

Yukawa Unification in Renormalizable SU(5) GUT with Froggatt-Nielsen mechanism

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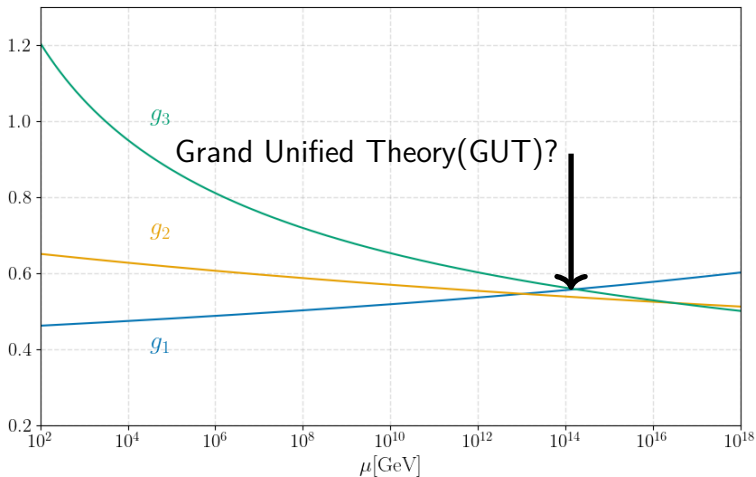
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Work in Progress...

Based on “Fermion Multiplicities at the GUT Scale: A Statistical Study of Unification and Proton Decay”
arXiv:2604.08237 with Akifumi Chitose, Masahiro Ibe and Satoshi Shirai

Gauge Coupling Unification?

Standard Model (SM) : $SU(3)_c \times SU(2)_L \times U(1)_Y$



Minimal SU(5) GUT

- $SU(5) \supset SU(3)_c \times SU(2)_L \times U(1)_Y$
- **Matter unification:** SM fermions embedded into $(\bar{\mathbf{5}} \oplus \mathbf{10}) \times 3$ **chiral representation**.

$$\bar{\mathbf{5}} = \begin{pmatrix} \bar{d} \\ \ell \end{pmatrix}, \quad \mathbf{10} = \begin{pmatrix} \bar{u} & q^T \\ -q & \bar{e} \end{pmatrix}$$

- $SU(5)$ is spontaneously broken to the SM by an adjoint scalar Σ ,

$$\langle \Sigma \rangle = V_{GUT} \times \begin{pmatrix} 2 & & & \\ & 2 & & \\ & & 2 & \\ & & & -3 \\ & & & & -3 \end{pmatrix}$$

Known Issues in Minimal SU(5) GUT

Two classic problems

- Gauge coupling non-unification
- Too short proton lifetime

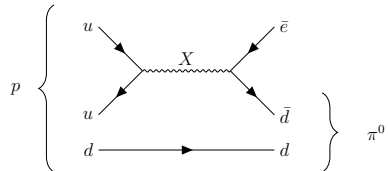
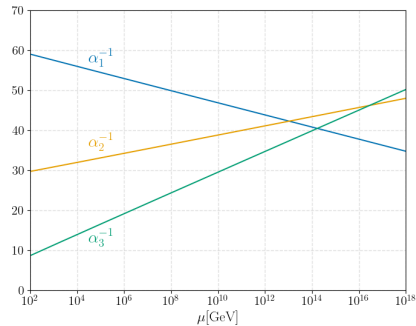
Lifetime estimate:

$$\tau(p \rightarrow \pi^0 + e^+) \simeq 5 \times 10^{26} \left(\frac{M_X/g_5}{10^{14} \text{ GeV}} \right)^4 \text{ year}$$

Current experimental limit (Super-Kamiokande):

