

Preparing Higgs Theory for Precision Measurements at Future Colliders

Mariko Kikuchi (Saga U.)

Collaborators:

Masashi Aiko (Miyakonojo NCT)

Shinya Kanemura (U. of Osaka)

Kodai Sakurai (Tsuruoka NCT)

Sora Taniguchi (U. of Osaka)

Kei Yagyu (Tokyo U. of Science)

JHEP05(2026)136 [2601.15983(hep-ph)]

Talk Plan

- Higgs Physics → New physics
- Future collider experiments
- Indirect test by Higgs precision measurements
Precise calculations for 125 Higgs decay processes
- Our current research results (*Complex Triplet Model*)
- Summary

Unknown things in the Higgs sector

The origin of the Higgs field

The origin of mass and flavour?

Why $\mu^2 < 0$?

Meta-stability of the vacuum

Detail of EW phase transition

Connecting New Physics to the Higgs Sector

Higgs field	I	Y
Φ	1/2	1/2

$$\mathcal{L}_{\text{SM Higgs}} = |D_\mu \Phi|^2$$

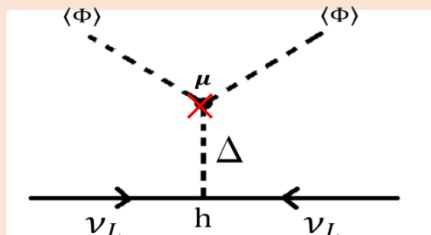
$$- y_f \bar{f}_L \Phi f_R + \text{h.c.}$$

$$- (\mu^2 |\Phi|^2 + \lambda |\Phi|^4)$$

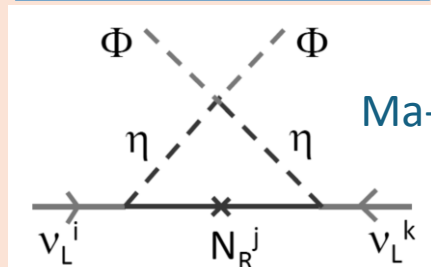
New Physics and Extended Higgs sectors

Neutrino and Higgs

Type II Seesaw



Radiative seesaw



Ma-model

DM and Higgs

Spin-0 WIMP $\rightarrow$ Additional scalar boson

$$A \rightarrow -A$$

✂ Strongly constrained by direct DM searches.
(In Singlet model, $m_{DM} \gtrsim 10\text{TeV}$ 共鳴領域以外)

EW Phase transition

Condition for 1st Order PT :
(sphaleron decoupling condition)

$$\frac{\varphi_C}{T_C} \gtrsim 1$$

Additional bosons loop can make sphaleron decoupling condition.

They are well-motivated !!

$\Phi + S$ (一重項場)

電弱一次相転移, 暗黒物質,
CPの破れ, ...

$\Phi + \Phi'$ (二重項場)

超対称模型, 暗黒物質, 電弱一次相転移,
CPの破れ, ニュートリノ質量(輻射シーソー), ...

$\Phi + \Delta$ (三重項場)

ニュートリノ質量(TypeIIシーソー),
電弱一次相転移, 暗黒物質, ...

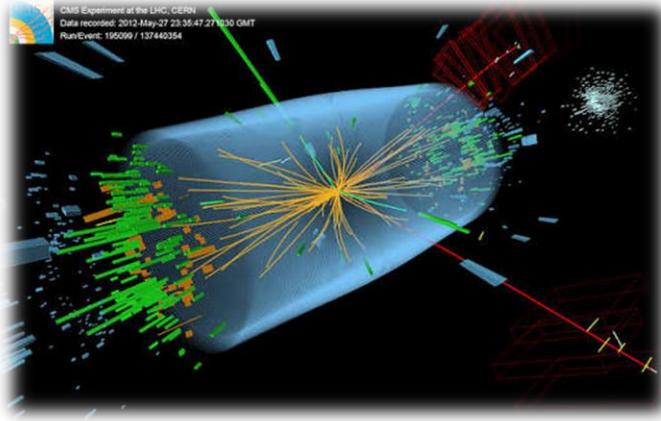


Higgs precision measurements

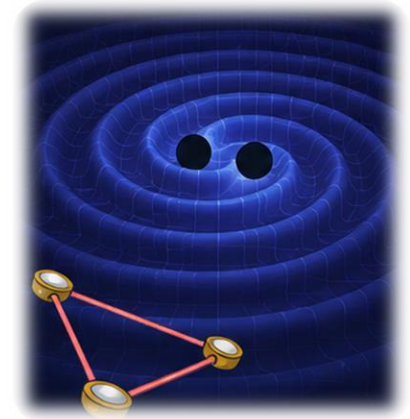
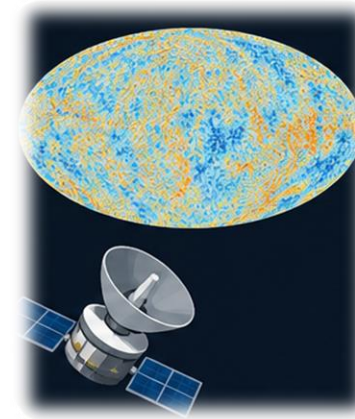
at future collider experiments

How do we approach these problems?

Collider experiments

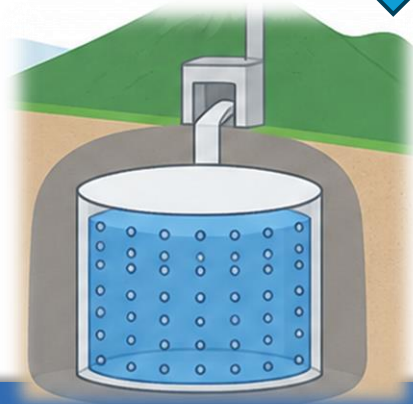


Cosmological survey(CMB, GW, DM, ...)

