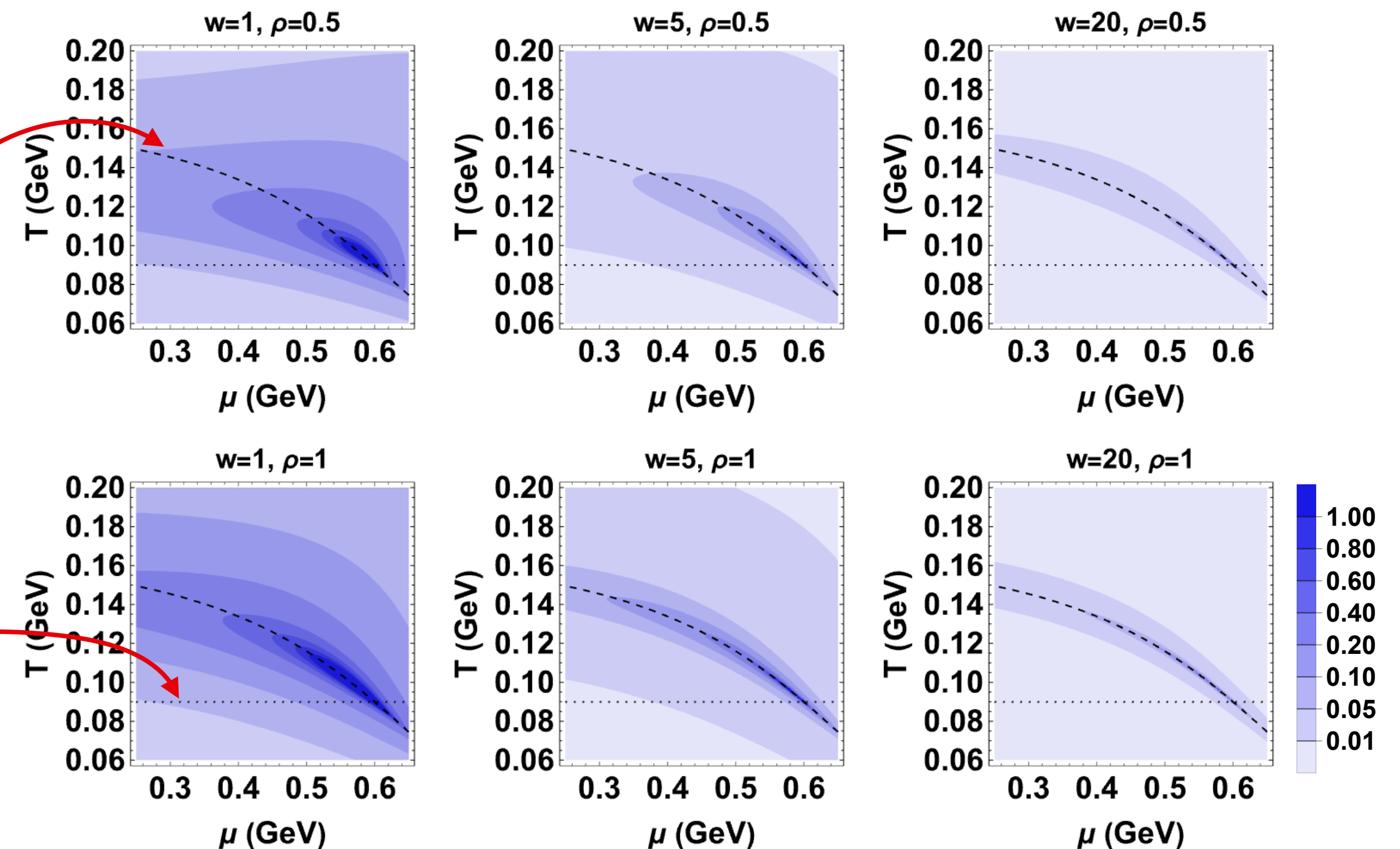
$$\langle \delta n_p^2 \rangle \propto g_p^2 \xi^{2-\eta}$$

Along  $r$ -axis:

$$\frac{\xi_{\text{QCD}}^2}{w^2} \propto \frac{1}{w^{2/3}} \left( \frac{\rho}{\mu_c^2 - \mu^2} \right)^{4/3}$$

Along  $h$ -axis:

$$\frac{\xi_{\text{QCD}}^2}{w^2} \propto \frac{1}{w^{6/5} (\mu_c^2 - \mu^2)^{4/5}}$$



$$\mu_c = 600 \text{ MeV}, \alpha_2 = 0^\circ (\alpha_1 = 16.6^\circ, T_c = 90 \text{ MeV})$$

JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Required Input for MaxEnt: EoS



- Maximum entropy does not require unknown parameters - takes only EoS as input
  - Because the EoS is not known yet, it has its own free parameters, e.g. mapping parameters from BEST EoS, for which we make a choice
    - $\mu_c = 600 \text{ MeV}, \alpha_2 = 0^\circ$
    - $(\alpha_1 = 16.6^\circ, T_c = 90 \text{ MeV})$
    - $w = \{1, 5, 20\} \rho = \{0.5, 1, 2\}$

# Required Input for MaxEnt: EoS

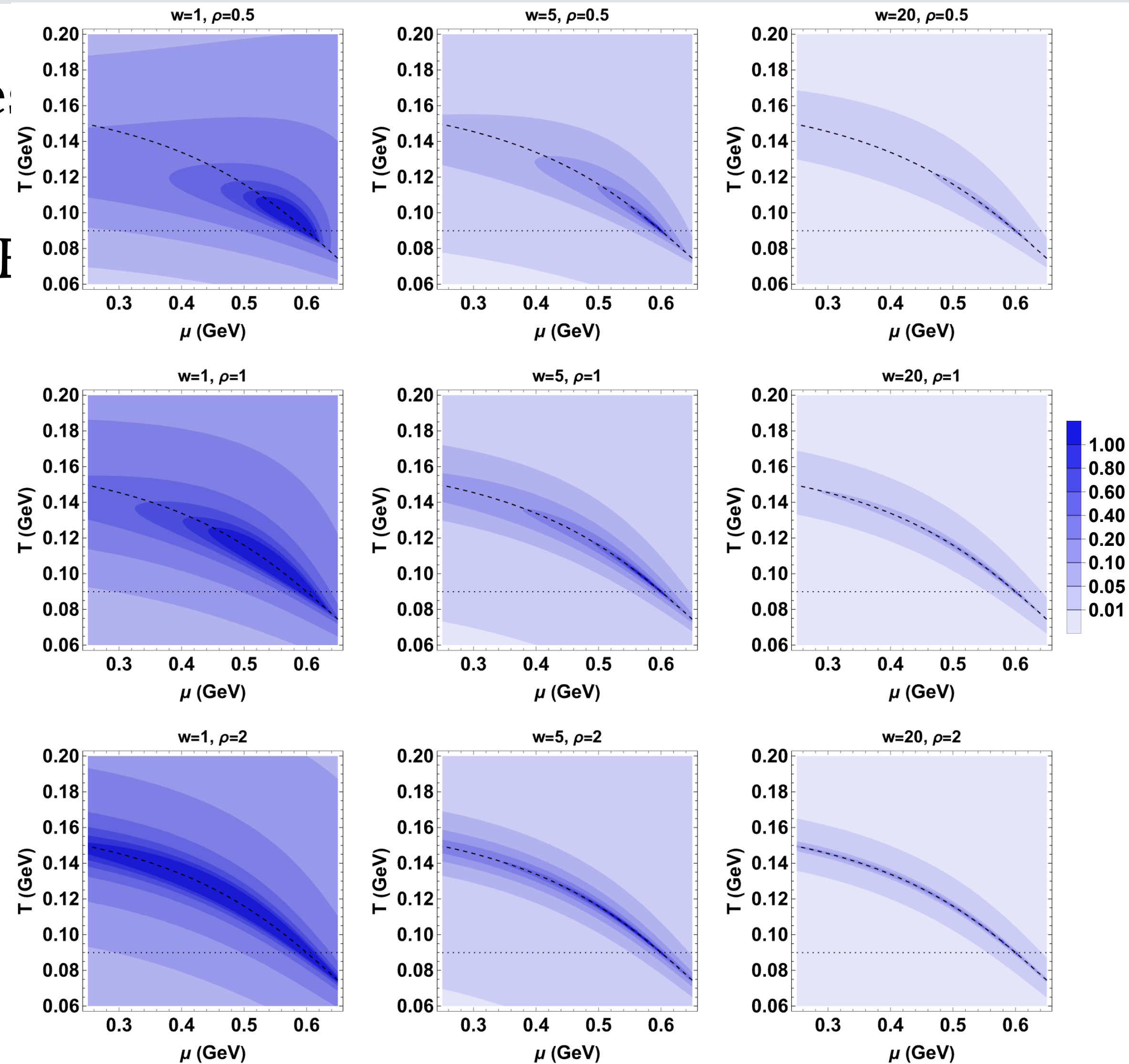
- Maximum entropy does not require any EoS as input
- Because the EoS is not known, we use a set of parameters from BI

$$\hat{\Delta}G_{A_1 \dots A_k} = \hat{\Delta}\mathcal{H}_{a_1 \dots a_k} \prod_{i=1}^k P_{A_i}^{a_i}$$

matches fluctuations at all orders

$$\Delta H_{kn} \equiv \langle \delta n^k \rangle$$

$$k = 2$$



only EoS as input  
 e.g. mapping  
 $(\mu, T) \rightarrow (T_c, \alpha_2)$   
 $(\mu, T) = (0.5, 0.15) \rightarrow (90 \text{ MeV}, \alpha_2 = 0)$   
 $(\mu, T) = (0.6, 0.1) \rightarrow (90 \text{ MeV}, \alpha_2 = 0.5)$   
 $(\mu, T) = (0.6, 0.07) \rightarrow (90 \text{ MeV}, \alpha_2 = 2)$

JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Required Input for MaxEnt: EoS

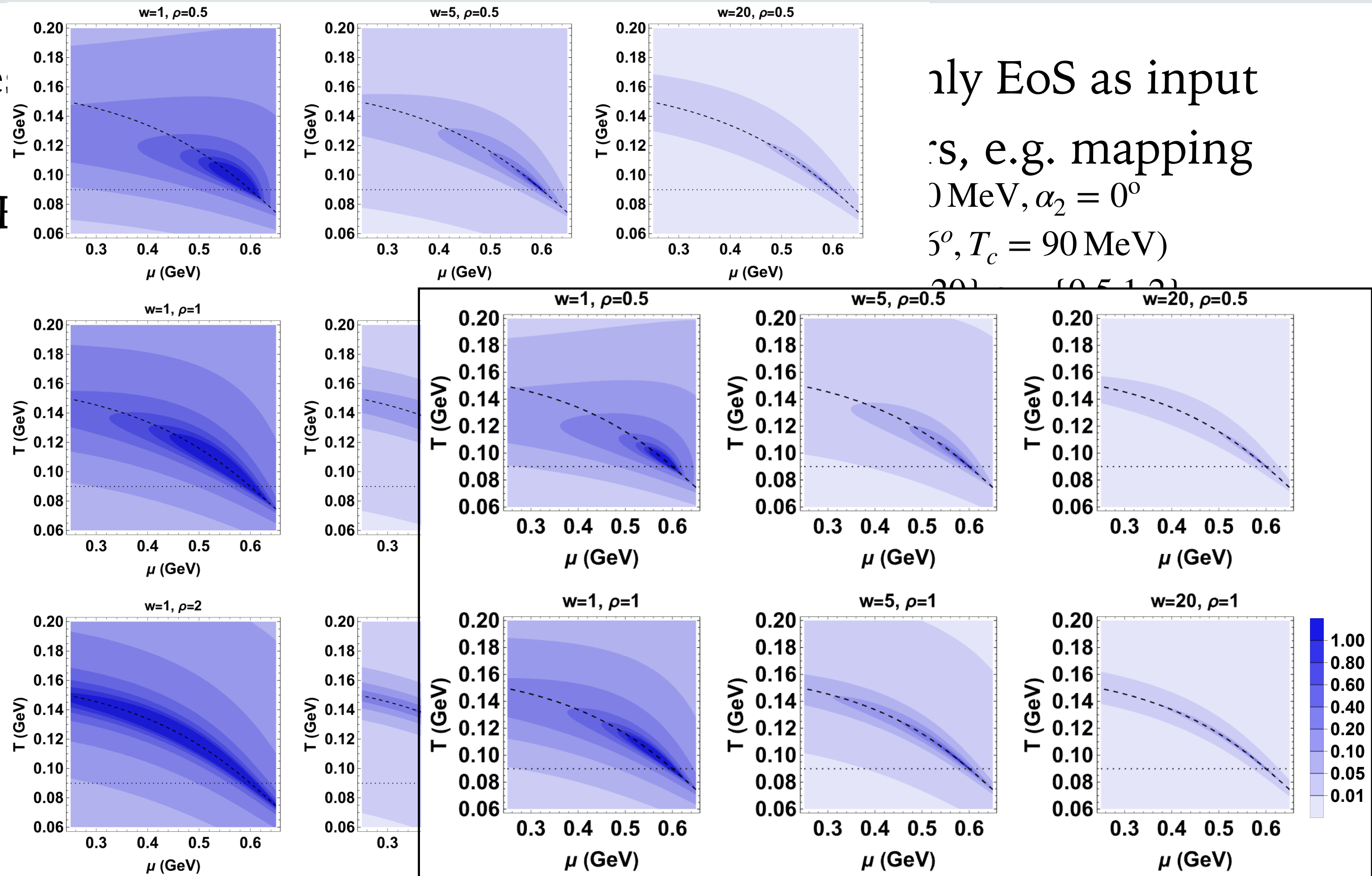
- Maximum entropy does not require any EoS as input
- Because the EoS is not known, e.g. mapping  $(\mu, T) \rightarrow (p, \epsilon)$ , e.g. mapping  $(\mu, T) \rightarrow (p, \epsilon)$  (e.g.  $\alpha_2 = 0^{\circ}$ ,  $\alpha_3 = 5^{\circ}$ ,  $T_c = 90$  MeV)

$$\hat{\Delta}G_{A_1 \dots A_k} = \hat{\Delta}\mathcal{H}_{a_1 \dots a_k} \prod_{i=1}^k P_{A_i}^{a_i}$$

matches fluctuations at all orders

$$\Delta H_{kn} \equiv \langle \delta n^k \rangle$$

$$k = 2$$



JMK, K. Rajagopal, M. Pradheep, M. Stephanov, T. Yin, PRD (2020, Editors' Suggestion)

# Required Input for MaxEnt: EoS

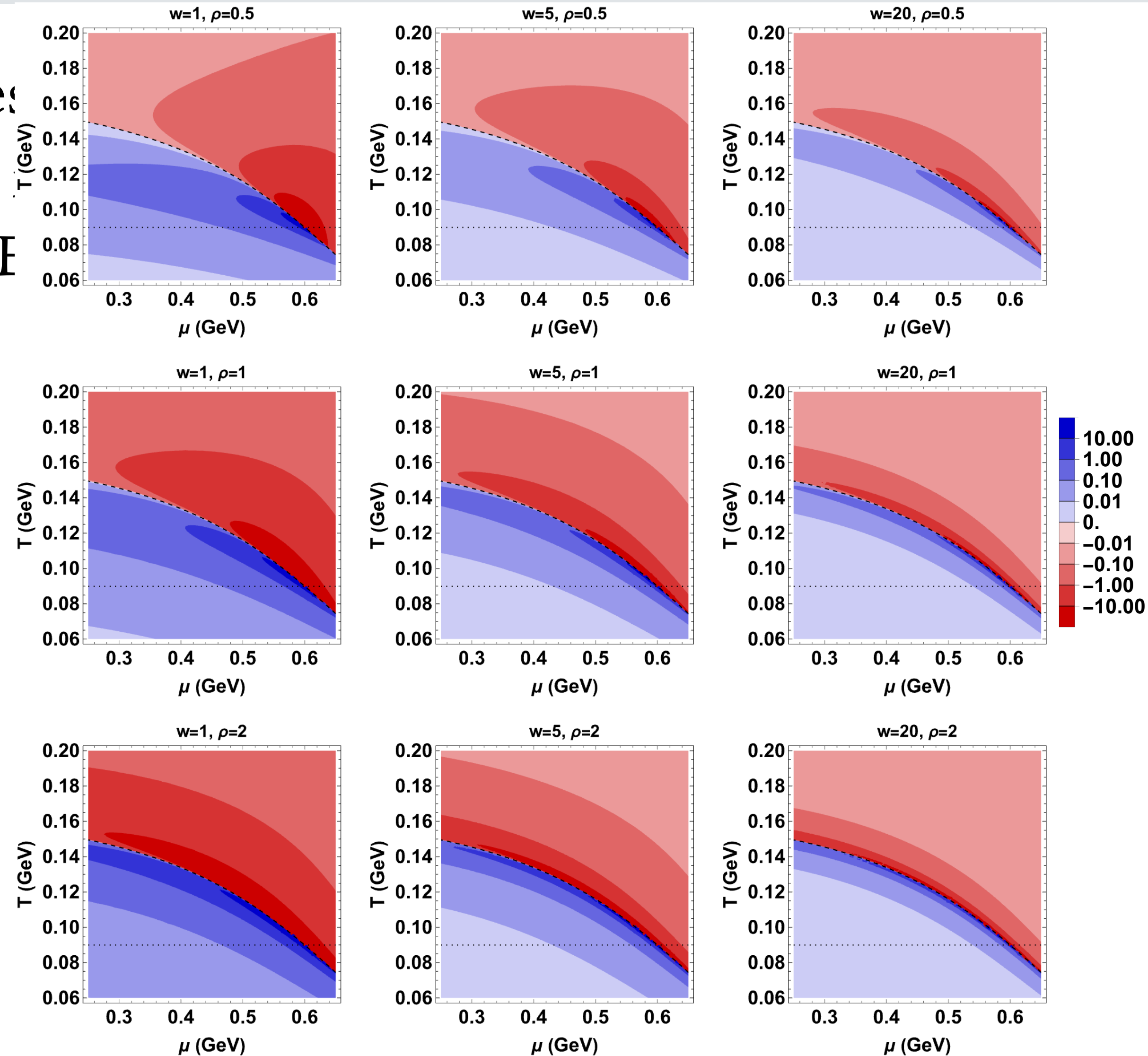
- Maximum entropy does
- Because the EoS is
- parameters from BE

$$\hat{\Delta}G_{A_1 \dots A_k} = \hat{\Delta}\mathcal{H}_{a_1 \dots a_k} \prod_{i=1}^k P_{A_i}^{a_i}$$

matches fluctuations at all orders

$$\Delta H_{kn} \equiv \langle \delta n^k \rangle$$

$$k = 3$$



ly EoS as input  
 s, e.g. mapping  
 MeV,  $\alpha_2 = 0^0$   
 ,  $T_c = 90$  MeV)  
 0}  $\rho = \{0.5, 1, 2\}$

JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Required Input for MaxEnt: EoS

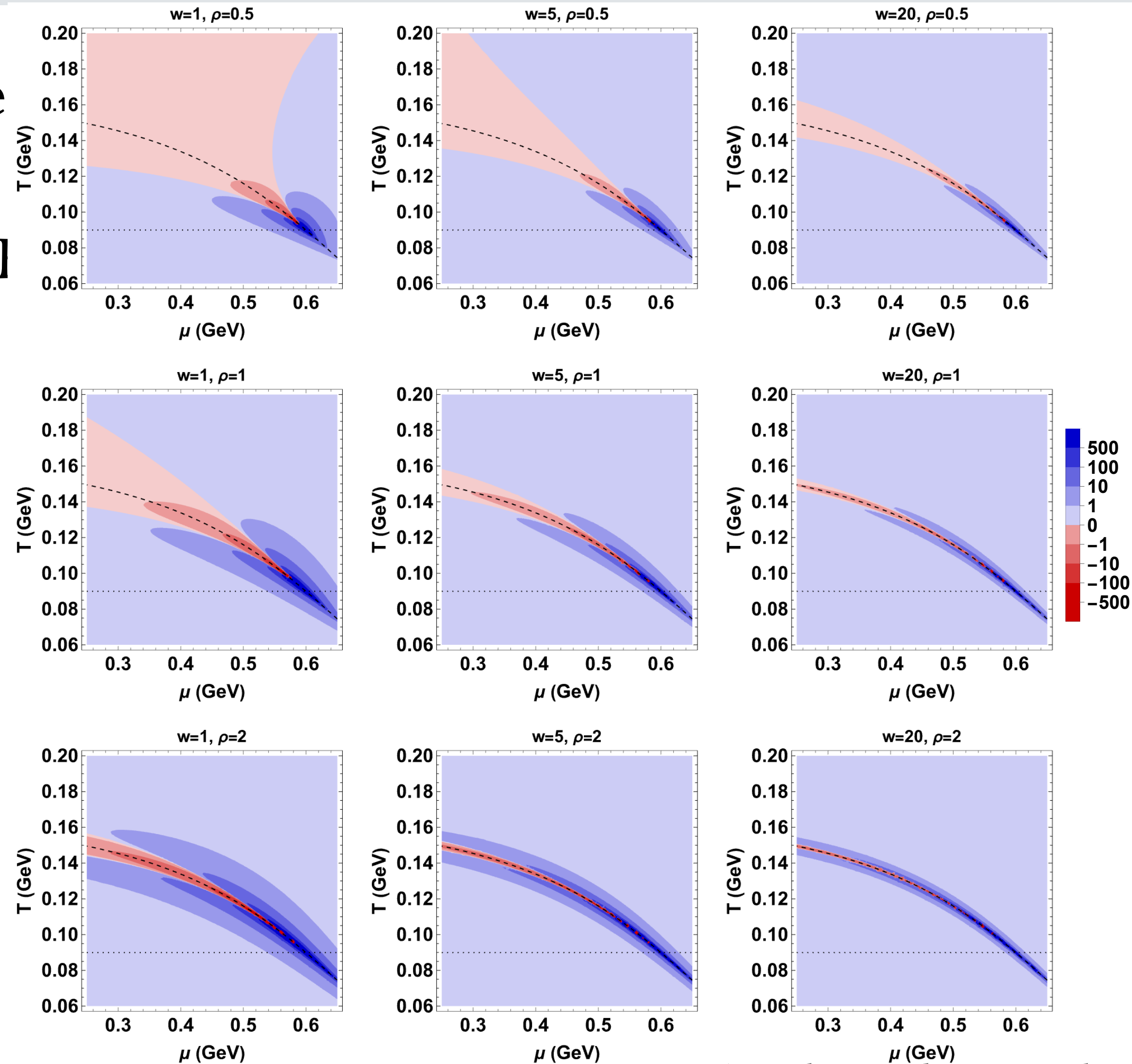
- Maximum entropy does not require any EoS as input
- Because the EoS is not known, we use parameters from BL

$$\hat{\Delta}G_{A_1 \dots A_k} = \hat{\Delta}\mathcal{H}_{a_1 \dots a_k} \prod_{i=1}^k P_{A_i}^{a_i}$$

matches fluctuations at all orders

$$\Delta H_{kn} \equiv \langle \delta n^k \rangle$$

$$k = 4$$



Maximum entropy does not require any EoS as input  
 Because the EoS is not known, we use parameters from BL  
 e.g. mapping  $(\mu, T_c)$  to  $(\mu, T_c)$   
 $(\mu, T_c) = (90 \text{ MeV}, 90 \text{ MeV})$   
 $\rho = \{0.5, 1, 2\}$

JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Quantifying Critical Fluctuation Signatures



- Extract critical contribution to proton multiplicity fluctuations from an Ising-like critical point & study influence of unknown EoS parameters
  - Focus on deviations from Hadron Resonance Gas background of EoS close to critical point where singular part dominates

$$P(\mu, T) = P^{\text{reg}}(\mu, T) + P^{\text{sing}}(\mu, T)$$

HRG

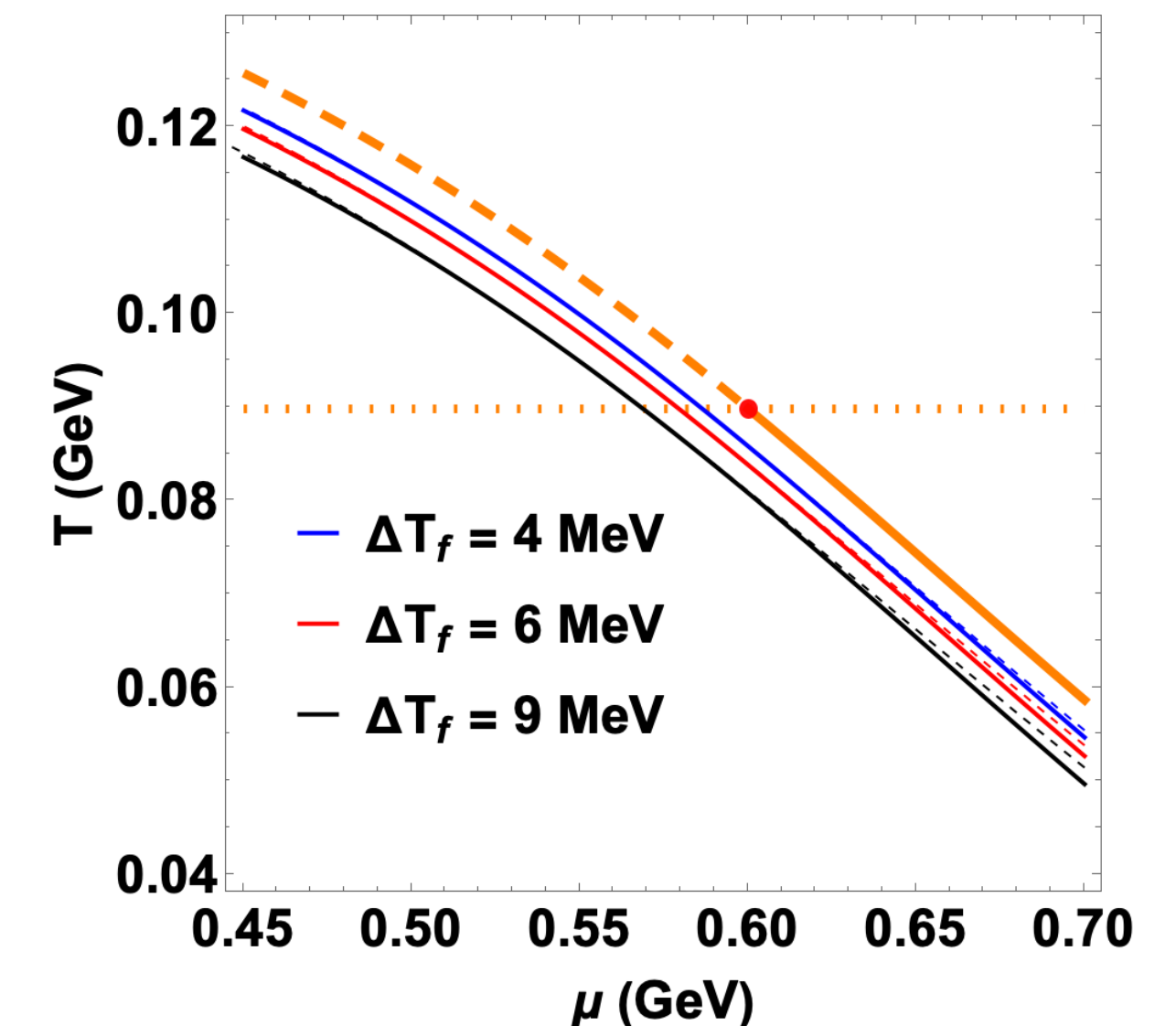
3D Ising model

- Use quadratic map for Ising transition line

*M. Kahangirwe et al, PRD (2024)*

- Calculate multiplicity fluctuations along freeze-out lines

$$T_f(\mu_B) = T_{\text{crossover}}(\mu_B) - \Delta T_f$$



*JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)*

# Proton Factorial Cumulants from MaxEnt



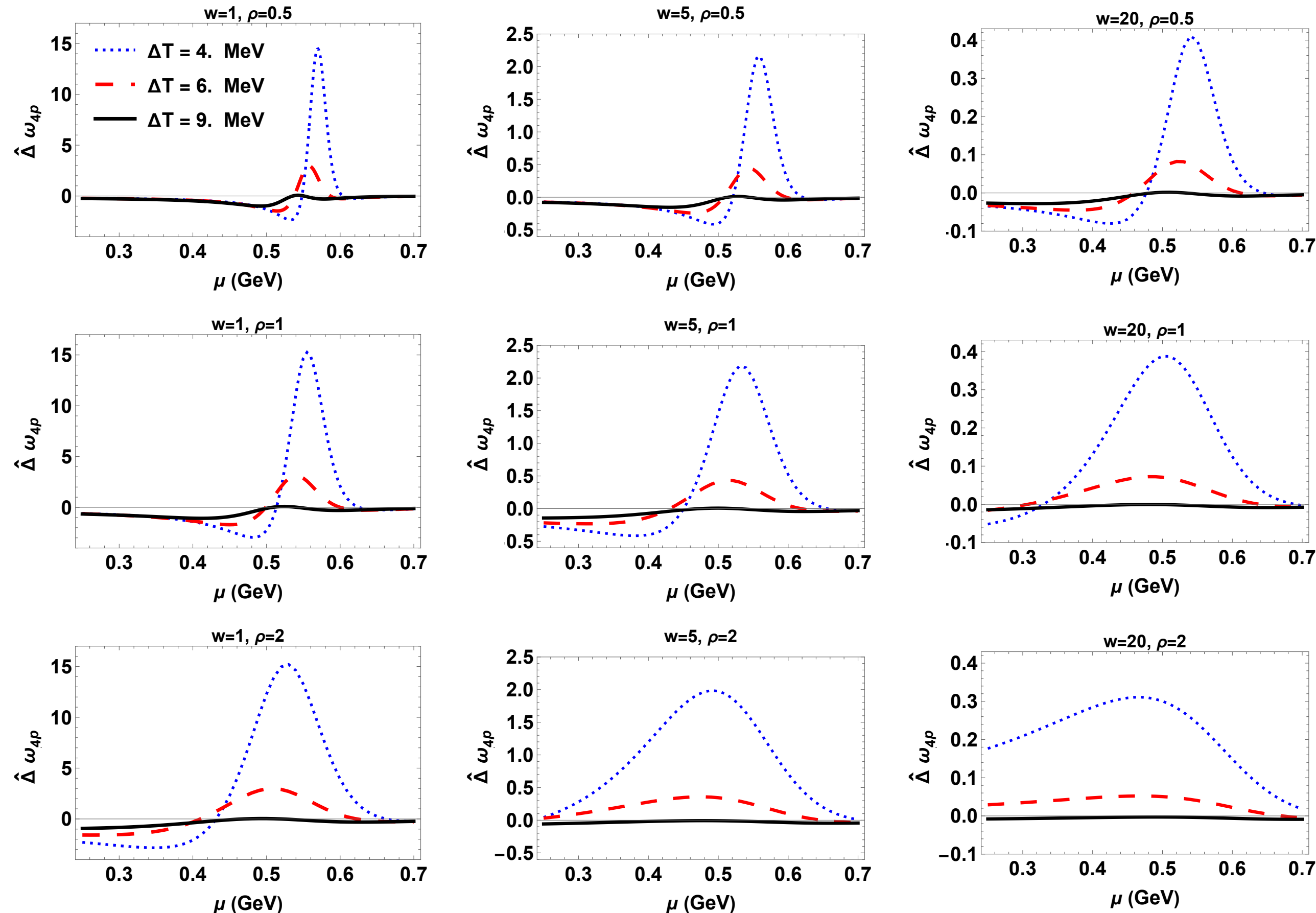
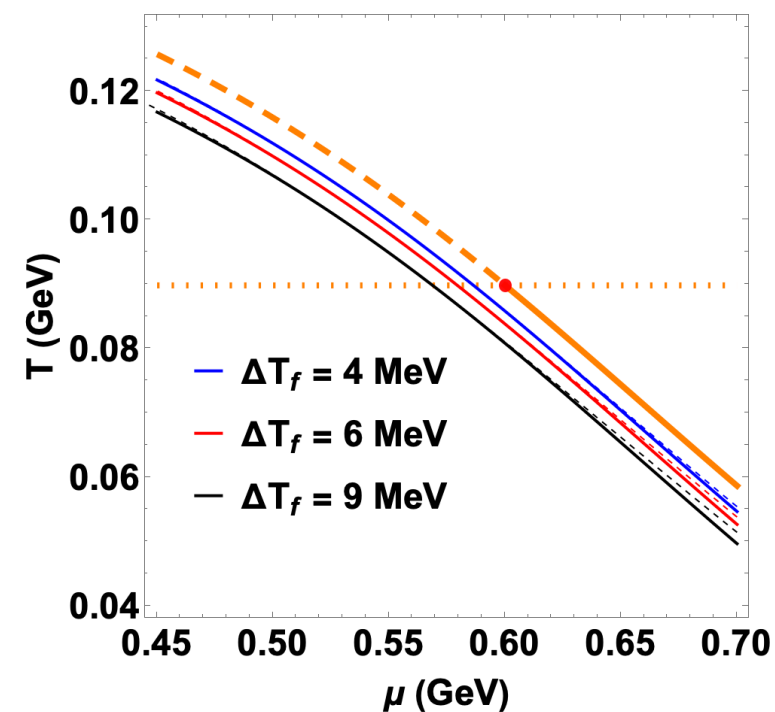
- Critical contribution to fourth order proton factorial cumulants in equilibrium for range of EoS parameters and freeze-out scenarios

