

International Workshop Led by Young Researchers "Quantum thermalization, Hydrodynamics and Gravity"

Active Matter, KPZ, and (possibly) Gravity

Kazuaki Takasan

Department of Applied Physics, The University of Tokyo



Chat with Masataka (One of the Organizers)

02/35



I don't really work on thermalization, hydrodynamics, or gravity.
What should I talk about ?



Our aim is to learn from different research fields.
Please share recent developments in nonequilibrium condensed matter physics.
How about active matter and KPZ ?



OK

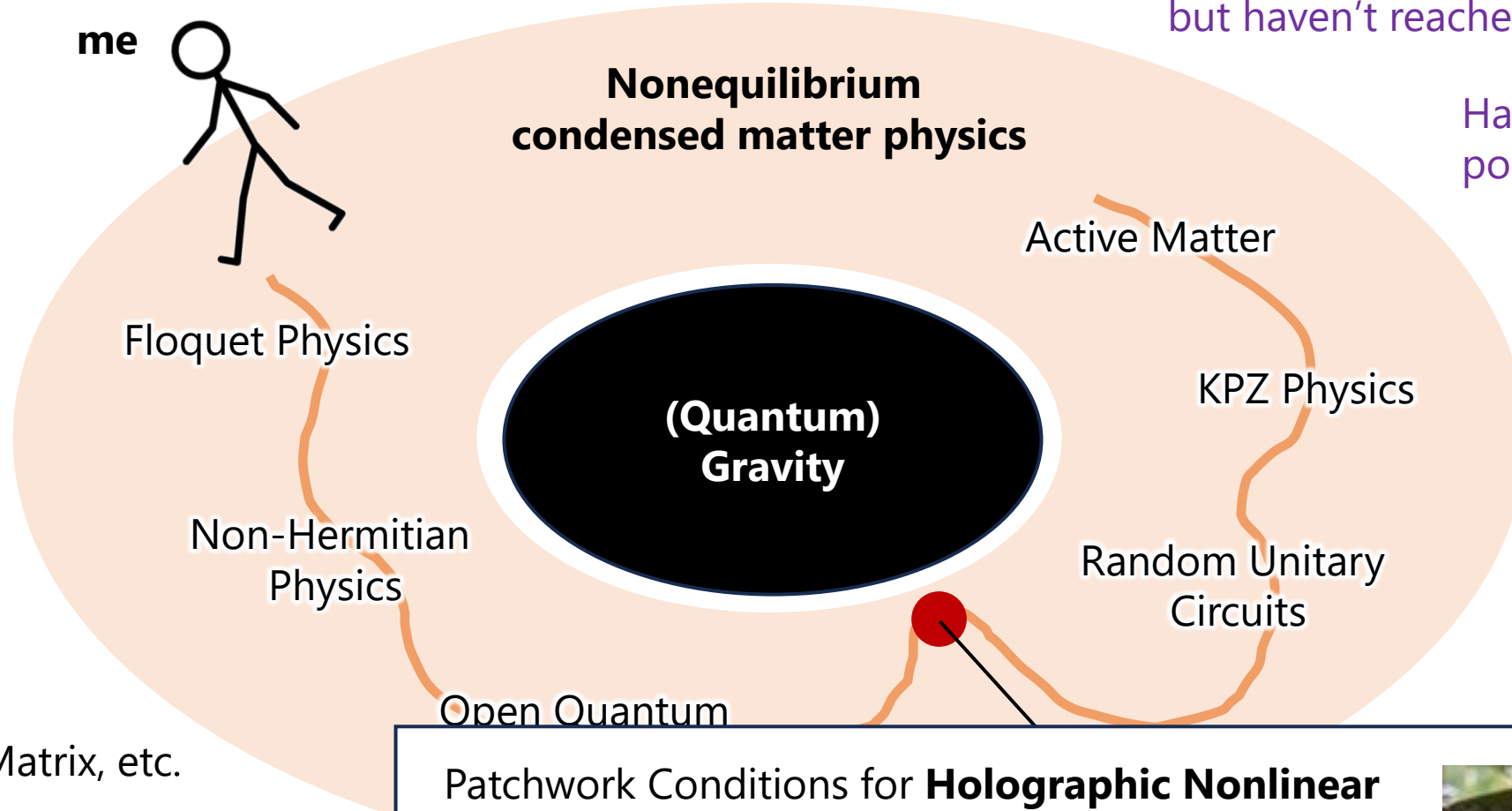
Why I put "gravity" in the title of this talk

03/35

Because I'm interested in.

I feel attracted to it,
but haven't reached the center yet

Happy to discuss
possible connections



Holography,
Entanglement,
SYK, Random Matrix, etc.

Holography can be a tool for
nonequilibrium systems
(applied holography)

Patchwork Conditions for **Holographic Nonlinear Responses**: A Computational Method for Electric Conductivity and Friction Coefficient

[S. Ishigaki](#), [S. Nakamura](#), [KT](#), [arXiv: 2303.02633](#)



Research Interests

Nonequilibrium Phases of Matter

Topic 1

Quantum Active Matter

Topic 2

KPZ in the Heisenberg Spin Chain

Collaborators



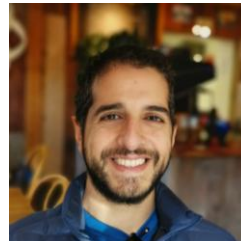
K. Adachi
RIKEN



K. Kawaguchi
UTokyo



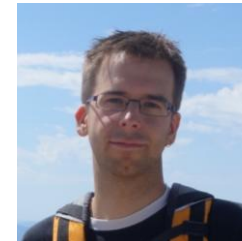
K. A. Takeuchi
UTokyo



O. Busani
Univ. Edinburgh



P. L. Ferrari
Bonn U

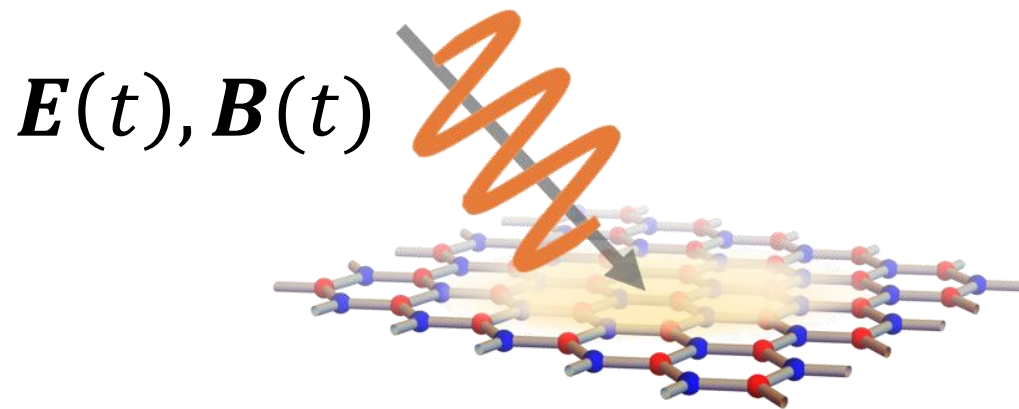


R. Vasseur
Univ. Geneva

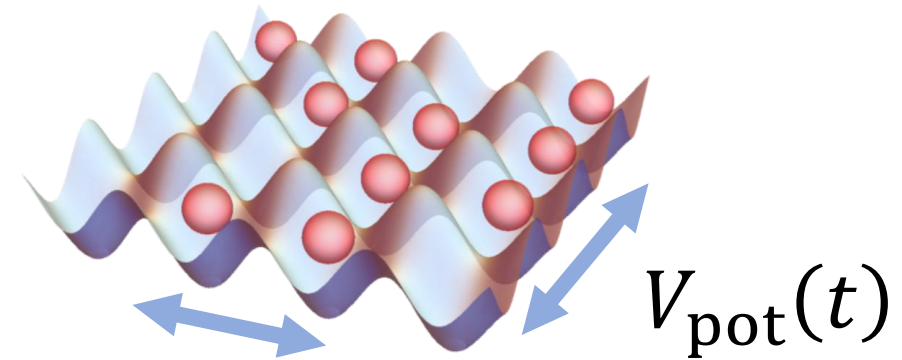


J. De Nardis
Cergy Paris Univ.

Stable (or metastable) quantum many-body states of matter realized far from equilibrium



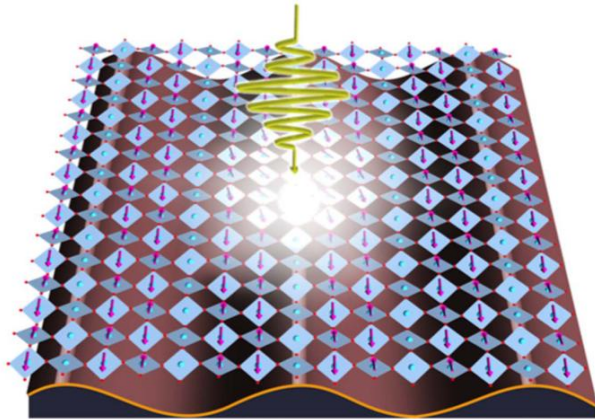
Solid-state systems



Ultracold atoms

Recent experimental developments allow us to investigate them

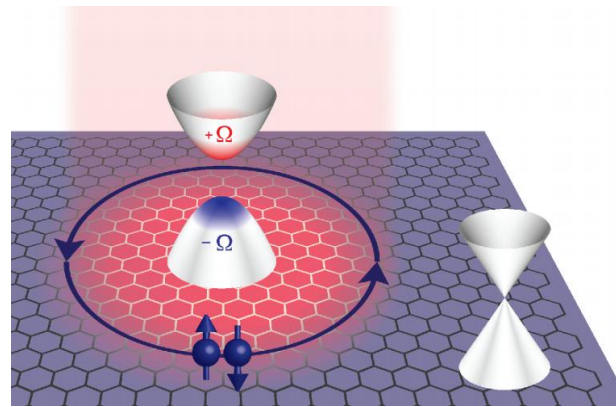
Light-induced superconductivity



Fausti *et al.* Science (2011), Kaiser *et al.* PRB (2014), Mitrano *et al.* Nature (2016), ...

SC-like response even in room temperature

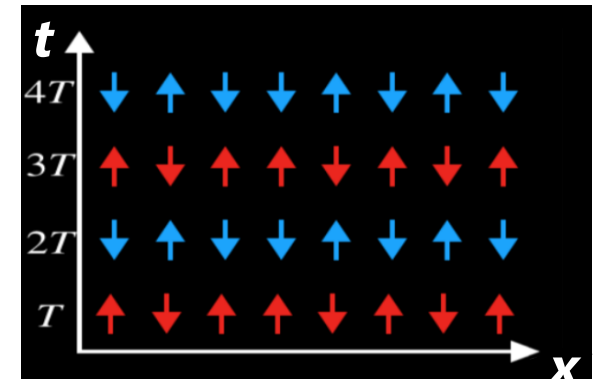
Floquet topological phases



Oka and Aoki PRB (2009), Jotzu *et al.* Nature (2014), McIver *et al.* Nat. Phys (2019)

Ultrafast (fs~ps) control of topological nature

Discrete time crystal



Else *et al.* PRL (2016), Khemani *et al.* PRL (2016), Zhang *et al.* Nature (2017), ...