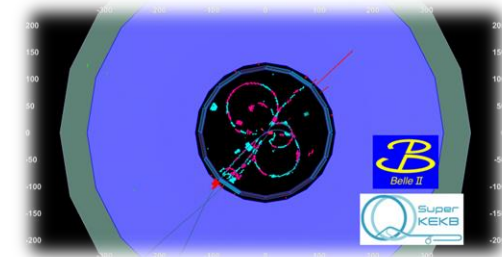
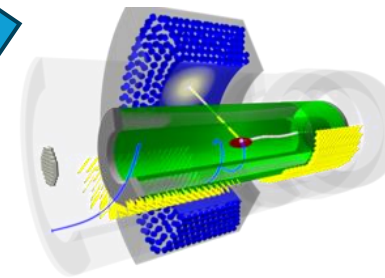


Higgs Physics

DM detection, ν -experiment



Flavor Physics Experiments



Current Status of Future Colliders

Near-term Timeline

- 2022–2026 : LHC Run 3
- 2026–2030 : Long Shutdown 3 (HL-LHC upgrade)
- 2030–2041 : High-Luminosity LHC (HL-LHC)

Candidate Future Colliders(e^+e^-)

- **FCC-ee → FCC-hh** **Preferred option in the 2026 European Strategy**
- CEPC Circular
- ILC Linear
- CLIC Linear

<https://home.cern/the-cern-council-decided-to-update-the-european-strategy-for-particle-physics/>

Expected Higgs coupling precisions

@ FCC-ee

Coupling	HL-LHC	FCC-ee	FCC-ee + FCC-hh
κ_Z (%)	1.3*	0.10	0.10
κ_W (%)	1.5*	0.29	0.25
κ_b (%)	2.5*	0.38 / 0.49	0.33 / 0.45
κ_g (%)	2*	0.49 / 0.54	0.41 / 0.44
κ_τ (%)	1.6*	0.46	0.40
κ_c (%)	—	0.70 / 0.87	0.68 / 0.85
κ_γ (%)	1.6*	1.1	0.30
$\kappa_{Z\gamma}$ (%)	10*	4.3	0.67
κ_t (%)	3.2*	3.1	0.75
κ_μ (%)	4.4*	3.3	0.42
$ \kappa_s $ (%)	—	$^{+29}_{-67}$	$^{+29}_{-67}$
Γ_H (%)	—	0.78	0.69
$\mathcal{B}_{\text{inv}} (<, 95\% \text{ CL})$	$1.9 \times 10^{-2} *$	5×10^{-4}	2.3×10^{-4}
$\mathcal{B}_{\text{unt}} (<, 95\% \text{ CL})$	$4 \times 10^{-2} *$	6.8×10^{-3}	6.7×10^{-3}

2504.00672 [hep-ex]

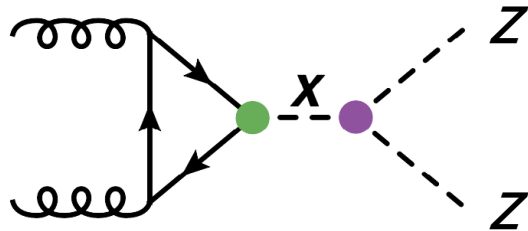
@ LCF(ILC)

	LCF, L[ab ⁻¹]	250 GeV 3, pol.	250 + 550 GeV 3+4, pol.	250 + 550 GeV 3+8, pol.
Higgs couplings				
$g(Hb\bar{b})$		0.72	0.46	0.38
$g(Hc\bar{c})$		1.45	1.04	0.87
$g(Hgg)$		1.31	0.86	0.70
$g(H\tau\tau)$		0.83	0.61	0.53
$g(HWW)$		0.34	0.24	0.22
$g(HZZ)$		0.34	0.24	0.22
$g(H\gamma\gamma)$		1.02	0.98	0.94
$g(H\gamma Z)$		7.51	5.75	5.41
$g(H\mu\mu)$		3.87	3.68	3.53
Higgs widths				
total width		1.39	0.97	0.85
width to invisible (95%CL)		0.36	0.33	0.31
width to unclassified (95%CL)		1.53	1.23	1.11

2503.19983 [hep-ex]