Normalized proton factorial cumulants:

$$\hat{\Delta}\omega_{kp} = \frac{\hat{\Delta}H_{a_1 \dots a_k} P_p^{a_1} \dots P_p^{a_k}}{\langle N_p \rangle}$$

$$= \frac{\mathcal{K}_{4p}}{\mathcal{K}_{1p}}$$

Along freeze-out curves:



JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Proton Factorial Cumulants from MaxEnt



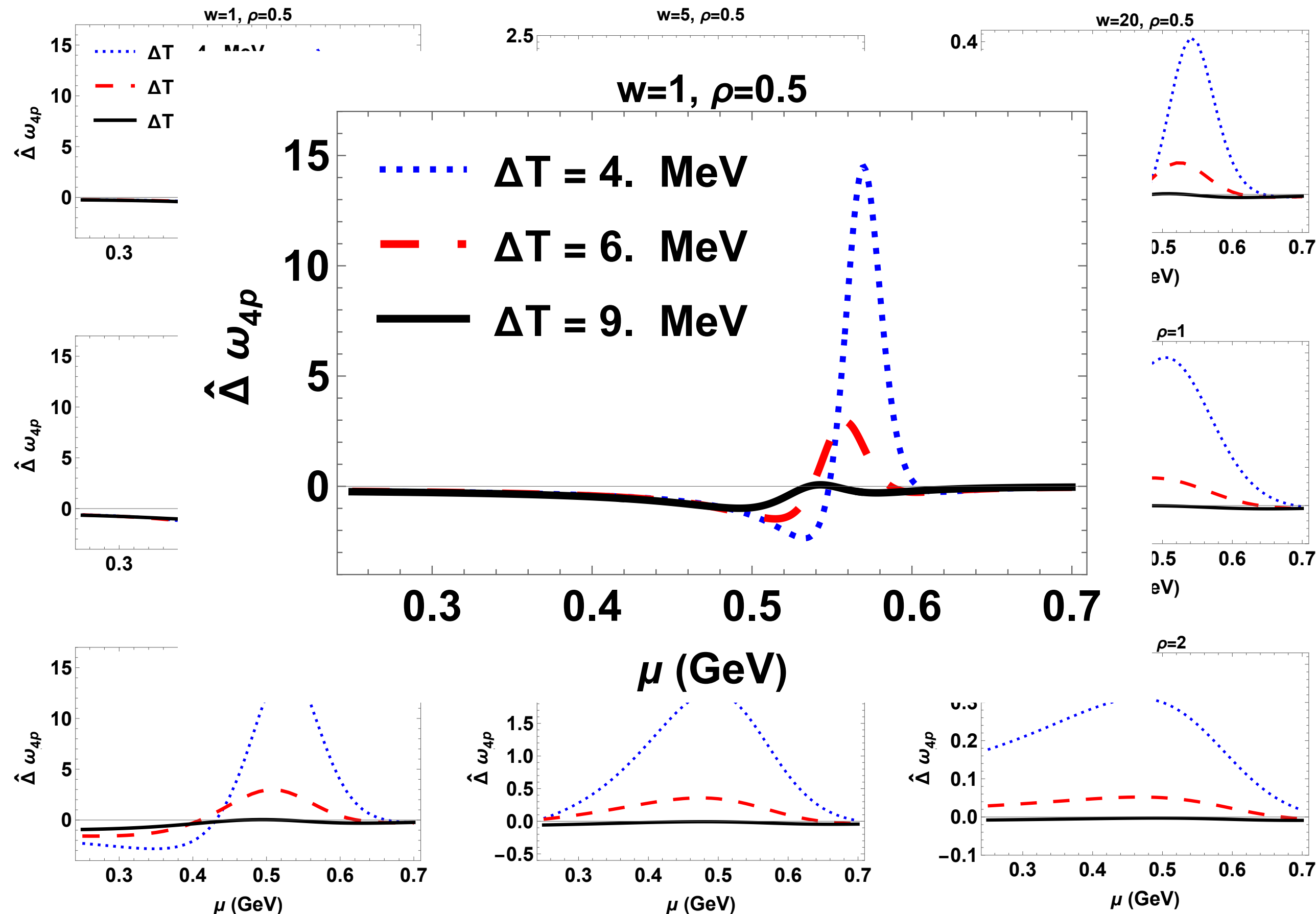
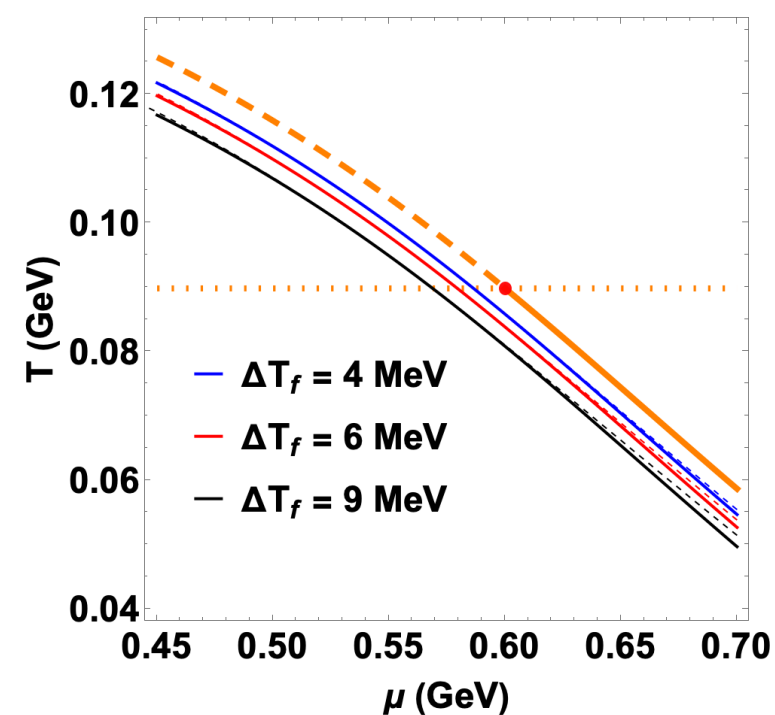
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$$\hat{\Delta}\omega_{kp} = \frac{\hat{\Delta}H_{a_1 \dots a_k} P_p^{a_1} \dots P_p^{a_k}}{\langle N_p \rangle}$$

$$= \frac{\kappa_{4p}}{\kappa_{1p}}$$

Along freeze-out curves:



JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Proton Factorial Cumulants from MaxEnt



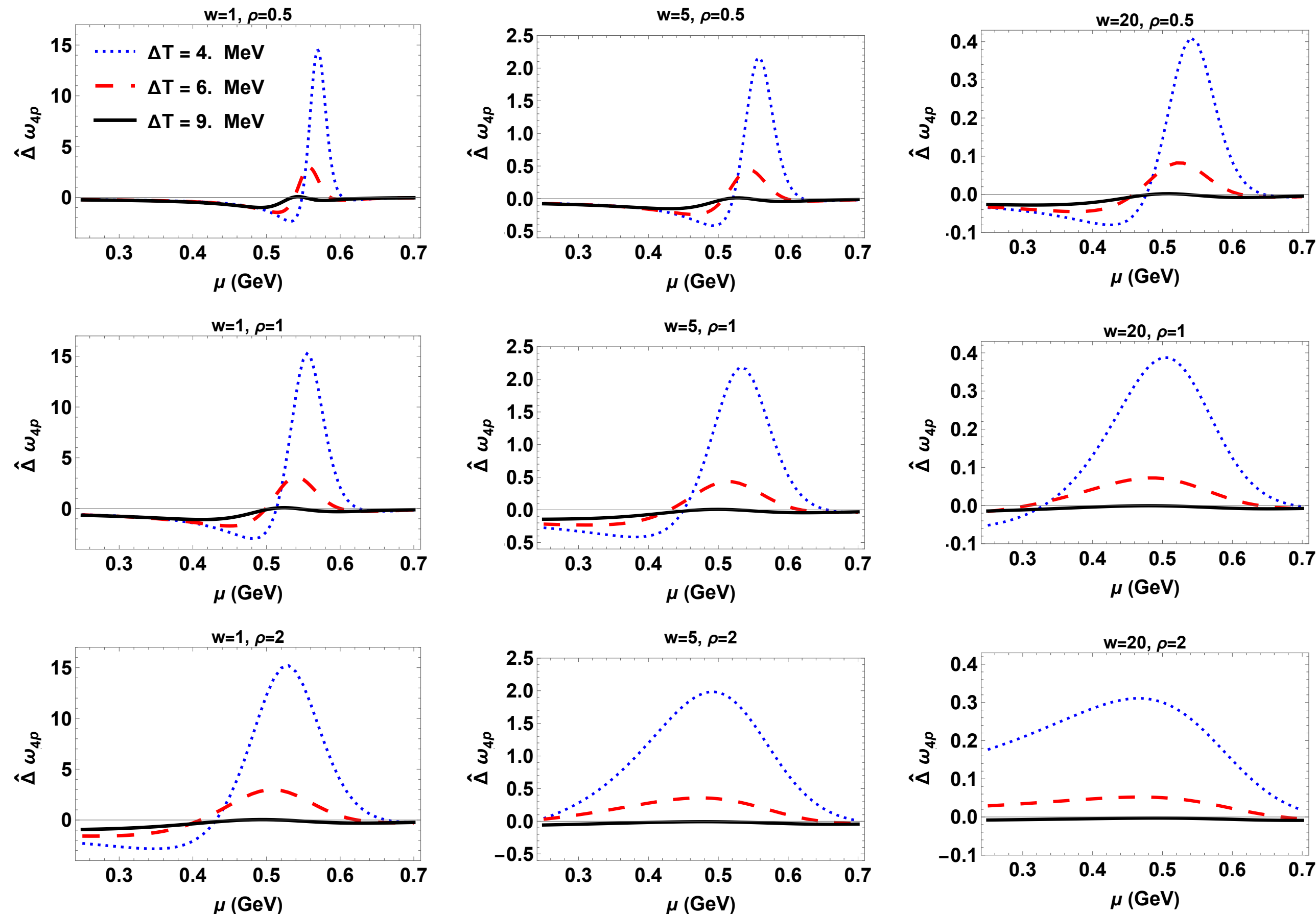
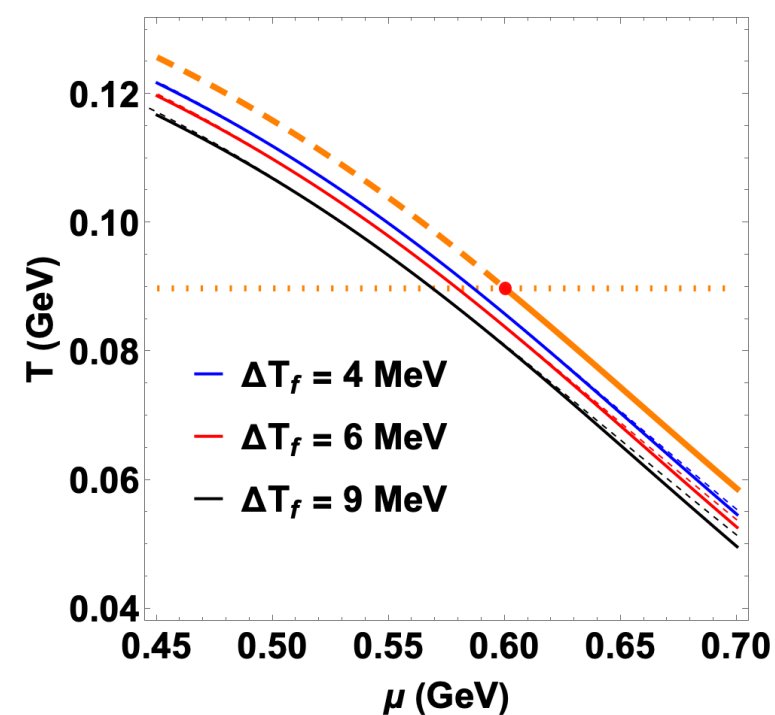
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$$\hat{\Delta}\omega_{kp} = \frac{\hat{\Delta}H_{a_1 \dots a_k} P_p^{a_1} \dots P_p^{a_k}}{\langle N_p \rangle}$$

$$= \frac{\kappa_{4p}}{\kappa_{1p}}$$

Along freeze-out curves:



Increasing  $\rho$   
increases peak width

JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Proton Factorial Cumulants from MaxEnt



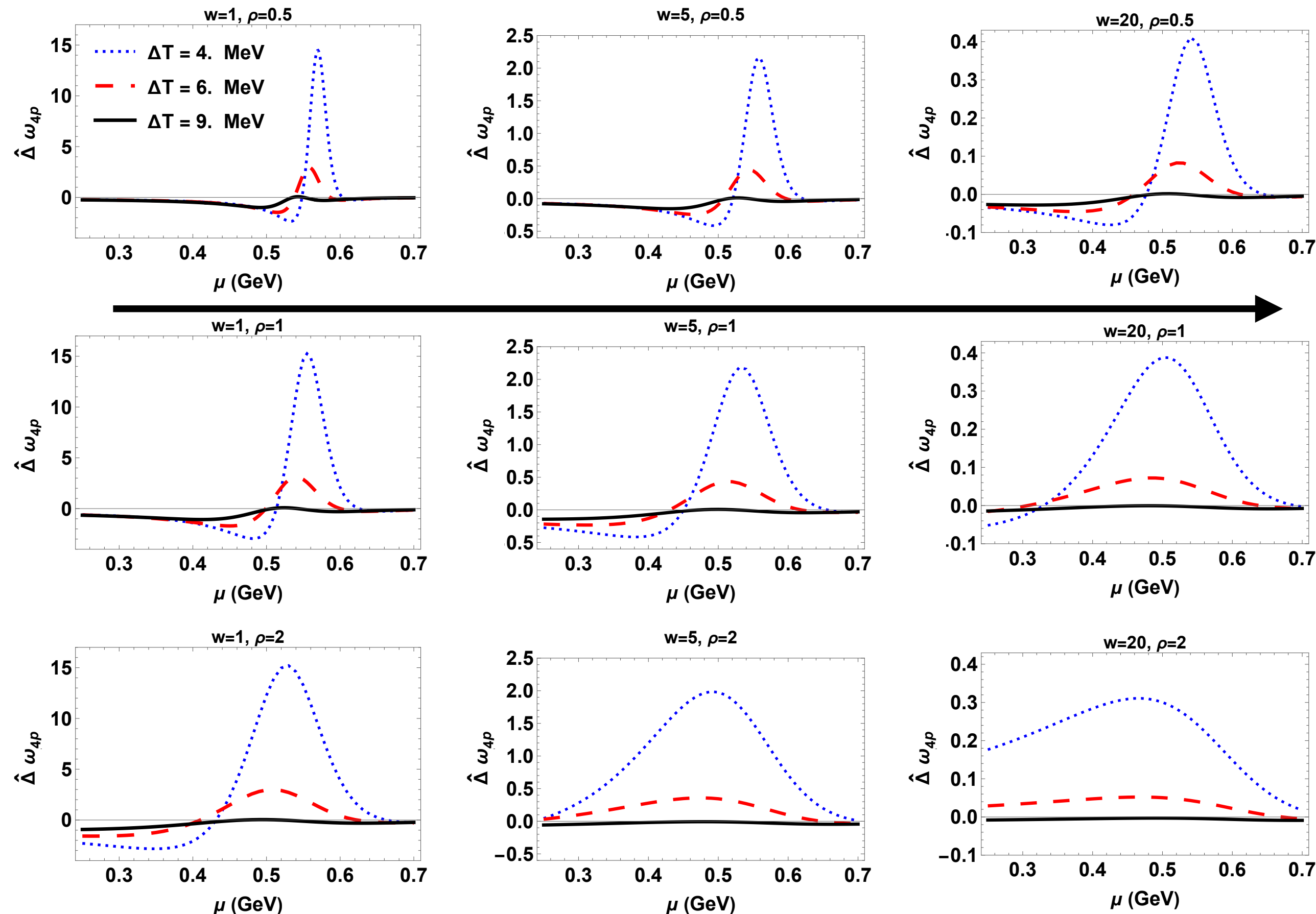
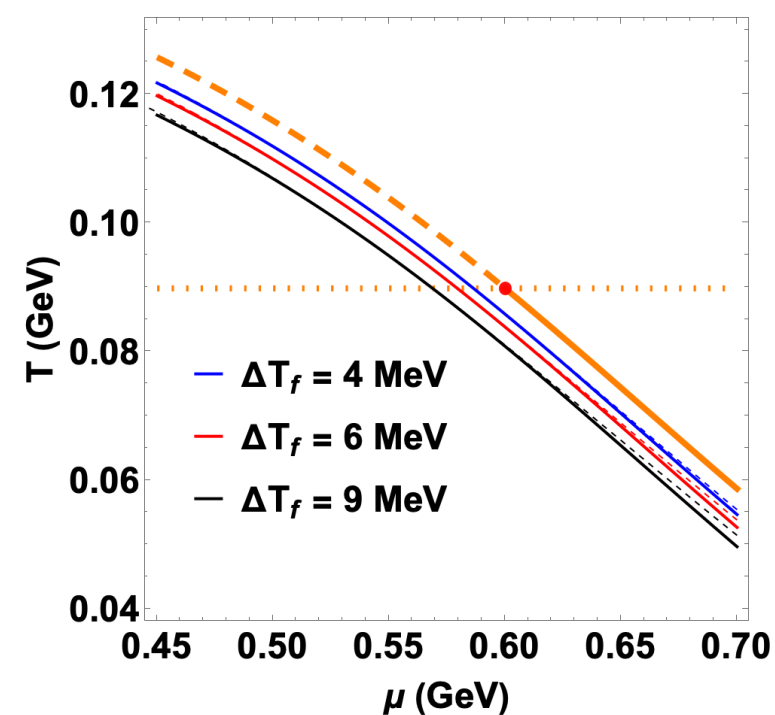
- Critical contribution to fourth order proton factorial cumulants in equilibrium for range of EoS parameters and freeze-out scenarios