SSB in time direction, prohibited in equilibrium

Watanabe and Oshikawa PRL (2015)

Nonequilibrium Phases of Matter

Not limited by the rule in equilibrium states

e.g., minimize free energy



Possibility to overcome "limitation" in current condensed matter physics

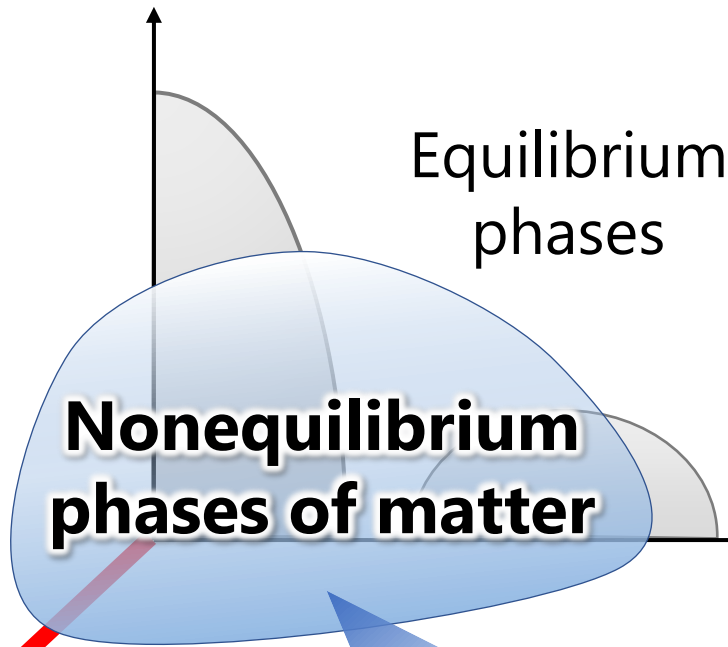
e.g., room-temp. SC



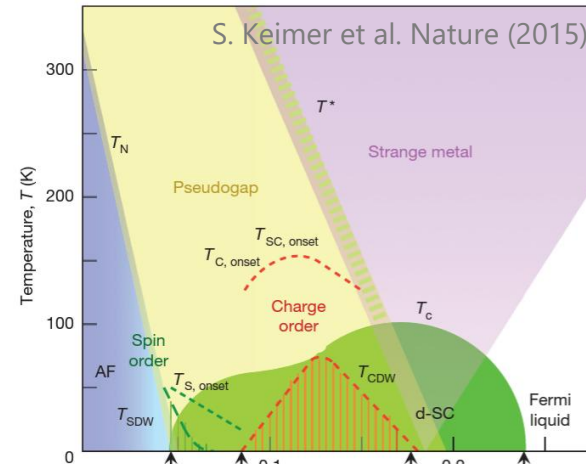
Deepen our understanding of phases of matter

Future novel technology

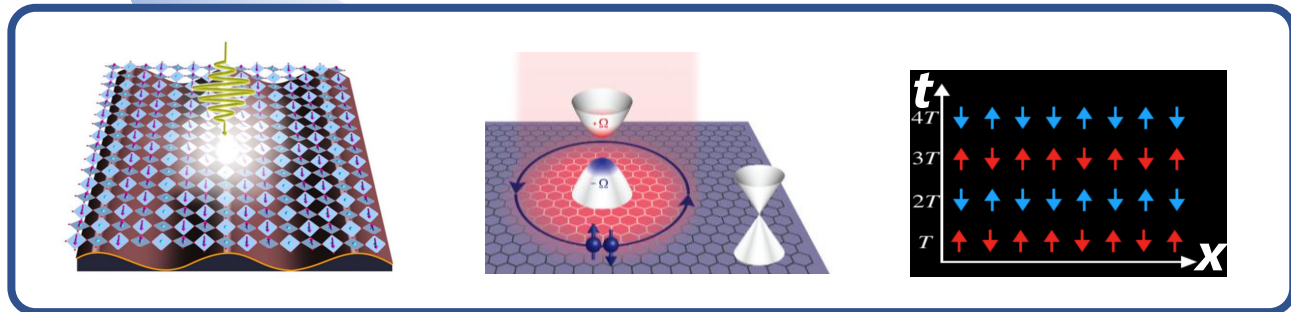
Temperature

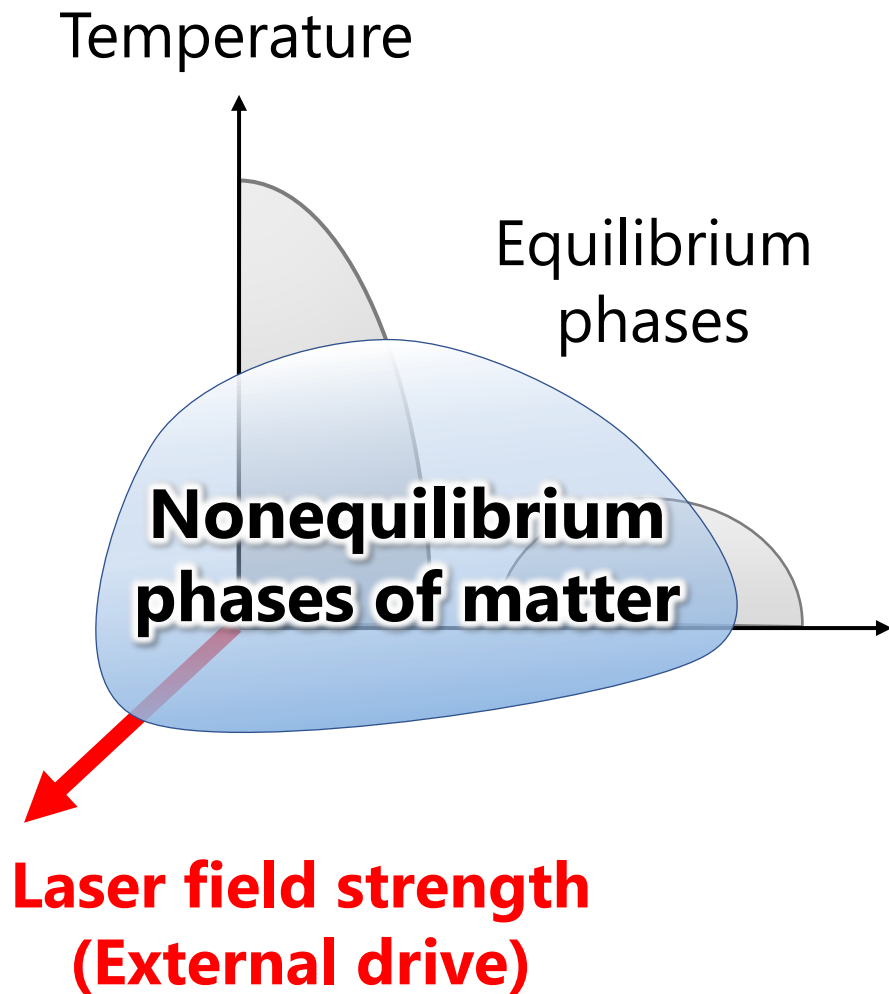


length
ive)



Pressure / Doping





Deepen our understanding of phases of matter

Future novel technology

There remain many challenges...

1. To explore new types of phases of matter

Quantum Active Matter (topic 1)

[KT*](#), K. Adachi*, K. Kawaguchi, Phys. Rev. Research **6**, 023096 (2024)

K. Adachi, [KT](#), K. Kawaguchi, Phys. Rev. Research **4**, 013194 (2022)

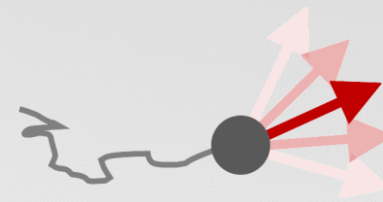
2. To understand their universal aspects

KPZ universality in spin chains (topic 2)

K. A. Takeuchi, [KT](#), O. Busani, P. L. Ferrari, R. Vasseur, J. De Nardis, Phys. Rev. Lett. **134**, 097104 (2025)

Research Interests	Nonequilibrium Phases of Matter
Topic 1	Quantum Active Matter
Topic 2	KPZ in the Heisenberg Spin Chain
Discussion	Possible Connections to Gravity

What is active matter ?



Self-propelled
particle

Collection of self-propelled particles = **Active matter**

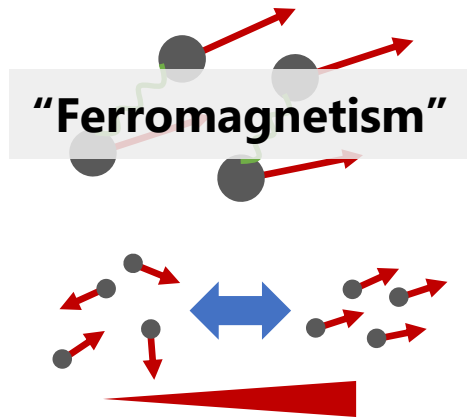
Nonequilibrium phase transitions not happening in static matter

Nonequilibrium phase transitions

Nonequilibrium phase transitions not happening in static matter



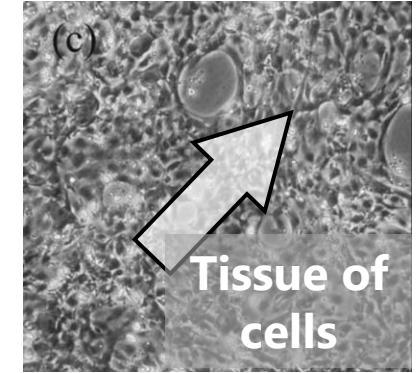
Flocking transition Flock = 群れ in Japanese



+ **Aligning interaction**

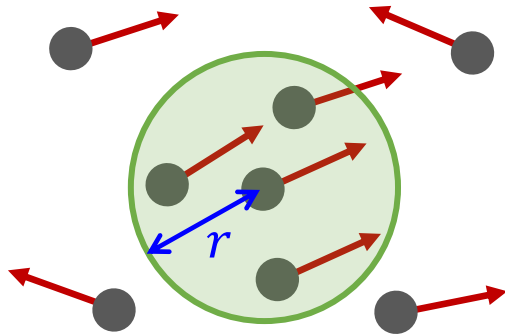
→ SSB of rotational symmetry

Can happen in 2D at finite-T,
PROHIBITED by the Mermin-Wagner theorem



B. Szabó et al. Phys. Rev. E **74**, 061908 (2006)

2D Vicsek model



Fixed velocity
Directional noise

T. Vicsek et al. PRL (1995)

High temperature

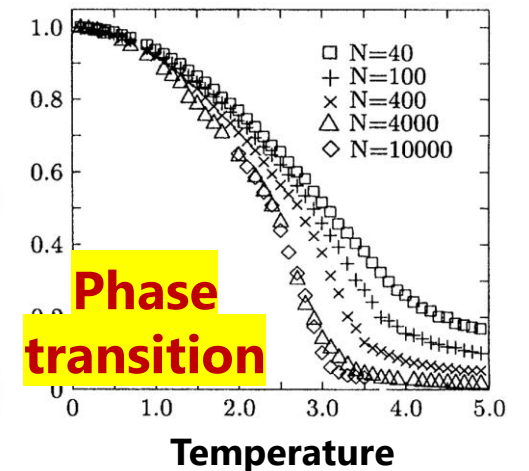


Low temperature



Video by K. Adachi

Order parameter

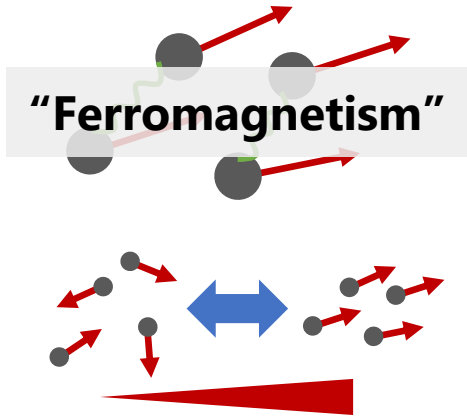


Nonequilibrium phase transitions

Nonequilibrium phase transitions not happening in static matter



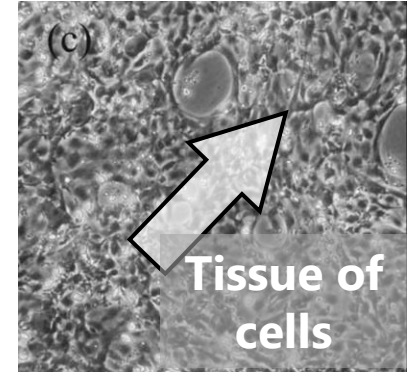
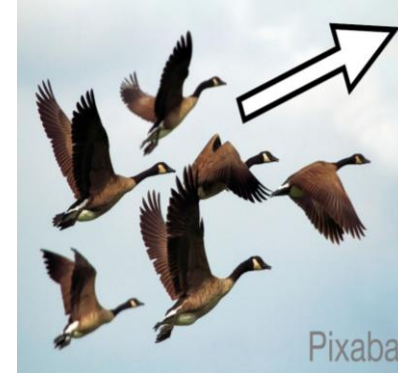
Flocking transition Flock = 群れ in Japanese



+ **Aligning interaction**

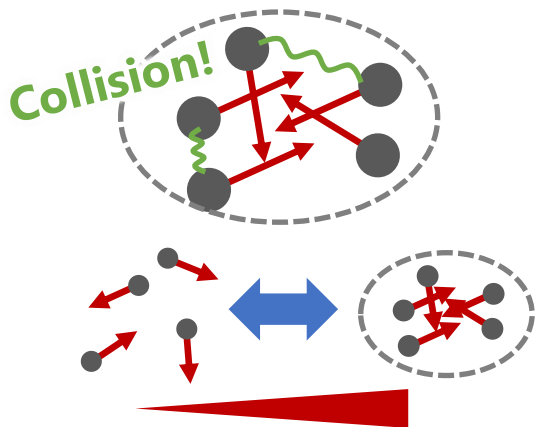
→ SSB of rotational symmetry

Can happen in 2D at finite-T,
PROHIBITED by the Mermin-Wagner theorem



B. Szabó et al. Phys. Rev. E **74**, 061908 (2006)

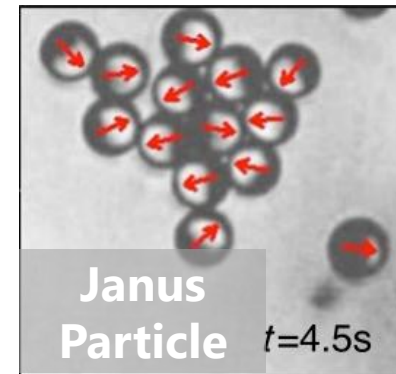
Motility-induced phase separation (MIPS)