How to search at colliders

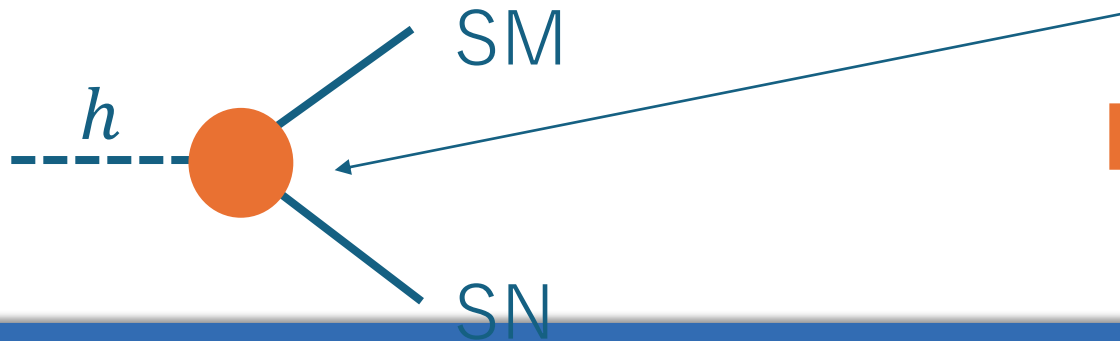
Direct searches of heavy Higgs



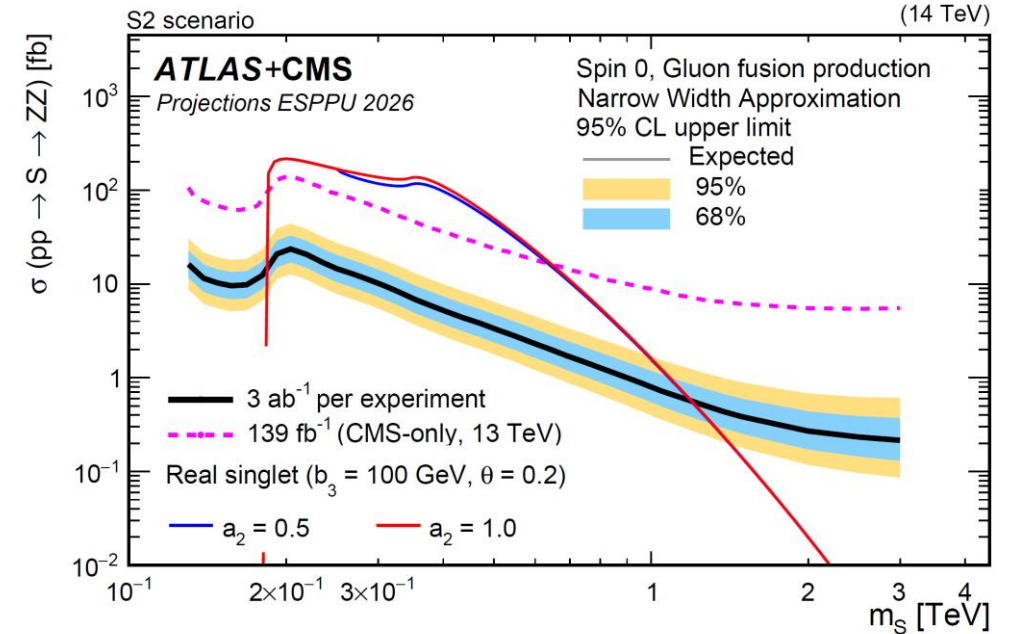
Indirect test by h -precision

In SM : couplings $\propto$ masses

Deviations from the SM predictions = New physics effects



Precision measurements of h
Higgs factory !!



2504.00672 [hep-ex]

Other potential of e⁺e⁻ colliders

① Z-pole

m_Z, Flavor Phys. (Z→bb)

② WW threshold (m_W)

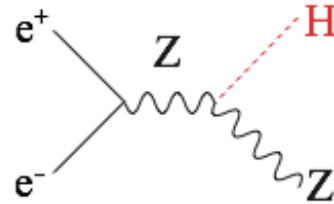
m_W, Γ_W, EW parameters

③ 240 GeV (HZ)

hZZ, Higgs couplings,
Γ_h

④ 365 GeV (tt)

Top mass

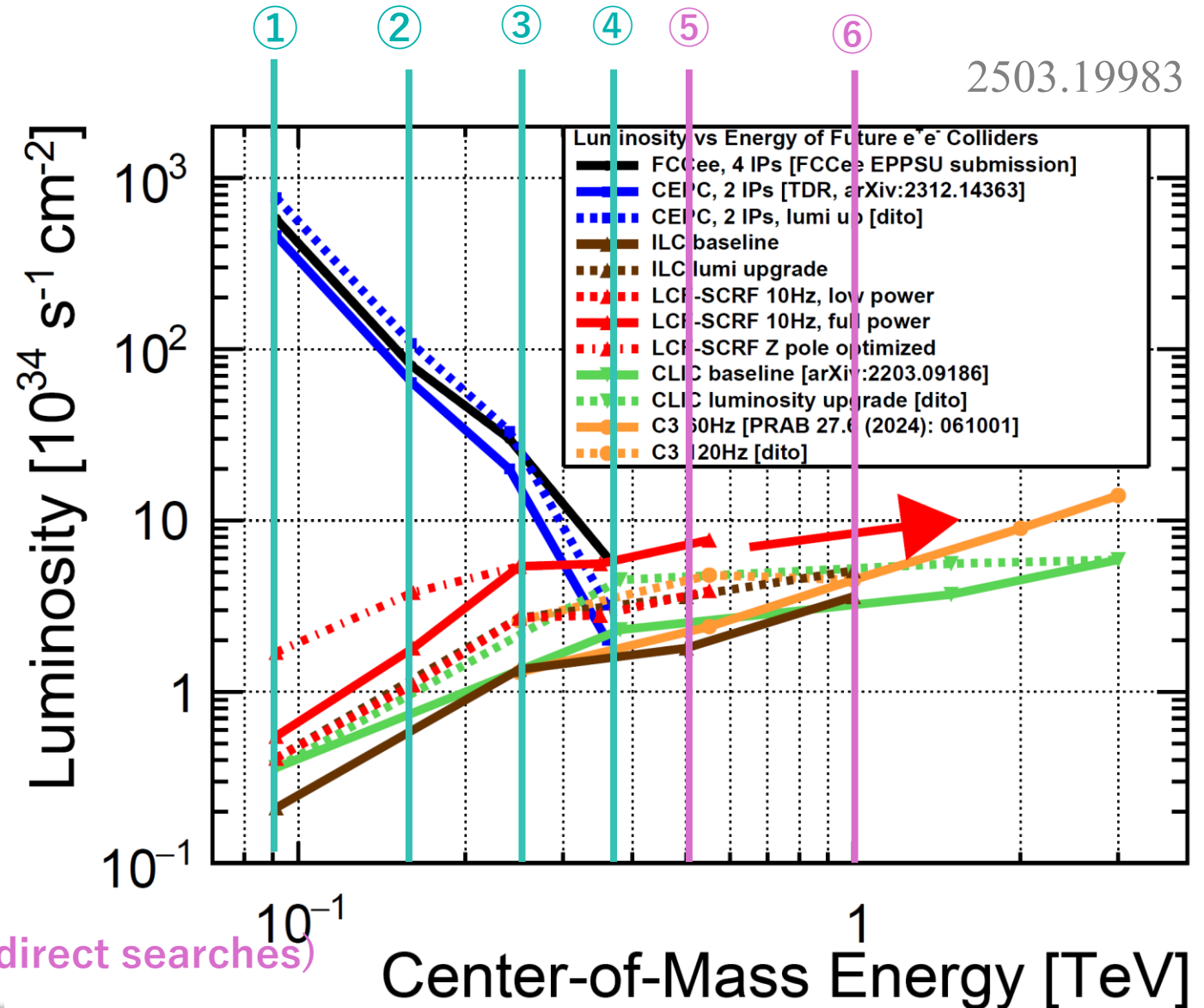


⑤ 500 GeV – 500 GeV

(Top-Yukawa, Double Higgs pro. , hhh)

⑥ 1TeV

(More precise measure hhh, Heavy Higgs direct searches)