Normalized proton factorial cumulants:

$$\hat{\Delta}\omega_{kp} = \frac{\hat{\Delta}H_{a_1 \dots a_k} P_p^{a_1} \dots P_p^{a_k}}{\langle N_p \rangle}$$

$$= \frac{\mathcal{K}_{4p}}{\mathcal{K}_{1p}}$$

Along freeze-out curves:



Increasing  $w$   
reduces peak height

*JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)*

# Proton Factorial Cumulants from MaxEnt

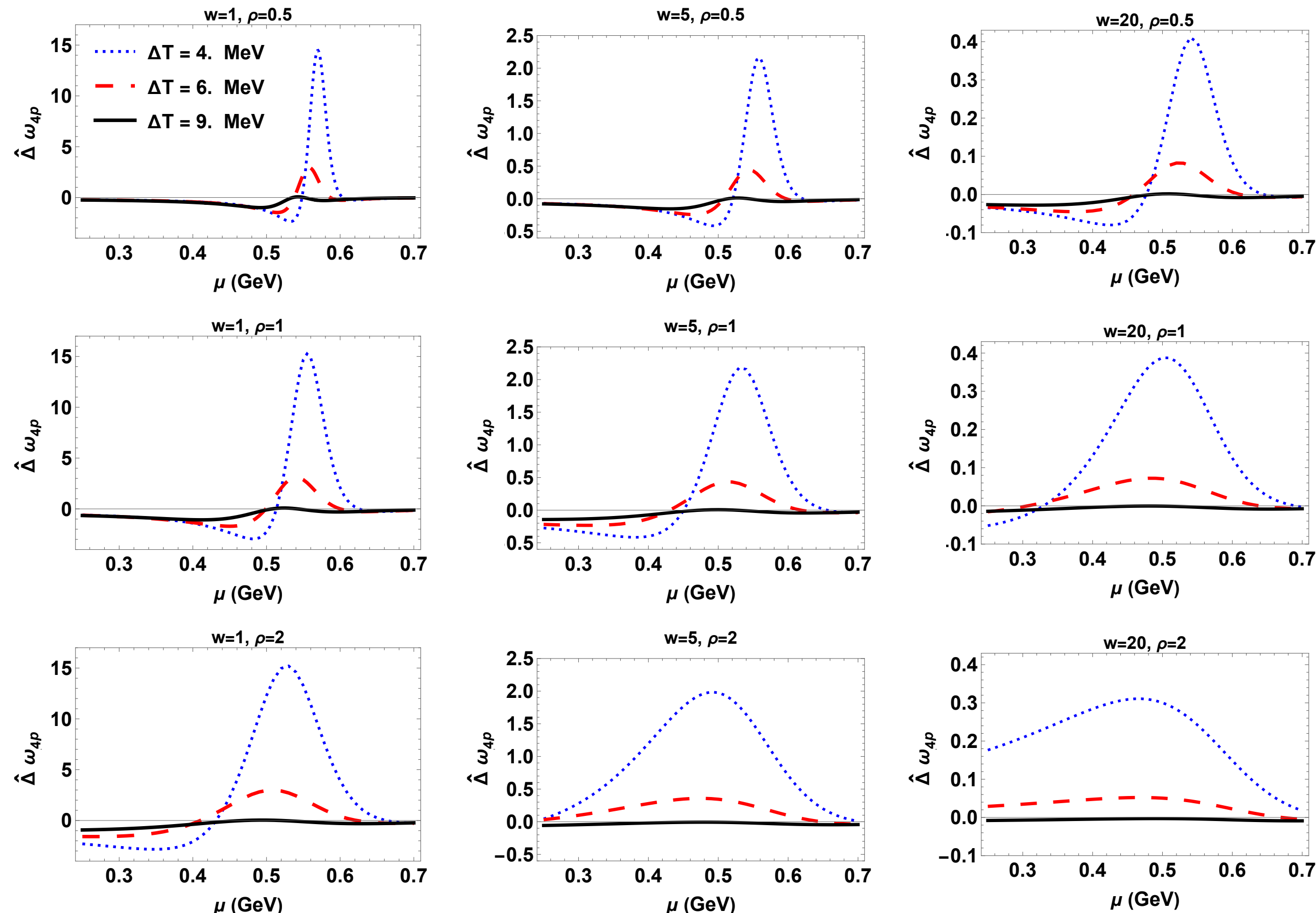
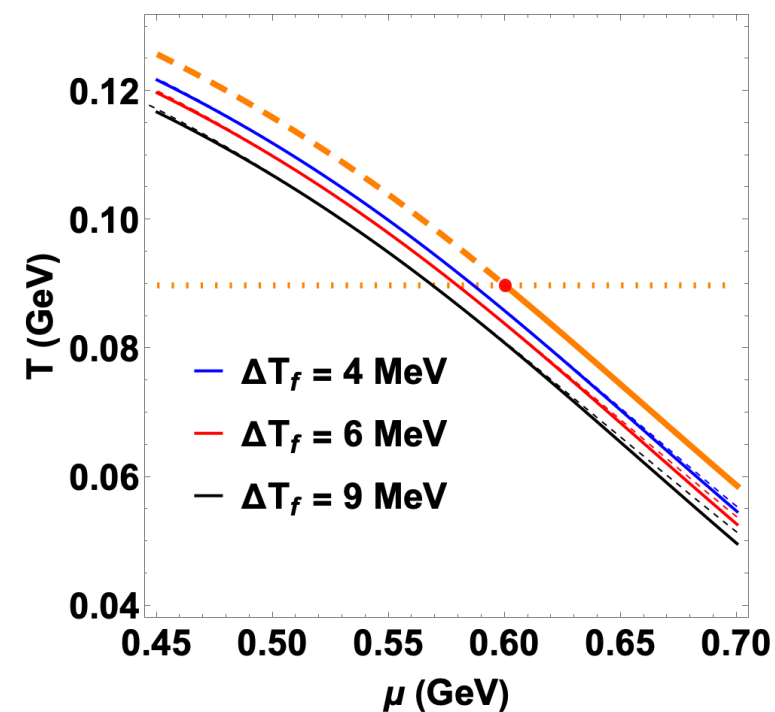


- Critical contribution to fourth order proton factorial cumulants in equilibrium for range of EoS parameters and freeze-out scenarios

Normalized proton factorial cumulants:

$$\hat{\Delta}\omega_{kp} = \frac{\hat{\Delta}H_{a_1 \dots a_k} P_p^{a_1} \dots P_p^{a_k}}{\langle N_p \rangle} = \frac{\kappa_{4p}}{\kappa_{1p}}$$

Along freeze-out curves:



Scaling of magnitude:

$$\hat{\Delta}\omega_p^k \sim w^{-1.2} \Delta T_f^{1.2-k} \sim \xi^{k(5-\eta)/2-3}$$

Uncertainty in  $\Delta T_f$ :  
factorial cumulants can vary by order of magnitude

JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Proton Factorial Cumulants from MaxEnt



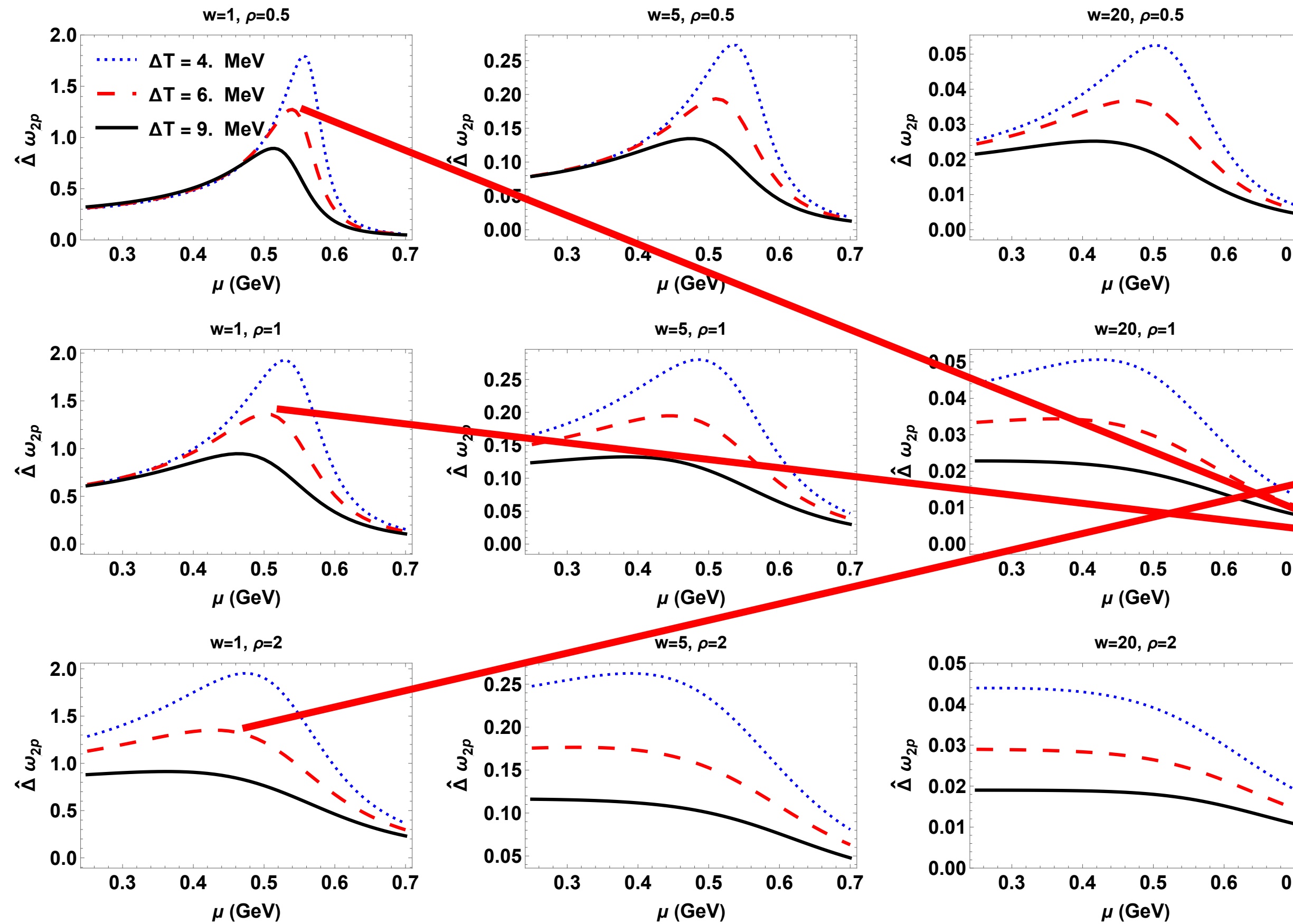
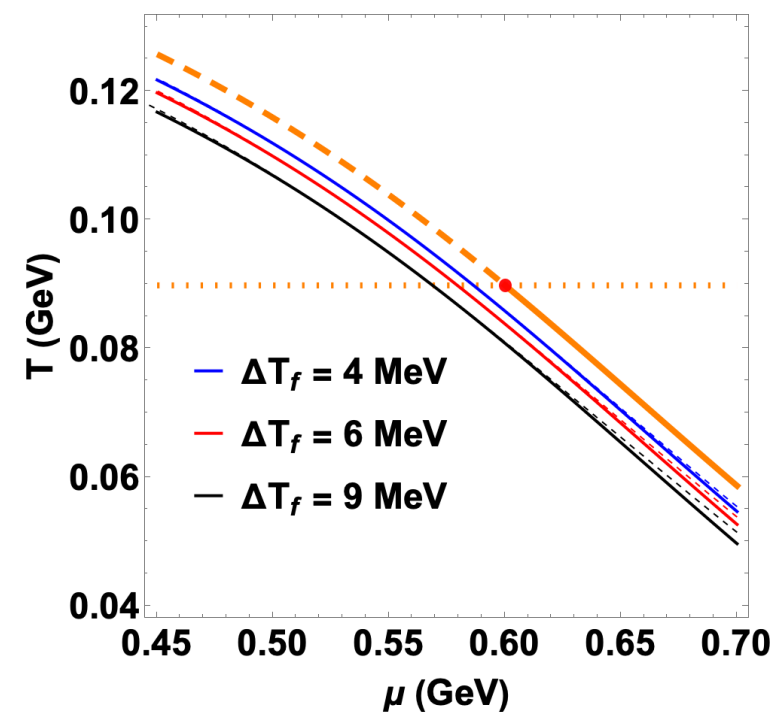
- Critical contribution to second order proton factorial cumulants in equilibrium for range of EoS parameters and freeze-out scenarios

Normalized proton factorial cumulants:

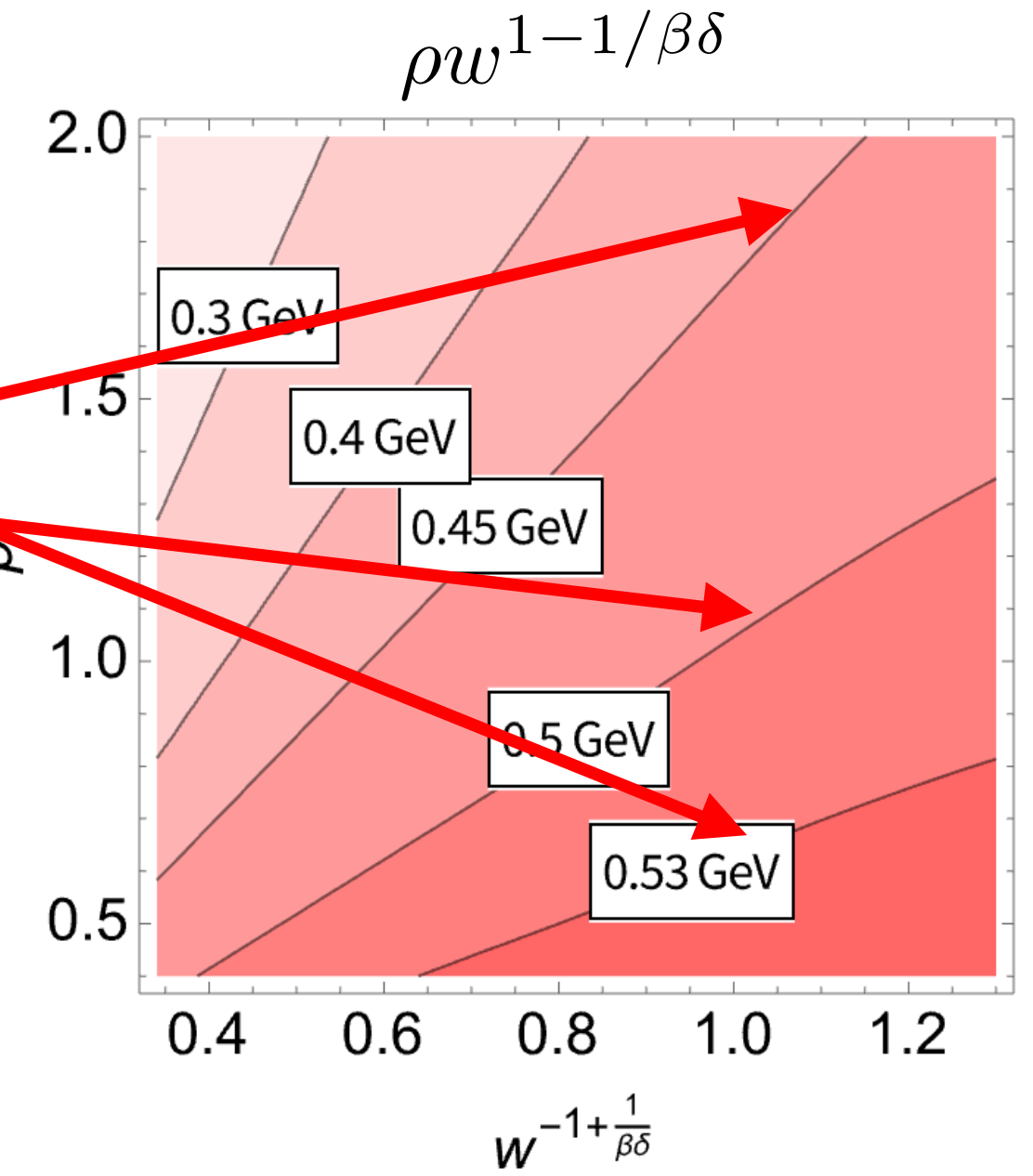
$$\hat{\Delta}\omega_{kp} = \frac{\hat{\Delta}H_{a_1 \dots a_k} P_p^{a_1} \dots P_p^{a_k}}{\langle N_p \rangle}$$

$$= \frac{\kappa_{4p}}{\kappa_{1p}}$$

Along freeze-out curves:



Scaling of location:

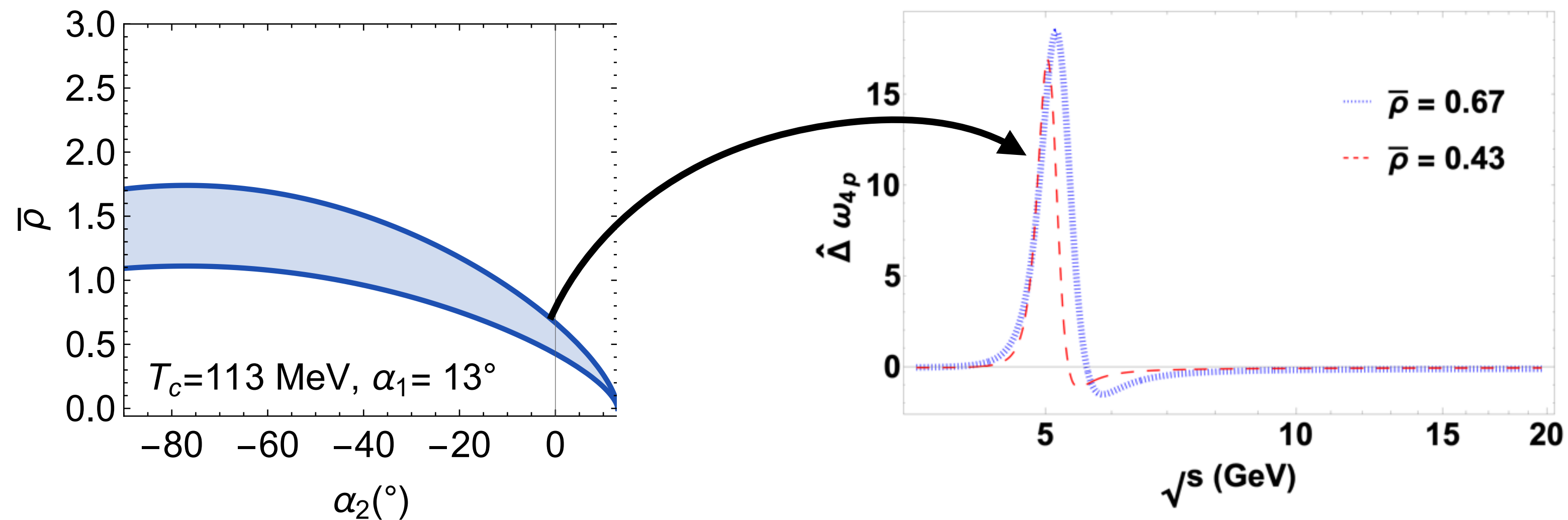


JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)

# Other Constraints on EoS from $\hat{\omega}_{kp}$ with MaxEnt

- Padé estimates for CP can also constrain (via imaginary part,  $c_2$ ) the combination of non-universal parameters  $\rightarrow$  fix  $T_c, \alpha_1, \alpha_2$  to determine  $\bar{\rho} = \rho_w^{1-1/(\beta\delta)}$

$$\bar{\rho} = \left( \frac{1}{c_2 T_c^{\beta\delta-1}} \frac{|\sin \alpha_{12}|}{|\sin \alpha_1|^{\beta\delta+1}} \right)^{1/\beta\delta}$$



Uncertainty on Lattice Padé estimates translates to spread of  $\bar{\rho}$   
with location of critical signal stable

# Future Work: Critical Effects Beyond Equilibrium



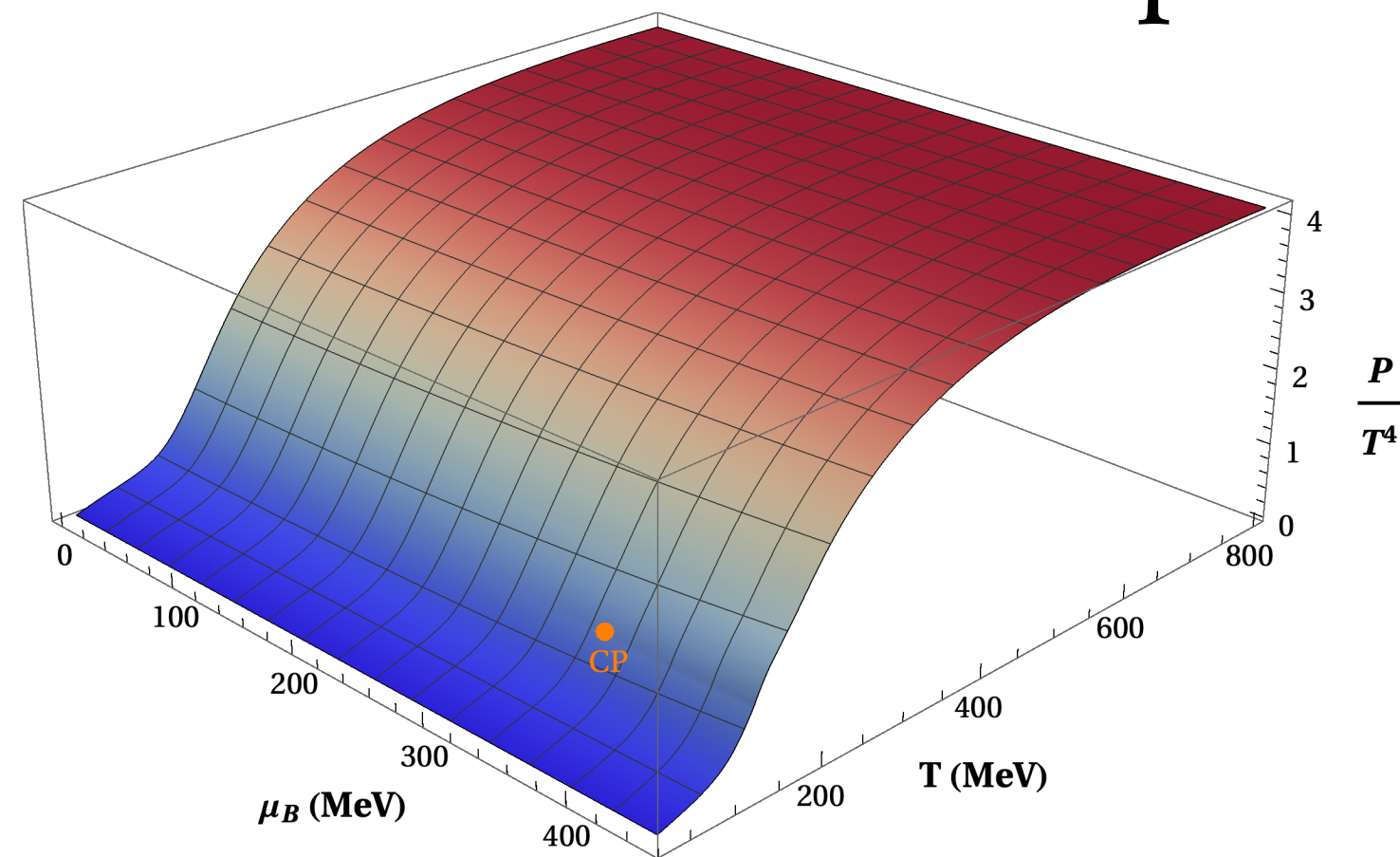
- Direct protons
  - ✓ Qualitative results persist with resonance feed-down (Appendix B)
- Current background EoS is HRG
  - ➔ Extend to more realistic EoS constrained by lattice
  - ➔ Compare to non-critical baseline *P. Braun-Munzinger, B. Friman et al, NPA (2021)*  
*V. Vovchenko, V. Koch & C. Shen, PRC (2022)*
- Assumption of thermal equilibrium
  - ➔ Full Hydro+ simulations *M. Stephanov and Y. Yin, PRD (2018)*  
*K. Rajagopal, G. Ridgway et al, PRD (2020)*  
*M. Pradeep, K. Rajagopal et al, PRD (2022)*
- Simultaneous particlization and chemical freeze-out
  - ➔ Couple to hadronic afterburner *Critical effects seen to survive in SMASH:*  
*J. Hammelman, M. Bluhm et al, PRC (2024)*

*JMK, K. Rajagopal, M. Pradeep, M. Stephanov, Y. Yin, PRD (2026, Editors' Suggestion)*

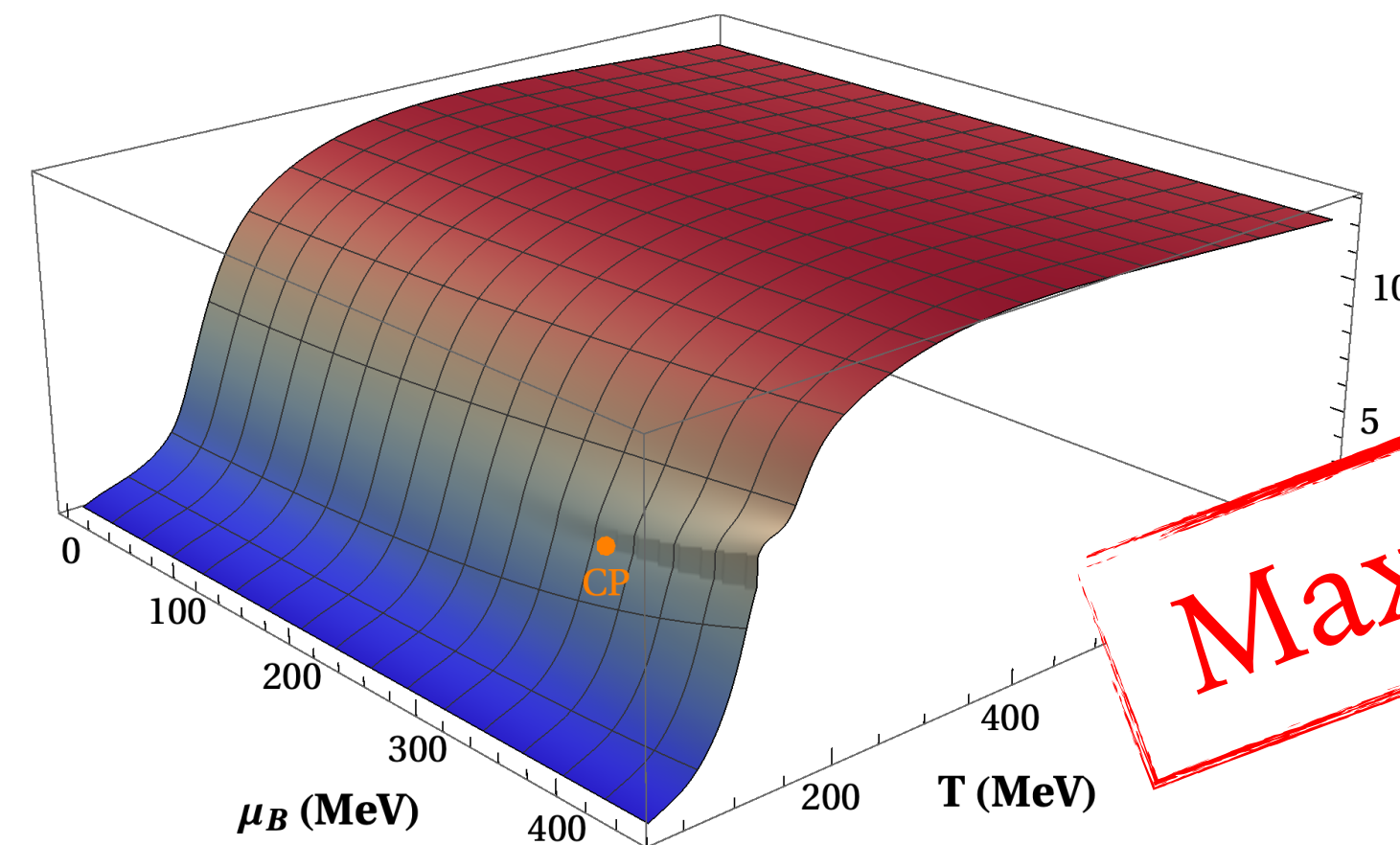
# How to use EoS to connect to experiment?