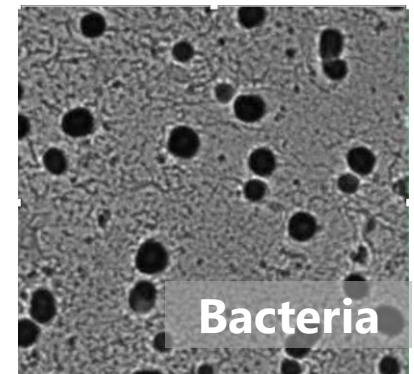
+ **Repulsive interaction**

→ "Droplet" is formed

Phase separation occurs
WITHOUT attractive force



I. Buttinoni et al., Phys. Rev. Lett. **110**, 238301 (2013)

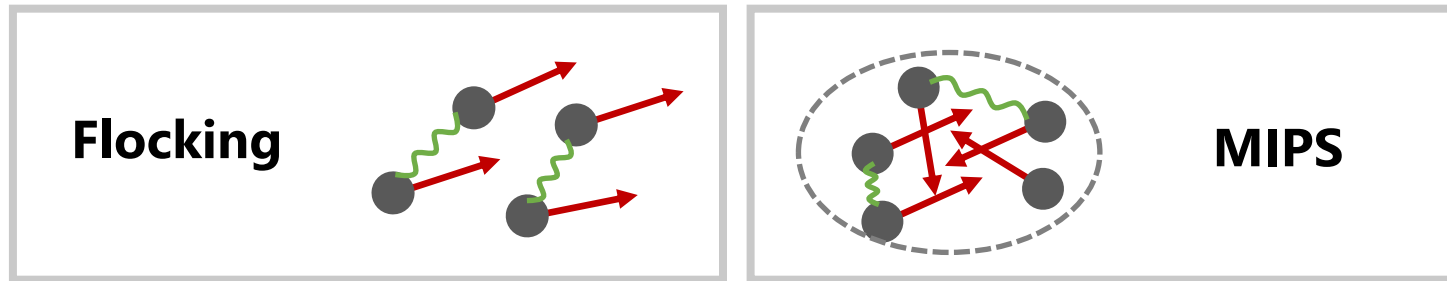


G. Liu et al. Phys. Rev. Lett. **122**, 248102 (2019)

Nonequilibrium phase transitions

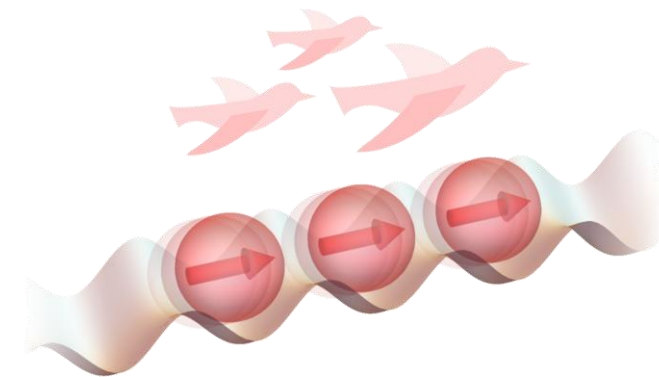
13/35

Nonequilibrium phase transitions not happening in static matter



Quantum analogue of these phase transitions ?

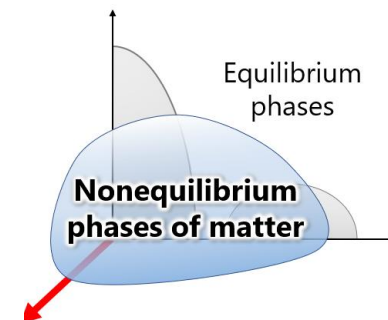
e.g., Increase activity (velocity) \rightarrow Ferromagnetism



It has **not been considered** when we have started this project...


\therefore The main target of active matter physics has been **living things** and it has been studied in the field of **biophysics**

I believe this is a natural question both for stat-phys and cond-mat



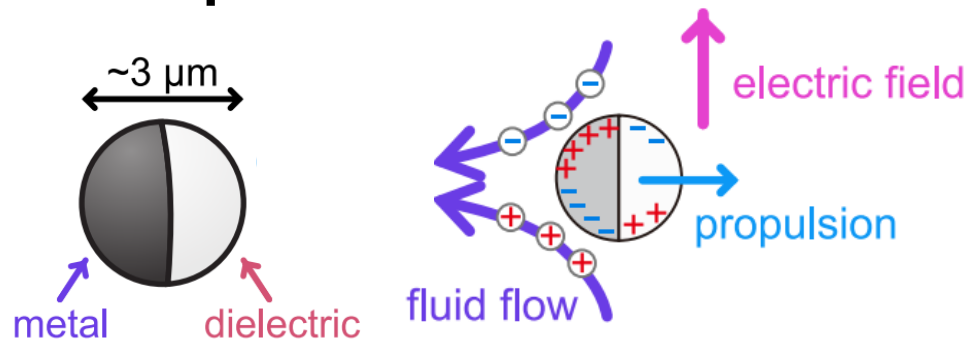
20/37

What is (can be) “quantum active matter” ?

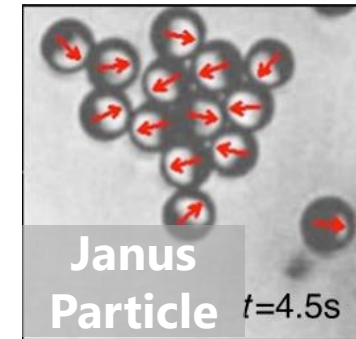
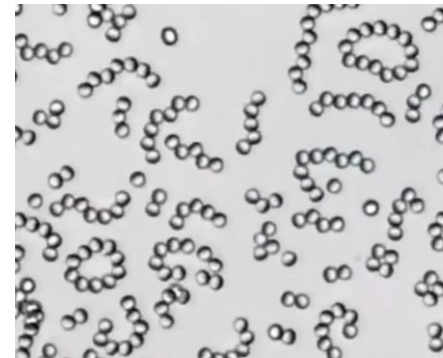
We don't know “quantum living things” |  \rangle

Our strategy: **Design** quantum system that behaves like active matter

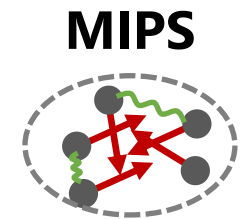
Ex. Janus particle



D. Nishiguchi, JPSJ, **92**, 121007 (2023)

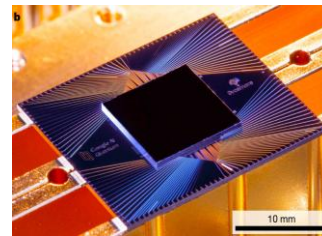


Janus Particle $t=4.5\text{s}$
I. Buttinoni et al., Phys. Rev. Lett. **110**, 238301 (2013)

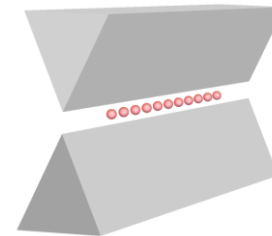


Now, we have various platforms for highly-tunable (programmable) quantum many-body systems

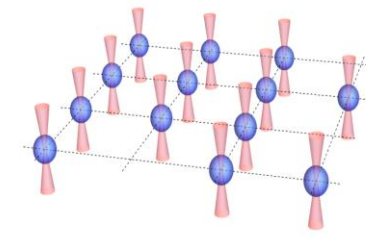
Let's use them to design quantum active matter!



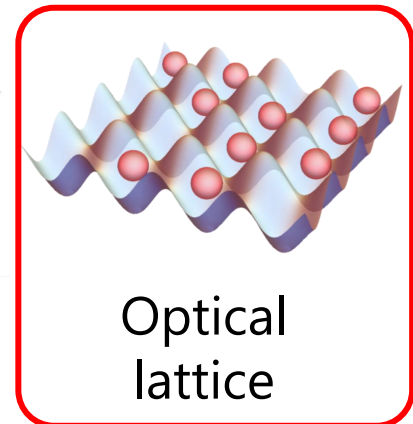
SC qubits



Trapped ions



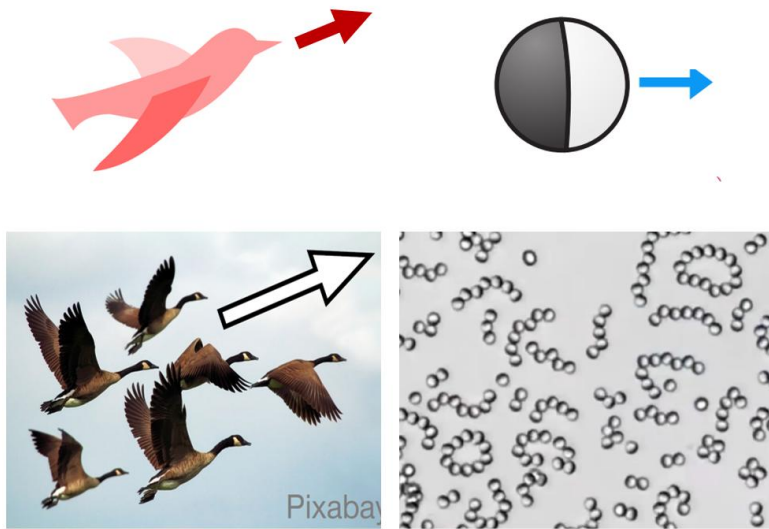
Optical tweezer



Optical lattice

Our guiding principles to construct quantum active matter

Our guiding principle to construct q. active matter

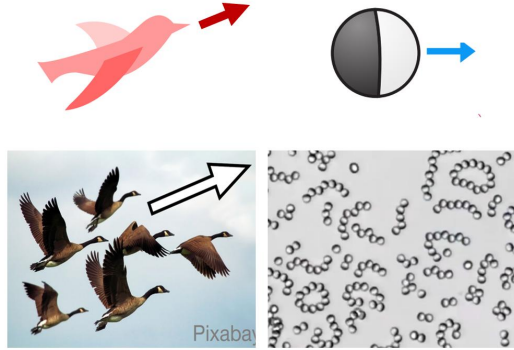


1. Governed by quantum mechanics
2. Nonequilibrium many-body system
3. Each element has own orientation
4. The orientation can change (no special direction)
- 5. The orientation determines the kinetic motion**

Remark: We do not claim that this is the definition

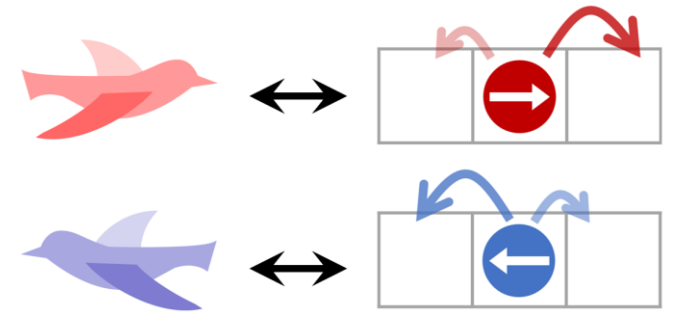
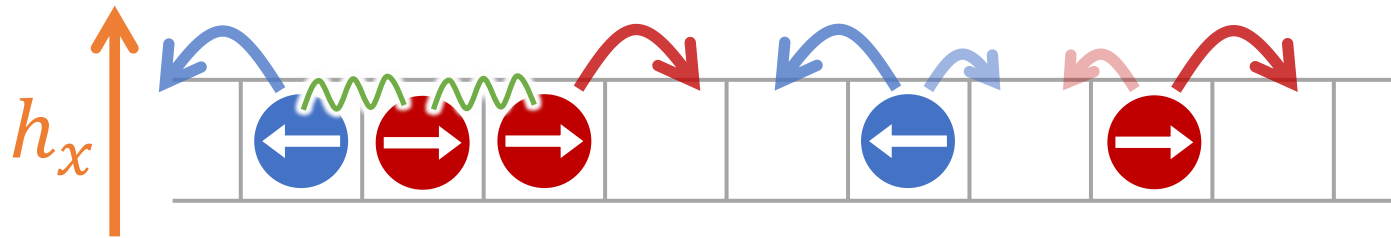
K. Adachi, [KT](#), K. Kawaguchi, Phys. Rev. Research **4**, 013194 (2022)
[KT*](#), K. Adachi*, K. Kawaguchi, Phys. Rev. Research **6**, 023096 (2024)