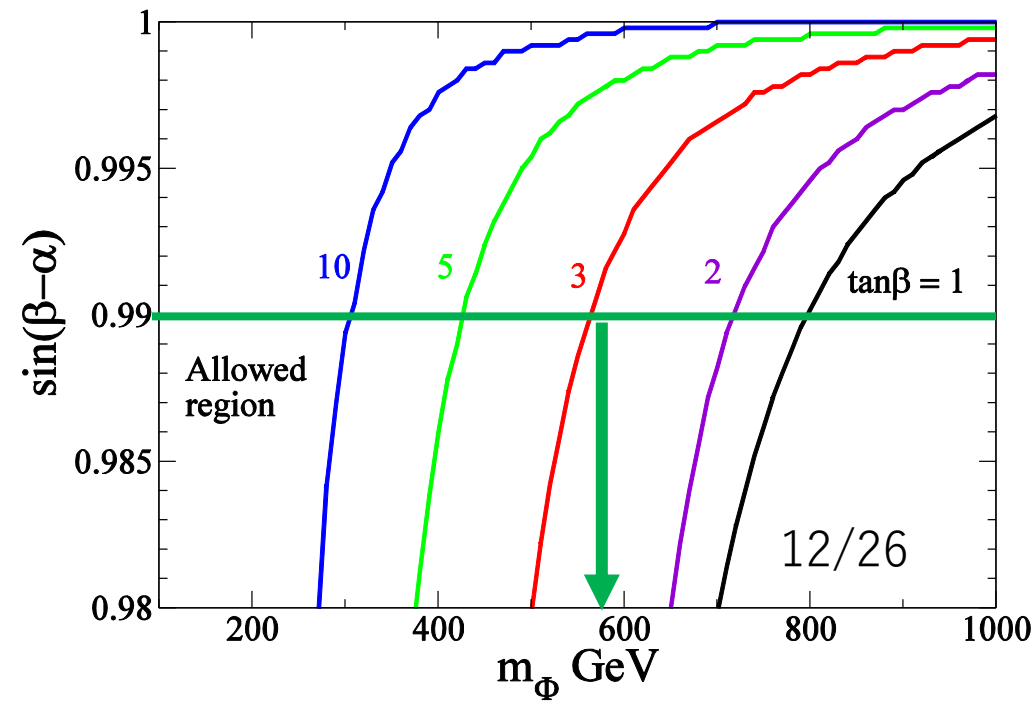




Precise calculations for 125 Higgs decay processes

What Do Deviations Tell Us? (1) New physics scale

Constraints from perturbative unitarity



例えば

$$\Delta\kappa_Z = -1\%$$

Deviations in Higgs couplings give upper bounds on Additional Higgs mass.

Characteristic deviation pattern

$\kappa_X = \frac{c_{hXX}}{c_{hXX}^{SM}}$	κ_Z	κ_d	κ_e
HSM	$\cos\alpha$	$\cos\alpha$	$\cos\alpha$
Type-I 2HDM	$\sin(\beta - \alpha) \left[= \sin\beta \cos(-\alpha) + \cos\beta \sin(-\alpha) \right]$	$\sin(\beta - \alpha) + \cot\beta \cos(\beta - \alpha)$	$\sin(\beta - \alpha) + \cot\beta \cos(\beta - \alpha)$
Type-II 2HDM	$\sin(\beta - \alpha)$	$\sin(\beta - \alpha) - \tan\beta \cos(\beta - \alpha)$	$\sin(\beta - \alpha) - \tan\beta \cos(\beta - \alpha)$
Type-X 2HDM	$\sin(\beta - \alpha)$	$\sin(\beta - \alpha) + \cot\beta \cos(\beta - \alpha)$	$\sin(\beta - \alpha) - \tan\beta \cos(\beta - \alpha)$
Triplet (Y=1)	$\cos\beta' \cos\alpha + 2\sin\beta' \sin\alpha$	$\cos\alpha / \cos\beta$	$\cos\alpha / \cos\beta$
	$\Delta\kappa_Z = \kappa_Z - 1 \ll 1$	$\Delta\kappa_d$	$\Delta\kappa_e$
HSM	$-x^2/2$	$-x^2/2$	$-x^2/2$
Type-I 2HDM	$-x^2/2$	$+\cot\beta x$	$+\cot\beta x$
Type-II 2HDM	$-x^2/2$	$-\tan\beta x$	$-\tan\beta x$
Type-X 2HDM	$-x^2/2$	$+\cot\beta x$	$-\tan\beta x$
Triplet (Y=1)	$-(\alpha^2 + \beta^2 - 2\sqrt{2}\alpha\beta)/2$	$-(\alpha^2 - \beta^2)/2$	$-(\alpha^2 - \beta^2)/2$