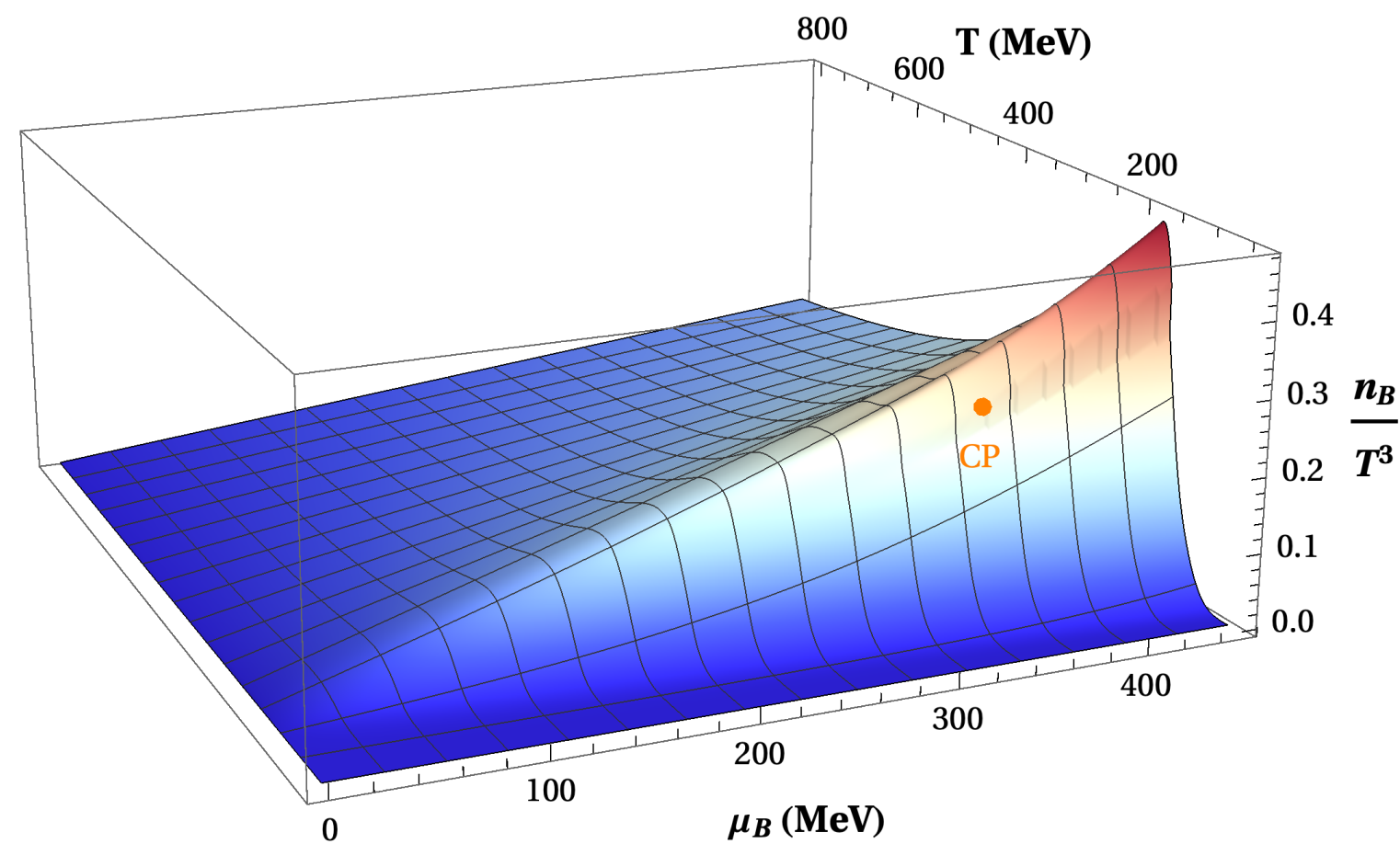
... in a model independent way



Pressure

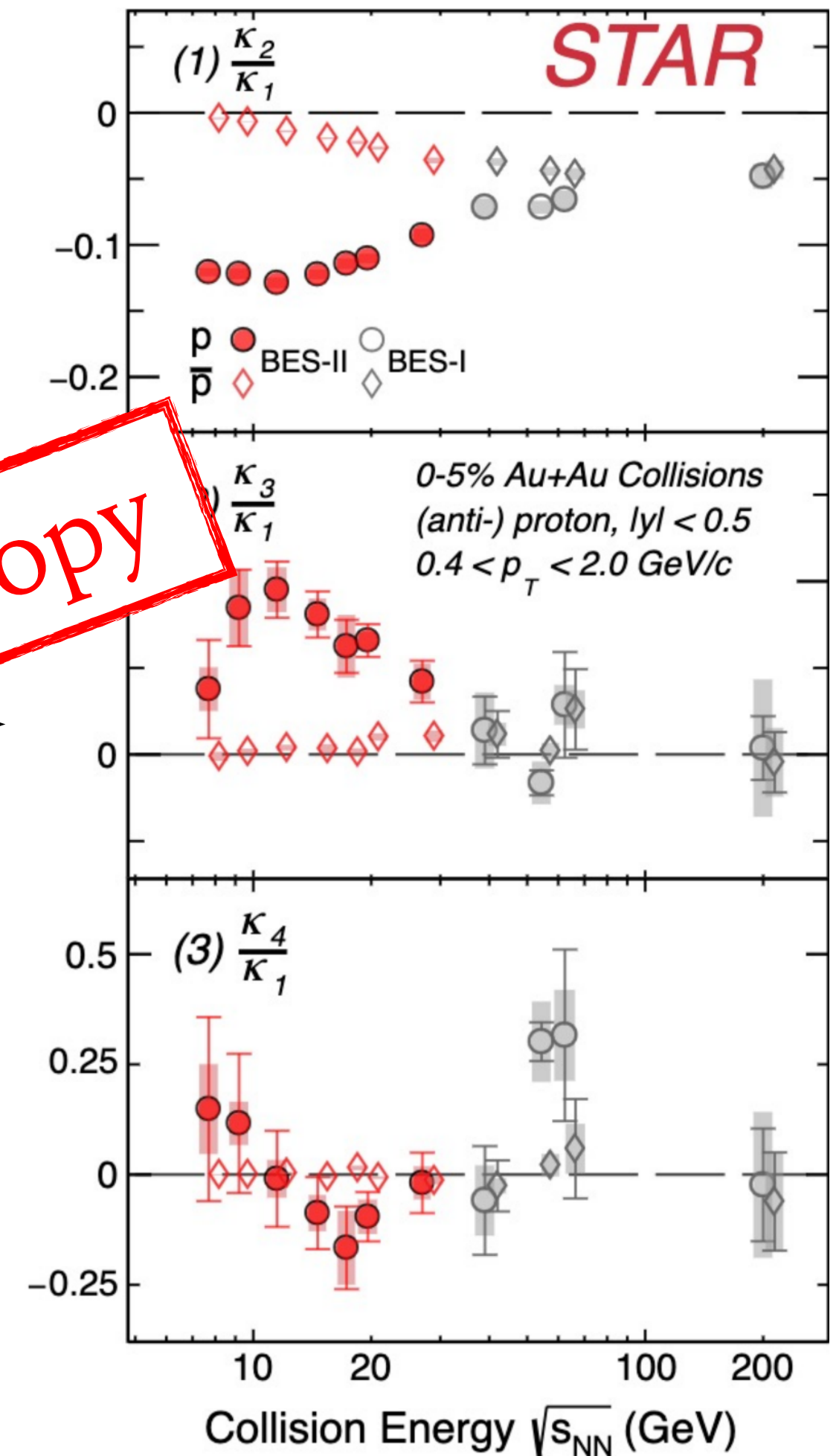


Energy density



Baryon density

**Maximum Entropy**

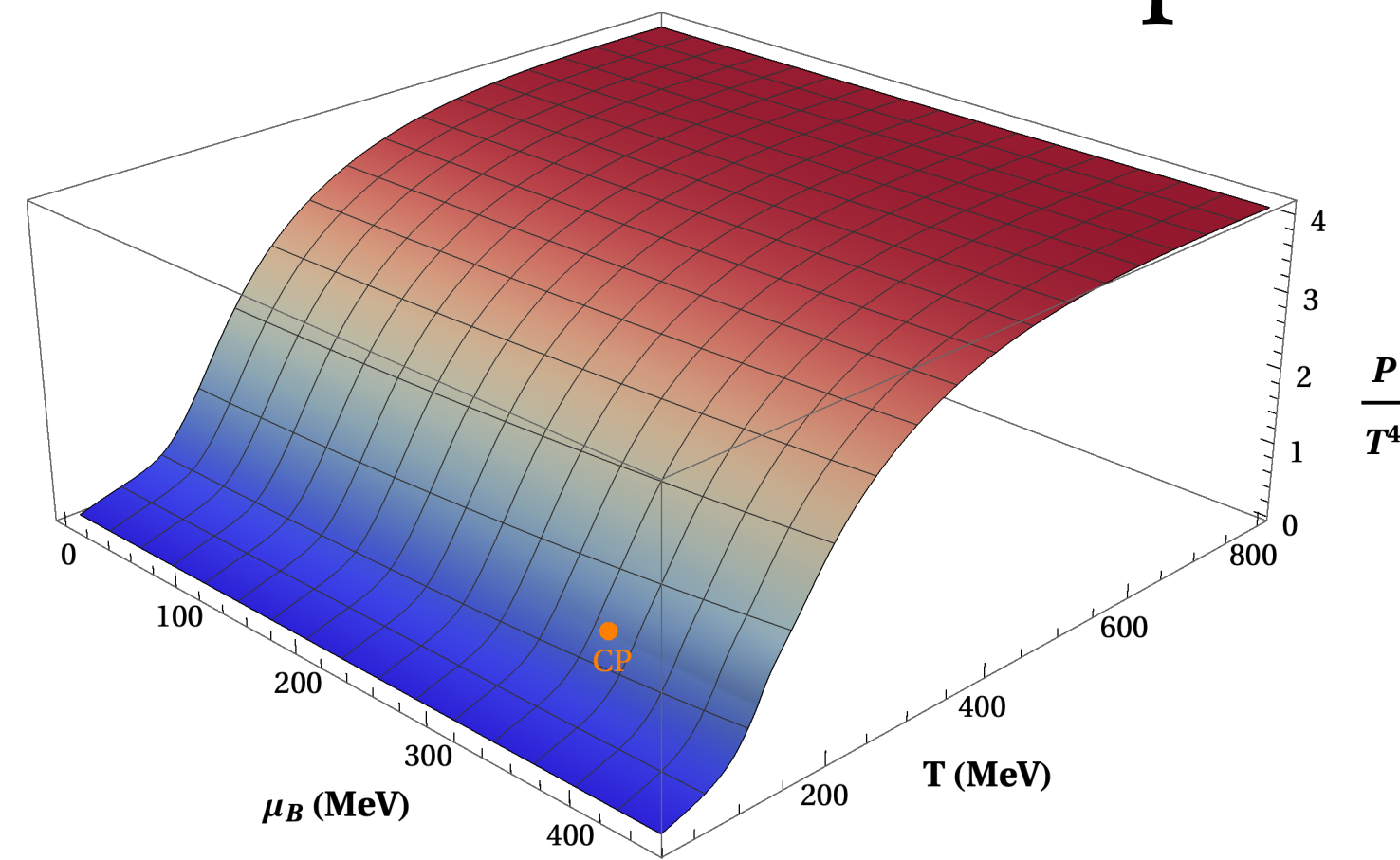


JMK et al, EPJ+ (2021)  
A. Pandav (STAR collaboration), CPOD 2024

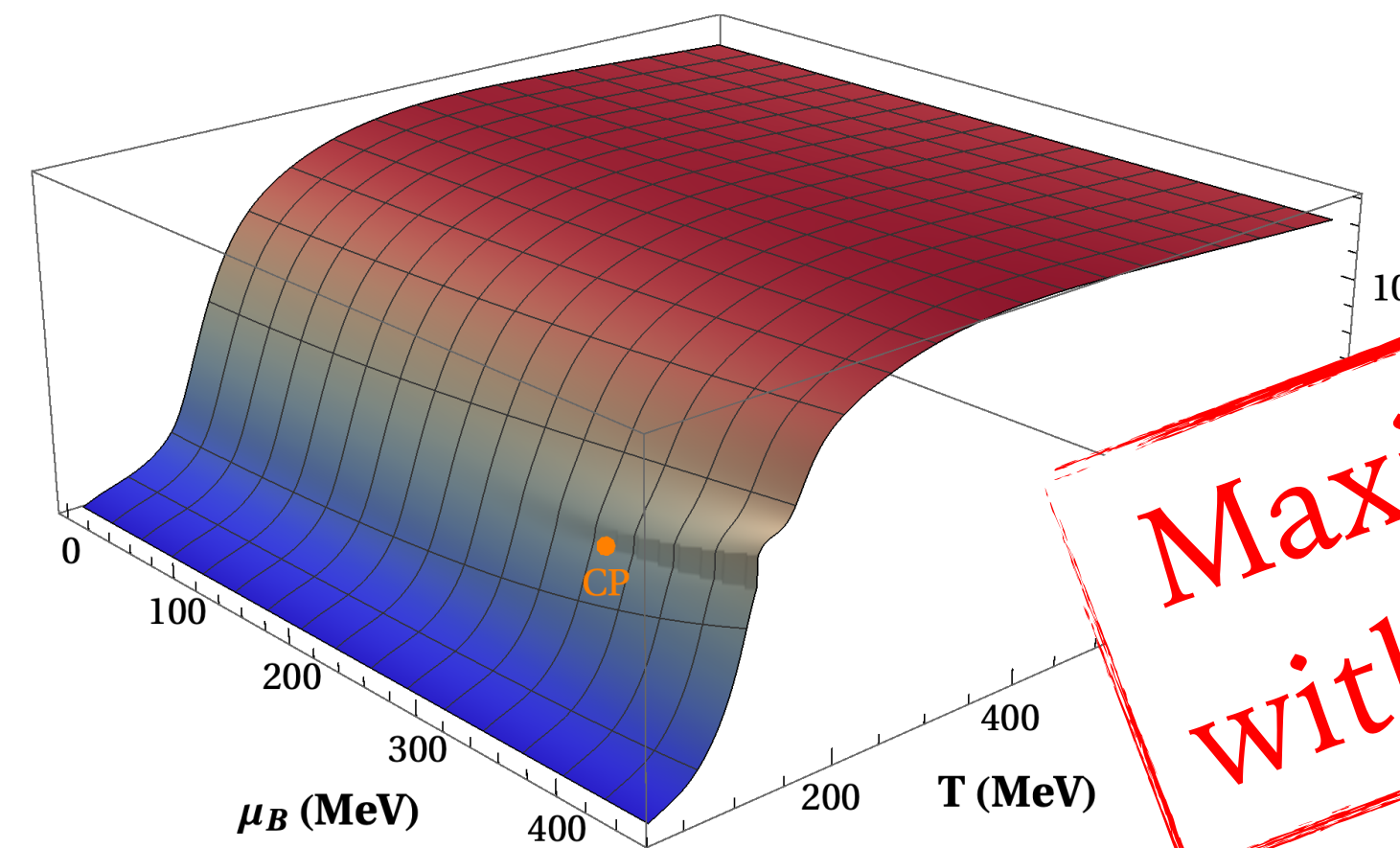
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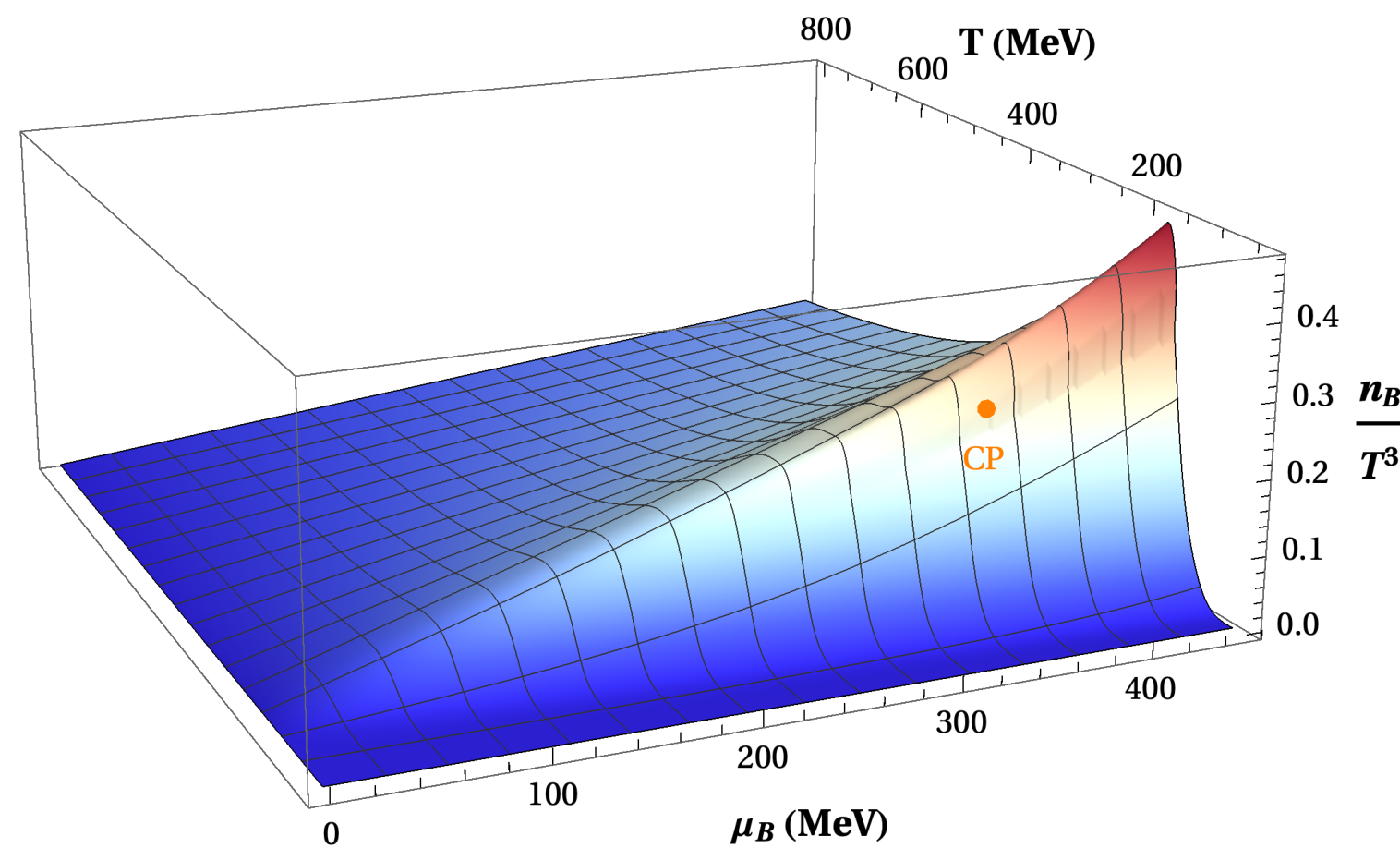
... in a model independent way