Our quantum active matter model

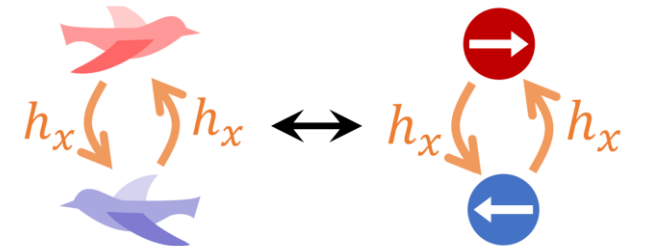


1. Governed by quantum mechanics
2. Nonequilibrium many-body system
3. Each element has own orientation
4. The orientation can change (no special direction)
- 5. The orientation determines the kinetic motion**

Two-component interacting bosons (fermions) in 1D

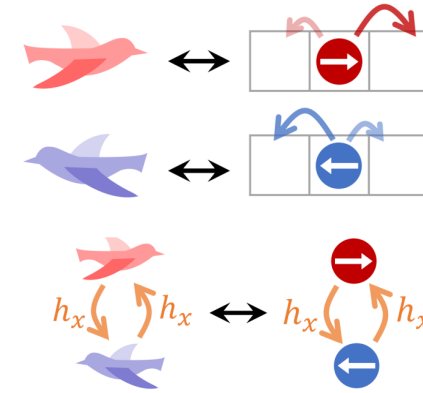
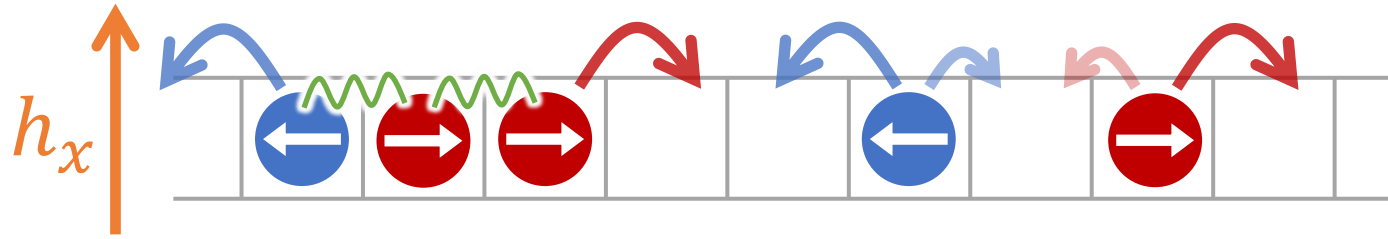


$$\hat{H} = -t \sum_{i=1}^L \sum_{s=\pm} \{ (1+s\varepsilon) \hat{a}_{i+1,s}^\dagger \hat{a}_{i,s} + (1-s\varepsilon) \hat{a}_{i,s}^\dagger \hat{a}_{i+1,s} - h_x \hat{a}_{i,s}^\dagger \hat{a}_{i,-s} \} + \hat{H}_{\text{int}}$$



One of the simplest models of "quantum active matter"

*not a unique option



$$\hat{H} = -t \sum_{i=1}^L \sum_{s=\pm} \{ (1+s\varepsilon) \hat{a}_{i+1,s}^\dagger \hat{a}_{i,s} + (1-s\varepsilon) \hat{a}_{i,s}^\dagger \hat{a}_{i+1,s} - h_x \hat{a}_{i,s}^\dagger \hat{a}_{i,-s} \} + \hat{H}_{\text{int}}$$

1. Relation to classical active matter

$$\frac{dp}{dt} = Hp$$

Master equation of Markov process for active lattice gas model

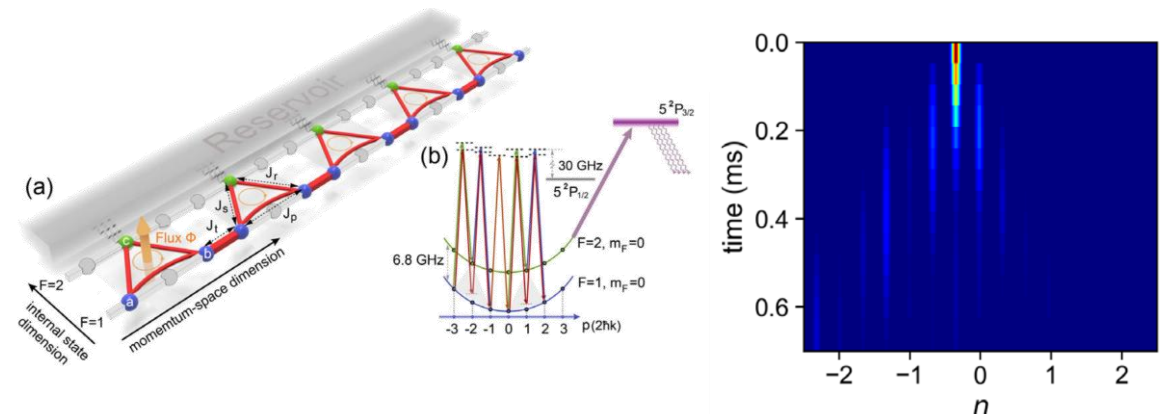
Asymmetric hopping \Leftrightarrow Asymmetric transition rate
= activity in the classical cases

Our construction is not just an analogy!
Formal reason to call this model "active matter"

K. Adachi, [KT](#), K. Kawaguchi, PRR **4**, 013194 (2022)

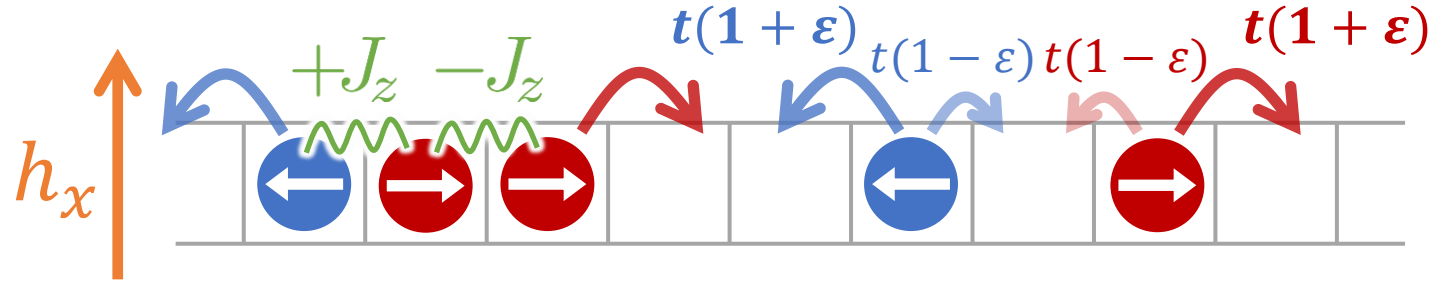
2. Experimental realization

Spin-**independent** one has been realized by ultracold atoms in a dissipative optical lattice



Q. Liang, ..., B. Yan, Phys. Rev. Lett. **129**, 070401 (2022)

Model: Non-Hermitian 2-component interacting bosons



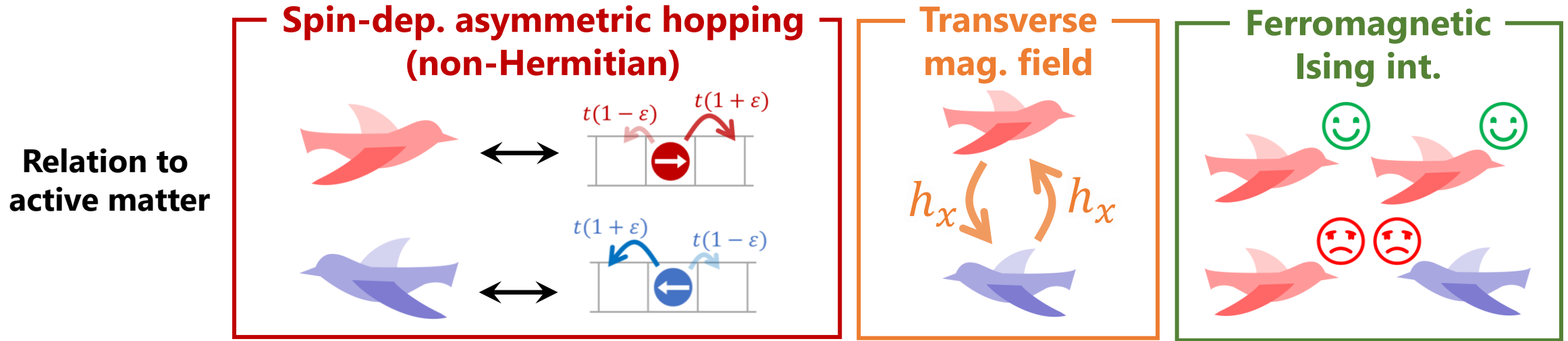
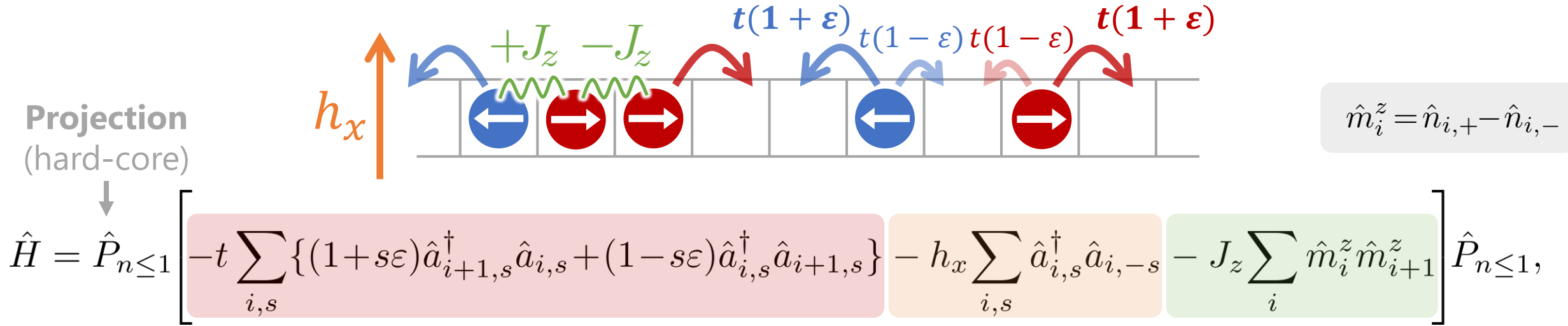
$$\hat{m}_i^z = \hat{n}_{i,+} - \hat{n}_{i,-}$$

$$\hat{H} = -t \sum_{i,s} \{ (1+s\epsilon) \hat{a}_{i+1,s}^\dagger \hat{a}_{i,s} + (1-s\epsilon) \hat{a}_{i,s}^\dagger \hat{a}_{i+1,s} \} - h_x \sum_{i,s} \hat{a}_{i,s}^\dagger \hat{a}_{i,-s} - J_z \sum_i \hat{m}_i^z \hat{m}_{i+1}^z$$

$$+U \sum_i \hat{n}_{i,+} \hat{n}_{i,-}$$

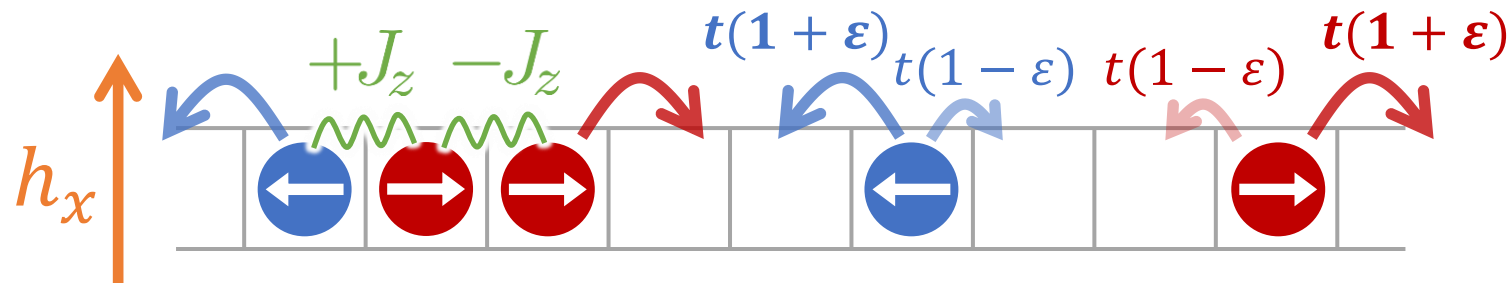
On-site repulsive interaction

Model: Non-Hermitian 2-component interacting bosons

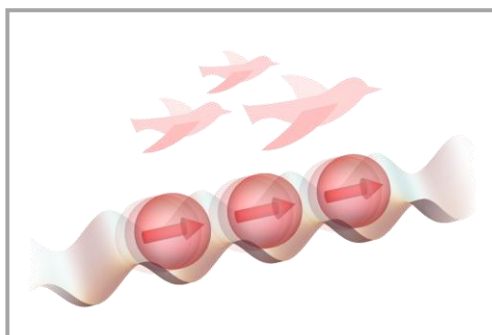


We study the ground state (= the eigenstate with the **smallest Re[E]**) of this model in the following
 Reasons: 1. Experimentally relevant, 2. Corresponds to the steady state in Markov process