If $\Delta\kappa_Z = -1\%$

-1%

+O(10)%

-O(10)%

+O(10)%

-1%

+O(10)%

-O(10)%

-O(10)%

Decoupling & non-decoupling

スカラー粒子のループ効果 も デカップリング 効果になり得る

$$m_{\Phi}^2 = \lambda_i v^2 + M^2$$

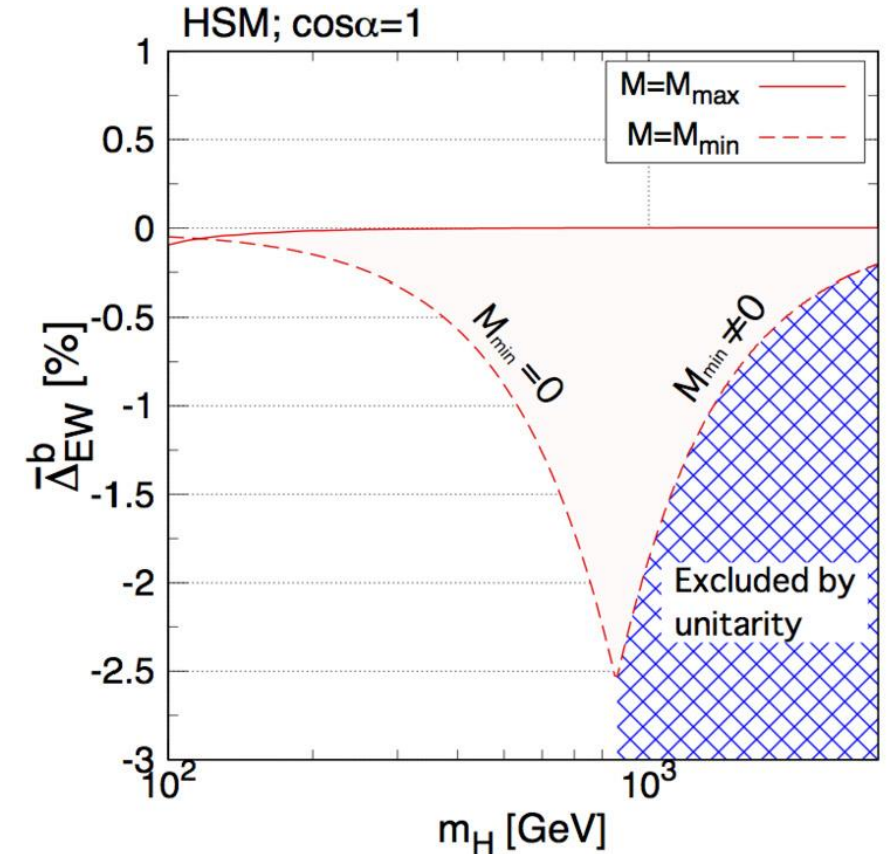
付加的ヒッグス場 ϕ の質量
は2種類のソース

- $M^2 \gg \lambda_i v^2$ デカップリングする
- $M^2 \sim \lambda_i v^2$ **ノンデカップリング効果**

例 $\Gamma(h \rightarrow f \bar{f}) = \Gamma_0(h \rightarrow f \bar{f}) \left[1 + \Delta_{EW}^f + \Delta_{QCD}^f \right]$

$$\bar{\Delta}_{EW}^X \simeq -\frac{1}{16\pi^2} \frac{1}{6} \sum_{\Phi} c_{\Phi} \frac{m_{\Phi}^2}{v^2} \left(1 - \frac{M^2}{m_{\Phi}^2} \right)^2$$

$$m_{\Phi}^2 \left(1 - \frac{M^2}{m_{\Phi}^2} \right)^2 \begin{cases} \propto \frac{1}{m_{\Phi}^2} & (M^2 \gg v^2) \quad \text{デカップリング} \\ \propto m_{\Phi}^2 & (M^2 \sim v^2) \quad \text{ノンデカップリング} \end{cases}$$



Kanemura, KM, Mawatari,
Sakurai, Yagy (2018)

Fortran program for **decay BRs** of **Higgs bosons** with **loop corrections** in **extended Higgs models**.

	Real Singlet	Softly broken Z2 2HDM (Type I)	IDM	Complex Triplet	SM
h_{125} -decays Couplings, BR, Γ	✓	✓	✓	New !!	✓
Heavy Higgs decays BR, Γ	✓	✓	✓		✓

- **Loop corrections** : EW and scalar-NLO
 - UV div. → On-shell renormalizations
 - Gauge dependence → Remove by Pinch technique
 - IR div. via photon loop → Cancel by real photon emission

QCD (NNLO, NLO) corrections ($\overline{\text{MS}}$ -bar scheme).

- **Constraints** : Vacuum stability(tree level, RGE improved), Tree-level unitarity, Triviality, True vacuum, ST parameters

A. Djouadi, Phys. Rept., 457, 1 (2008),
 B. M. Spira, Prog. Part. Nucl. Phys., 95, 98 (2017),
 K. G. Chetyrkin, A. Kwiatkowski, Nucl. Phys., B461, 3 (1996)



H decays with loop corrections in Complex Triplet Model

Complex Triplet Higgs model

	SU(2) _L	U(1) _Y
Φ	2	1/2
Δ	3	1

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$$

$$\Delta = \begin{pmatrix} \frac{\Delta^+}{\sqrt{2}} & \Delta^{++} \\ \Delta^0 & \frac{\Delta^-}{\sqrt{2}} \end{pmatrix}$$

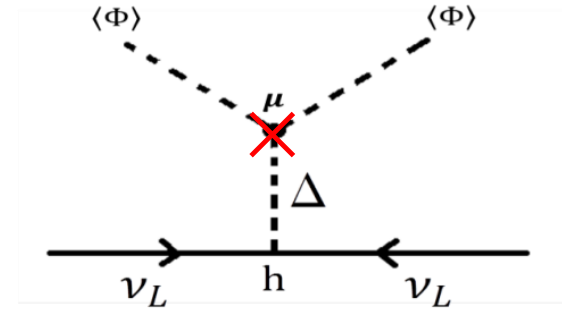
$$\phi^0 = \frac{1}{\sqrt{2}}(\phi + v_\phi + i\chi)$$

$$\Delta^0 = \frac{1}{\sqrt{2}}(\delta + v_\Delta + i\eta)$$

- Neutrino mass (Type-II seesaw)

$$\mathcal{L}_\nu = h_{ij} \overline{L}_L^i c i\tau_2 \Delta L_L^j + hc$$

$$M_\nu^{ij} = \frac{h^{ij} \mu v_\Phi^2}{M_\Delta^2}$$



Cheng, Li (1980),
Mohapatra,
Senjanovic(1981)

- Mass eigenstates

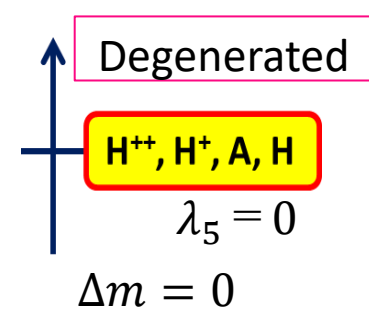
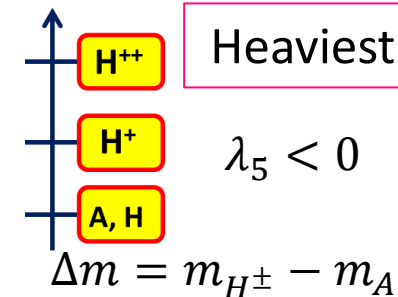
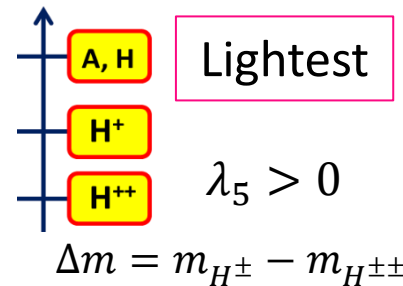
h

SM-like Higgs boson

$H^{\pm\pm}, H^\pm, A, H$

triplet-like Higgs bosons

$$\Delta m = m_{H^\pm} - m_{\text{lightest}}$$



- EW rho parameter

$$\rho_{\text{Tree}} = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = \frac{v_\phi^2 + 2v_\Delta^2}{v_\phi^2 + 4v_\Delta^2} \neq 1$$

Constraints from ρ_{exp} $v_\Delta^2 \ll v_\phi^2$

$\rho = 1.00031 \pm 0.00019$
(PDG2024)

Coupling deviations

- hVV couplings can be larger than SM values.