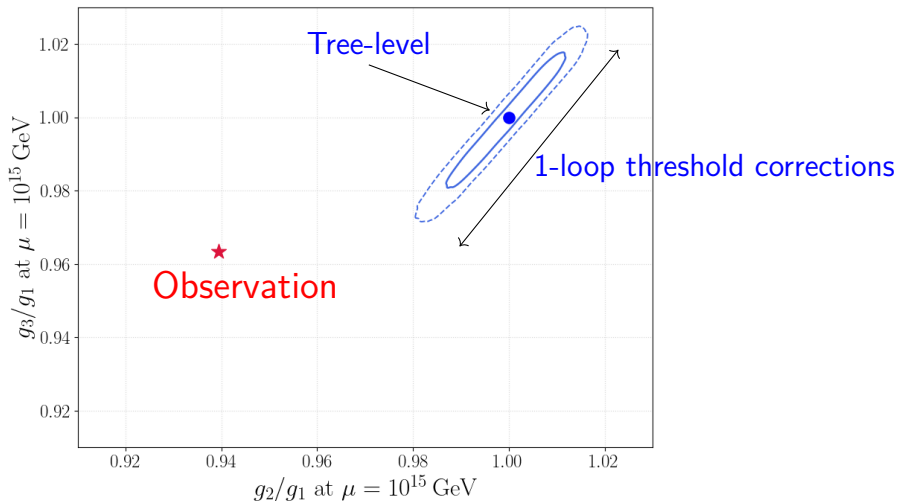
$$\tau(p \rightarrow \pi^0 + e^+) \gtrsim 2.4 \times 10^{34} \text{ year}$$

Ref.: A. Takenaka *et al.* (Super-Kamiokande), Phys. Rev. D **102** (2020) 112011



Gauge Coupling Prediction for Minimal SU(5) GUT

Free Parameters: M_X , M_{H_c} , M_Σ , g_5



Is Minimality Really Well-Motivated?

BSM Physics

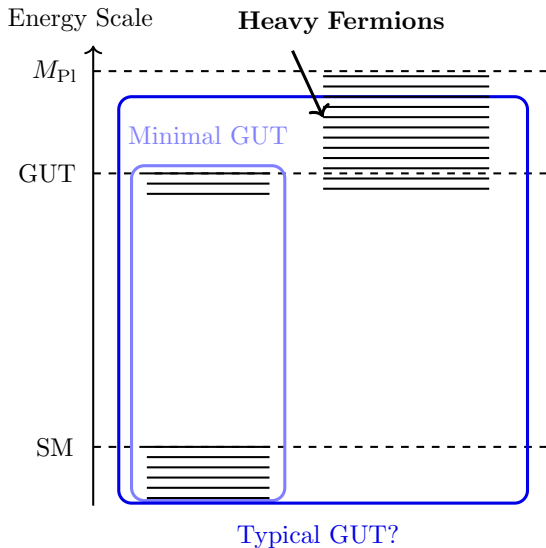
- Quantum Gravity?
- Extra Dimensions?
- Axion?
- Seesaw Mechanism?
- Flavor Models?

GUT may be connected to other BSM physics.

Minimality is not essential.

What Does a Typical GUT Look Like?

- Many additional fermions at the GUT scale?



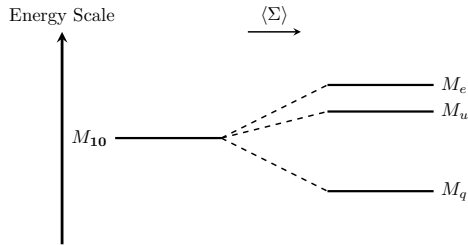
Model Setup

Matter Content (contain three chiral generations)

$$(\bar{\mathbf{5}}, \mathbf{10}) \times 3 + (\mathbf{5}, \bar{\mathbf{5}}) \times n + (\mathbf{10}, \bar{\mathbf{10}}) \times n$$

Vector-like fermions have a generally complex mass matrix:

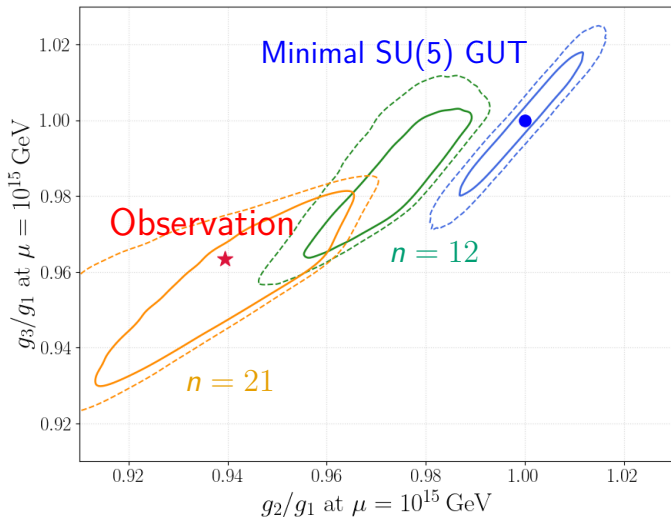
$$\mathcal{L}_M = -\mathbf{10} (M_{10} + Y_{10} \Sigma) \bar{\mathbf{10}} - \bar{\mathbf{5}} (M_5 + Y_5 \Sigma) \mathbf{5} + \text{h.c.} .$$



- After GUT breaking, the masses of GUT multiplets are split.
- Mass splittings affect the threshold corrections.

Result: Gauge Coupling Unification

Case with random $(M_{10,5})_{ij} = \mathcal{O}(1) \times M_0$, $(Y_{10,5})_{ij} = \mathcal{O}(1)$, $M_0 = 0.7V_{\text{GUT}}$



Origin of SM Fermions

Mass Term: $\mathcal{L}_M = -\mathbf{10} (M_{10} + Y_{10} \Sigma) \overline{\mathbf{10}} - \overline{\mathbf{5}} (M_5 + Y_5 \Sigma) \mathbf{5} + \text{h.c.}$

$\text{rank}(M_{10,5} + Y_{10,5} \Sigma) = n \Rightarrow$ Only 3 generations remain massless (SM fermions)

- SM fermions are linear combinations of multiple fermions.

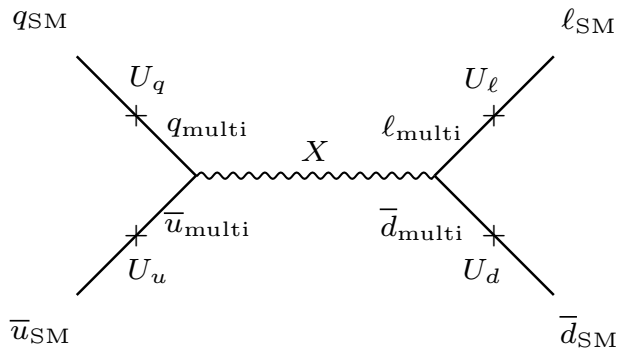
$$q_{\text{SM}} = U_q q_{\text{multi}}, \quad \bar{u}_{\text{SM}} = U_u \bar{u}_{\text{multi}}, \quad \bar{e}_{\text{SM}} = U_e \bar{e}_{\text{multi}}, \dots$$

- After GUT breaking, the unitary matrices differ among SM fermions.