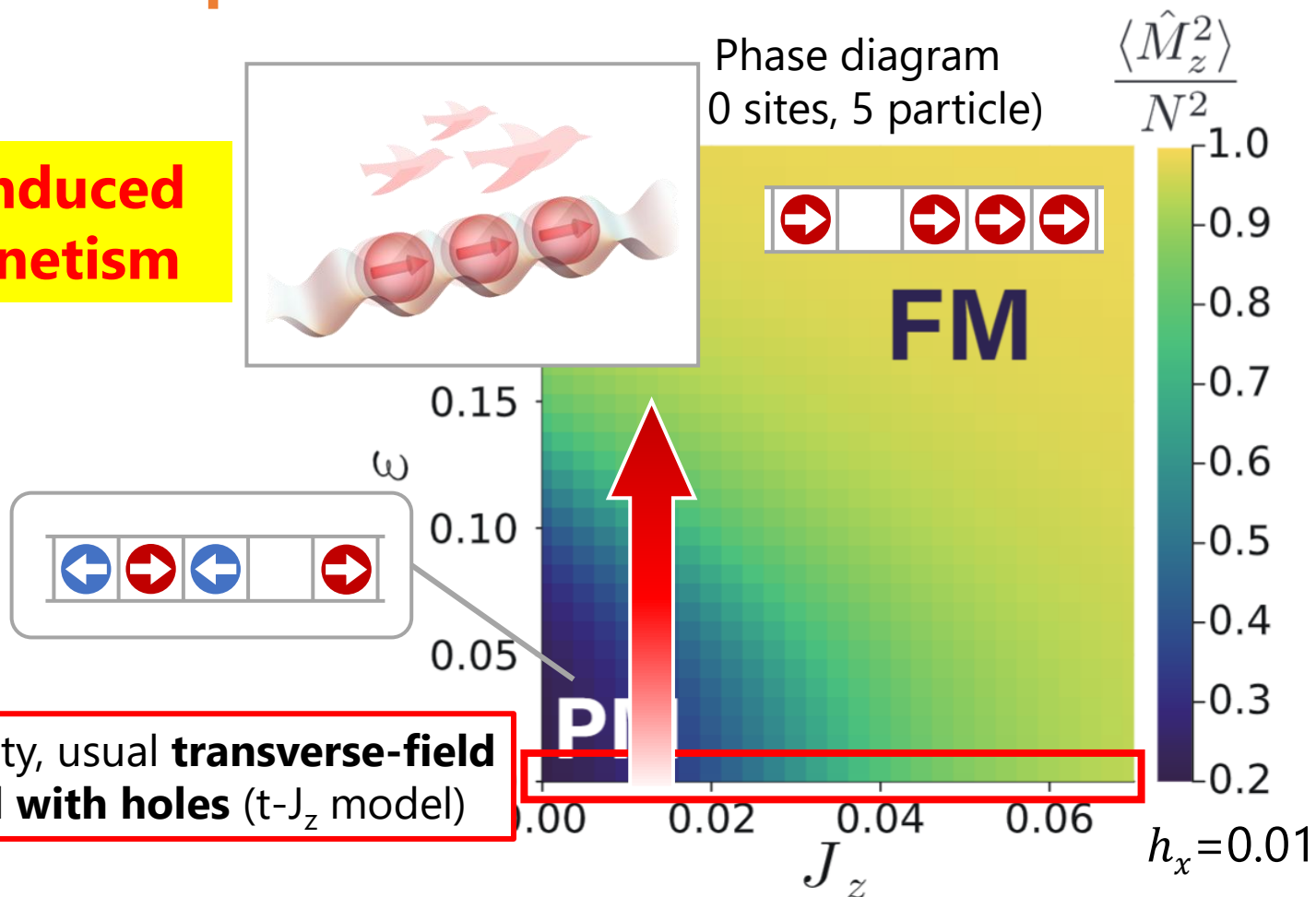
Numerical results: Exact diagonalization



Activity-induced ferromagnetism



Phase diagram (0 sites, 5 particle)



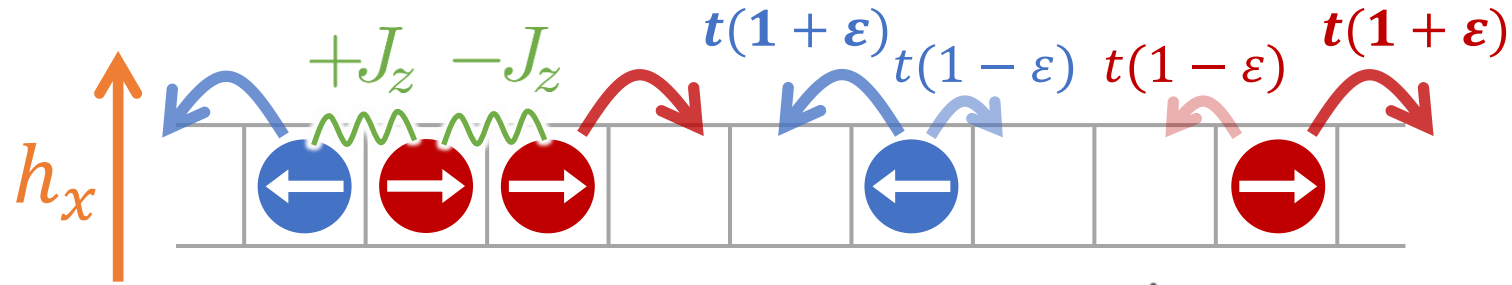
Without activity, usual **transverse-field Ising model with holes** (t - J_z model)

cf. Active Ising model (classical)

No FM order in 1D

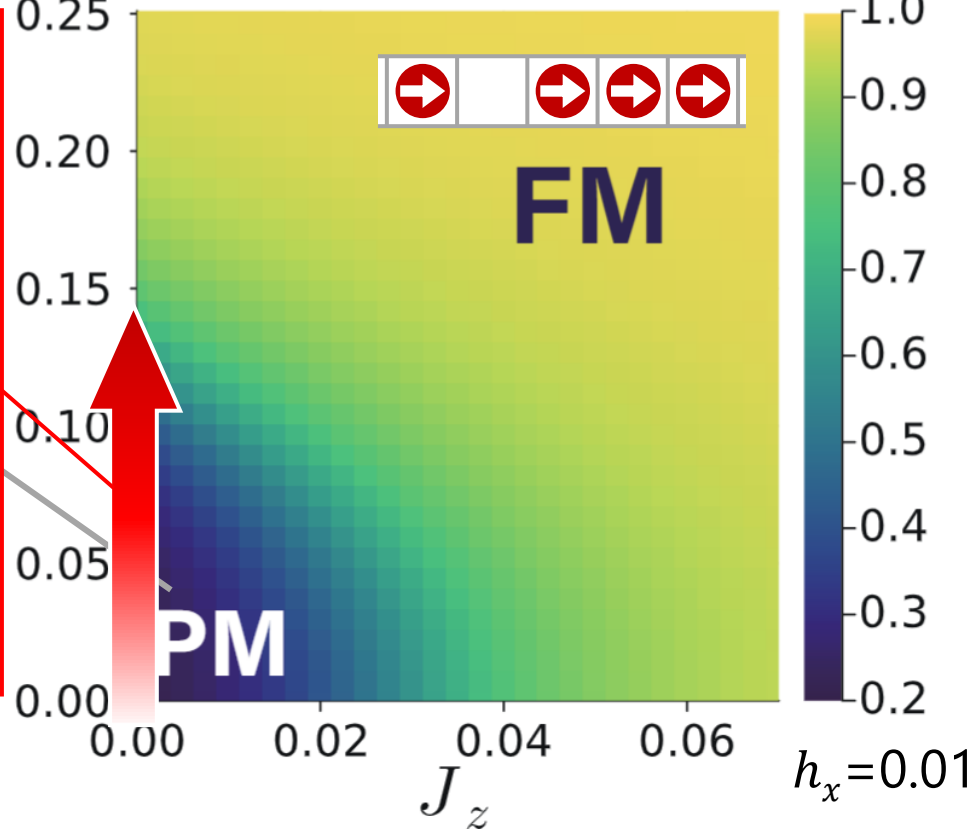
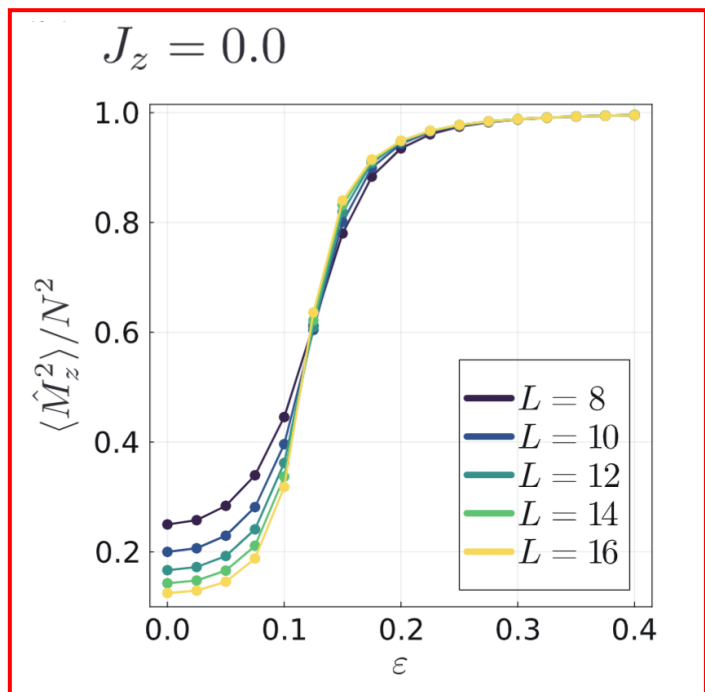
A. P. Solon, J. Tailleur, Phys. Rev. Lett. (2013); Phys. Rev. E (2015)

Numerical results: Exact diagonalization

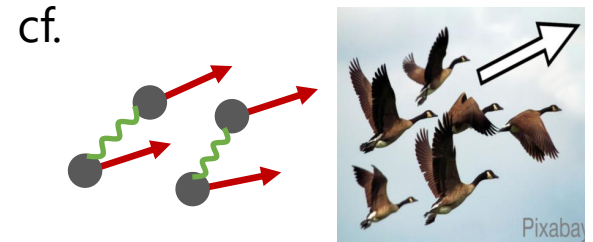


Phase diagram
(10 sites, 5 particle)

$$\frac{\langle \hat{M}_z^2 \rangle}{N^2}$$

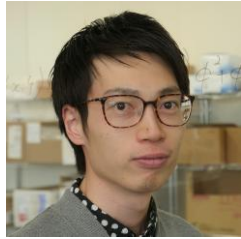


Ferromagnetic order emerges without the Ising interaction

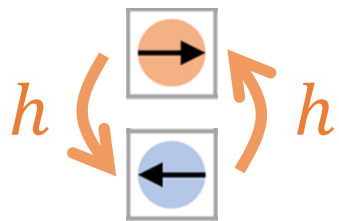
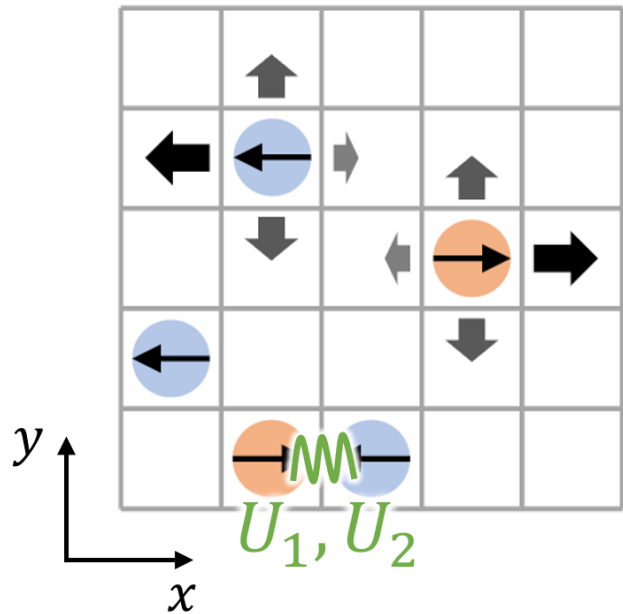