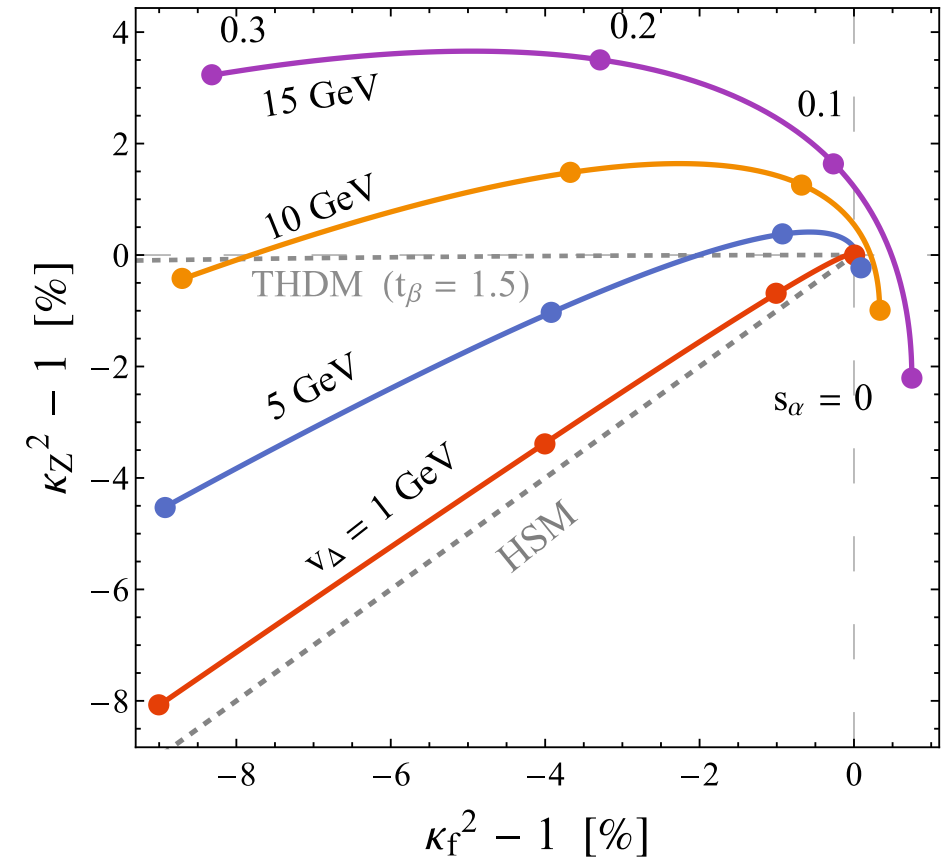
$$K_{hZZ} = \cos\beta' \cos\alpha + 2\sin\beta' \sin\alpha$$

$$K_{hWW} = \cos\beta \cos\alpha + \sqrt{2}\sin\beta \sin\alpha$$

$$\left(\tan\beta = \frac{\sqrt{2}v_\Delta}{v_\phi}, \tan\beta' = \frac{2v_\Delta}{v_\phi} \right)$$

- A larger value of v_Δ leads to larger deviations in hVV couplings.

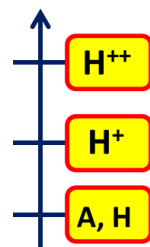
$$K_{hZZ} \simeq \cos\alpha + \frac{4v_\Delta}{v} \sin\alpha$$



Does the characteristic deviation pattern appear even with radiative correction?

$$K_{hff} = \cos\alpha / \cos\beta$$

Constrains



Heaviest

$$\lambda_5 < 0$$

$$\Delta m = m_{H^\pm} - m_A$$

◆ EW precision observables

$$(m_W^{\text{rem}})^2 = \frac{m_Z^2(1 + c_{\beta'}^2)}{4} \left[1 + \sqrt{1 - \frac{8\pi\alpha_{\text{em}}}{\sqrt{2}G_F m_Z^2(1 + c_{\beta'}^2)(1 - \Delta r)}} \right]$$

$$\Delta r \supset -\frac{1}{16\pi^2} \ln(m_{H^\pm} - m_A)^2 \quad \Rightarrow \quad \text{Large } \Delta m \text{ is required.}$$

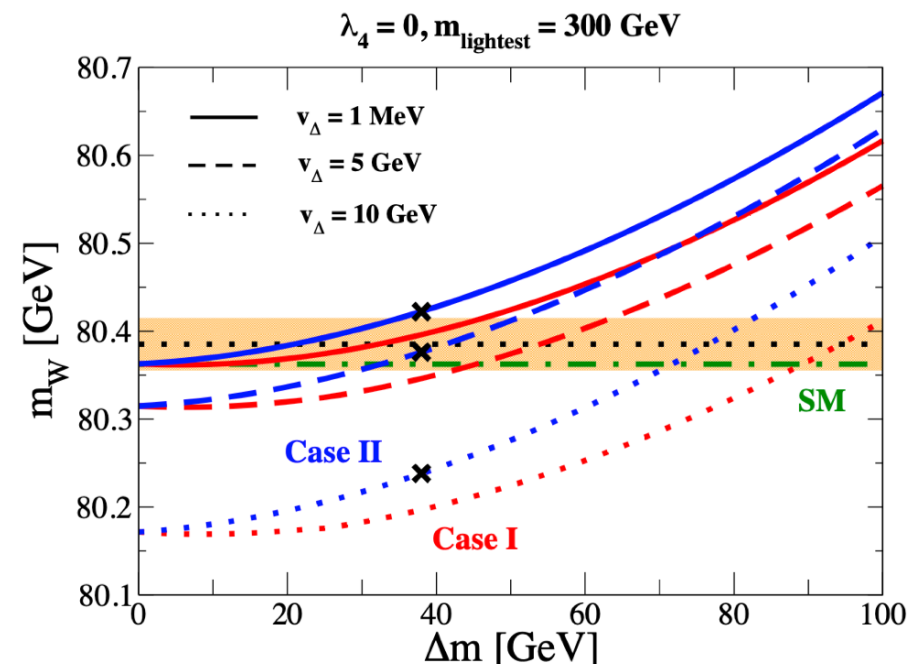
◆ $h \rightarrow \gamma\gamma$

$$\Gamma(h \rightarrow \gamma\gamma)_{\text{HTM}} = \left| \text{t-loop} + \text{w}^+ \text{ loop} + \text{H}^{++}, \text{H}^+ \text{ loop} + \text{H}^{++}, \text{H}^+ \text{ triangle} \right|^2$$

$\Rightarrow$ Larger $m_{H^{++}}$ is preferred.