$$U_q \neq U_u \neq U_e, \quad U_d \neq U_\ell$$

Suppression of Proton Decay

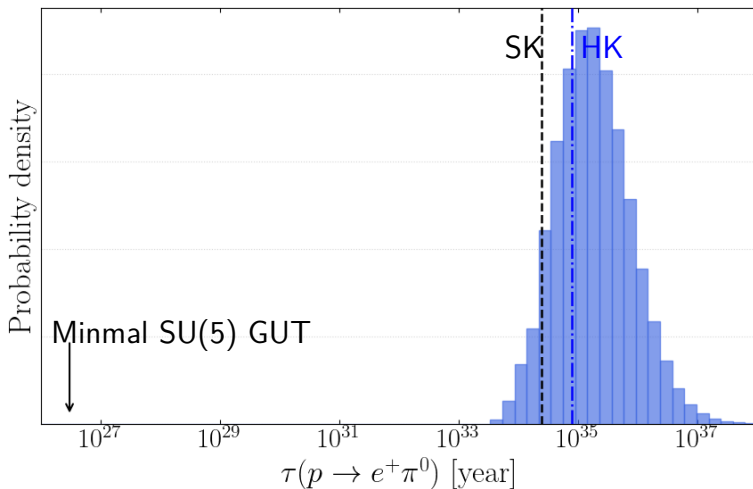
For large n ,



$$\propto (U_q^\dagger U_u) \times (U_\ell^\dagger U_d) \sim \frac{1}{n}$$

For an $n \times n$ random unitary matrix, $U_{ij} = \mathcal{O}\left(\frac{1}{\sqrt{n}}\right)$, $(U^\dagger U')_{ij} = \mathcal{O}\left(\frac{1}{\sqrt{n}}\right)$

Result: Proton Lifetime Predictions



$$\text{SK : } \tau(p \rightarrow \pi^0 + e^+) \gtrsim 2.4 \times 10^{34} \text{ yr}, \quad \text{HK : } \tau(p \rightarrow \pi^0 + e^+) \gtrsim 7.8 \times 10^{34} \text{ yr}$$

Short Summary

- Minimality is not essential.
- Typical GUT can address the problems of minimal $SU(5)$, including gauge unification and proton decay.

How about Yukawa?

- Similar sizes of down-type quark and lepton Yukawa couplings:

$$y_b \sim y_\tau, \quad y_s \sim y_\mu, \quad y_d \sim y_e$$

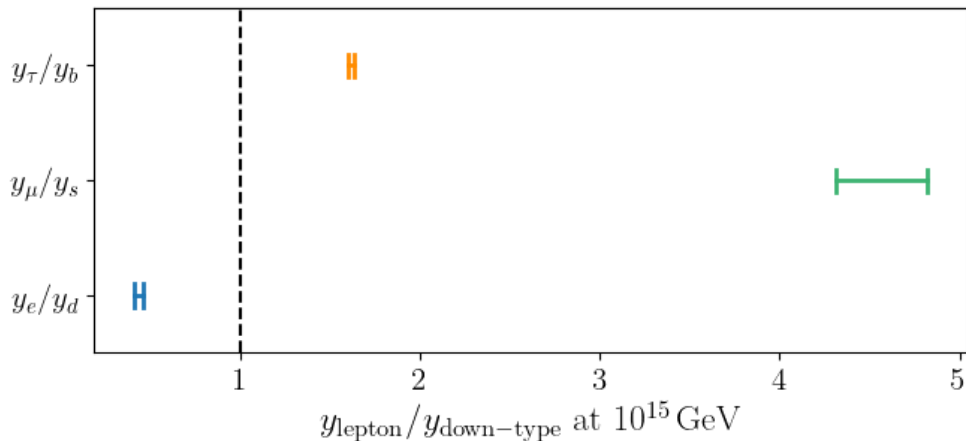
- In usual GUTs, they can originate from a common Yukawa interaction:

$$\mathcal{L}_Y = y_5 H_5^\dagger \mathbf{10} \bar{\mathbf{5}} + \text{h.c.} \quad \Rightarrow \quad y_{\text{down}} = y_{\text{lepton}} = y_5$$



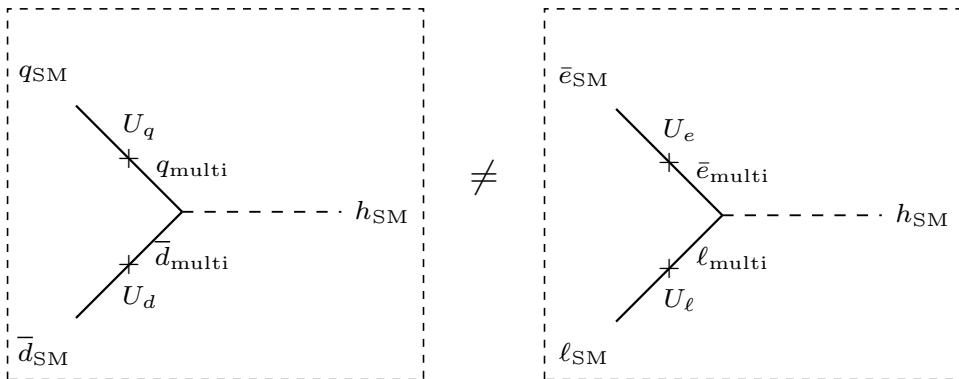
A hint of Grand Unification?