Flocking without aligning interaction is a **unique feature of quantum?**

Two-dimensional quantum active matter

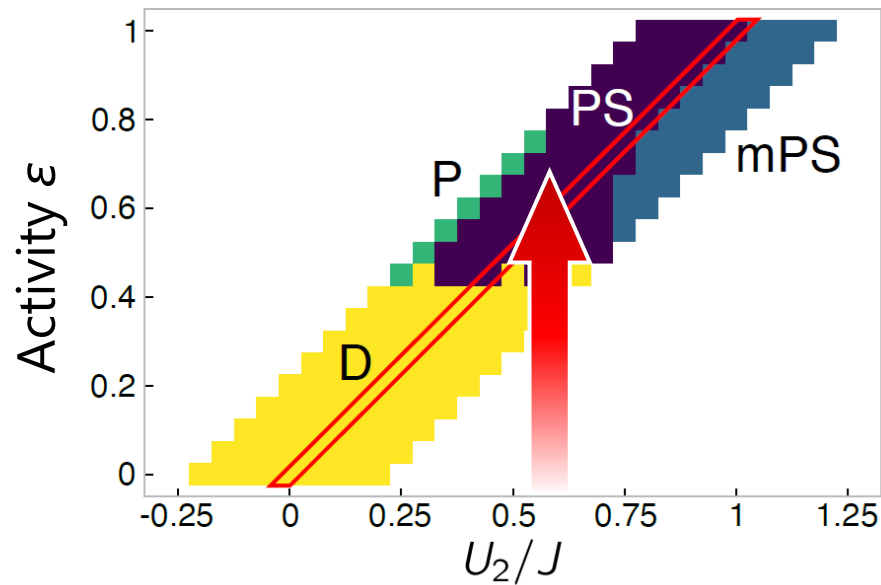


two-component hardcore bosons on the square lattice



Quantum Monte Carlo simulation
(Diffusion Monte Carlo)

$$U_1 = 2J$$



Phase separation



Disorder

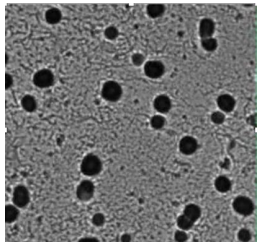
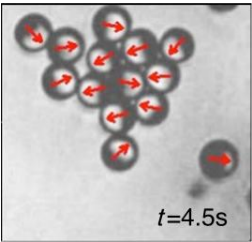
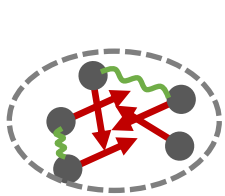


MIPS in quantum active matter

micro PS (mPS)

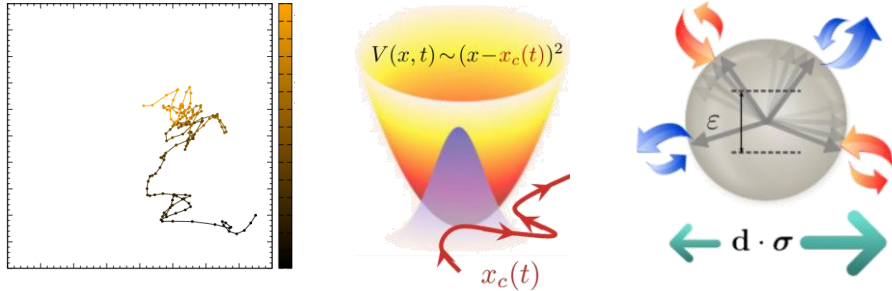


Polar order (ferromag.)



stabilized by quantum effect

Active quantum particle



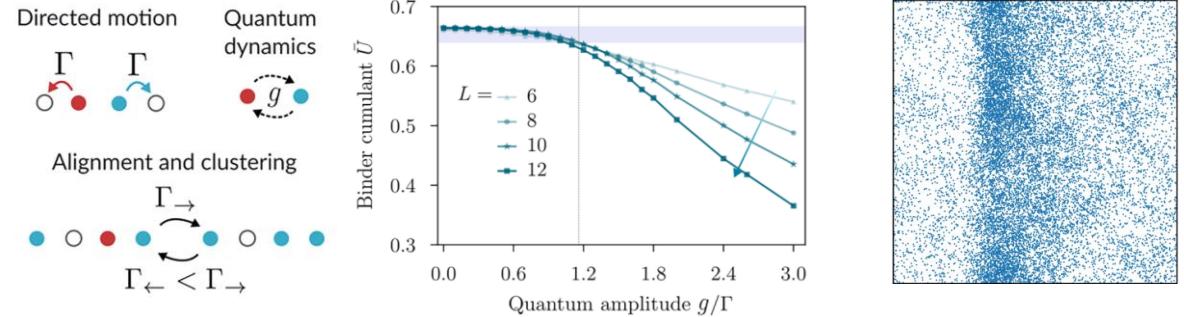
Yamagishi, Hatano, Obuse, Sci. Rep. (2024)

Antonov, Zheng, Liebchen, Löwen, PRR (2025)

Penner, Viotti, Fazio, Arrachea, von Oppen, PRB (2025)

J. Gipouloux, M. Brunelli, L. Cugliandolo, R. Fazio, M. Schirò, arXiv (2026)

Flocks in different quantum models

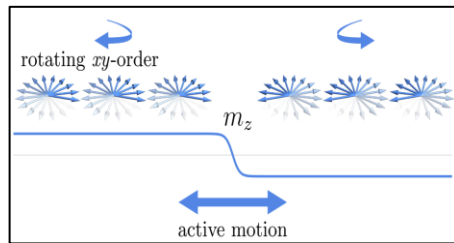


Khasseh, Wald, Moessner, Weber, Heyl, PRL (2025)

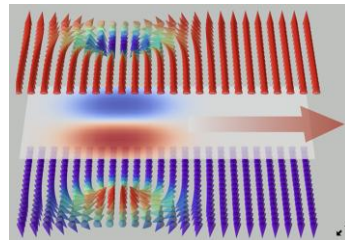
Yuan, Cui, Chen, Sun, arXiv (2024)

Steiner, von Oppen, Egger arXiv (2026)

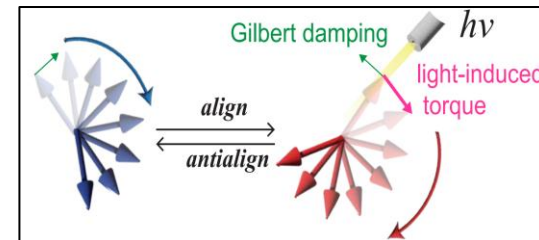
Topological spin texture behaving like active matter Nonreciprocal phase transitions



Hardt, Doostani, Diehl, del Ser, Rosch, Nat. Comm. (2025)



de Souza Silva, Correia, Pina Velasquez PRL (2025)



Hanai, Ootsuki, Tazai, Nat. Comm. (2025)

Challenges: 1. Properties unique to "quantum" active matter? , 2. Universal description, 3. Experiments.

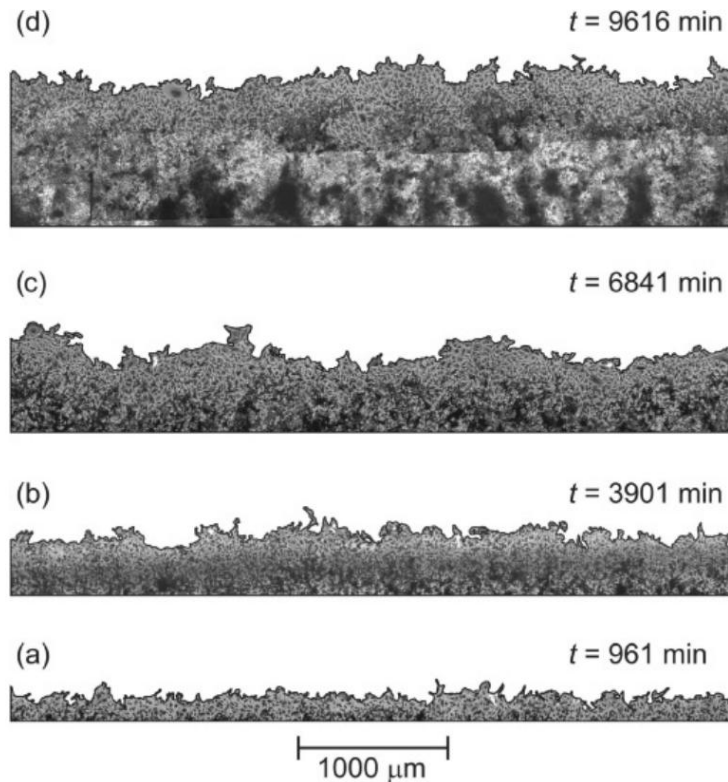
Research Interests	Nonequilibrium Phases of Matter
Topic 1	Quantum Active Matter
Topic 2	KPZ in the Heisenberg Spin Chain
Discussion	Possible Connections to Gravity

Kardar-Parisi-Zhang (KPZ) Universality

= One of the universality classes for nonequilibrium fluctuations

Interface Growth

Cancer cells



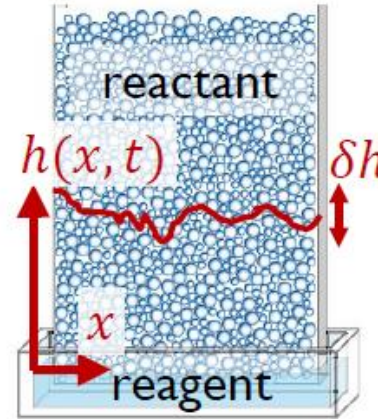
M. A. C. Huergo et al. PRE (2015)

Fire front



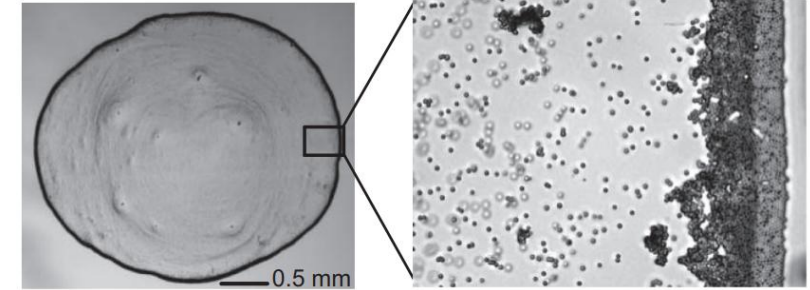
J. Maunuksela et al PRL (1997)

Chemical reaction



S. Atis et al PRL (2013)

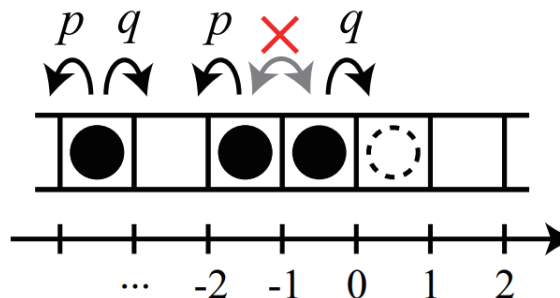
Particle deposition



P.J. Yunker et al. Nature (2011)