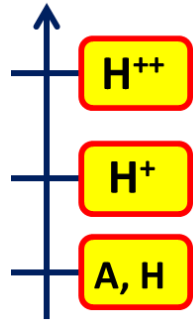
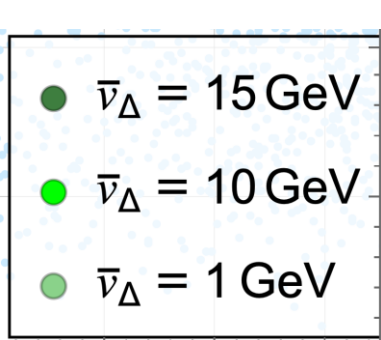
◆ Perturbative unitarity

$\Rightarrow$ Upper bound for $m_{H^{++}}$



Aoki, Kanemura, MK, Yagyu (2012)

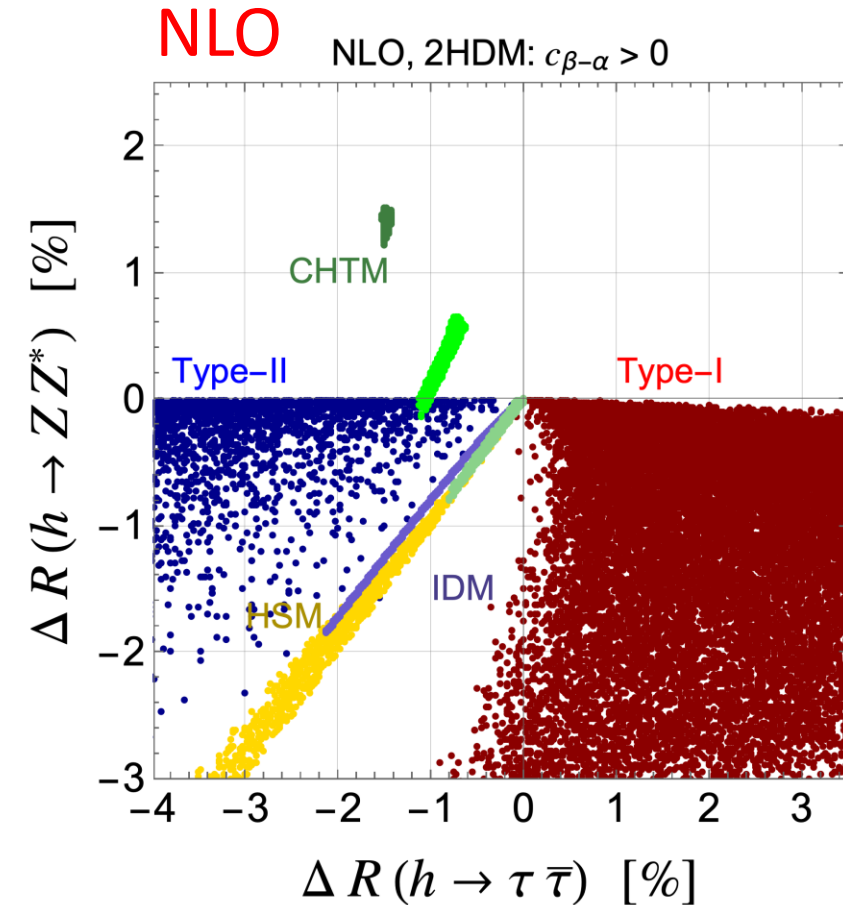
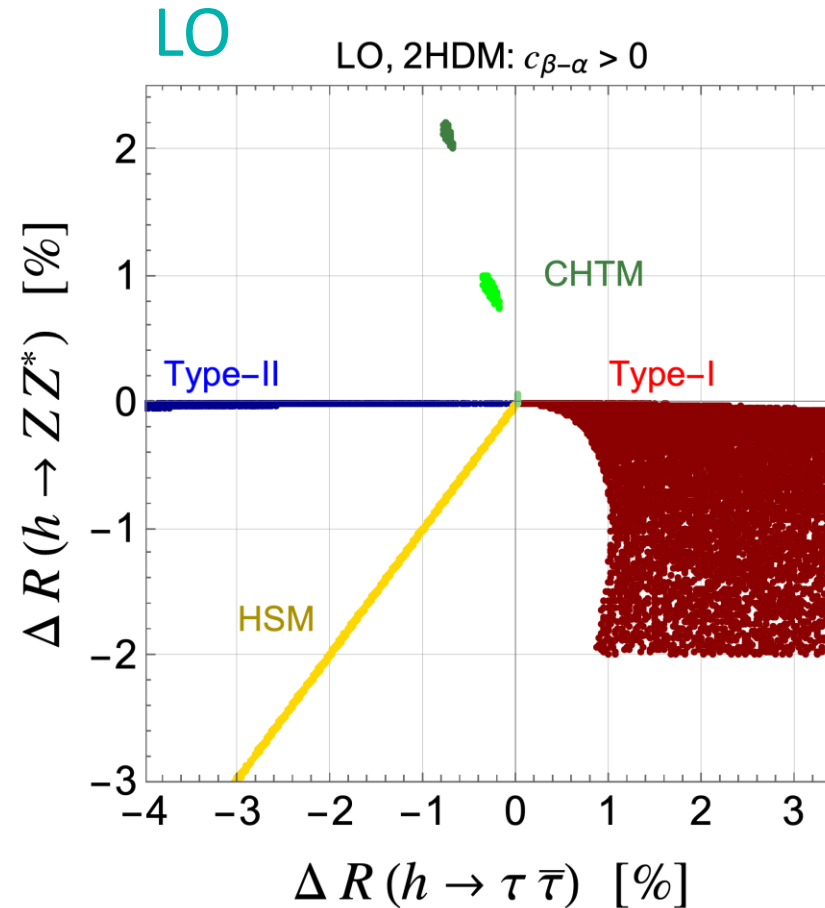
Decays of $h \rightarrow ZZ^*$ and $h \rightarrow \tau\tau$ with NLO corr.