No Yukawa Unification in Minimal SU(5)

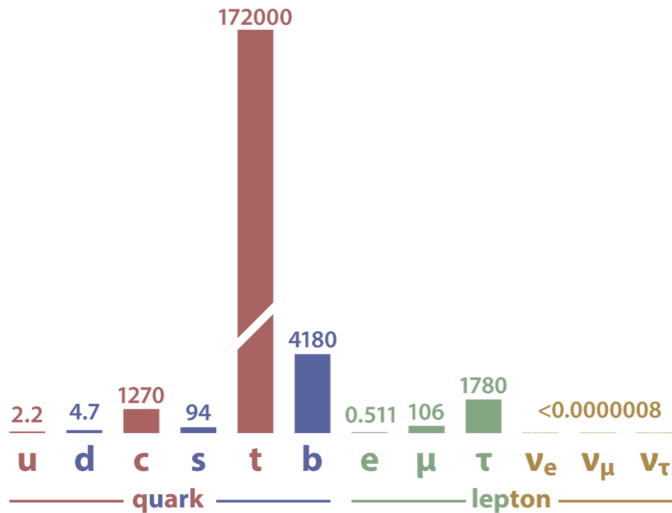


Typical GUT avoid Yukawa Unification

- SM Fermions as Linear Combinations of Multiple fermions
- Down-type quark and lepton Yukawa couplings need not be unified



Yukawa Hierarchy



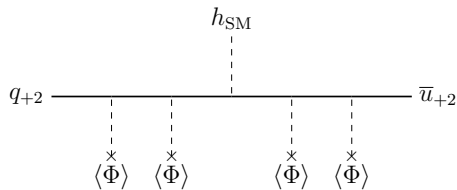
⇒ Flavor Symmetry

Froggatt-Nielsen Mechanism

- Global $U(1)$ flavor symmetry
- SM fermions have generation-dependent $U(1)$ charges (FN charges)
- $U(1)_{\text{FN}}$ broken by a complex scalar field Φ (FN charge = -1)
- Heavy fermions in the UV theory (often omitted...)
- Yukawa unification cannot be solved by FN.

Ref.: C. D. Froggatt and H. B. Nielsen, Nucl. Phys. B **147** (1979) 277

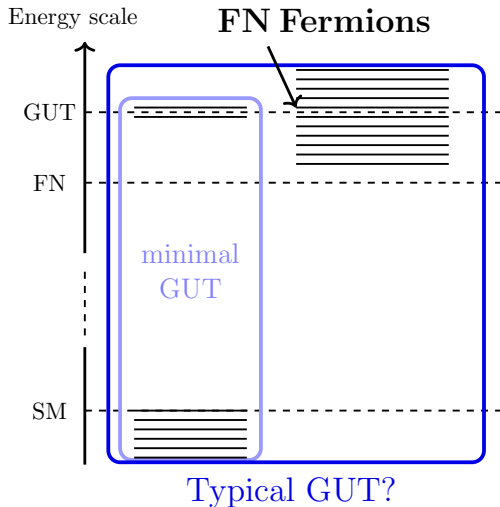
e.g. Yukawa interaction of q_{+2} and $\bar{u}_{+2}$



$$\mathcal{L} = \mathcal{O}(1) \times \left(\frac{\langle \Phi \rangle}{M} \right)^{(2+2)} q_{+2} \bar{u}_{+2} h_{\text{SM}}$$