Particle Transport

ASEP



Nonlinear Fluctuating Hydrodynamics

Anharmonic chain

$$H = \sum_{j=1}^N \left(\frac{1}{2} p_j^2 + V(q_{j+1} - q_j) \right)$$

Caveat: Not all interface growth are KPZ

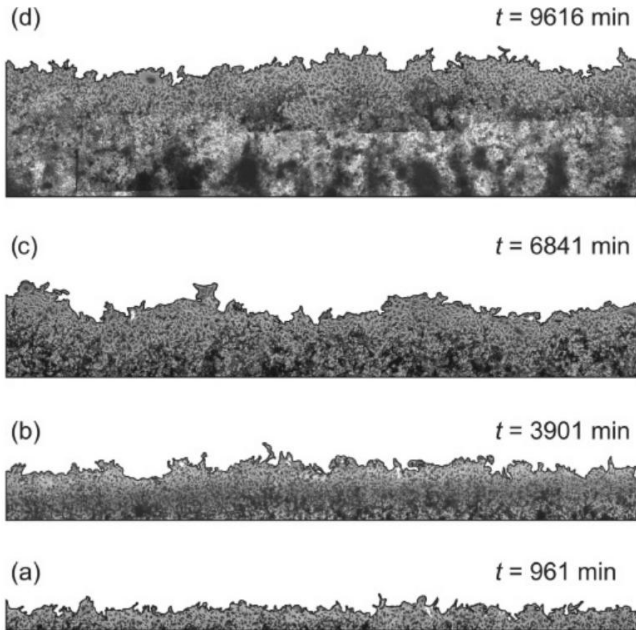
Review: K. A. Takeuchi, Physica A 504, 77-105 (2018)

Kardar-Parisi-Zhang (KPZ) Universality

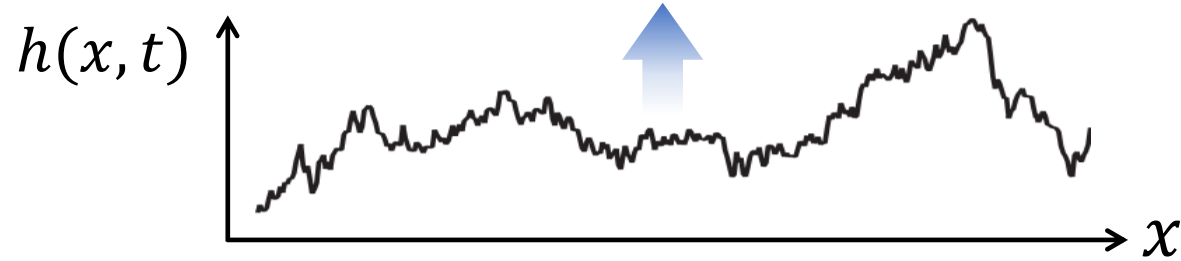
= One of the universality classes for nonequilibrium fluctuations

Interface Growth

Cancer cells



M. A. C. Huergo et al. PRE (2015)



KPZ equation $\partial_t h = v_0 + v \partial_x^2 h + \frac{\lambda}{2} (\partial_x h)^2 + \sqrt{D} \eta$
 K.-P.-Z. PRL (1986)

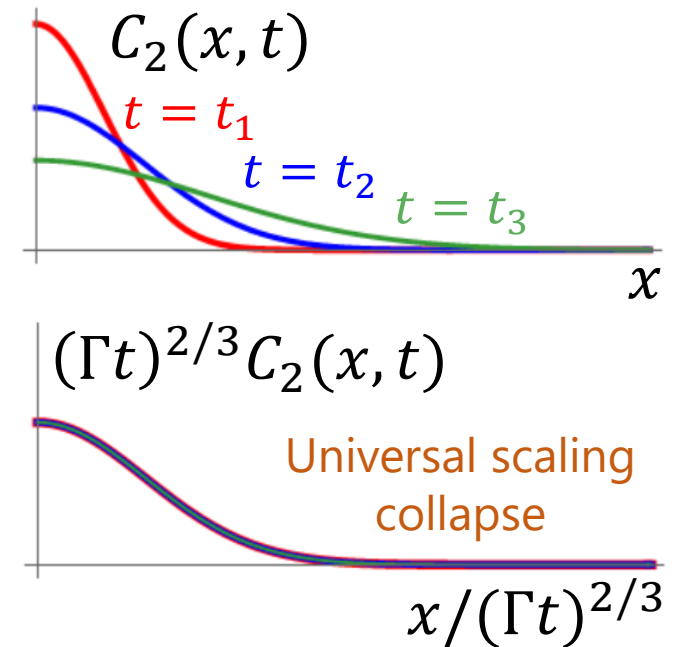
Universal Exponent

$h(x_0, t) \sim v_\infty t + (\Gamma t)^{1/3} \chi_{St}$
 Universal Distribution

Correlation function

$C_2(x, t) = \langle \partial_x h(x, t) \partial_x h(0, 0) \rangle$
 $\approx \frac{A}{(\Gamma t)^{2/3}} f_{KPZ} \left(\frac{x}{(\Gamma t)^{2/3}} \right)$

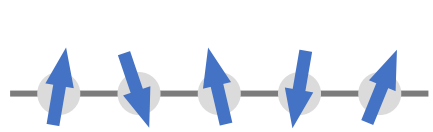
f_{KPZ} : Prähofer–Spohn scaling function



*This talk only treats the stationary subclass

Emergence of the KPZ universality in a spin chain

Heisenberg spin chain



$$\hat{H} = J \sum_{i=1}^{L-1} \hat{S}_i \cdot \hat{S}_{i+1}$$

Two-point function at $T=\infty$

$$C_2(x, t) = \frac{1}{2L} \text{Tr}[\hat{S}_x^z(t) \hat{S}_0^z(0)]$$

$$\left[\hat{S}_x^z(t) = e^{i\hat{H}t} \hat{S}_x^z e^{-i\hat{H}t} \right]$$

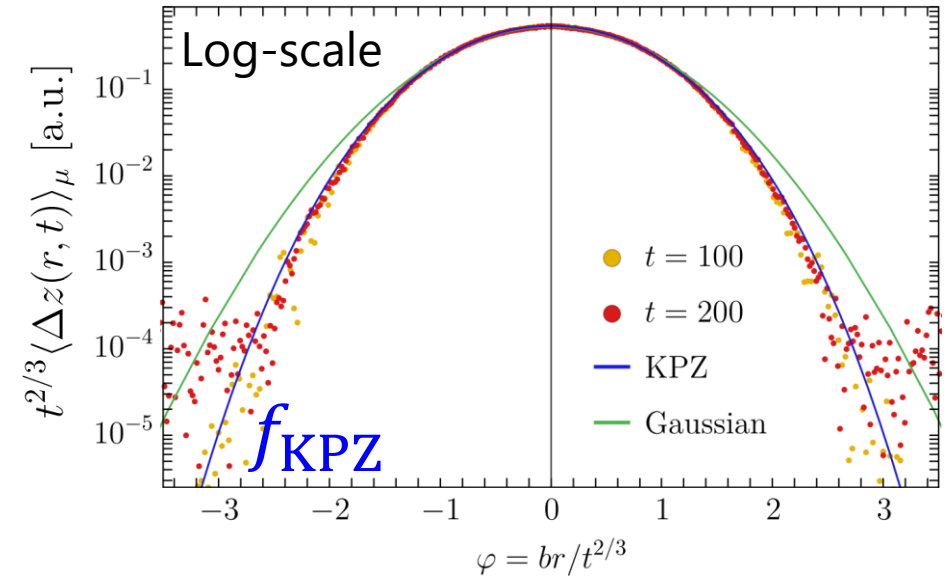
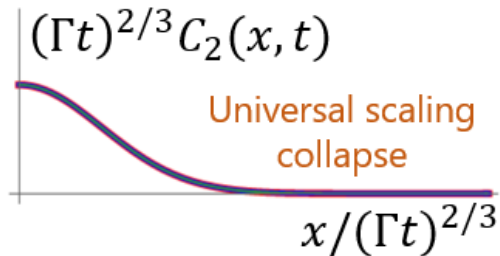
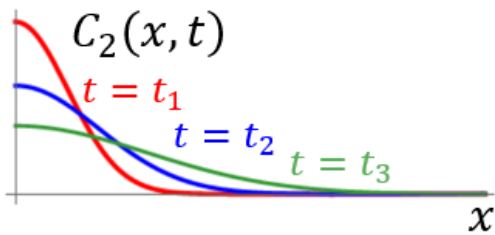
The KPZ scaling function f_{KPZ} emerges in the infinite temperature two-point function !

PHYSICAL REVIEW LETTERS 122, 210602 (2019)

Kardar-Parisi-Zhang Physics in the Quantum Heisenberg Magnet

Marko Ljubotina, Marko Žnidarič, and Tomaž Prosen

Physics Department, Faculty of Mathematics and Physics, University of Ljubljana, 1000 Ljubljana, Slovenia



Surprise

KPZ behavior emerges in deterministic quantum dynamics, whereas it is conventionally associated with classical stochastic systems.

Remark: The KPZ behavior can also appear in the classical Heisenberg-type spin chain.

Discrepancy with the KPZ predictions

Dictionary

$$C_2(x, t) = \langle \partial_x h(x, t) \partial_x h(0, 0) \rangle \simeq \langle S_x^z(t) S_0^z(0) \rangle_{T=\infty}$$

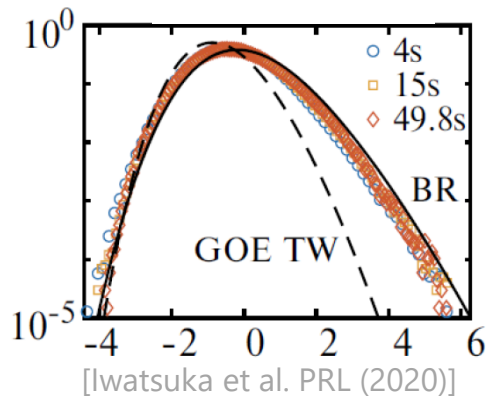
$$\text{Slope } \partial_x h(x, t) \leftrightarrow -\hat{S}_x^z(t)$$

$$\text{Height } h(x, t) \leftrightarrow \int_0^t \hat{j}_{x+1/2}^z(t) dt =: \hat{Q}(t) \quad \text{Integrated spin current}$$

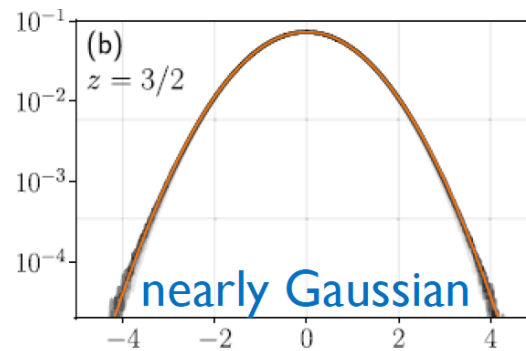
$$\left[\text{Spin current op. } \hat{j}_{x+1/2} = \frac{ij}{2} (S_x^+ S_{x+1}^- - S_x^- S_{x+1}^+) \right]$$

Other quantities not matching with the KPZ predictions

One-point distribution of $h(x, t)$ vs $\hat{Q}_x(t)$



[Iwatsuka et al. PRL (2020)]



[Krajnik et al. PRL (2024)]

$h(x, t)$: Baik-Rains Distribution
Finite Skewness & Kurtosis

$\hat{Q}_x(t)$: Nearly Gaussian
Skewness = 0 & Kurtosis \simeq 0

Confirmed via

numerical simulation

Krajnik et al. Phys. Rev. Lett. 132, 017101 (2024)

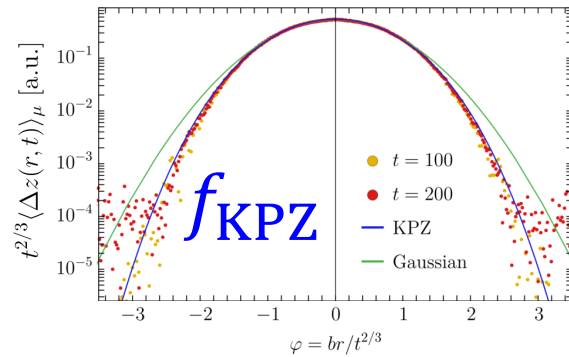
and **q. comp. simulation**

Rosenberg et al. (Google) Science 384, 48 (2024)

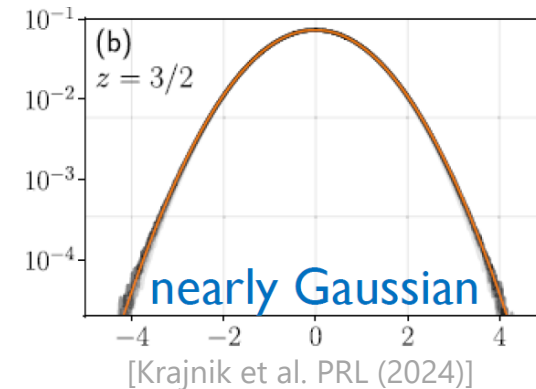
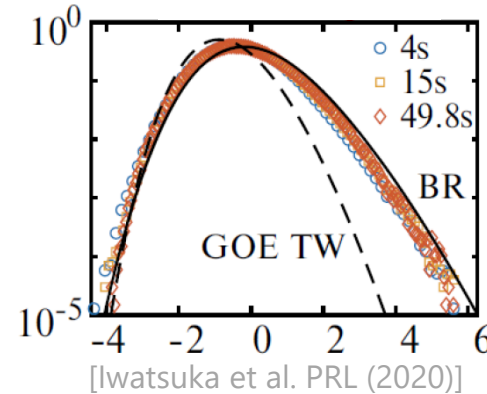
Clear discrepancy with the KPZ predictions

Then, is it really KPZ ?