Pressure

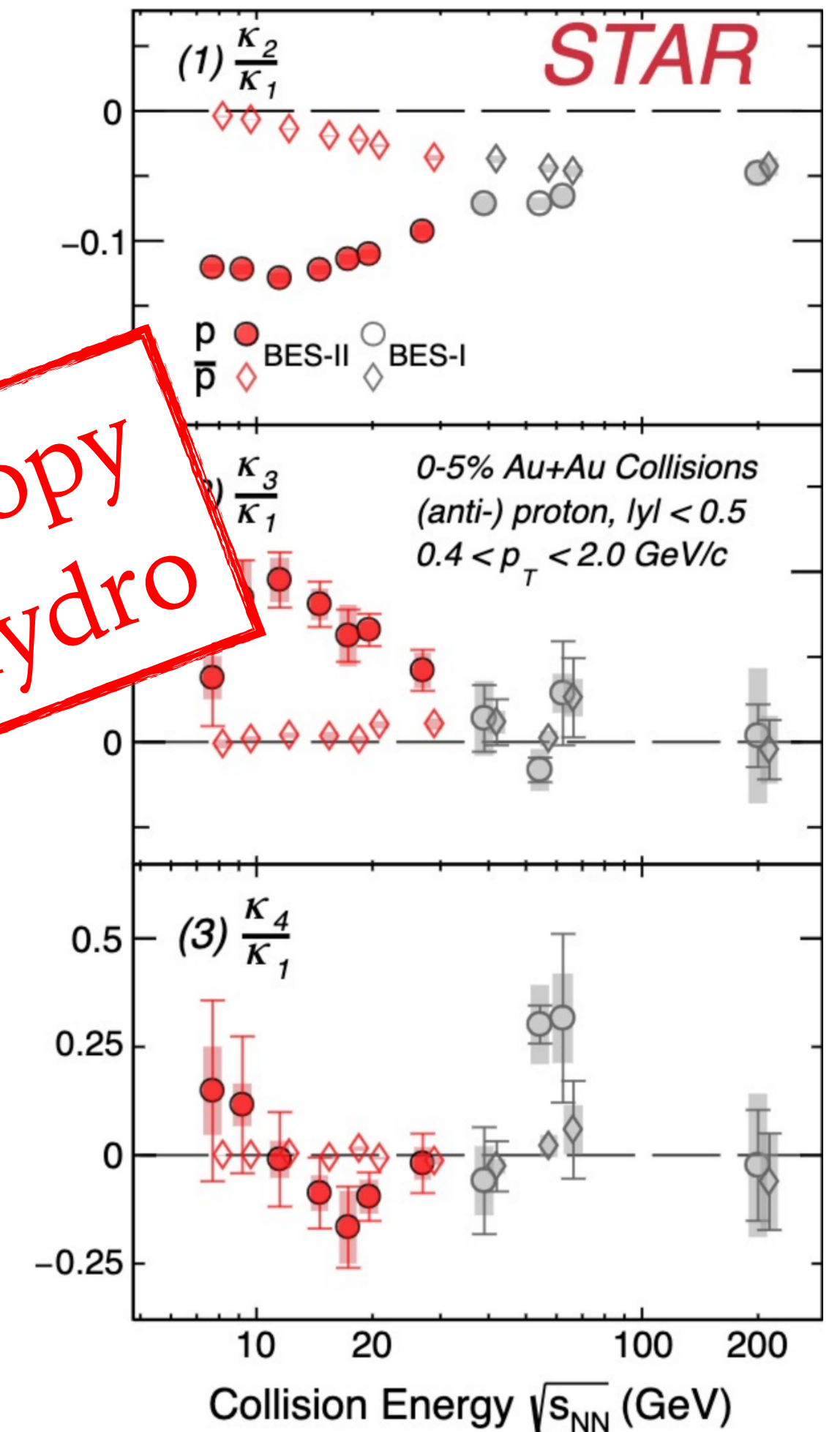


Energy density



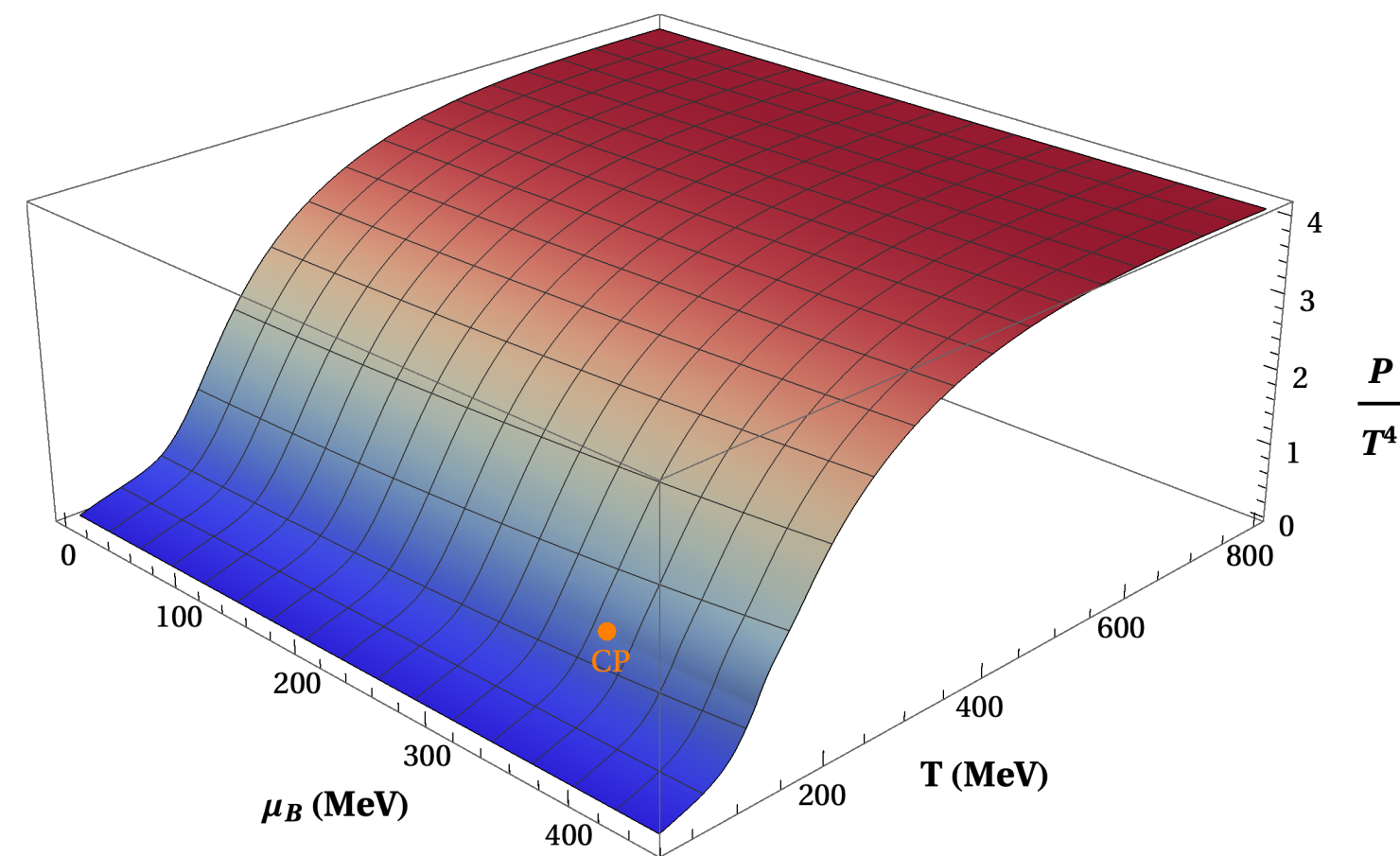
Baryon density

Maximum Entropy  
with realistic hydro

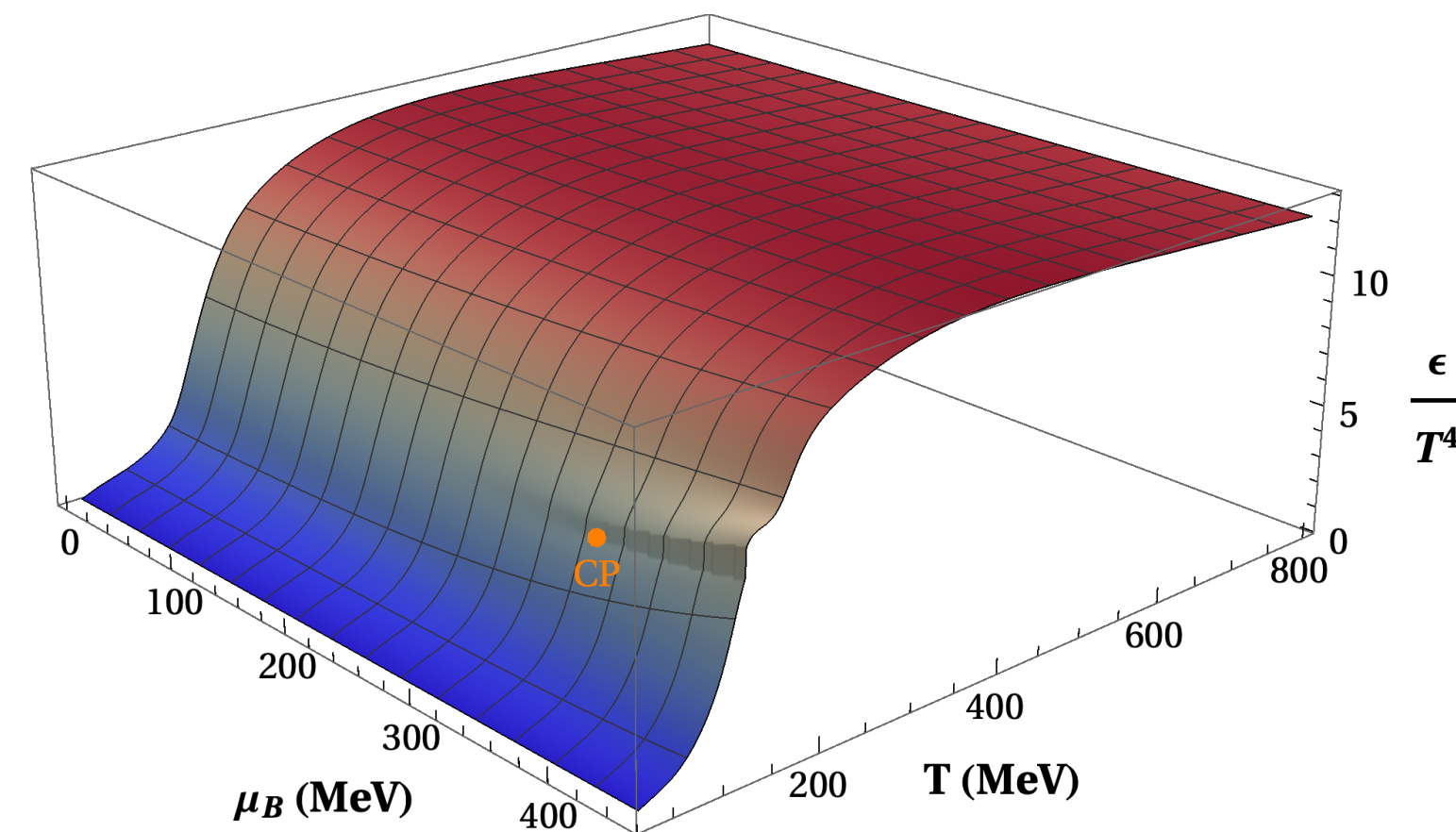


JMK et al, EPJ+ (2021)  
A. Pandav (STAR collaboration), CPOD 2024

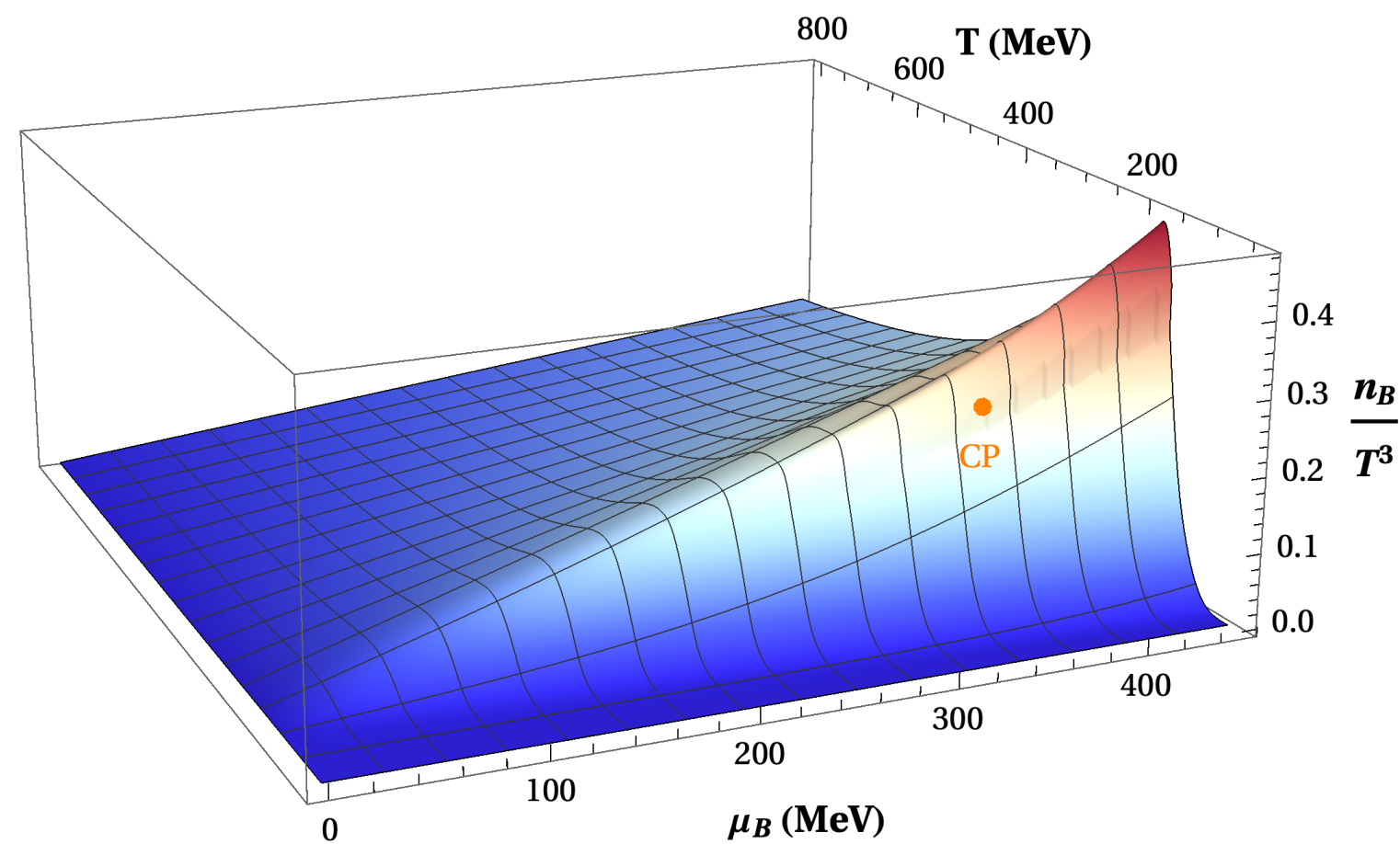
# How to extract EoS from experiment?



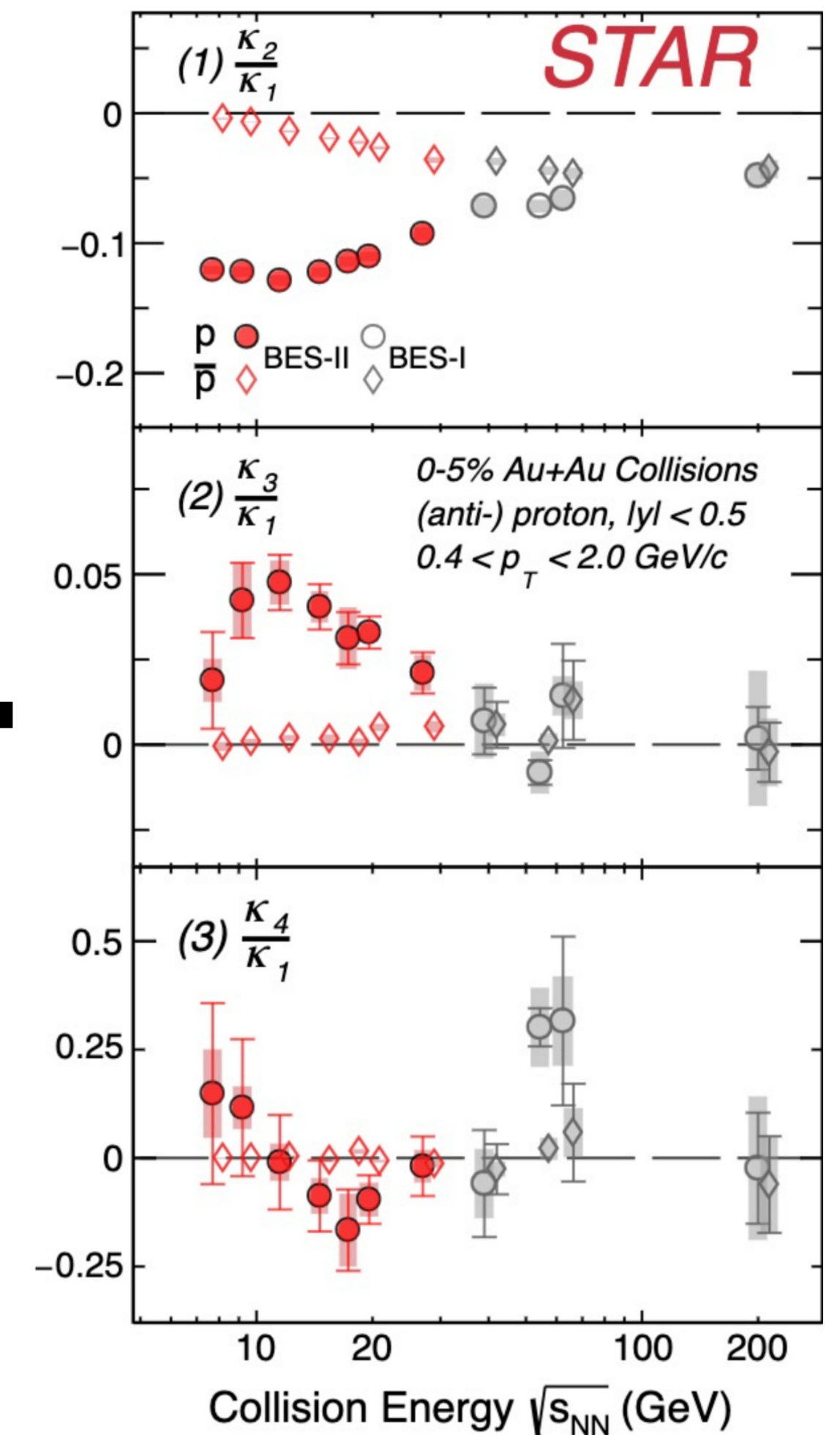
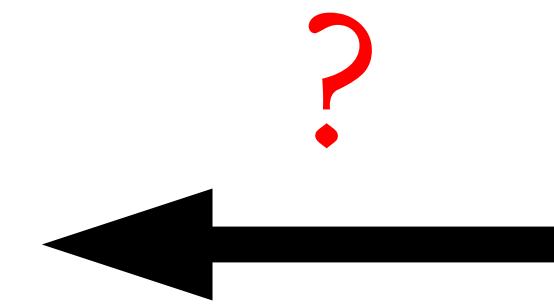
Pressure