$$300 \text{ GeV} \leq m_A \leq 1 \text{ TeV},$$

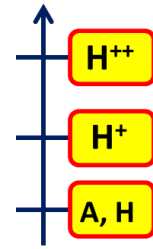
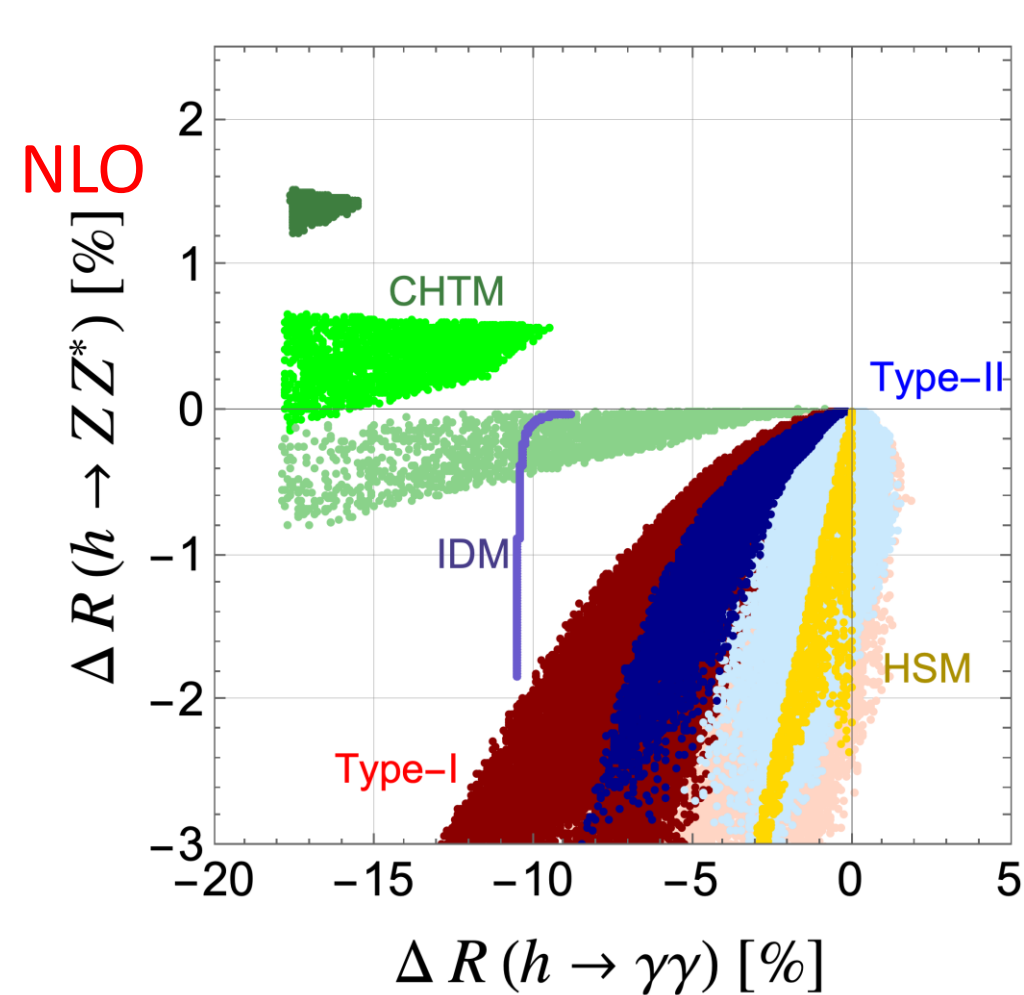
$$0 \leq \Delta m \leq 150 \text{ GeV}$$

$$0 \leq \lambda_4 \leq 4\pi$$



At NLO level, $\Delta R(h \rightarrow ZZ^*) \gtrsim 1\%$, by which CHTM is separable from other models.

Decays of $h \rightarrow ZZ^*$ and $h \rightarrow \gamma\gamma$



$h \rightarrow \gamma\gamma$ can be significantly deviated from SM.

$\Delta R(h \rightarrow \gamma\gamma) \simeq -20\%$ in $m_{H^{++}} > m_{H^+} > m_H$

$\Delta R(h \rightarrow \gamma\gamma) \simeq +10\%$ in $m_H > m_{H^+} > m_{H^{++}}$

Current bound : $|\Delta\mu_{\gamma\gamma}| < 0.18$ at 2σ ,

Future measurements :

$\Gamma(h \rightarrow \gamma\gamma) \sim 2.6\%$ @ HL-LHC [2504.00672]

$\Gamma(h \rightarrow ZZ^*) \sim 0.5\%$ @ ILC250+550 [2503.19983]

The characteristic deviation patterns are expected to be detected at HL LHC or future e+e- Collider

Summary

- Future e+e- colliders will provide not only Higgs precision measurements, but also EW and flavour precision.
- Precise theoretical calculations are essential to the future precision measurements.
- H-COUP provides loop-corrected Higgs decay predictions in a variety of extended Higgs models.
- Comparing characteristic patterns of Higgs coupling deviations enables us to discriminate among different Higgs sectors.
- Current our result → Complex Triple model

$$\Delta R(h \rightarrow VV^*) \sim +1\%, \Delta R(h \rightarrow \gamma\gamma) \sim -20\% \text{ [H++ Heaviest scenario]}$$
$$\Delta R(h \rightarrow \gamma\gamma) \sim +10\% \text{ [H++ Lightest scenario]}$$

Indirect searches can probe heavy Higgs bosons beyond the direct-search reach !!