Two- point spin-spin correlation function



One-point distribution of height / Integrated spin current



It is not a conventional emergence of the KPZ universality.
But, what's going on here ? We need more information.

Our work K. A. Takeuchi, [KT](#), O. Busani, P. L. Ferrari, R. Vasseur, J. De Nardis, PRL 134, 097104 (2025)

Based on the detailed numerical simulation, we have examined other correlators:

- **All the two-point correlators match with KPZ, including the numerical prefactor**
- **Higher-order / non-two-point quantities do not match with KPZ**

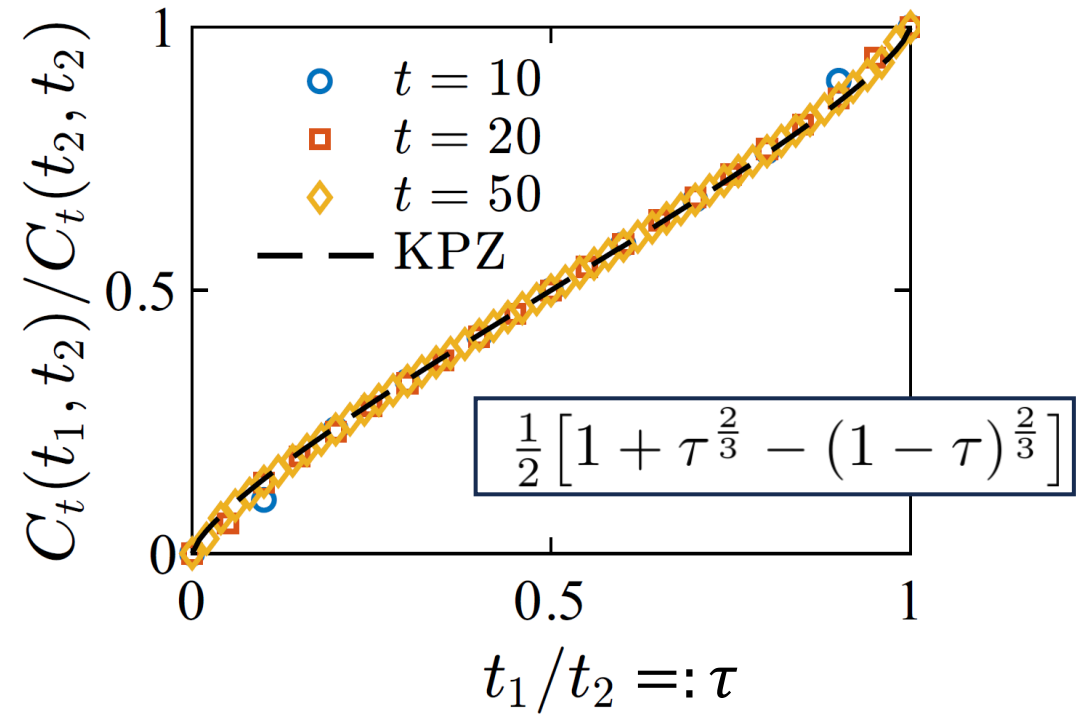
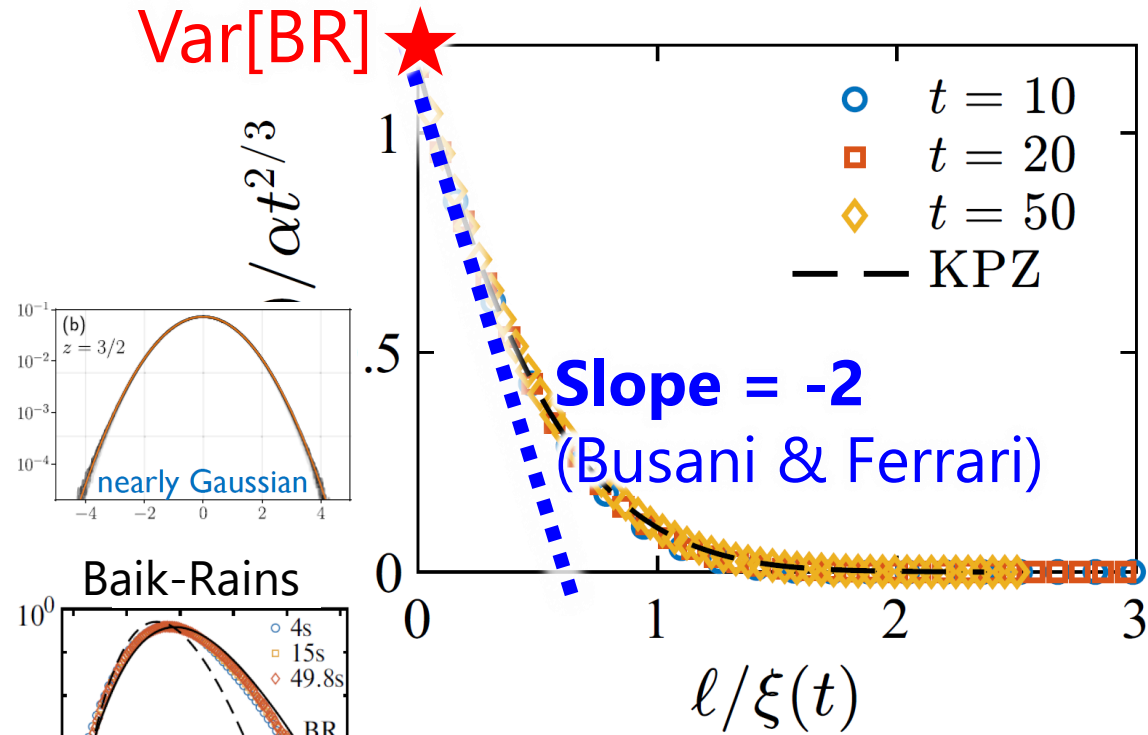
Then, we have concluded that this is a **partial yet definite** emergence of the KPZ universality

Our results: Partial yet definite emergence of KPZ

Spatial correlation of integrated spin current $C_s(\ell, t) = \langle h(x, t)h(x + \ell, t) \rangle - \langle h(x, t) \rangle^2$

Temporal correlation of integrated spin current $C_t(t_1, t_2) = \langle h(x, t_1)h(x, t_2) \rangle - \langle h(x, t_1) \rangle \langle h(x, t_2) \rangle$

[c.f. Previous works: two-point spin-spin correlation $C_2(x, t) = \langle \partial_x h(x, t) \partial_x h(0, 0) \rangle$]



The **variance** of the BR distribution appears !

1. Extending the quantum lattice models showing the KPZ-like behavior

i. **Integrability** and ii. **Non-Abelian symmetry** play an important role

e.g., One-dimensional Hubbard model C. P. Moca et al. PRB 108, 235139 (2023)

Remark: The dynamical exponent $z = 3/2$ (superdiffusion) is not enough to claim KPZ.
The **scaling function** needs to be examined.

$$\xi(t) \sim t^{1/z}$$

Sharp conditions to obtain the KPZ behavior are not obtained yet.

2. Theoretical/Analytical understanding of the KPZ behavior is still an open problem

Two-component nonlinear fluctuating hydrodynamics

J. De Nardis, S. Gopalakrishnan, R. Vasseur, Phys. Rev. Lett. **131**, 197102 (2023)

Related review: S. Gopalakrishnan, R. Vasseur, Ann. Rev. Cond. Matt. Phys. 15, 159 (2024)

Their hydrodynamics does not reproduce the kurtosis correctly

Further theoretical investigation for why there is the partial emergence is necessary

Outline

Research Interests

Nonequilibrium Phases of Matter

Topic 1

Quantum Active Matter

Topic 2

KPZ in the Heisenberg Spin Chain

Discussion

Possible Connections to Gravity

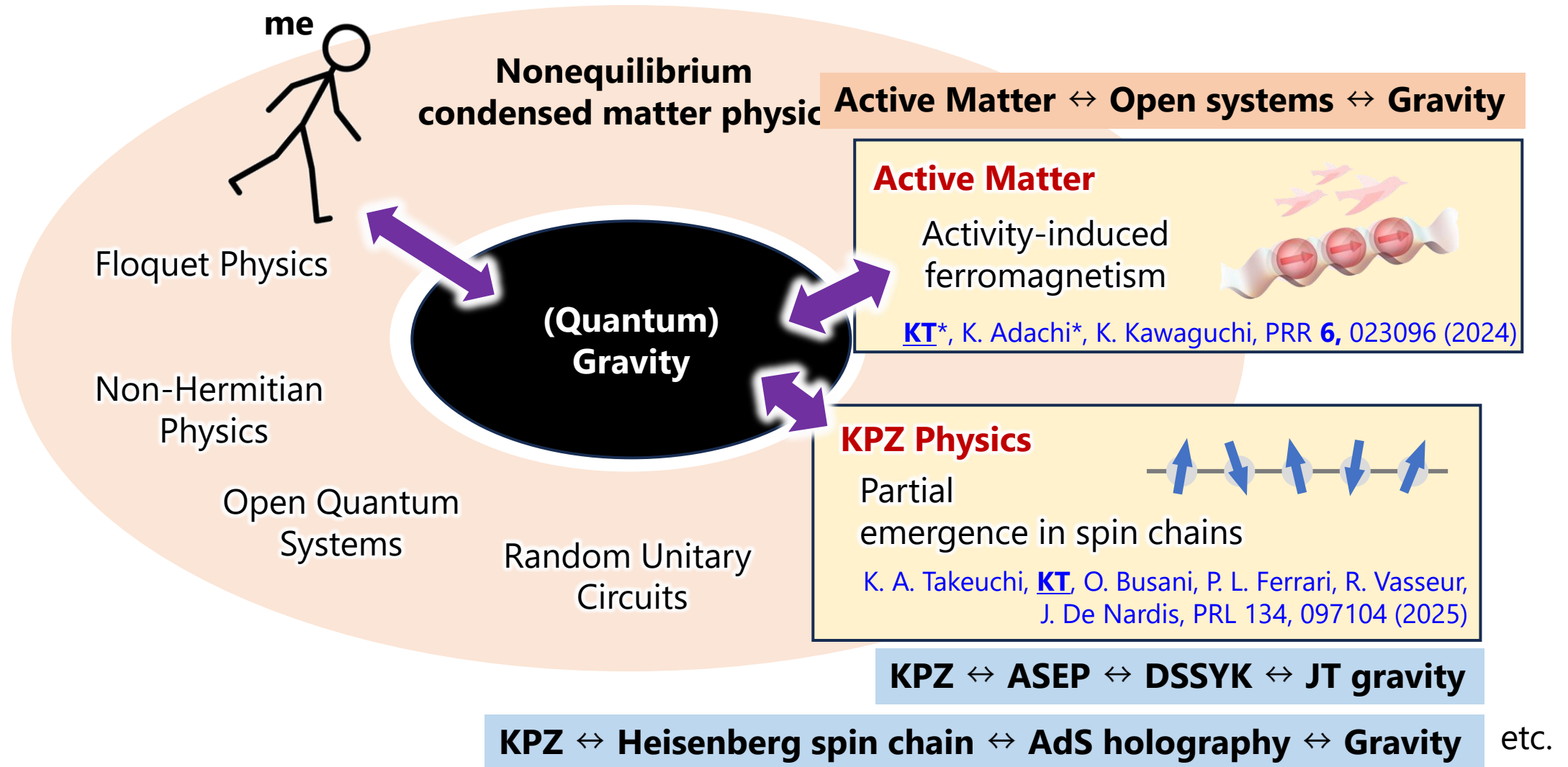
Discussions: Possible Connections to Gravity

Omitted in the shared version

Discussions: Possible Connections to Gravity

Omitted in the shared version

Summary: Active Matter, KPZ, and (possibly) Gravity



I would be happy to discuss this with you,
hear your ideas, and hopefully explore a possible collaboration.

Kazuaki Takasan (Univ. of Tokyo)
E-mail: takasan@ap.t.u-tokyo.ac.jp