Energy density

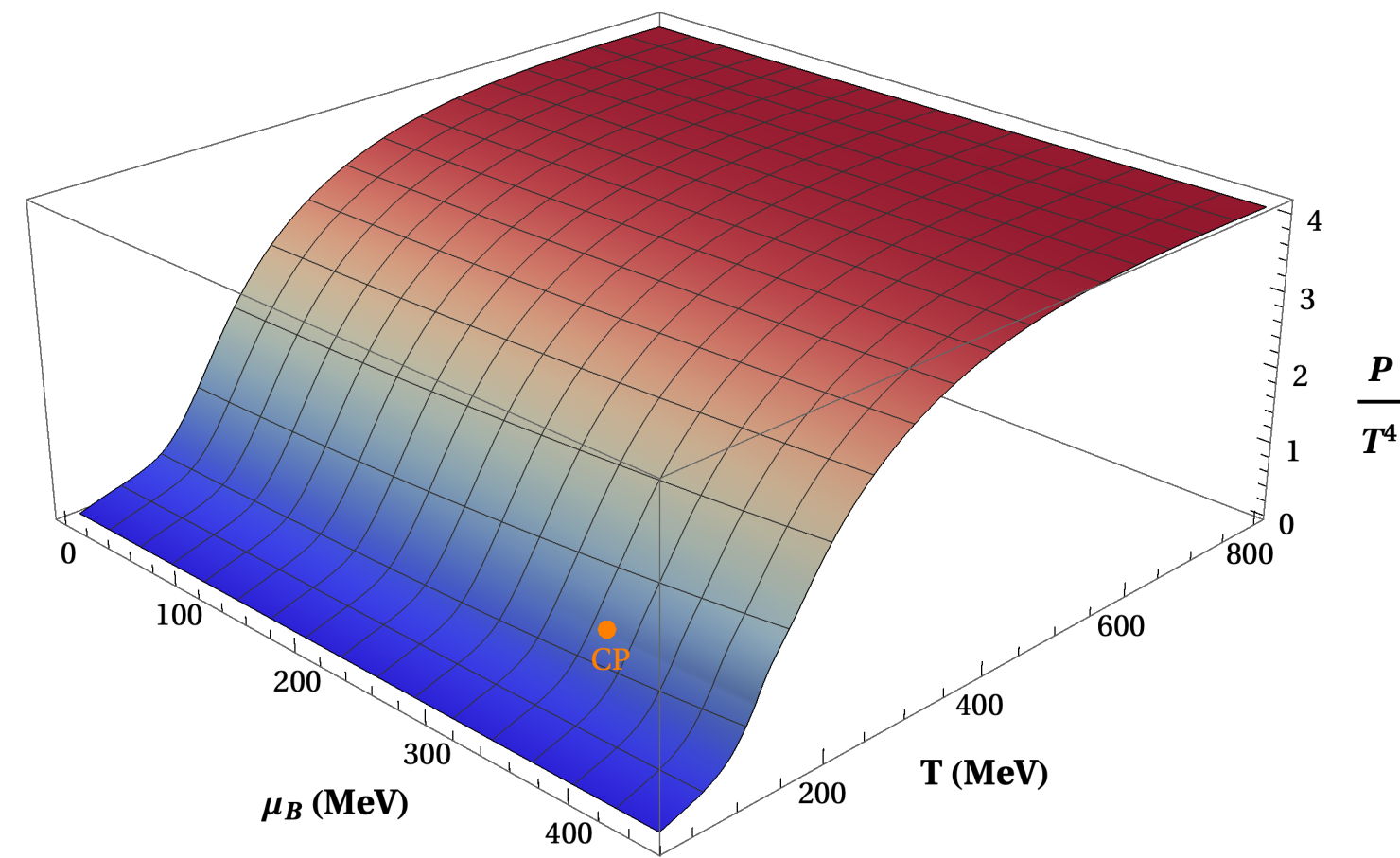


Baryon density

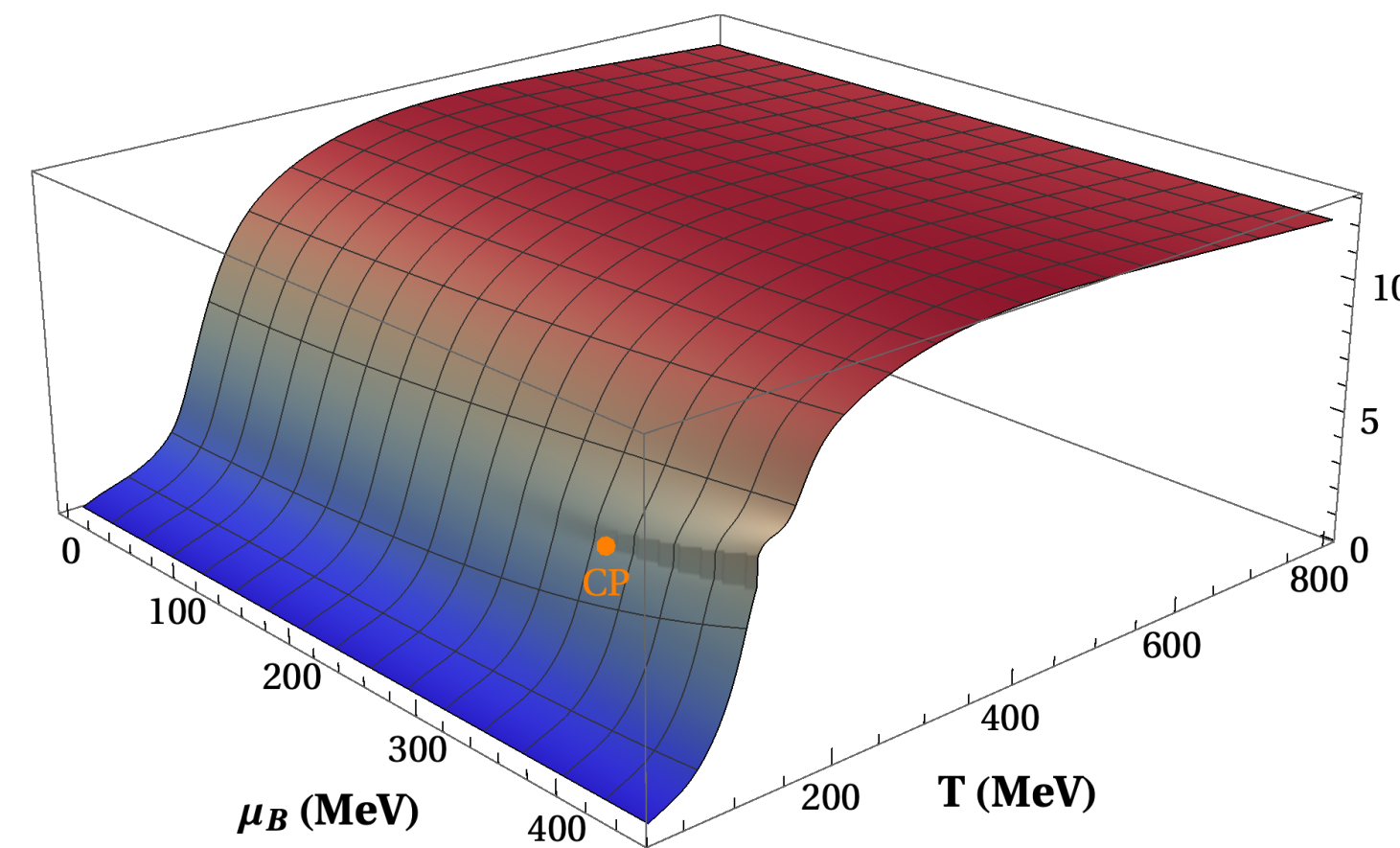


JMK et al, EPJ+ (2021)  
A. Pandav (STAR collaboration), CPOD 2024

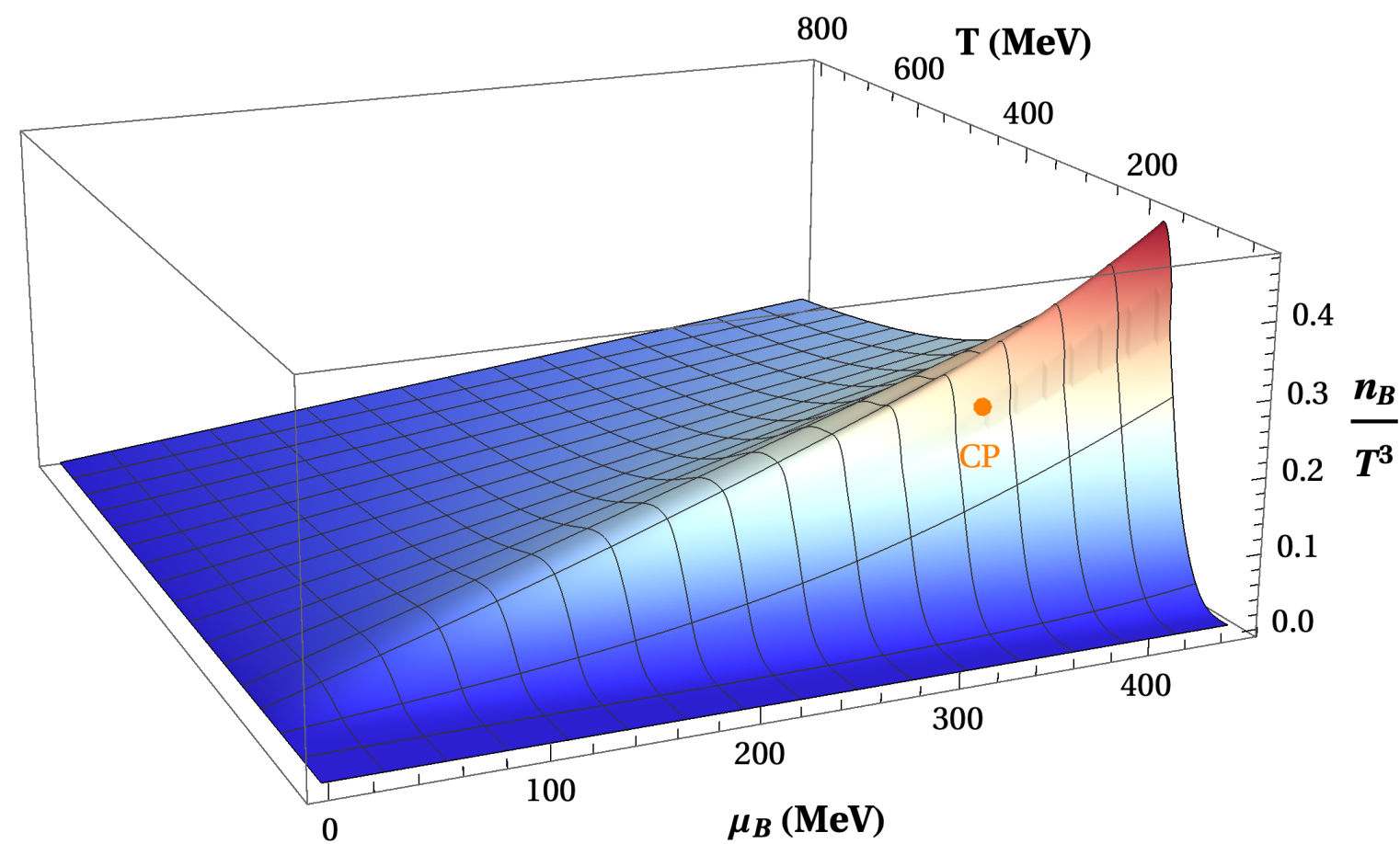
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Pressure

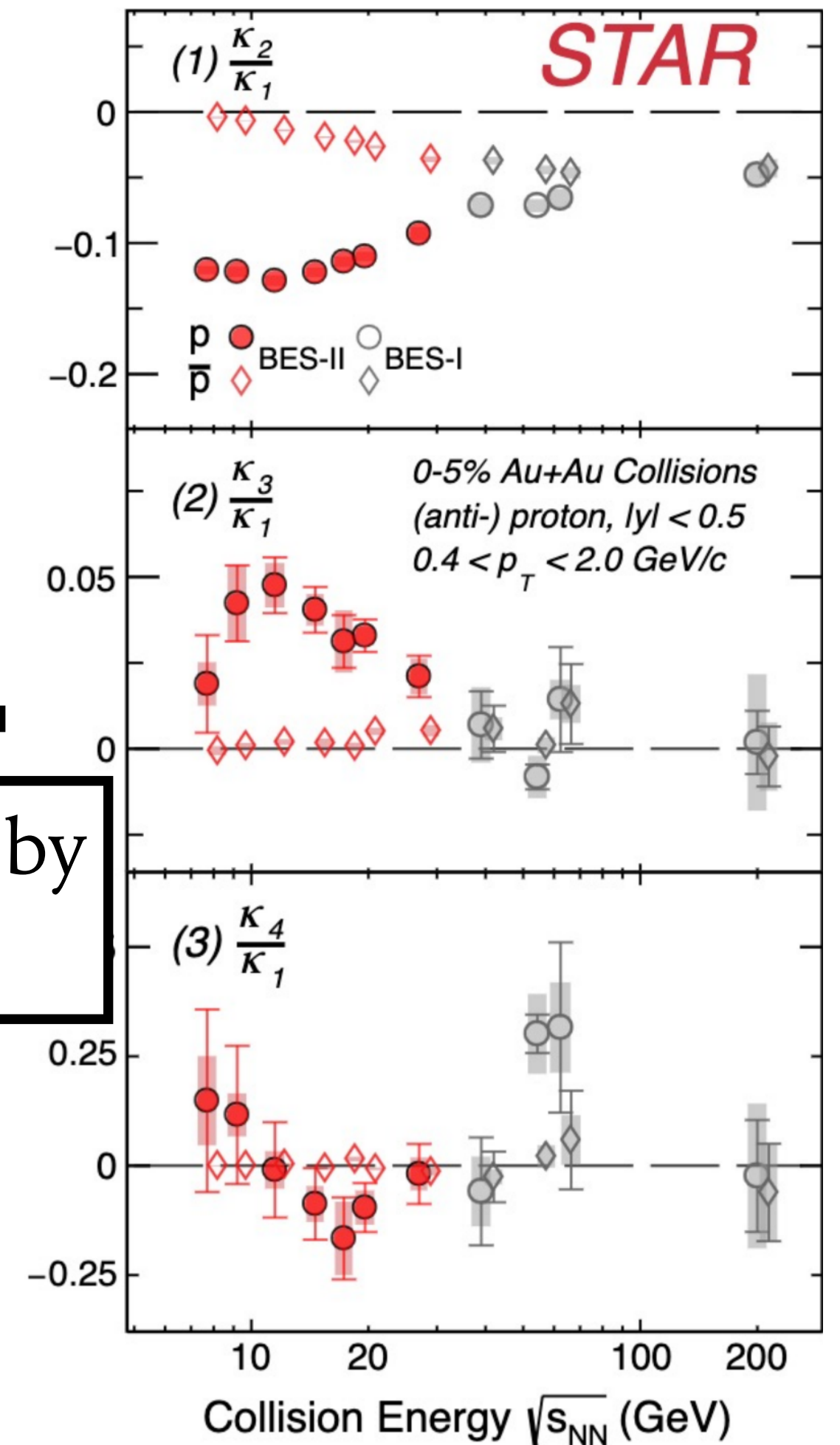


Energy density



Baryon density

← ?  
See next talk by Grégoire!



JMK et al, EPJ+ (2021)  
A. Pandav (STAR collaboration), CPOD 2024

- The search for the critical point/first order transition continues...
- Criticality incorporated within the original BEST EoS
  - Updated to strangeness neutrality, higher coverage of phase diagram, first order features + ongoing work: 4D EoS
- Max Ent provides a framework that takes the EoS as input and calculates the particle multiplicity fluctuations in the least biased way
  - Dependence of proton factorial cumulants on EoS parameters derived
  - Future work: implement these new updates into Hydro+/freeze-out
- Next talk: inverse problem - use experimental data to obtain EoS inputs