$$\epsilon \equiv \frac{\langle \Phi \rangle}{M} = \mathcal{O}(0.1)$$

Vector-like as FN fermions



- $\frac{\text{FN scale}}{\text{GUT scale}} \sim 0.1$
- FN fermions at GUT scale

- FN charges are assigned such that, after integrating out the heavy fermions,

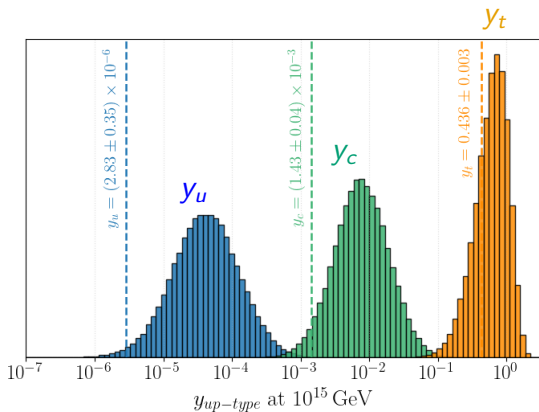
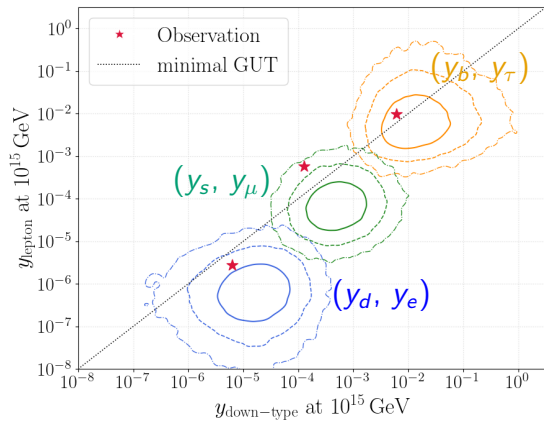
$$Q^{(10)} = (2, 1, 0), \quad Q^{(\bar{5})} = (2, 2, 2).$$

- This gives the hierarchical Yukawa textures

$$y_u \sim \begin{pmatrix} \epsilon^4 & \epsilon^3 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & \epsilon^1 \\ \epsilon^2 & \epsilon^1 & \epsilon^0 \end{pmatrix}, \quad y_{d,\ell} \sim \begin{pmatrix} \epsilon^4 & \epsilon^4 & \epsilon^4 \\ \epsilon^3 & \epsilon^3 & \epsilon^3 \\ \epsilon^2 & \epsilon^2 & \epsilon^2 \end{pmatrix}.$$

$$\epsilon = 0.06 \quad \Leftrightarrow \quad \text{Observation?}$$

Successful Yukawa Unification



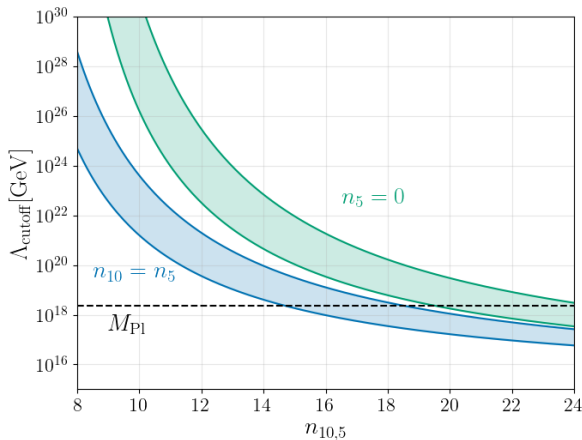
Summary

- Typical GUT model can achieve gauge coupling unification and sufficiently long proton lifetime.
- Yukawa unification (Bottom-tau problem) can be avoided.
- Heavy fermions can be interpreted as FN fermions.
- Typical GUT with the Froggatt-Nielsen mechanism can explain observed Yukawa couplings.

End

Upper Limit on the Number of Heavy Fermions

Good unification $\neq$ unlimited fermions: **Landau poles** set an upper bound.



- Many fermions accelerate the **divergence of the coupling g_5** .
- For $n_5 = n_{10} \lesssim 18$ or $n_5 = 0, n_{10} \lesssim 24$, no landau pole appears.

$$g_5(\Lambda_{\text{cutoff}}) = 4\pi$$
$$g_5(10^{15} \text{ GeV}) = 0.48\text{--}0.55$$